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# The Signature Method for DAEs arising in the Modeling of Electrical Circuits

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## Abstract

We consider the Signature Method ( $\Sigma$ -method) for the structural analysis of differential-algebraic equations (DAEs) that arise in the modeling and simulation of electrical circuits. Different formulations of the set of model equations are considered. We show that for some formulations the structural approach may fail for certain circuit topologies, while other formulations are better suited for a structural analysis. In particular, we show that for the branch-oriented model equations the Signature Method always succeeds with a structural index that corresponds to the differentiation index of the system. The results are illustrated by a number of examples.

*Keywords:* Differential-algebraic equation, structural analysis, signature method, modified nodal analysis, MNA, modified loop analysis, MLA, branch-oriented model.

*2010 MSC:* 34A09, 65L80, 94C05

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## 1. Introduction

Modeling and simulation of dynamical systems is an important issue in the development of technical innovations. More and more equation-based object-oriented modeling environments such as Dymola, MapleSim or 20sim are used as tools for automatized modeling and simulation of multi-physical systems. An important problem class are electrical circuits or electrical components that are

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embedded in multi-physical applications. Circuit equations typically lead to large-scale systems of differential-algebraic equations (DAEs). For such DAE systems it is well-known that, due to the occurrence of hidden constraints, an *index reduction* or *regularization* is required for a robust numerical integration, e.g., to avoid drift-off, instabilities or artificial oscillations, see [1, 2, 3, 4]. In most modeling environments a structural analysis of the model equations is used to perform a regularization, i.e., to determine a formulation that is suited for the numerical integration. Here, usually *Pantelides Algorithm* [5] or the *Signature Method* ( $\Sigma$ -method) [6] are used in combination with the dummy derivative approach [7]. These structural methods are powerful tools since they are computationally very efficient. In [6] it has been shown that the structural analysis based on the  $\Sigma$ -method works reliable for many classes of DAEs (including mechanical multibody systems or systems in Hessenberg form). However, it is also known that structural approaches may fail for certain problems [8, 9]. In this paper, we will consider the structural analysis for commonly used circuit equations. We will present a number of examples that show that the structural approach may fail for certain circuit topologies or certain formulations of the model equations. Fortunately, an appropriate formulation of the circuit equations using branch-oriented model equations allows the secure application of the  $\Sigma$ -method.

The paper is organized as follows. In Section 2 we collect some preliminaries. Then, in Section 3, we consider different formulations of the model equations for electrical circuits. Next, in Section 4, we recapitulate the basic ideas of the  $\Sigma$ -method and apply the approach to the different circuit equations. We will see that for some formulations the  $\Sigma$ -method may fail depending on the topology of the circuit, while other formulations are better suited for the structural analysis. We end with some concluding remarks in Section 5. Some additional graph theoretical results are given in Appendix A.

## 2. Preliminaries

In general, a DAE is given as a nonlinear system

$$F(t, x, \dot{x}) = 0, \quad (1)$$

where  $F : \mathbb{I} \times \mathbb{D}_x \times \mathbb{D}_{\dot{x}} \rightarrow \mathbb{R}^n$  is a continuous function,  $\mathbb{D}_x, \mathbb{D}_{\dot{x}} \subset \mathbb{R}^n$  are open,  $\mathbb{I} = [t_0, t_f] \subset \mathbb{R}$  and  $x : \mathbb{I} \rightarrow \mathbb{D}_x$  is a continuously differentiable unknown function. For a differentiable time depending function  $x$ , the  $i$ -th derivative of  $x$  with respect to  $t$  is denoted by  $x^{(i)}(t) = d^i x(t)/dt^i$  for  $i \in \mathbb{N}$ , using the convention  $x^{(1)}(t) = \dot{x}(t)$ , and  $x^{(2)}(t) = \ddot{x}(t)$ .

For a matrix  $A \in \mathbb{R}^{m,n}$ ,  $\text{im } A$  denotes the image of  $A$ ,  $\ker A$  denotes the kernel of  $A$ , and  $\text{rank } A$  denotes the rank of  $A$ . Furthermore, a square matrix  $A$  that is positive definite or positive semi-definite is denoted by  $A > 0$  or  $A \geq 0$ , respectively.

**Definition 1.** A function  $x : \mathbb{I} \rightarrow \mathbb{D}_x$  is said to be a solution of the DAE (1) if it is continuously differentiable for all  $t \in \mathbb{I}$  and (1) is fulfilled pointwise for all  $t \in \mathbb{I}$ . A function  $x : \mathbb{I} \rightarrow \mathbb{D}_x$  is called a solution of the initial value problem (1) and  $x(t_0) = x_0$  with  $x_0 \in \mathbb{D}_x$  if it is a solution of (1) and satisfies the initial condition  $x(t_0) = x_0$ . An initial value  $x_0 \in \mathbb{D}_x$  is called consistent, if the initial value problem (1) and  $x(t_0) = x_0$  has a solution.

The original equation (1) and its derivatives up to order  $\ell$  can be gathered into a so-called *derivative array*

$$\mathcal{F}_\ell(t, x, \dot{x}, w) = \begin{bmatrix} F(t, x, \dot{x}) \\ \frac{d}{dt} F(t, x, \dot{x}) \\ \vdots \\ (\frac{d}{dt})^\ell F(t, x, \dot{x}) \end{bmatrix}, \quad (2)$$

where

$$w = [\ddot{x}, \dots, x^{(\ell+1)}].$$

**Definition 2** ([10]). Suppose that (1) is solvable. If  $v$  is uniquely determined by  $(t, x)$  and  $\mathcal{F}_\ell(t, x, v, w) = 0$  for all consistent values and  $\nu_d$  is the least such

integer  $\ell$  that this holds for, we call  $v_d$  the differentiation index (d-index) of the DAE (1).

If the d-index is well-defined, one can extract a so-called *underlying ODE*  $\dot{x}(t) = \Phi(x(t), t)$  from the derivative array with the property that every solution of the DAE (1) also solves the underlying ODE.

In the following, we will also need some basic results of graph theory. Let  $G = (V, B, \Psi)$  denote a *directed graph* with  $V = \{v_1, v_2, \dots\}$  the finite set of nodes (or vertices) and  $B = \{b_1, b_2, \dots\}$  the finite set of branches (or edges).  $\Psi : B \rightarrow V \times V$  denotes the *incidence map* that maps every branch  $b \in B$  onto some ordered pair of nodes  $(v_1, v_2)$ , i.e.,  $\Psi(b) = (v_1, v_2)$  for some  $v_1, v_2 \in V$ , and we say that  $b$  is *directed from*  $v_1$  *to*  $v_2$ . Here,  $v_1$  is called the *initial node* and  $v_2$  the *terminal node* of  $b$ . For every branch  $b \in B$ , we define an additional branch  $-b$  being directed from the terminal to the initial node of  $b$ , that is  $\Psi(-b) = (v_2, v_1)$ , and the set  $\tilde{B} = \{b, -b \mid b \in B\}$ . A tuple  $(b_1, \dots, b_r)$  of branches  $b_i \in \tilde{B}$ ,  $i = 1, \dots, r$ , and  $\Psi(b_i) = (v_{k_i}, v_{k_{i+1}})$  for all  $i = 1, \dots, r-1$  is called *path from*  $v_{k_1}$  *to*  $v_{k_r}$ . A path is called *simple*, if  $v_{k_1}, \dots, v_{k_{r-1}}$  are distinct. A *loop* is a simple path such that  $v_{k_1} = v_{k_r}$ . A directed graph  $G$  is *connected* if for every pair of nodes there exists a path between them. A *subgraph*  $G' := (V', B', \Psi|_{B'})$  of a connected graph  $G$  is a graph such that  $V' \subset V$ ,  $B' \subset B|_{V'} := \{b \in B \mid \Psi(b) \in V' \times V'\}$ .

A *cutset* is a set  $B_c$  of branches of a connected graph  $G$  such that the graph  $G_c$  that results when the branches in  $B_c$  are deleted from  $G$  is disconnected, and adding any branch in  $B_c$  to  $G_c$  would result again in a connected graph.

A (*spanning*) *tree* in a connected graph is a connected subgraph which contains all nodes and has no loops. We will also use the term *tree* to refer to the set of branches contained in this subgraph. Once a tree has been chosen the branches in the tree are called *twigs*, whereas the remaining ones are called *links*. The set of links defines the *cotree*. Let  $n_\eta$  be the number of nodes and  $n_b$  denote the number of branches in the connected graph. Then, any tree defines  $n_\eta - 1$  twigs and  $n_b - n_\eta + 1$  links. For more details we refer to [11, 12] and to Appendix A.

### 85 3. Model Equations for Electrical Circuits

We consider lumped electrical circuits containing (possibly nonlinear) resistors, capacitors, and inductors, as well as voltage sources and current sources. The modeling of the dynamical behavior of such electrical circuits is based on Kirchhoff's laws together with the constitutive relations for the electrical components.  
 90 However, there are different ways to set up the set of model equations. We will present four different approaches to model electrical circuits in this section.

#### 3.1. The Modified Nodal Analysis

A common way for the modeling of electrical circuits is the *Modified Nodal Analysis (MNA)* [13]. The circuit is modeled as a directed graph whose branches correspond to the circuit elements and whose nodes correspond to the interconnections of these elements. The topological structure of such a graph with  $n_\eta$  nodes and  $n_b$  branches can be described by an incidence matrix  $A_0 \in \mathbb{R}^{n_\eta, n_b}$  with entries

$$a_{ij} = \begin{cases} 1 & \text{if branch } j \text{ leaves node } i, \\ -1 & \text{if branch } j \text{ enters node } i, \\ 0 & \text{if branch } j \text{ is not incident with node } i. \end{cases}$$

If the network graph is connected the rows of  $A_0$  are linearly dependent and we can choose one arbitrary node as reference node. By eliminating the corresponding row in the incidence matrix we obtain the *reduced incidence matrix*  $A \in \mathbb{R}^{n_\eta-1, n_b}$  that then has full row rank. Kirchhoff's current law and Kirchhoff's voltage law takes the form

$$A\iota = 0, \quad \nu = A^T\eta, \tag{3}$$

where  $\iota$  denotes the vector of all branch currents,  $\nu$  denotes the vector of all branch voltages, and  $\eta$  denotes the vector of all node potentials (excepting the  
 95 reference node for which the node potential is set to zero). Assuming that the

branches are ordered by the type of component, we can split  $A$  into element-related incidence matrices

$$A = \begin{bmatrix} A_C & A_L & A_{\mathcal{R}} & A_{\mathcal{V}} & A_I \end{bmatrix}, \quad (4)$$

such that  $A_C \in \mathbb{R}^{n_\eta-1, n_C}$ ,  $A_L \in \mathbb{R}^{n_\eta-1, n_L}$ ,  $A_{\mathcal{R}} \in \mathbb{R}^{n_\eta-1, n_{\mathcal{R}}}$ ,  $A_{\mathcal{V}} \in \mathbb{R}^{n_\eta-1, n_{\mathcal{V}}}$ , and  $A_I \in \mathbb{R}^{n_\eta-1, n_I}$ . Here,  $n_{\mathcal{V}}$  denotes the number of voltage sources,  $n_I$  the number of current sources,  $n_C$  the number of capacitors,  $n_L$  the number of inductors, and  $n_{\mathcal{R}}$  the number of resistors in the circuit, respectively. The vectors  $\nu$  and  $\iota$  are split accordingly into

$$\nu = \begin{bmatrix} \nu_C^T & \nu_L^T & \nu_{\mathcal{R}}^T & \nu_{\mathcal{V}}^T & \nu_I^T \end{bmatrix}^T, \quad \iota = \begin{bmatrix} \iota_C^T & \iota_L^T & \iota_{\mathcal{R}}^T & \iota_{\mathcal{V}}^T & \iota_I^T \end{bmatrix}^T.$$

Furthermore, we use the constitutive element relations

$$\nu_L = \frac{d}{dt}\phi(\iota_L), \quad \iota_C = \frac{d}{dt}q(\nu_C), \quad \iota_{\mathcal{R}} = g(\nu_{\mathcal{R}}), \quad \iota_I = I_s(t), \quad \nu_{\mathcal{V}} = \mathcal{V}_s(t) \quad (5)$$

for inductors, capacitors, resistors, current and voltage sources, where  $g : \mathbb{R}^{n_{\mathcal{R}}} \rightarrow \mathbb{R}^{n_{\mathcal{R}}}$  is the conductance function,  $q : \mathbb{R}^{n_C} \rightarrow \mathbb{R}^{n_C}$  is the charge function and  $\phi : \mathbb{R}^{n_L} \rightarrow \mathbb{R}^{n_L}$  is the flux function. Here, we restrict to the case of independent current and voltage sources described by the source functions  $I_s(t)$  and  $\mathcal{V}_s(t)$ , respectively. In general, also controlled sources are possible, see [14]. Altogether, we get a DAE system of the form

$$A_C \frac{d}{dt}q(A_C^T \eta) + A_L \iota_L + A_{\mathcal{R}} g(A_{\mathcal{R}}^T \eta) + A_{\mathcal{V}} \nu_{\mathcal{V}} + A_I I_s = 0, \quad (6a)$$

$$\frac{d}{dt}\phi(\iota_L) - A_L^T \eta = 0, \quad (6b)$$

$$A_{\mathcal{V}}^T \eta - \mathcal{V}_s = 0. \quad (6c)$$

The system (6) consists of  $(n_\eta - 1) + n_L + n_{\mathcal{V}}$  equations in the  $(n_\eta - 1) + n_L + n_{\mathcal{V}}$  unknowns  $[\eta, \iota_L, \nu_{\mathcal{V}}]$  and is also known as the *MNA equations* of the electrical circuit. Note that in (6) we have omitted the dependency on time for better readability. For details on the constitutive element relations and on the derivation of the MNA equations, see also [12, 14, 15, 16].

We say that the DAE system (6) is *well-posed* if it satisfies the following assumptions.

(A1) The circuit contains no  $\mathcal{V}$ -loops, i.e.,  $A_{\mathcal{V}}$  has full column rank.

(A2) The circuit contains no  $I$ -cutsets, i.e.,  $[A_C \ A_L \ A_{\mathcal{R}} \ A_{\mathcal{V}}]$  has full row rank.

(A3) The charge function  $q : \mathbb{R}^{n_c} \rightarrow \mathbb{R}^{n_c}$  is continuously differentiable and the Jacobian

$$C(\nu_C) := \frac{\partial}{\partial \nu_C} q(\nu_C)$$

120 is symmetric and pointwise positive definite.

(A4) The flux function  $\phi : \mathbb{R}^{n_L} \rightarrow \mathbb{R}^{n_L}$  is continuously differentiable and the Jacobian

$$L(\iota_L) := \frac{\partial}{\partial \iota_L} \phi(\iota_L),$$

is symmetric and pointwise positive definite.

(A5) The conductance function  $g : \mathbb{R}^{n_{\mathcal{R}}} \rightarrow \mathbb{R}^{n_{\mathcal{R}}}$  is continuously differentiable and the Jacobian

$$\mathcal{G}(\nu_{\mathcal{R}}) := \frac{\partial}{\partial \nu_{\mathcal{R}}} g(\nu_{\mathcal{R}}),$$

is symmetric and pointwise positive definite.

A  $\mathcal{V}$ -loop is defined as a loop in the circuit graph that consists only of branches corresponding to voltage sources. In the same way, a  $\mathcal{CV}$ -loop means a loop  
 125 that consists only of branches corresponding to capacitances and/or voltage sources. Likewise, an  $I$ -cutset is a cutset in the circuit graph that consists only of branches corresponding to current sources, and an  $\mathcal{LI}$ -cutset is a cutset that consists only of branches corresponding to inductances and/or current sources. Assumption (A1) implies that there are no short-circuits. In a similar manner,  
 130 the occurrence of  $I$ -cutsets may lead to contradictions in the Kirchhoff laws (source functions may not sum up to zero), which is excluded by assumption (A2). The assumptions (A3), (A4) and (A5) imply that all circuit elements are passive, i.e., they do not generate energy.

Due to the special structure of (6), it is possible to determine the index by graph  
 135 theoretical considerations.



**Theorem 3** ([14, 16]). *Consider an electrical circuit with circuit equations as in (6). Assume that the assumptions (A1)-(A5) hold.*

1. *The following statements are equivalent:*

- *the MNA equations (6) are of d-index  $\nu_d = 0$ ;*
- *the circuit contains neither voltage sources nor  $\mathcal{RLI}$ -cutsets;*
- *$n_v = 0$  and  $\text{rank } A_C = n_\eta - 1$ .*

2. *The following statements are equivalent:*

- *the MNA equations (6) are of d-index  $\nu_d = 1$ ;*
- *the circuit contains neither  $\mathcal{LI}$ -cutsets nor  $\mathcal{CV}$ -loops (except for pure  $\mathcal{C}$ -loops);*
- *$\text{rank}[A_C, A_{\mathcal{R}}, A_v] = n_\eta - 1$  and  $\ker[A_C, A_v] = \ker A_C \times \{0\}$ .*

3. *The following statements are equivalent:*

- *the MNA equations (6) are of d-index  $\nu_d = 2$ ;*
- *the circuit contains  $\mathcal{LI}$ -cutsets or  $\mathcal{CV}$ -loops which are no pure  $\mathcal{C}$ -loops;*
- *$\text{rank}[A_C, A_{\mathcal{R}}, A_v] < n_\eta - 1$  or  $\ker[A_C, A_v] \neq \ker A_C \times \{0\}$ .*

In the case that  $\nu_d = 2$  the MNA equations (6) contain hidden constraints that can be revealed by differentiating certain parts of the system, see [14]. In particular, the source functions that belong to  $\mathcal{CV}$ -loops or  $\mathcal{LI}$ -cutsets have to be differentiable if the DAE has d-index  $\nu_d = 2$ . Note that under the given assumptions, the MNA equations (6) will always have a d-index  $\nu_d \leq 2$ .

**Remark 4.** The formulation using the MNA equations (6) belong to the class of *nodal methods* that are characterized by the use of node potentials as fundamental modal variables together with some branch variables. The MNA equations (6) are often used in circuit simulation programs (e.g. in SPICE or TITAN), because their compact form allows for efficient numerical computations [14, 15, 17].

### 3.2. The Modified Loop Analysis

An alternative way to model electrical circuits is the *Modified Loop Analysis* (MLA) [18]. Here, in order to describe the topology of the circuit graph, instead of the incidence matrix  $A$  one uses the so-called *loop matrix*  $B_0 \in \mathbb{R}^{n_\ell, n_b}$  with entries defined as follows

$$b_{ij} = \begin{cases} 1 & \text{if branch } j \text{ belongs to loop } i \text{ and has the same orientation,} \\ -1 & \text{if branch } j \text{ belongs to loop } i \text{ and has the contrary orientation,} \\ 0 & \text{if branch } j \text{ does not belong to loop } i. \end{cases}$$

Here,  $n_\ell$  denotes the number of all oriented loops in the directed graph. By removing all linearly dependent rows in the loop matrix we get the *reduced loop matrix*  $B \in \mathbb{R}^{n_b - n_\eta + 1, n_b}$  of full row rank (for details see Appendix A). Now, Kirchhoff's current law and Kirchhoff's voltage law takes the form

$$B\nu = 0, \quad i = B^T j,$$

where  $j(t) \in \mathbb{R}^{n_b - n_\eta + 1}$  is the vector of all loop currents. If we split  $B$  into element-related loop matrices

$$B = \begin{bmatrix} B_C & B_L & B_{\mathcal{R}} & B_{\mathcal{V}} & B_I \end{bmatrix}$$

and insert the constitutive element relations we get a DAE system of the form

$$\begin{aligned} B_L \mathcal{L}(B_L^T j) B_L^T \frac{d}{dt} j + B_{\mathcal{R}} r(B_{\mathcal{R}}^T j) + B_C \nu_C + B_I \nu_I + B_{\mathcal{V}} \mathcal{V}_s &= 0, \\ C(\nu_C) \frac{d}{dt} \nu_C - B_C^T j &= 0, \\ B_I^T j + I_s &= 0. \end{aligned} \quad (7)$$

Here,  $r : \mathbb{R}^{n_{\mathcal{R}}} \rightarrow \mathbb{R}^{n_{\mathcal{R}}}$  denotes the resistance function<sup>1</sup> and system (7) consists of  $n_b - n_\eta + 1 + n_c + n_i$  equations in the  $n_b - n_\eta + 1 + n_c + n_i$  unknowns  $[j, \nu_C, \nu_I]$ . The equations (7) are also known as the *MLA equations*. Again, the index of the MLA equations (7) can be determined based on the topology of the circuit.

<sup>1</sup>If the resistance function  $r$  is continuously differentiable with  $\mathcal{R} = \frac{\partial}{\partial i_{\mathcal{R}}} r$ , then  $\mathcal{R}$  is the pointwise inverse of  $\mathcal{G}$ .

**Theorem 5** ([16]). *Consider an electrical circuit with MLA equations (7). Assume that (A1)-(A5) hold.*

- 170 1. *The following statements are equivalent:*
  - *the MLA equations (7) are of d-index  $\nu_d = 0$ ;*
  - *the circuit contains neither current sources nor  $\mathcal{CRV}$ -loops;*
  - *$n_l = 0$  and  $\text{rank } B_L = n_b - n_\eta + 1$ .*
2. *The following statements are equivalent:*
  - 175 • *the MLA equations (7) are of d-index  $\nu_d = 1$ ;*
  - *the circuit contains neither  $\mathcal{CV}$ -loops nor  $\mathcal{LI}$ -cutsets (except for pure  $\mathcal{L}$ -cutsets);*
  - *$\text{rank}[B_L, B_R, B_I] = n_b - n_\eta + 1$  and  $\ker[B_L, B_I] = \ker B_L \times \{0\}$ .*
3. *The following statements are equivalent:*
  - 180 • *the MLA equations (7) are of d-index  $\nu_d = 2$ ;*
  - *the circuit contains  $\mathcal{CV}$ -loops or  $\mathcal{LI}$ -cutsets which are no pure  $\mathcal{L}$ -cutsets;*
  - *$\text{rank}[B_L, B_R, B_I] < n_b - n_\eta + 1$  or  $\ker[B_L, B_I] \neq \ker B_L \times \{0\}$ .*

**Remark 6.** Note that for a given electrical circuit the MNA equations (6) and  
 185 the MLA equations (7) may have different d-index depending on the topology of the circuit, in particular, if the circuit contains pure  $\mathcal{C}$ -loops or pure  $\mathcal{L}$ -cutsets, cf. Example 6.

### 3.3. Branch-Oriented Model Equations

In branch-oriented model formulations Kirchhoff's current and voltage laws are stated as

$$A\iota = 0, \tag{8}$$

$$B\nu = 0. \tag{9}$$

Together with the constitutive element relations (5) we obtain a DAE system of the form

$$C(\nu_C) \frac{d}{dt} \nu_C = \iota_C, \quad (10a)$$

$$\mathcal{L}(\iota_L) \frac{d}{dt} \iota_L = \nu_L, \quad (10b)$$

$$0 = \iota_{\mathcal{R}} - g(\nu_{\mathcal{R}}), \quad (10c)$$

$$0 = A_C \iota_C + A_{\mathcal{R}} \iota_{\mathcal{R}} + A_{\mathcal{L}} \iota_L + A_I \iota_I + A_{\mathcal{V}} \iota_{\mathcal{V}}, \quad (10d)$$

$$0 = B_C \nu_C + B_{\mathcal{R}} \nu_{\mathcal{R}} + B_{\mathcal{L}} \nu_L + B_I \nu_I + B_{\mathcal{V}} \nu_{\mathcal{V}}, \quad (10e)$$

$$0 = \iota_I - I_s, \quad (10f)$$

$$0 = \nu_{\mathcal{V}} - \mathcal{V}_s, \quad (10g)$$

which consists of  $2n_b = 2(n_C + n_{\mathcal{R}} + n_{\mathcal{L}} + n_I + n_{\mathcal{V}})$  equations in the unknown branch currents  $\iota_*$  and branch voltages  $\nu_*$  for  $* \in \{C, \mathcal{R}, \mathcal{L}, I, \mathcal{V}\}$ . We will call (10) the *branch-oriented model equations* of the electrical circuit.

**Theorem 7.** *Consider an electrical circuit with branch-oriented model equations (10). Assume that (A1)-(A5) hold.*

1. *The following statements are equivalent:*

- the branch-oriented model equations (10) are of  $d$ -index  $\nu_d = 1$ ;
- $\text{rank}[A_{\mathcal{R}}, A_C, A_{\mathcal{V}}] = n_{\eta} - 1$  and  $\ker[A_C, A_{\mathcal{V}}] = \{0\}$ ;
- $\text{rank}[B_{\mathcal{L}}, B_{\mathcal{R}}, B_I] = n_b - n_{\eta} + 1$  and  $\ker[B_{\mathcal{L}}, B_I] = \{0\}$ ;
- the circuit contains neither  $\mathcal{LI}$ -cutsets nor  $\mathcal{CV}$ -loops (including pure  $\mathcal{C}$ -loops and pure  $\mathcal{L}$ -cutsets).

2. *The following statements are equivalent:*

- the branch-oriented model equations (10) are of  $d$ -index  $\nu_d = 2$ ;
- $\text{rank}[A_{\mathcal{R}}, A_C, A_{\mathcal{V}}] < n_{\eta} - 1$  or  $\ker[A_C, A_{\mathcal{V}}] \neq \{0\}$ ;
- $\text{rank}[B_{\mathcal{L}}, B_{\mathcal{R}}, B_I] < n_b - n_{\eta} + 1$  or  $\ker[B_{\mathcal{L}}, B_I] \neq \{0\}$ ;
- the circuit contains  $\mathcal{LI}$ -cutsets or  $\mathcal{CV}$ -loops (including pure  $\mathcal{C}$ -loops and pure  $\mathcal{L}$ -cutsets).

*Proof.* For the geometric index the proof is given in [12]. Except for differences in the smoothness requirements (which do not apply here) the geometric index is equal to the differentiation index [10, 19].  $\square$

Next, we introduce the *cutset matrix*  $Q_0$  of a connected directed graph. The removal of any cutset in the graph results in a directed graph with two connected components  $C_1$  and  $C_2$ . Given a branch in the cutset, one terminal node must be in  $C_1$  and the other one in  $C_2$ . Thus, we can define two different orientations in this cutset: from  $C_1$  to  $C_2$  or from  $C_2$  to  $C_1$ . With this the cutset matrix  $Q_0 = [q_{ij}]$  can be defined as follows

$$q_{ij} = \begin{cases} 1 & \text{if branch } j \text{ is in cutset } i \text{ with the same orientation,} \\ -1 & \text{if branch } j \text{ is in cutset } i \text{ with the opposite orientation,} \\ 0 & \text{if branch } j \text{ is not in cutset } i. \end{cases}$$

Again,  $n_\eta - 1$  linearly independent rows of  $Q_0$  define a reduced cutset matrix  $Q \in \mathbb{R}^{n_\eta - 1, n_b}$ . Furthermore, it holds that  $BA^T = BQ^T = 0$  and  $\text{im } B^T = \ker A = \ker Q$ , see Theorem 24. Thus, we have that  $Q = MA$  for some nonsingular matrix  $M$  and Kirchhoff's current law (8) can be replaced by the relation  $Q\iota = 0$ . In any circuit graph, we can always choose a tree such that all voltage sources correspond to twigs and all current sources correspond to links. This choice of a tree leads to a set of *fundamental cutsets* and *fundamental loops* (each one uniquely defined by a twig or a link, respectively), and choosing the orientations of these cutsets and loops coherently with the orientations of the corresponding twig or link, the matrices  $Q$  and  $B$  take the form

$$Q = \begin{bmatrix} I_r & -F^T \end{bmatrix}, \quad B = \begin{bmatrix} F & I_s \end{bmatrix}$$

for a certain matrix  $F \in \mathbb{R}^{s,r}$  with  $s = n_b - n_\eta + 1$  and  $r = n_\eta - 1$ . This results from the orthogonality property  $BQ^T = 0$ , see Theorem 24 and the details in Appendix A. The first block matrix in  $B$  and  $Q$  is associated with the twigs of the tree, whereas the second block is associated with the links.

With this Kirchhoff's current law (8) and Kirchhoff's voltage law (9) can be represented as

$$\mathbf{i}_1 = F^T \mathbf{i}_2, \quad (11a)$$

$$\mathbf{v}_2 = -F \mathbf{v}_1, \quad (11b)$$

where the subscript **1** denotes the tree elements while the subscript **2** denotes the cotree elements. We can split  $\mathbf{v}_1$  and  $\mathbf{v}_2$  into

$$\mathbf{v}_1 = [\nu_{C1}^T, \nu_{L1}^T, \nu_{q'}^T, \nu_{R1}^T]^T,$$

225 consisting of twig capacitors, twig inductors, voltage sources, and twig resistors, and

$$\mathbf{v}_2 = [\nu_{L2}^T, \nu_{C2}^T, \nu_I^T, \nu_{R2}^T]^T,$$

consisting of link inductors, link capacitors, current sources, and link resistors. Note again that the tree is chosen in such a way that voltage sources always belong to the tree elements while current sources always belong to the cotree elements. A similar splitting can be done for the vectors  $\mathbf{i}_1$  and  $\mathbf{i}_2$  into

230

$$\mathbf{i}_1 = [i_{C1}^T \ i_{L1}^T \ i_{q'}^T \ i_{R1}^T]^T, \quad \mathbf{i}_2 = [i_{L2}^T \ i_{C2}^T \ i_I^T \ i_{R2}^T]^T,$$

and the matrix  $F$  into

$$F = \begin{bmatrix} F_{11} & F_{12} & F_{13} & F_{14} \\ F_{21} & F_{22} & F_{23} & F_{24} \\ F_{31} & F_{32} & F_{33} & F_{34} \\ F_{41} & F_{42} & F_{43} & F_{44} \end{bmatrix}. \quad (12)$$

A *proper tree* in a connected circuit graph is a tree which contains all voltage sources and all capacitances as well as (possibly) some resistors, but neither current sources nor inductors. Likewise, a *normal tree* in a connected circuit graph is a tree which contains all voltage sources, no current sources, as many capacitors as possible, and as few inductors as possible; it may also contain some resistors. It can be shown that a connected circuit graph has neither  $\mathcal{CV}$ -loops nor  $\mathcal{LI}$ -cutsets if and only if it contains a proper tree, see [12]. Note that due

to Lemma 27 there exists no loops just defined by twigs, and no cutsets just defined by links. Thus, in a normal tree the fundamental cutsets defined by each twig inductor only have link inductors and currents sources. Analogously, in a normal tree the fundamental loops defined by each link capacitor only involve twig capacitors and voltage sources. Thus, for a normal tree we have  $F_{22} = 0$ ,  $F_{24} = 0$  and  $F_{42} = 0$  in (12). Summarizing these results we get from (10) by replacing (10d),(10e) by (11a),(11b) that

$$\begin{aligned}
 \mathcal{C}(\nu_{C1}, \nu_{C2}) \frac{d}{dt} \begin{bmatrix} \nu_{C1} \\ \nu_{C2} \end{bmatrix} &= \begin{bmatrix} \imath_{C1} \\ \imath_{C2} \end{bmatrix}, \\
 \mathcal{L}(\imath_{L1}, \imath_{L2}) \frac{d}{dt} \begin{bmatrix} \imath_{L1} \\ \imath_{L2} \end{bmatrix} &= \begin{bmatrix} \nu_{L1} \\ \nu_{L2} \end{bmatrix}, \\
 0 &= \begin{bmatrix} \imath_{R1} \\ \imath_{R2} \end{bmatrix} - g(\nu_{R1}, \nu_{R2}), \\
 \begin{bmatrix} \nu_{L2} \\ \nu_{C2} \\ \nu_I \\ \nu_{R2} \end{bmatrix} &= - \begin{bmatrix} F_{11} & F_{12} & F_{13} & F_{14} \\ F_{21} & 0 & F_{23} & 0 \\ F_{31} & F_{32} & F_{33} & F_{34} \\ F_{41} & 0 & F_{43} & F_{44} \end{bmatrix} \begin{bmatrix} \nu_{C1} \\ \nu_{L1} \\ \nu_{\mathcal{V}} \\ \nu_{R1} \end{bmatrix}, \\
 \begin{bmatrix} \imath_{C1} \\ \imath_{L1} \\ \imath_{\mathcal{V}} \\ \imath_{R1} \end{bmatrix} &= \begin{bmatrix} F_{11}^T & F_{21}^T & F_{31}^T & F_{41}^T \\ F_{12}^T & 0 & F_{32}^T & 0 \\ F_{13}^T & F_{23}^T & F_{33}^T & F_{43}^T \\ F_{14}^T & 0 & F_{34}^T & F_{44}^T \end{bmatrix} \begin{bmatrix} \imath_{L2} \\ \imath_{C2} \\ \imath_I \\ \imath_{R2} \end{bmatrix}, \\
 0 &= \nu_{\mathcal{V}} - \mathcal{V}_s \\
 0 &= \imath_I - I_s.
 \end{aligned} \tag{13}$$

Note that the DAEs (10) and (13) have the same d-index, since the performed transformations do not change the analytical properties of the system. Thus, the statements of Theorem 7 hold also for (13).

### 3.4. Port-Hamiltonian Circuit Equations

235 If we consider the electrical circuit as a power-based network model of inter-connected subsystems that mutually influence each other via energy flow, this directly leads us to the starting point of port-Hamiltonian systems theory, see [20, 21]. The formulation of a physical system as *port-Hamiltonian system* has many advantages as e.g. the preservation of energy, the preservation of passivity  
240 or stability, see e.g. [21].

Here, we consider *linear port-Hamiltonian DAEs (pHDAE)* of the form

$$\begin{aligned} E\dot{x} &= (J - R)Qx - EKx + Bu, \\ y &= B^T Qx, \end{aligned} \quad (14)$$

where  $J, R, K \in C(\mathbb{I}, \mathbb{R}^{n,n})$ ,  $E, Q \in C^1(\mathbb{I}, \mathbb{R}^{n,n})$  with  $R = R^T \geq 0$ , and  $E^T Q = Q^T E \geq 0$  pointwise on  $\mathbb{I}$ , satisfying

$$\frac{d}{dt}(Q^T E) = Q^T(EK - JQ) + (EK - JQ)^T Q, \quad (15)$$

as well as  $B \in C(\mathbb{I}, \mathbb{R}^{n,p})$ . Note that also more general formulations are possible, see [22, 23]. Here,  $u$  denotes the  $p$ -dimensional input of the system and  $y$  denotes the  $p$ -dimensional output of the system; together they define the (external) ports. The matrix  $J$  can be seen as the interconnection matrix,  $R$  is the  
245 resistance matrix, and  $E^T Q$  describes the total energy of the system represented by the *Hamiltonian*  $\mathcal{H}(x) = \frac{1}{2}x^T E^T Qx$ . The matrix  $K$  is required to describe equivalence transformations in the time-varying setting, for constant coefficients one may set  $K$  to zero.

**Remark 8.** Kirchhoff's current law and Kirchhoff's voltage law define a (separable) *Dirac structure*  $\mathcal{D} = \{(\iota, \nu) | A\iota = 0, \nu = A^T \eta \text{ for some } \eta \in \mathbb{R}^n\}$  that  
250 describes the underlying geometric structure of the port-Hamiltonian DAE. The currents through the electrical components are the flows and the voltages across the electrical components are the efforts defining the port variables of the Dirac structure.



In order to formulate the circuit equations as a pHDAE, we start by considering the system of equations consisting of Kirchhoff laws (3) together with the constitutive element relations (5). We insert the relations  $\nu_{\mathcal{V}} = \mathcal{V}_s$ ,  $\nu_L = A_L^T \eta$  and  $\iota_I = I_s$ , and, assuming linear element relations for the resistances, capacitances and inductances, a reordering of the equations yields

$$\mathcal{C}(t) \frac{d}{dt} \nu_C = \iota_C, \quad (16a)$$

$$\mathcal{L}(t) \frac{d}{dt} \iota_L = A_L^T \eta, \quad (16b)$$

$$0 = -\nu_C + A_C^T \eta, \quad (16c)$$

$$0 = -\nu_{\mathcal{R}} + A_{\mathcal{R}}^T \eta, \quad (16d)$$

$$0 = A_{\mathcal{V}}^T \eta - \mathcal{V}_s, \quad (16e)$$

$$0 = \iota_{\mathcal{R}} - \mathcal{G}(t) \nu_{\mathcal{R}}, \quad (16f)$$

$$0 = -A_L \iota_L - A_C \iota_C - A_{\mathcal{R}} \iota_{\mathcal{R}} - A_{\mathcal{V}} \iota_{\mathcal{V}} - A_I I_s, \quad (16g)$$

$$0 = -\nu_I + A_I^T \eta. \quad (16h)$$

Considering the last equation (16h) as an output equation, we obtain a pHDAE of the form (14) with

$$E = \left[ \begin{array}{cc|cccccc} \mathcal{C}(t) & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \mathcal{L}(t) & 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{array} \right], \quad R = \left[ \begin{array}{cc|cccccc} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \mathcal{G}(t) & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{array} \right],$$

$$J = \left[ \begin{array}{cc|cccccc} 0 & 0 & I & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & A_L^T \\ \hline -I & 0 & 0 & 0 & 0 & 0 & 0 & A_C^T \\ 0 & 0 & 0 & 0 & 0 & -I & 0 & A_{\mathcal{R}}^T \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & A_{\mathcal{V}}^T \\ 0 & 0 & 0 & I & 0 & 0 & 0 & 0 \\ 0 & -A_L & -A_C & -A_{\mathcal{R}} & -A_{\mathcal{V}} & 0 & 0 & 0 \end{array} \right], \quad B = \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \\ -I & 0 \\ 0 & 0 \\ 0 & -A_I \end{bmatrix}, \quad Q = I,$$

255 and with state variables

$$x = [\nu_C^T, \iota_L^T, \iota_C^T, \iota_R^T, \iota_{\mathcal{V}}^T, \nu_R^T, \eta^T]^T,$$

and (external) port variables

$$u = \begin{bmatrix} \mathcal{V}_s \\ I_s \end{bmatrix}, \quad y = \begin{bmatrix} -\iota_{\mathcal{V}} \\ -\nu_I \end{bmatrix}.$$

The Hamiltonian is given by

$$\mathcal{H}(x) = \frac{1}{2}x^T E^T Q x = \frac{1}{2}\nu_C^T \mathcal{C}(t)\nu_C + \frac{1}{2}\iota_L^T \mathcal{L}(t)\iota_L$$

and describes the total energy of the system. In the following, we call system (16) the *port-Hamiltonian Circuit equations (pHC equations)*.

**Remark 9.** Condition (15) is met by defining  $K = \text{diag}(K_{11}, K_{22}, 0)$ , where  $K_{11} = K_{11}^T$  and  $K_{22} = K_{22}^T$  are given by the unique solutions of the matrix Lyapunov equations

$$\begin{aligned} CK_{11} + K_{11}C &= \dot{C}, \\ LK_{22} + K_{22}L &= \dot{L}. \end{aligned}$$

260 Since we can always find a variable transformation for  $x$  that eliminates  $K$  in (14), we can omit these parts in the formulation of the port-Hamiltonian circuit equations (16) (cf. also [22]).

**Theorem 10.** Consider an electrical circuit with pHC equations (16). Assume that (A1)-(A5) hold.

1. The following statements are equivalent:

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- the pHC equations (16) are of  $d$ -index  $\nu_d = 1$ ;
  - $\text{rank}[A_R, A_C, A_{\mathcal{V}}] = n_{\eta} - 1$  and  $\ker[A_C, A_{\mathcal{V}}] = \{0\}$ ;
  - the circuit contains neither  $\mathcal{LI}$ -cutsets nor  $\mathcal{CV}$ -loops (including pure  $\mathcal{C}$ -loops).

2. The following statements are equivalent:

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- the  $pHC$  equations (16) are of  $d$ -index  $\nu_d = 2$ ;
- $\text{rank}[A_{\mathcal{R}}, A_C, A_{\mathcal{V}}] < n_\eta - 1$  or  $\ker[A_C, A_{\mathcal{V}}] \neq \{0\}$ ;
- the circuit contains  $\mathcal{LI}$ -cutsets or  $\mathcal{CV}$ -loops.

In order to prove Theorem 10 we make use of the following Lemma.

**Lemma 11.** *Let  $A \in \mathbb{R}^{n_1, m}$ ,  $B \in \mathbb{R}^{n_1, n_2}$ , and  $C \in \mathbb{R}^{m, m}$  be symmetric positive*  
 275 *definite. Then for*

$$M = \begin{bmatrix} ACA^T & B \\ -B^T & 0 \end{bmatrix}$$

*it holds that  $\ker M = \ker [A, B]^T \times \ker B$ . In particular,  $M$  is invertible if and only if  $\ker [A, B]^T = \{0\}$  and  $\ker B = \{0\}$ .*

*Proof.* For  $v \in \ker [A, B]^T \times \ker B$  it follows immediately that  $v \in \ker M$ . For  
 the converse, let  $v \in \ker M$  and partition  $v = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix}$  according to the block  
 280 structure of  $M$ . Then

$$v^T M v = v_1^T A C A^T v_1 + v_1 B v_2 - v_2^T B^T v_1 = 0 \iff v_1^T A C A^T v_1 = 0.$$

Thus,  $v_1 \in \ker A^T$  and from  $Mv = 0$  we get that  $Bv_2 = 0$  and  $-B^T v_1 = 0$ , and consequently  $v \in \ker [A, B]^T \times \ker B$ .  $\square$

In the following, we denote by  $Q_C$ ,  $Q_{\mathcal{V}-C}$ ,  $Q_{\mathcal{R}-C\mathcal{V}}$ ,  $Q_{\mathcal{V}}$ ,  $\bar{Q}_C$ , and  $\bar{Q}_{\mathcal{V}-C}$  projectors onto  $\ker A_C^T$ ,  $\ker A_{\mathcal{V}}^T Q_C$ ,  $\ker A_{\mathcal{R}}^T Q_C Q_{\mathcal{V}-C}$ ,  $\ker A_{\mathcal{V}}^T$ ,  $\ker A_C$ , and  $\ker Q_C^T A_{\mathcal{V}}$ ,  
 285 respectively. The complementary projectors will be denoted by  $P := I - Q$ , with the corresponding sub-index. In order to shorten notations, we use the abbreviation  $Q_{C\mathcal{R}\mathcal{V}} := Q_C Q_{\mathcal{V}-C} Q_{\mathcal{R}-C\mathcal{V}}$ .

**Lemma 12.** [14] If  $C$ ,  $\mathcal{L}$  and  $\mathcal{G}$  are positive definite, then the matrices

$$\begin{aligned} H_1 &:= A_C C A_C^T + Q_C^T Q_C \\ H_2 &:= Q_C^T A_{\mathcal{V}} A_{\mathcal{V}}^T Q_C + Q_{\mathcal{V}-C}^T Q_{\mathcal{V}-C} \\ H_3 &:= A_{\mathcal{V}}^T Q_C Q_C^T A_{\mathcal{V}} + \bar{Q}_{\mathcal{V}-C}^T \bar{Q}_{\mathcal{V}-C} \\ H_4 &:= Q_{\mathcal{V}-C}^T Q_C^T A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T Q_C Q_{\mathcal{V}-C} + Q_{\mathcal{R}-C}^T Q_{\mathcal{R}-C} \\ H_5 &:= Q_{C\mathcal{R}\mathcal{V}}^T A_{\mathcal{L}} \mathcal{L}^{-1} A_{\mathcal{L}}^T Q_{C\mathcal{R}\mathcal{V}} + P_{C\mathcal{R}\mathcal{V}}^T P_{C\mathcal{R}\mathcal{V}} \\ H_6 &:= \bar{Q}_{\mathcal{V}-C}^T A_{\mathcal{V}}^T H_1^{-1} A_{\mathcal{V}} \bar{Q}_{\mathcal{V}-C} + \bar{P}_{\mathcal{V}-C}^T \bar{P}_{\mathcal{V}-C} \\ H_7 &:= A_C^T A_C + \bar{Q}_C^T \bar{Q}_C \end{aligned}$$

are nonsingular.

*Proof of Theorem 10.* Omitting the output equations, the pHc equations (16) are of the form

$$\begin{bmatrix} E_{11} & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & A_{12} \\ -A_{12}^T & A_{22} \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ f_2 \end{bmatrix}, \quad (17)$$

where  $E_{11}$  is (pointwise) nonsingular. The DAE (17) is of d-index 1 if and only if  $A_{22}$  is invertible. In this case, differentiation of the constraints yields

$$\begin{aligned} 0 &= -A_{12}^T \dot{x}_1 + A_{22} \dot{x}_2 + \dot{f}_2 \\ \implies \dot{x}_2 &= A_{22}^{-1} A_{12}^T \dot{x}_1 - A_{22}^{-1} \dot{f}_2 = A_{22}^{-1} A_{12}^T E_{11}^{-1} A_{12} x_2 - A_{22}^{-1} \dot{f}_2. \end{aligned}$$

It holds that

$$A_{22} = \begin{bmatrix} 0 & 0 & 0 & 0 & A_C^T \\ 0 & 0 & 0 & -I & A_{\mathcal{R}}^T \\ 0 & 0 & 0 & 0 & A_{\mathcal{V}}^T \\ 0 & I & 0 & -\mathcal{G} & 0 \\ -A_C & -A_{\mathcal{R}} & -A_{\mathcal{V}} & 0 & 0 \end{bmatrix} \sim \left[ \begin{array}{cc|ccc} 0 & -I & 0 & 0 & 0 \\ I & -\mathcal{G} & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & A_C^T \\ 0 & 0 & 0 & 0 & A_{\mathcal{V}}^T \\ 0 & 0 & -A_C & -A_{\mathcal{V}} & -A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T \end{array} \right]$$

by basic row and column transformations. Due to Lemma 11, the matrix  $A_{22}$  is invertible if and only if

$$\ker[A_{\mathcal{R}}, A_C, A_{\mathcal{V}}]^T = \{0\} \quad \text{and} \quad \ker[A_C, A_{\mathcal{V}}] = \{0\},$$

i.e., if and only if the matrix  $[A_C \ A_{\mathcal{R}} \ A_{\mathcal{V}}]$  has full row rank and  $\begin{bmatrix} A_C^T \\ A_{\mathcal{V}}^T \end{bmatrix}$  has  
 290 full column rank. Due to Lemma 26 the first condition is equivalent to the  
 absence of  $\mathcal{LI}$ -cutsets in the circuit, while the second condition is equivalent to  
 the absence of  $\mathcal{CV}$ -loops (including pure  $\mathcal{C}$ -loops).

If the matrix  $A_{22}$  is not invertible, then there are hidden constraints contained  
 in the system. We follow the lines of the proof of Theorem 5.1 in [14]. If we  
 differentiate the algebraic constraints in (16) we get

$$\begin{aligned} 0 &= -\frac{d}{dt}\nu_C + A_C^T \frac{d}{dt}\eta, \\ 0 &= -\frac{d}{dt}\nu_{\mathcal{R}} + A_{\mathcal{R}}^T \frac{d}{dt}\eta, \\ 0 &= A_{\mathcal{V}}^T \frac{d}{dt}\eta + \frac{d}{dt}\mathcal{V}_s, \\ 0 &= \frac{d}{dt}\iota_{\mathcal{R}} - \mathcal{G} \frac{d}{dt}\nu_{\mathcal{R}}, \\ 0 &= -A_C \frac{d}{dt}\iota_C - A_{\mathcal{R}} \frac{d}{dt}\iota_{\mathcal{R}} - A_I \frac{d}{dt}I_s - A_{\mathcal{V}} \frac{d}{dt}\iota_{\mathcal{V}} - A_{\mathcal{L}} \frac{d}{dt}\iota_{\mathcal{L}}. \end{aligned}$$

Inserting the expression for  $\frac{d}{dt}\nu_C$  and  $\frac{d}{dt}\iota_{\mathcal{L}}$  from (16a) and (16b), as well as  
 $\frac{d}{dt}\iota_{\mathcal{R}} = \mathcal{G} A_{\mathcal{R}}^T \frac{d}{dt}\eta$  yields

$$\begin{aligned} 0 &= -\mathcal{C}^{-1}\iota_C + A_C^T \frac{d}{dt}\eta, \\ 0 &= A_{\mathcal{V}}^T \frac{d}{dt}\eta + \frac{d}{dt}\mathcal{V}_s, \\ 0 &= -A_C \frac{d}{dt}\iota_C - A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T \frac{d}{dt}\eta - A_I \frac{d}{dt}I_s - A_{\mathcal{V}} \frac{d}{dt}\iota_{\mathcal{V}} - A_{\mathcal{L}} \mathcal{L}^{-1} A_{\mathcal{L}}^T \eta. \end{aligned} \tag{18}$$

In order to extract the underlying ODE, we need to derive expressions for  $\frac{d}{dt}\iota_C$ ,  
 $\frac{d}{dt}\iota_{\mathcal{V}}$  and  $\frac{d}{dt}\eta$  from (18). Then we also get representations for  $\frac{d}{dt}\iota_{\mathcal{R}}$  and  $\frac{d}{dt}\nu_{\mathcal{R}}$   
 from the above relations. We consider the following splitting

$$\begin{aligned} \frac{d}{dt}\eta &= P_C \frac{d}{dt}\eta + Q_C P_{\mathcal{V}-C} \frac{d}{dt}\eta + Q_C Q_{\mathcal{V}-C} P_{\mathcal{R}-C} \frac{d}{dt}\eta + Q_C \mathcal{R} \frac{d}{dt}\eta \\ &=: \frac{d}{dt}\eta_1 + \frac{d}{dt}\eta_2 + \frac{d}{dt}\eta_3 + \frac{d}{dt}\eta_4 = \dot{\eta}_1 + \dot{\eta}_2 + \dot{\eta}_3 + \dot{\eta}_4, \\ \frac{d}{dt}\iota_{\mathcal{V}} &= \bar{P}_{\mathcal{V}-C} \frac{d}{dt}\iota_{\mathcal{V}} + \bar{Q}_{\mathcal{V}-C} \frac{d}{dt}\iota_{\mathcal{V}} =: \frac{d}{dt}\iota_1 + \frac{d}{dt}\iota_2, \\ \frac{d}{dt}\iota_C &= \bar{P}_{C-\mathcal{V}} \frac{d}{dt}\iota_C + \bar{Q}_{C-\mathcal{V}} \frac{d}{dt}\iota_C =: \frac{d}{dt}\iota_3 + \frac{d}{dt}\iota_4, \end{aligned}$$

giving

$$0 = -\mathcal{C}^{-1}\iota_C + A_C^T \dot{\eta}_1, \quad (19a)$$

$$0 = A_{\mathcal{V}}^T(\dot{\eta}_1 + \dot{\eta}_2) + \dot{\mathcal{V}}_s, \quad (19b)$$

$$0 = -A_C \frac{d}{dt} \iota_C - A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T (\dot{\eta}_1 + \dot{\eta}_2 + \dot{\eta}_3) - A_I \dot{I}_s - A_{\mathcal{V}} \frac{d}{dt} \iota_{\mathcal{V}} - A_{\mathcal{L}} \mathcal{L}^{-1} A_{\mathcal{L}}^T \eta. \quad (19c)$$

Multiplication of (19c) with  $Q_C^T$ ,  $Q_{\mathcal{V}-C}^T Q_C^T$  and  $Q_{C\mathcal{R}\mathcal{V}}^T$  yields

$$0 = -Q_C^T \left[ A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T (\dot{\eta}_1 + \dot{\eta}_2 + \dot{\eta}_3) + A_I \dot{I}_s + A_{\mathcal{V}} \frac{d}{dt} \iota_{\mathcal{V}} + A_{\mathcal{L}} \mathcal{L}^{-1} A_{\mathcal{L}}^T \eta \right], \quad (20a)$$

$$0 = -Q_{\mathcal{V}-C}^T Q_C^T \left[ A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T (\dot{\eta}_1 + \dot{\eta}_2 + \dot{\eta}_3) + A_I \dot{I}_s + A_{\mathcal{L}} \mathcal{L}^{-1} A_{\mathcal{L}}^T \eta \right], \quad (20b)$$

$$0 = -Q_{C\mathcal{R}\mathcal{V}}^T \left[ A_I \dot{I}_s + A_{\mathcal{L}} \mathcal{L}^{-1} A_{\mathcal{L}}^T \eta \right]. \quad (20c)$$

If we multiply (19a) with  $H_1^{-1} A_C \mathcal{C}$  we get

$$\dot{\eta}_1 = P_C \frac{d}{dt} \eta = H_1^{-1} A_C \iota_C, \quad (21)$$

and inserting (21) into (19b) gives

$$0 = A_{\mathcal{V}}^T (H_1^{-1} A_C \iota_C + \dot{\eta}_2) + \dot{\mathcal{V}}_s. \quad (22)$$

Moreover, multiplication with  $H_2^{-1} Q_C^T A_{\mathcal{V}}$  yields

$$0 = H_2^{-1} Q_C^T A_{\mathcal{V}} A_{\mathcal{V}}^T (H_1^{-1} A_C \iota_C + \dot{\eta}_2) + \dot{\mathcal{V}}_s,$$

giving  $P_{\mathcal{V}-C} \frac{d}{dt} \eta = f(\iota_C, \dot{\mathcal{V}}_s)$  as a function depending on  $\iota_C$  and  $\dot{\mathcal{V}}_s$ , such that  $\dot{\eta}_2 = Q_C f(\iota_C, \dot{\mathcal{V}}_s)$ . Inserting the expressions for  $\dot{\eta}_1$  and  $\dot{\eta}_2$  into (20b) and multiplication of the resulting system with  $H_4^{-1}$  gives a representation of  $P_{\mathcal{R}-C\mathcal{V}} \frac{d}{dt} \eta$  such that the multiplication with  $Q_C Q_{\mathcal{V}-C}$  gives  $\dot{\eta}_3 = g(\eta, \iota_C, \dot{\mathcal{V}}_s, \dot{I}_s)$  as function of  $\eta, \iota_C, \dot{\mathcal{V}}_s, \dot{I}_s$ . Furthermore, inserting the expressions for  $\dot{\eta}_1$ ,  $\dot{\eta}_2$  and  $\dot{\eta}_3$  into (20a) and multiplying the resulting equation with  $H_3^{-1} A_{\mathcal{V}}^T Q_C$  yields a representation for  $\bar{P}_{\mathcal{V}-C} \frac{d}{dt} \iota_{\mathcal{V}} = \frac{d}{dt} \iota_1$  as function of  $\eta, \iota_C, \dot{\mathcal{V}}_s, \dot{I}_s$ .

Multiplication of (22) with  $\bar{Q}_{\mathcal{V}-C}^T$  yields

$$0 = \bar{Q}_{\mathcal{V}-C}^T A_{\mathcal{V}}^T H_1^{-1} A_C \iota_C + \bar{Q}_{\mathcal{V}-C}^T \dot{\mathcal{V}}_s. \quad (23)$$

If we differentiate (20c) and (23) a second time, we get

$$0 = -Q_{\mathcal{C}\mathcal{R}\mathcal{V}}^T \left[ A_I \ddot{I}_s + A_L \mathcal{L}^{-1} A_L^T (\dot{\eta}_1 + \dot{\eta}_2 + \dot{\eta}_3 + \dot{\eta}_4) \right], \quad (24)$$

$$0 = \bar{Q}_{\mathcal{V}-\mathcal{C}}^T A_{\mathcal{V}}^T H_1^{-1} A_C \frac{d}{dt} \iota_C + \bar{Q}_{\mathcal{V}-\mathcal{C}}^T \ddot{\mathcal{V}}_s. \quad (25)$$

Inserting the expressions for  $\dot{\eta}_1$ ,  $\dot{\eta}_2$  and  $\dot{\eta}_3$  and multiplying (24) with  $H_5^{-1}$  gives

$$0 = -H_5^{-1} Q_{\mathcal{C}\mathcal{R}\mathcal{V}}^T A_L \mathcal{L}^{-1} A_L^T \dot{\eta}_4 + \tilde{f}(\iota_C, \eta, \dot{\mathcal{V}}_s, \dot{I}_s, \ddot{I}_s),$$

which gives an expression for  $\dot{\eta}_4$  as function of  $\iota_C, \eta, \dot{\mathcal{V}}_s, \dot{I}_s, \ddot{I}_s$ . Multiplication of (25) with  $H_6^{-1}$  yields

$$0 = H_6^{-1} \bar{Q}_{\mathcal{V}-\mathcal{C}}^T A_{\mathcal{V}}^T H_1^{-1} A_C \frac{d}{dt} \iota_C + H_6^{-1} \bar{Q}_{\mathcal{V}-\mathcal{C}}^T \ddot{\mathcal{V}}_s,$$

and by inserting  $A_C \frac{d}{dt} \iota_C$  from (19c) we get  $\frac{d}{dt} \iota_2$  as function of  $\iota_C, \eta, \dot{\mathcal{V}}_s, \ddot{\mathcal{V}}_s, \dot{I}_s$ . Inserting the obtained expressions for  $\dot{\eta}_1, \dot{\eta}_2, \dot{\eta}_3, \frac{d}{dt} \iota_1$  and  $\frac{d}{dt} \iota_2$  into (19c) yields

$$A_C (\bar{P}_C \frac{d}{dt} \iota_C + \bar{Q}_C \frac{d}{dt} \iota_C) = F(\eta, \iota_C, \dot{\mathcal{V}}_s, \dot{I}_s, \ddot{\mathcal{V}}_s).$$

Finally, multiplication with  $H_7^{-1} A_C^T$  gives

$$\bar{P}_C \frac{d}{dt} \iota_C = H_7^{-1} A_C^T F(\eta, \iota_C, \dot{\mathcal{V}}_s, \dot{I}_s, \ddot{\mathcal{V}}_s),$$

and from the derivatives of (19a) and (21), and inserting the expression for  $\bar{P}_C \frac{d}{dt} \iota_C$  we get

$$\bar{Q}_C \frac{d}{dt} \iota_C = [C A_C^T H_1^{-1} A_C - I] H_7^{-1} A_C^T F(\eta, \iota_C, \dot{\mathcal{V}}_s, \dot{I}_s, \ddot{\mathcal{V}}_s).$$

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□

**Remark 13.** Equation (20c) is a hidden constraint on  $\eta$  (in particular on  $\eta_4$ ) and corresponds to the hidden constraints of the MNA equations (29) in case of the existence of  $\mathcal{L}I$ -cutsets, see [14]. In particular,  $Q_{\mathcal{C}\mathcal{R}\mathcal{V}} = 0$  if and only if the circuit does not contain  $\mathcal{L}I$ -cutsets. Equation (23) poses a hidden constraint on  $\iota_C$  and it holds that  $\bar{Q}_{\mathcal{V}-\mathcal{C}} = 0$  if and only if the circuit contains no  $\mathcal{C}\mathcal{V}$ -loops.

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**Remark 14.** In contrast to the MNA formulation (6) or the MLA formulation (7) the pHF formulation (16) will never be a DAE of d-index  $\nu_d = 0$ , i.e., it will

never take the form of an ODE. Otherwise, the conditions for d-index  $\nu_d = 1$  and  $\nu_d = 2$  are similar to the conditions for the MNA equation (6) with the exception of pure  $\mathcal{C}$ -loops. If a circuit contains no  $\mathcal{LI}$ -cutsets and only pure  $\mathcal{C}$ -loops the MNA equations (6) are of d-index  $\nu_d = 1$  while the pH equations (16) are of d-index  $\nu_d = 2$  (cf. Example 6).

#### 4. Structural Analysis of the Circuit Equations

For the analysis and regularization of DAEs structural approaches are widely used in equation-based modeling environments. In this paper, we focus on the *Signature method* ( $\Sigma$ -method) [6] for the structural analysis of DAEs. Another popular structural method is the *Pantelides algorithm* [5] that is related to the  $\Sigma$ -method, see [6]. First, we will review the basic steps of the  $\Sigma$ -method. More details can be found in [6]. Then, in Sections 4.2, 4.3, 4.4 and 4.5 we apply the  $\Sigma$ -method to the different formulations of the circuit equations derived in the previous sections. Although, the four formulations describe the same dynamical behavior of the circuit, there are small differences in the analytical properties (as e.g. the index of the system) and, as we will see, there may be huge differences when it comes to the applicability of the  $\Sigma$ -method. Some of the formulations are suited for a structural analysis while others are not depending on the topology of the circuit.

##### 4.1. The Signature Method for DAEs

The  $\Sigma$ -method can be applied to regular nonlinear DAEs of arbitrary high order  $p$  of the form

$$F(t, x, \dot{x}, \dots, x^{(p)}) = 0, \quad (26)$$

with  $F : \mathbb{I} \times \mathbb{R}^n \times \dots \times \mathbb{R}^n \rightarrow \mathbb{R}^n$  sufficiently smooth. We denote by  $F_i$  the  $i$ th component of the vector-valued function  $F$  and by  $x_j$  the  $j$ th component of the vector  $x$ . Then, the  $\Sigma$ -method consists of the following steps:



1. Building the *signature matrix*  $\Sigma = [\sigma_{ij}]_{i,j=1,\dots,n}$  with

$$\sigma_{ij} = \begin{cases} \text{highest order of derivative of } x_j \text{ in } F_i, \\ -\infty \text{ if } x_j \text{ does not occur in } F_i. \end{cases}$$

2. Finding a *highest value transversal (HVT)* of  $\Sigma$ , i.e., a *transversal*  $T$  of  $\Sigma$

$$T = \{(1, j_1), (2, j_2), \dots, (n, j_n)\},$$

where  $(j_1, \dots, j_n)$  is a permutation of  $(1, \dots, n)$ , with maximal *value*

$$\text{Val}(T) := \sum_{(i,j) \in T} \sigma_{ij}.$$

3. Computing the *offset vectors*  $c = [c_i]_{i=1,\dots,n}$  and  $d = [d_j]_{j=1,\dots,n}$  with  $c_i \geq 0, d_j \geq 0$  such that

$$d_j - c_i \geq \sigma_{ij} \quad \text{for all } i, j = 1, \dots, n, \quad (27a)$$

$$d_j - c_i = \sigma_{ij} \quad \text{for all } (i, j) \in T. \quad (27b)$$

4. Forming the  $\Sigma$ -*Jacobian*  $\mathfrak{J} = [\mathfrak{J}_{ij}]_{i,j=1,\dots,n}$ , with

$$\mathfrak{J}_{ij} := \begin{cases} \frac{\partial F_i}{\partial x_j^{(\sigma_{ij})}} & \text{if } d_j - c_i = \sigma_{ij}, \\ 0 & \text{otherwise.} \end{cases}$$

5. Building the *reduced derivative array*

$$\mathcal{F}(t, \mathcal{X}) = \begin{bmatrix} F_1(t, x, \dot{x}, \dots, x^{(p)}) \\ \frac{d}{dt} F_1(t, x, \dot{x}, \dots, x^{(p)}) \\ \vdots \\ \frac{d^{c_1}}{dt^{c_1}} F_1(t, x, \dot{x}, \dots, x^{(p)}) \\ \vdots \\ F_n(t, x, \dot{x}, \dots, x^{(p)}) \\ \frac{d}{dt} F_n(t, x, \dot{x}, \dots, x^{(p)}) \\ \vdots \\ \frac{d^{c_n}}{dt^{c_n}} F_n(t, x, \dot{x}, \dots, x^{(p)}) \end{bmatrix} = 0, \quad (28)$$

with

$$\mathcal{X} = \begin{bmatrix} x_1 & \dot{x}_1 & \dots & x_1^{(d_1)} & \dots & x_n & \dot{x}_n & \dots & x_n^{(d_n)} \end{bmatrix}^T.$$

6. Success check: if  $\mathcal{F}(t, \mathcal{X}) = 0$ , considered locally as an algebraic system, has a solution  $(t^*, \mathcal{X}^*) \in \mathbb{I} \times \mathbb{R}^{n + \sum_{i=1}^n d_i}$  and  $\mathfrak{J}$  is nonsingular at  $(t^*, \mathcal{X}^*)$ , then  $(t^*, \mathcal{X}^*)$  is a consistent point and the method succeeds.

335 If the  $\Sigma$ -method succeeds, it allows to determine the *structural index* of the DAE as

$$\nu_S := \max_i c_i + \begin{cases} 0 & \text{if all } d_j > 0, \\ 1 & \text{if some } d_j = 0, \end{cases}$$

and  $\text{Val}(\Sigma)$ , defined as the value of the highest value transversal  $T$ , corresponds to the number of degrees of freedom of the system. We call  $\mathfrak{J}$  the  $\Sigma$ -Jacobian since it is in general not equal to the Jacobian  $\frac{\partial F}{\partial x}$  or  $\frac{\partial F}{\partial \dot{x}}$ , but defined by the offset  
 340 vectors. Note that the success check of the  $\Sigma$ -method is performed locally at a fixed point  $(t^*, \mathcal{X}^*)$ , such that the result may hold only locally in a neighborhood of a consistent point. The HVT defines a mapping of maximal value between variables and equations, but it is usually not uniquely determined. A HVT, as well as the offset vectors, can be computed by solving a linear assignment  
 345 problem (LAP), see [6]. That means,  $\Sigma$  is the matrix of the LAP, where each assignment is specified by a transversal. This LAP (as a special kind of a linear programming problem) also has a dual problem, and the offset vectors  $c$  and  $d$  are the corresponding solutions of the dual problem. Note that the offset vectors  $c$  and  $d$  are not uniquely defined by the conditions (27), since for any feasible  
 350 solution  $c$  and  $d$ , also the vectors  $[c_i + \theta]_i$  and  $[d_j + \theta]_j$  form a solution for any  $\theta > 0$ . However, since there exists a unique element-wise smallest solution of the dual problem, these so-called *canonical offsets* are uniquely determined and independent of the chosen HVT, see [6, Theorem 3.6].

The crucial step in the  $\Sigma$ -method is the success check, i.e., the verification of  
 355 regularity of the  $\Sigma$ -Jacobian at a consistent point. Systems for which the  $\Sigma$ -Jacobian is singular for all points  $(t, \mathcal{X})$  that solve the enlarged system (28), or

systems for which there exists no HVT, are called *structurally singular*. Accordingly, we call systems for which the  $\Sigma$ -method succeeds *structurally regular*.

It has been shown in [6] that the  $\Sigma$ -method works successfully for certain (structured) classes of DAE systems, among others for systems in Hessenberg form including semi-explicit systems of d-index 2 and the equations of motion of constrained multibody systems of d-index 3. For such systems the  $\Sigma$ -method succeeds (locally at a consistent point) with  $\nu_S = \nu_d$ . In general, if the  $\Sigma$ -method succeeds, the structural index gives an upper bound for the d-index of the system, see [6], and the obtained structural information can be used to determine a regularization of the system, see [24]. However, there are also cases where the  $\Sigma$ -method fails, see [8, 9]. In particular, this can happen for circuit equations as we will see in the following sections.

#### 4.2. The Signature Method for the MNA equations

We want to apply the  $\Sigma$ -method to the MNA equations (6). To start with, we consider the following two examples.

**Example 1.** Consider the circuit given in Figure 1 consisting of a current source with source function  $I_s$ , two inductors with inductances  $\mathcal{L}_1$  and  $\mathcal{L}_2$ , a capacitor with capacitance  $\mathcal{C}$  and a resistor with conductance  $\mathcal{G}$ . The directed graph corresponding to the circuit is given in Figure 2.

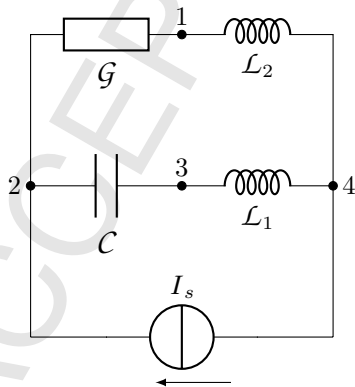


Figure 1:  $\mathcal{LRCI}$ -circuit

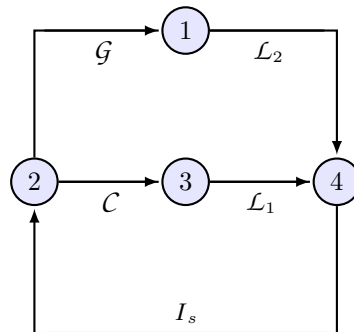


Figure 2: Graph for  $\mathcal{LRCI}$ -circuit

(a) If we select node 1 as reference node, the reduced incidence matrix is given by

$$A = [A_{\mathcal{R}}, A_C, A_L, A_I] = \left[ \begin{array}{c|c|c|c|c} 1 & 1 & 0 & 0 & -1 \\ 0 & -1 & 1 & 0 & 0 \\ 0 & 0 & -1 & -1 & 1 \end{array} \right]$$

and the MNA equations are given by

$$0 = C(\dot{\eta}_2 - \dot{\eta}_3) + G\eta_2 - I_s,$$

$$0 = -C(\dot{\eta}_2 - \dot{\eta}_3) + v_{L_1},$$

$$0 = -v_{L_1} - v_{L_2} + I_s,$$

$$0 = \mathcal{L}_1 i_{L_1} - (\eta_3 - \eta_4),$$

$$0 = \mathcal{L}_2 i_{L_2} + \eta_4,$$

with unknowns  $x = [\eta_2, \eta_3, \eta_4, v_{L_1}, v_{L_2}]^T$ . The signature matrix corresponding to this DAE system is given by

$$\Sigma = \left[ \begin{array}{ccc|cc} \boxed{1} & 1 & - & - & - \\ 1 & \boxed{1} & - & 0 & - \\ - & - & - & \boxed{0} & 0 \\ \hline - & 0 & \boxed{0} & 1 & - \\ - & - & 0 & - & \boxed{1} \end{array} \right],$$

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where one possible HVT is marked by gray boxes. Here, the entry  $-$  stands for  $-\infty$ . The canonical offset vectors are  $c = [0, 0, 1, 0, 0]$  and  $d = [1, 1, 0, 1, 1]$ , and  $\text{Val}(\Sigma) = 3$ . The corresponding  $\Sigma$ -Jacobian is given by

$$\mathfrak{J} = \left[ \begin{array}{ccc|cc} \mathcal{C} & -\mathcal{C} & 0 & 0 & 0 \\ -\mathcal{C} & \mathcal{C} & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & -1 \\ \hline 0 & 0 & 1 & \mathcal{L}_1 & 0 \\ 0 & 0 & 1 & 0 & \mathcal{L}_2 \end{array} \right],$$

and we see immediately that the success check of the  $\Sigma$ -method fails, since  $\mathfrak{J}$  is singular.

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- (b) If we select node 2 as reference node, the reduced incidence matrix is given by

$$A = [A_{\mathcal{R}}, A_C, A_{\mathcal{L}}, A_I] = \left[ \begin{array}{c|c|c|c|c} -1 & 0 & 0 & 1 & 0 \\ 0 & -1 & 1 & 0 & 0 \\ 0 & 0 & -1 & -1 & 1 \end{array} \right]$$

and, in this case, the MNA equations take the form

$$\begin{aligned} 0 &= \mathcal{G}\eta_1 + \iota_{\mathcal{L}_2}, \\ 0 &= \mathcal{C}\dot{\eta}_3 + \iota_{\mathcal{L}_1}, \\ 0 &= -\iota_{\mathcal{L}_1} - \iota_{\mathcal{L}_2} + I_s, \\ 0 &= \mathcal{L}_1 \dot{\iota}_{\mathcal{L}_1} - \eta_3 + \eta_4, \\ 0 &= \mathcal{L}_2 \dot{\iota}_{\mathcal{L}_2} - \eta_1 + \eta_4, \end{aligned}$$

with unknowns  $x = [\eta_1, \eta_3, \eta_4, \iota_{\mathcal{L}_1}, \iota_{\mathcal{L}_2}]^T$ . The signature matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{ccc|cc} 0 & - & - & - & 0 \\ - & 1 & - & 0 & - \\ - & - & - & 0 & 0 \\ \hline - & 0 & 0 & 1 & - \\ 0 & - & 0 & - & 1 \end{array} \right], \quad \mathfrak{J} = \left[ \begin{array}{ccc|cc} \mathcal{G} & 0 & 0 & 0 & 0 \\ 0 & \mathcal{C} & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & -1 \\ \hline 0 & 0 & 1 & \mathcal{L}_1 & 0 \\ -1 & 0 & 1 & 0 & \mathcal{L}_2 \end{array} \right]$$

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with canonical offsets  $c = [0, 0, 1, 0, 0]$  and  $d = [0, 1, 0, 1, 1]$ , and  $\text{Val}(\Sigma) = 2$ . Now, the  $\Sigma$ -method succeeds, since  $\mathfrak{J}$  is regular, and the structural index is determined to be  $\nu_S = 2$ .

- (c) If we select node 3 as reference node, the MNA equations take the form

$$\begin{aligned} 0 &= -\mathcal{G}(-\eta_1 + \eta_2) + \iota_{\mathcal{L}_2}, \\ 0 &= \mathcal{C}\dot{\eta}_2 + \mathcal{G}(-\eta_1 + \eta_2) - I_s, \\ 0 &= -\iota_{\mathcal{L}_1} - \iota_{\mathcal{L}_2} + I_s, \\ 0 &= \mathcal{L}_1 \dot{\iota}_{\mathcal{L}_1} + \eta_4, \\ 0 &= \mathcal{L}_2 \dot{\iota}_{\mathcal{L}_2} - \eta_1 + \eta_4, \end{aligned}$$

with unknowns  $x = [\eta_1, \eta_2, \eta_4, \iota_{L_1}, \iota_{L_2}]^T$ . The signature matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{ccc|cc} 0 & 0 & - & - & 0 \\ 0 & 1 & - & - & - \\ - & - & - & 0 & 0 \\ \hline - & - & 0 & 1 & - \\ 0 & - & 0 & - & 1 \end{array} \right], \quad \mathfrak{J} = \left[ \begin{array}{ccc|cc} \mathcal{G} & 0 & 0 & 0 & 0 \\ -\mathcal{G} & \mathcal{C} & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & -1 \\ \hline 0 & 0 & 1 & \mathcal{L}_1 & 0 \\ -1 & 0 & 1 & 0 & \mathcal{L}_2 \end{array} \right]$$

395 with canonical offsets  $c = [0, 0, 1, 0, 0]$  and  $d = [0, 1, 0, 1, 1]$  and  $\text{Val}(\Sigma) = 2$ . Again, the  $\Sigma$ -method succeeds with  $\nu_S = 2$ .

(d) If we select node 4 as reference node, the MNA equations take the form

$$\begin{aligned} 0 &= -\mathcal{G}(-\eta_1 + \eta_2) + \iota_{L_2}, \\ 0 &= \mathcal{C}(\dot{\eta}_2 - \dot{\eta}_3) + \mathcal{G}(-\eta_1 + \eta_2) - I_s, \\ 0 &= -\mathcal{C}(\dot{\eta}_2 - \dot{\eta}_3) + \iota_{L_1}, \\ 0 &= \mathcal{L}_1 \dot{\iota}_{L_1} - \eta_3, \\ 0 &= \mathcal{L}_2 \dot{\iota}_{L_2} - \eta_1, \end{aligned}$$

with unknowns  $x = [\eta_1, \eta_2, \eta_3, \iota_{L_1}, \iota_{L_2}]^T$ . The signature matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{ccc|cc} 0 & 0 & - & - & 0 \\ 0 & 1 & 1 & - & - \\ - & 1 & 1 & 0 & - \\ \hline - & - & 0 & 1 & - \\ 0 & - & - & - & 1 \end{array} \right], \quad \mathfrak{J} = \left[ \begin{array}{ccc|cc} \mathcal{G} & 0 & 0 & 0 & 0 \\ -\mathcal{G} & \mathcal{C} & -\mathcal{C} & 0 & 0 \\ 0 & -\mathcal{C} & \mathcal{C} & 0 & 0 \\ \hline 0 & 0 & 0 & \mathcal{L}_1 & 0 \\ -1 & 0 & 0 & 0 & \mathcal{L}_2 \end{array} \right]$$

with canonical offsets  $c = [0, 0, 0, 0, 0]$  and  $d = [0, 1, 1, 1, 1]$  and  $\text{Val}(\Sigma) = 4$ .

400 In this case, the success check fails due to the singularity of the  $\Sigma$ -Jacobian.

◁

In Example 1 we can observe that failure or success of the  $\Sigma$ -method can depend on the selection of the reference node. If we select node 1 or node 4, the  $\Sigma$ -method fails due to a singular  $\Sigma$ -Jacobian, if we select node 2 or node 3, the

405  $\Sigma$ -method succeeds with  $\nu_S = 2$ . Note that due to the occurrence of an  $\mathcal{LI}$ -cutset in the circuit given in Example 1, the MNA equations have d-index  $\nu_d = 2$  and this analytical property is independent of the chosen reference node.

**Example 2.** Consider the circuit given in Figure 3 with corresponding circuit graph given in Figure 4. If we select node 1 as reference node, we get the reduced

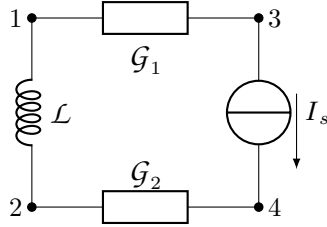


Figure 3:  $\mathcal{RLI}$ -circuit

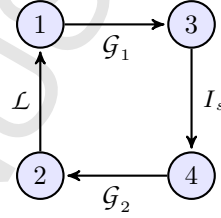


Figure 4: Graph for  $\mathcal{RLI}$ -circuit

410 incidence matrix

$$A = [A_{\mathcal{R}}, A_{\mathcal{L}}, A_I] = \left[ \begin{array}{cc|c|c} 0 & -1 & 1 & 0 \\ -1 & 0 & 0 & 1 \\ 0 & 1 & 0 & -1 \end{array} \right]$$

and the MNA equations take the form

$$0 = -\mathcal{G}_2(-\eta_2 + \eta_4) + \iota_{\mathcal{L}},$$

$$0 = \mathcal{G}_1\eta_3 + I_s,$$

$$0 = \mathcal{G}_2(-\eta_2 + \eta_4) - I_s,$$

$$0 = \mathcal{L}i_{\mathcal{L}} - \eta_2.$$

The corresponding signature matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{ccc|c} 0 & - & 0 & 0 \\ - & 0 & - & - \\ 0 & - & 0 & - \\ \hline 0 & - & - & 1 \end{array} \right], \quad \mathfrak{J} = \left[ \begin{array}{ccc|c} \mathcal{G}_2 & 0 & -\mathcal{G}_2 & 0 \\ 0 & \mathcal{G}_1 & 0 & 0 \\ -\mathcal{G}_2 & 0 & \mathcal{G}_2 & 0 \\ \hline -1 & 0 & 0 & \mathcal{L} \end{array} \right]$$

with canonical offsets  $c = [0, 0, 0, 0]$  and  $d = [0, 0, 0, 1]$  and  $\text{Val}(\Sigma) = 1$ . The success check of the  $\Sigma$ -method fails due to singularity of the  $\Sigma$ -Jacobian. If we

select node 2,3 or 4 as reference node, the  $\Sigma$ -method also fails due to singularity  
of  $\mathfrak{J}$ . ◁

Example 2 shows that there are cases where the  $\Sigma$ -method will always fail, independently of the selection of the reference node. Note again, that the circuit given in Example 2 contains an  $\mathcal{LI}$ -cutset and, thus, the MNA equations have d-index  $\nu_d = 2$ .

In order to explain the problems that show up in the examples we first rewrite the MNA equations (6) in the following form

$$\begin{aligned} A_C C(A_C^T \eta) A_C^T \frac{d}{dt} \eta + A_L \iota_L + A_{\mathcal{R}} g(A_{\mathcal{R}}^T \eta) + A_{\mathcal{V}} \iota_{\mathcal{V}} + A_I I_s &= 0, \\ \mathcal{L}(\iota_L) \frac{d}{dt} \iota_L - A_L^T \eta &= 0, \\ A_{\mathcal{V}}^T \eta - \mathcal{V}_s &= 0, \end{aligned} \quad (29)$$

using the definition of the Jacobians  $\mathcal{C}(\nu_C)$  and  $\mathcal{L}(\iota_L)$  in **(A3)** and **(A4)**. For such a system, we get a signature matrix of the form

$$\Sigma = \begin{bmatrix} \Sigma_{\mathfrak{C}\mathfrak{S}} & \Sigma_{A_L} & \Sigma_{A_{\mathcal{V}}} \\ \Sigma_{A_L}^T & \Sigma_{\mathcal{L}} & - \\ \Sigma_{A_{\mathcal{V}}}^T & - & - \end{bmatrix}, \quad (30)$$

where  $\Sigma_{\mathfrak{C}\mathfrak{S}} = \Sigma_{\mathfrak{C}\mathfrak{S}}^T$  is of size  $n_\eta - 1 \times n_\eta - 1$  with entries in  $\{-\infty, 0, 1\}$ ,  $\Sigma_{\mathcal{L}} = \Sigma_{\mathcal{L}}^T$  is of size  $n_L \times n_L$  with entries in  $\{-\infty, 0, 1\}$ ,  $\Sigma_{A_L}$  is of size  $n_\eta - 1 \times n_L$  with entries in  $\{-\infty, 0\}$  and  $\Sigma_{A_{\mathcal{V}}}$  is of size  $n_\eta - 1 \times n_{\mathcal{V}}$  with entries in  $\{-\infty, 0\}$ . The corresponding  $\Sigma$ -Jacobian has a block structure according to (30) with

$$\mathfrak{J} = \begin{bmatrix} \mathfrak{J}_{11} & \mathfrak{J}_{12} & \mathfrak{J}_{13} \\ \mathfrak{J}_{21} & \mathfrak{J}_{22} & 0 \\ \mathfrak{J}_{31} & 0 & 0 \end{bmatrix}. \quad (31)$$

Note that due to the symmetry of  $\Sigma$ , also the positions of the HVT will be symmetric on  $\Sigma$ , i.e., if  $(i, j) \in T$ , then also  $(j, i) \in T$ . In contrast to (30) the  $\Sigma$ -Jacobian (31) is not symmetric. If  $n_L = 0$  or  $n_{\mathcal{V}} = 0$ , then some blocks in (30)



and (31) may be void. Furthermore, under assumption **(A4)** we know that

$$\Sigma_{\mathcal{L}} = \begin{bmatrix} 1 & & \leq 1 \\ & \ddots & \\ \leq 1 & & 1 \end{bmatrix},$$

420 i.e., there are ones on the diagonal and entries  $\leq 1$  (i.e.,  $-\infty$ , 0 or 1) on the off-diagonal of  $\Sigma_{\mathcal{L}}$ .

*Case 1: no capacitances and no voltage sources.* At first, we assume that  $n_c = 0$  and  $n_v = 0$  (cf. Example 2). In this case, the MNA equations (29) take the form

$$\begin{aligned} A_{\mathcal{R}} g(A_{\mathcal{R}}^T \eta) + A_{\mathcal{L}} \iota_{\mathcal{L}} + A_I I_s &= 0, \\ \mathcal{L}(\iota_{\mathcal{L}}) \frac{d}{dt} \iota_{\mathcal{L}} - A_{\mathcal{L}}^T \eta &= 0, \end{aligned} \quad (32)$$

and the signature matrix (30) reduces to

$$\Sigma = \begin{bmatrix} \Sigma_{\mathfrak{G}} & \Sigma_{A_{\mathcal{L}}} \\ \Sigma_{A_{\mathcal{L}}}^T & \Sigma_{\mathcal{L}} \end{bmatrix}, \quad (33)$$

with  $\Sigma_{\mathfrak{G}} = \Sigma_{\mathfrak{G}}^T$  of size  $n_{\eta} - 1 \times n_{\eta} - 1$  and entries in  $\{-\infty, 0\}$ .

If we assume that we can find a HVT on the diagonal of (33), we get the canonical offsets  $c = [0, \dots, 0]$  and  $d = [0, \dots, 0, 1, \dots, 1]$  and the  $\Sigma$ -Jacobian

$$\mathfrak{J} = \begin{bmatrix} A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T & 0 \\ -A_{\mathcal{L}}^T & \mathcal{L} \end{bmatrix}.$$

425 Thus,  $\mathfrak{J}$  is nonsingular if and only if  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  is nonsingular, and in this case the  $\Sigma$ -method will succeed with structural index  $\nu_S = 1$ . Due to **(A5)** the matrix  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  is nonsingular if and only if  $A_{\mathcal{R}}^T$  has full column rank. In particular, the following statements are equivalent (see [14, 16]):

- $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  is nonsingular;

- 430
- $\ker A_{\mathcal{R}}^T = \{0\}$ ;

- the graph corresponding to the circuit contains no  $\mathcal{CLVI}$ -cutsets.<sup>1</sup>

Note that  $\mathcal{CLVI}$ -cutsets include  $\mathcal{LI}$ -cutsets as special case. Thus, the  $\Sigma$ -method will fail for systems of the form (32) and HVT on the diagonal whenever  $\nu_d > 1$ . But also pure  $I$ -cutsets (which are excluded by Assumption **(A2)**) or  
 435 pure  $\mathcal{L}$ -cutsets will lead to failure of the  $\Sigma$ -method in this setting.

If there exists no HVT on the diagonal (but the system is structurally well-posed), then necessarily  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  contains zero rows/columns. In this case, we can reorder the rows of  $A_{\mathcal{R}}$  such that

$$\Pi_{\mathcal{R}} A_{\mathcal{R}} = \begin{bmatrix} \tilde{A}_{\mathcal{R}} \\ 0 \end{bmatrix} \begin{matrix} \tilde{r} \\ n_{\eta} - 1 - \tilde{r} \end{matrix}$$

where  $\Pi_{\mathcal{R}}$  is a permutation such that  $\tilde{A}_{\mathcal{R}}$  contains no zero rows. Then, the corresponding permutation of (32) yields

$$\begin{aligned} \begin{bmatrix} \tilde{A}_{\mathcal{R}} \\ 0 \end{bmatrix} g(\tilde{A}_{\mathcal{R}}^T \tilde{\eta}_1) + \begin{bmatrix} \tilde{A}_{\mathcal{L},1} \\ \tilde{A}_{\mathcal{L},2} \end{bmatrix} \iota_{\mathcal{L}} + \begin{bmatrix} \tilde{A}_{I,1} \\ \tilde{A}_{I,2} \end{bmatrix} I_s = 0, \\ \mathcal{L}(\iota_{\mathcal{L}}) \frac{d}{dt} \iota_{\mathcal{L}} - \begin{bmatrix} \tilde{A}_{\mathcal{L},1}^T & \tilde{A}_{\mathcal{L},2}^T \end{bmatrix} \begin{bmatrix} \tilde{\eta}_1 \\ \tilde{\eta}_2 \end{bmatrix} = 0, \end{aligned} \quad (34)$$

440 where

$$\Pi_{\mathcal{R}} \eta =: \begin{bmatrix} \tilde{\eta}_1 \\ \tilde{\eta}_2 \end{bmatrix}, \quad \Pi_{\mathcal{R}} A_{\mathcal{L}} =: \begin{bmatrix} \tilde{A}_{\mathcal{L},1} \\ \tilde{A}_{\mathcal{L},2} \end{bmatrix} \quad \Pi_{\mathcal{R}} A_I =: \begin{bmatrix} \tilde{A}_{I,1} \\ \tilde{A}_{I,2} \end{bmatrix}.$$

The  $\Sigma$ -matrix for this permuted system takes the form

$$\Sigma = \begin{bmatrix} \tilde{\Sigma}_{\mathfrak{G}} & - & \Sigma_{\tilde{A}_{\mathcal{L},1}} \\ - & - & \Sigma_{\tilde{A}_{\mathcal{L},2}} \\ \Sigma_{\tilde{A}_{\mathcal{L},1}}^T & \Sigma_{\tilde{A}_{\mathcal{L},2}}^T & \Sigma_{\mathcal{L}} \end{bmatrix} \begin{matrix} \tilde{r} \\ n_{\eta} - 1 - \tilde{r} \\ n_{\mathcal{L}} \end{matrix}$$

with  $\tilde{\Sigma}_{\mathfrak{G}} = \tilde{\Sigma}_{\mathfrak{G}}^T$  of size  $\tilde{r} \times \tilde{r}$  and entries in  $\{-\infty, 0\}$  having zero-entries on the diagonal (which can be assumed w.l.o.g due to Assumption **(A5)**). Since the

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<sup>1</sup>Note that  $A_{\mathcal{R}}^T$  can only be of full column rank if  $n_{\eta} - 1 \leq n_{\mathcal{R}}$ . If  $n_{\eta} - 1 > n_{\mathcal{R}}$ , then there will be  $\mathcal{CLVI}$ -cutsets in the circuit.

system is structurally well-posed, there must be  $\tilde{s} := n_\eta - 1 - \tilde{r}$  positions for the HVT be taken from the block  $\Sigma_{\tilde{A}_{L,2}}$ , which can w.l.o.g. assumed to be positioned on the diagonal of the first  $\tilde{s}$ -by- $\tilde{s}$  block of  $\Sigma_{\tilde{A}_{L,2}}$ . This results in a  $\Sigma$ -matrix of the form

$$\left[ \begin{array}{cc|cc} \boxed{0} & & \leq 0 & \\ & \ddots & & \\ \leq 0 & & \boxed{0} & \\ \hline & & & \boxed{0} \\ & - & & \ddots \\ & & & \boxed{0} \\ \hline & \boxed{0} & 1 & \\ & & \ddots & \leq 1 \\ \leq 0 & & \boxed{0} & 1 \\ & & & \boxed{1} \\ & \leq 0 & \leq 1 & \ddots \\ & & & \boxed{1} \end{array} \right]$$

with canonical offsets

$$c = \left[ 0, \dots, 0 \mid 1, \dots, 1 \mid 0, \dots, 0 \right], \quad d = \left[ 0, \dots, 0 \mid 0, \dots, 0 \mid 1, \dots, 1 \right],$$

and the  $\Sigma$ -Jacobian takes the form

$$\mathfrak{J} = \begin{bmatrix} \tilde{A}_{\mathcal{R}} \mathcal{G} \tilde{A}_{\mathcal{R}}^T & 0 & 0 \\ 0 & 0 & -\tilde{A}_{L,2} \\ -\tilde{A}_{L,1}^T & -\tilde{A}_{L,2}^T & \mathcal{L} \end{bmatrix}.$$

We see that  $\mathfrak{J}$  can only be nonsingular if  $\tilde{A}_{\mathcal{R}} \mathcal{G} \tilde{A}_{\mathcal{R}}^T$  is nonsingular. As before,  $\tilde{A}_{\mathcal{R}} \mathcal{G} \tilde{A}_{\mathcal{R}}^T$  is nonsingular if and only if  $\ker \tilde{A}_{\mathcal{R}}^T = \{0\}$  or, equivalently, if the subgraph consisting of the nodes corresponding to  $\tilde{A}_{\mathcal{R}}$  contains no  $\mathcal{CLVI}$ -cutsets.

*Case 2: no resistors and no voltage sources.* Next we assume that  $n_{\mathcal{R}} = 0$  and  $n_{\mathcal{V}} = 0$ . In this case, the MNA equations (29) take the form

$$\begin{aligned} A_C C (A_C^T \eta) A_C^T \frac{d}{dt} \eta + A_L i_L + A_I I_s &= 0, \\ \mathcal{L}(i_L) \frac{d}{dt} i_L - A_L^T \eta &= 0, \end{aligned} \quad (35)$$

and the signature matrix (30) reduces to

$$\Sigma = \begin{bmatrix} \Sigma_{\mathfrak{C}} & \Sigma_{A_L} \\ \Sigma_{A_L}^T & \Sigma_{\mathcal{L}} \end{bmatrix},$$

445 where  $\Sigma_{\mathfrak{C}} = \Sigma_{\mathfrak{C}}^T$  is of size  $n_{\eta} - 1 \times n_{\eta} - 1$  with entries in  $\{-\infty, 0, 1\}$ . Assuming that there is a HVT on the diagonal we get

$$\Sigma = \left[ \begin{array}{c|c} \begin{matrix} \boxed{1} & & \\ & \ddots & \\ & & \boxed{1} \end{matrix} & \begin{matrix} & & \leq 0 \end{matrix} \\ \hline \begin{matrix} & & \leq 0 \end{matrix} & \begin{matrix} \boxed{1} & & \\ & \ddots & \\ & & \boxed{1} \end{matrix} \end{array} \right],$$

with  $c = [0, \dots, 0]$ ,  $d = [1, \dots, 1]$  and the  $\Sigma$ -Jacobian takes the form

$$\mathfrak{J} = \begin{bmatrix} A_C C A_C^T & 0 \\ 0 & \mathcal{L} \end{bmatrix}.$$

We see that  $\mathfrak{J}$  is nonsingular if and only if  $A_C C A_C^T$  is nonsingular. Since  $C$  is assumed to be symmetric positive definite,  $A_C C A_C^T$  will be nonsingular if and only if  $A_C^T$  has full column rank. Similar as before the following statements are  
450 equivalent (see [14, 16]):

- $A_C C A_C^T$  is nonsingular;
- $\ker A_C^T = \{0\}$ ;
- the graph corresponding to the circuit contains no  $\mathcal{RLV}$ -cutsets.<sup>2</sup>

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<sup>2</sup>Similar as before,  $A_C^T$  can only be of full column rank if  $n_{\eta} - 1 \leq n_C$ .

Again,  $\mathcal{LI}$ -cutsets are included as special case. Thus, the  $\Sigma$ -method will fail for  
 455 the MNA equations (35) whenever  $\nu_d > 1$ . But, again, also pure  $I$ -cutsets or  
 pure  $\mathcal{L}$ -cutsets will lead to a singular  $\Sigma$ -Jacobian in this setting.

If there is no HVT on the diagonal, then  $A_C C A_C^T$  contains zero rows/columns  
 and we can proceed in a similar manner as before using permutations of the  
 system such that

$$\Pi_C A_C = \begin{bmatrix} \tilde{A}_C \\ 0 \end{bmatrix}. \quad (36)$$

In this case the success check will fail whenever there are  $\mathcal{RLV}I$ -cutsets (in  
 particular  $\mathcal{LI}$ -cutsets) in the subgraph consisting of the nodes corresponding to  
 $\tilde{A}_C$ .

460 **Remark 15.** We have seen that  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  and  $A_C C A_C^T$  can be singular depend-  
 ing on the topology of the circuit. However, from Example 1 (cases (b) and  
 (c)) we can observe that  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  and  $A_C C A_C^T$  can both be singular and nev-  
 ertheless the  $\Sigma$ -method can succeed with  $\nu_S = \nu_d$ , even for higher index DAEs.  
 On the other hand, singular blocks in the  $\Sigma$ -Jacobian can also result from a  
 465 combination of the matrices  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  and  $A_C C A_C^T$  as can be seen in Example  
 1, case (d).

The previous discussion might suggest that the failure of the  $\Sigma$ -method is related  
 to the occurrence of  $\mathcal{LI}$ -cutsets. That this is not the case can be seen in the  
 following example.

**Example 3.** Consider the circuit given in Figure 5 with corresponding directed  
 graph given in Figure 6. If we select node 3 as reference node, the MNA equa-  
 tions take the form

$$\begin{bmatrix} C_1 + C_2 & -C_2 & -C_1 \\ -C_2 & C_2 & 0 \\ -C_1 & 0 & C_1 \end{bmatrix} \begin{bmatrix} \dot{\eta}_1 \\ \dot{\eta}_2 \\ \dot{\eta}_4 \end{bmatrix} + \begin{bmatrix} \mathcal{G}_1 & 0 & 0 \\ 0 & \mathcal{G}_2 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \eta_1 \\ \eta_2 \\ \eta_4 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ -1 \end{bmatrix} I_s = 0,$$

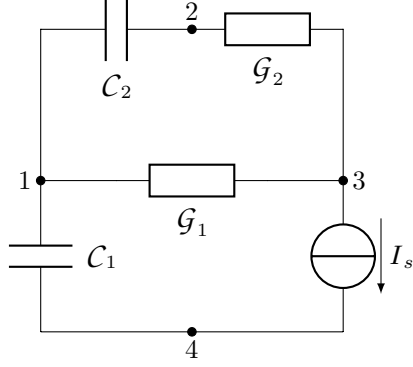


Figure 5:  $\mathcal{RCI}$ -circuit

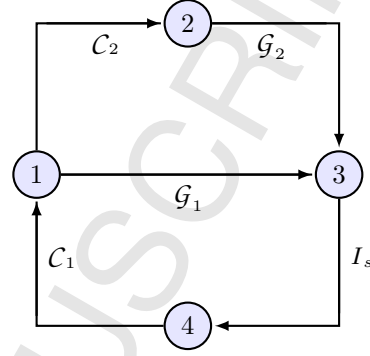


Figure 6: Graph for  $\mathcal{RCI}$ -circuit

470 and the signature matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \begin{bmatrix} \boxed{1} & 1 & 1 \\ 1 & \boxed{1} & - \\ 1 & - & \boxed{1} \end{bmatrix}, \quad \mathfrak{J} = \begin{bmatrix} C_1 + C_2 & -C_2 & -C_1 \\ -C_2 & C_2 & 0 \\ -C_1 & 0 & C_1 \end{bmatrix}.$$

The success check fails, since  $\mathfrak{J}$  is singular. If we choose one of the other nodes as reference node, then the  $\Sigma$ -method succeeds with structural index  $\nu_S = 1 = \nu_d$ .

◁

In Example 3 we have  $n_q = 0$  and  $n_L = 0$  such that (30) reduces to  $\Sigma = \Sigma_{\mathcal{CG}}$ .

475 Now, if  $A_C C A_C^T$  is singular (in Example 3 due to the occurrence of an  $\mathcal{RI}$ -cutset), and as reference a node is chosen that is not adjacent to a capacitor, we get  $\mathfrak{J} = A_C C A_C^T$  and the  $\Sigma$ -method fails. Example 3 shows that the  $\Sigma$ -method can fail also for MNA equations (29) of d-index  $\nu_d = 1$ .

We might consider more general forms of the MNA equations (29), but still, combining the ideas of Case 1 and Case 2 above, singular blocks of the form  $\tilde{A}_R G \tilde{A}_R^T$  or  $\tilde{A}_C C \tilde{A}_C^T$  can occur, corresponding to the occurrence of  $\mathcal{CLV}I$ -cutsets or  $\mathcal{RLV}I$ -cutsets in certain (sub)graphs. A complete characterization of all circuit configurations that lead to failure of the  $\Sigma$ -method for the MNA equations (29) has proven to be quite tedious. Since we will see later on that the  $\Sigma$ -method will always succeed for another formulation of the circuit equations  
485 we refrain from presenting a more detailed discussion of circuit topologies for

which the  $\Sigma$ -method will always give the right structural information at this point. At least we can formulate the following result.

**Theorem 16.** *Consider the system of MNA equations (29) and assume that (A1)-(A5) hold.*

1. *If the system (29) is of d-index  $\nu_d = 0$ , then the  $\Sigma$ -method succeeds with  $c = [0, \dots, 0]$ ,  $d = [1, \dots, 1]$  and structural index  $\nu_S = 0$ .*
2. *If  $n_c = 0$  and  $n_{\nu} = 0$  and there are no  $\mathcal{LI}$ -cutsets (including pure  $\mathcal{L}$ -cutsets) in the circuit, then the  $\Sigma$ -method for (29) succeeds with structural index  $\nu_S = 1$ .*
3. *If  $n_{\mathcal{R}} = 0$  and  $n_{\nu} = 0$  and there are no  $\mathcal{LI}$ -cutsets (including pure  $\mathcal{L}$ -cutsets) in the circuit, then the  $\Sigma$ -method for (29) succeeds with structural index  $\nu_S = 1$ .*

*Proof.* If the MNA equations (29) have d-index  $\nu_d = 0$ , then from Theorem 3 we know that  $n_{\nu} = 0$  and  $\text{rank } A_C = n_{\eta} - 1$ . Thus, the matrix  $A_C C A_C^T$  is regular and, consequently, also the  $\Sigma$ -Jacobian

$$\mathfrak{J} = \begin{bmatrix} A_C C A_C^T & 0 \\ 0 & \mathcal{L} \end{bmatrix}$$

is regular. As a result the  $\Sigma$ -method succeeds with  $c = [0, \dots, 0]$  and  $d = [1, \dots, 1]$ , such that  $\nu_S = \max_i c_i = 0$ . The other statements follow directly from the previous discussion.  $\square$

In the previous examples we have observed that failure or success of the  $\Sigma$ -method for the MNA equations (29) can depend on the selection of the reference node. However, in some cases the  $\Sigma$ -method always fails independently of the chosen reference (see Example 2, where  $A_{\mathcal{R}} G A_{\mathcal{R}}^T$  is a singular block in  $\mathfrak{J}$  for all possible choices of the reference node). Note again, that the selection of the reference node does not influence the analytical properties of the system (e.g. the index or the hidden constraints), and, in particular, the regularity of the matrices  $A_C C A_C^T$  and  $A_{\mathcal{R}} G A_{\mathcal{R}}^T$  is independent of the selection of the reference

510 node. Can we nevertheless give a characterization for a “good” choice of a reference node?

Consider  $Y = [y_{ij}] = A_K \mathcal{K} A_K^T$  with  $\mathcal{K} = \text{diag}(\mathcal{K}_1, \dots, \mathcal{K}_\ell) > 0$  and incidence matrix  $A_K$ . Then, we have  $y_{kk} = \sum_{j=1}^{\ell} a_{kj}^2 \mathcal{K}_j$  with

$$a_{kj}^2 = \begin{cases} 1 & \text{if branch } j \text{ is adjacent to node } k, \\ 0 & \text{else,} \end{cases}$$

i.e., for all  $k$  the diagonal entry  $y_{kk}$  of  $Y$  is the sum of the weights  $\mathcal{K}_i$ 's of all branches that are adjacent to node  $k$ . Furthermore, we have  $y_{ki} = \sum_{j=1}^{\ell} a_{kj} a_{ij} \mathcal{K}_j$ , with

$$a_{kj} a_{ij} = \begin{cases} -1 & \text{if branch } j \text{ connects the nodes } k \text{ and } i, \\ 0 & \text{else,} \end{cases}$$

for  $k \neq i$ , i.e., each off-diagonal entry  $y_{ki}$  of  $Y$  is the negative sum of the weights  $\mathcal{K}_i$ 's of all branches that connect node  $k$  with node  $i$ . If we replace  $Y$  by  $A_C C A_C^T$  or  $A_{\mathcal{R}} G A_{\mathcal{R}}^T$ , respectively, we see that the characteristic values of components  
 515 that are connected to the reference node appear only on the diagonal of  $Y$ , while the characteristic values of elements that are not connected to the reference node appear both on the diagonal and off-diagonal terms. In particular,  $A_C C A_C^T$  is only nonsingular if there exists a capacitive tree in the circuit graph. These observations (and numerous examples) show that if there exists no tree in the  
 520 circuit graph that consists only of capacitive and resistive branches, then as reference a node that is adjacent to a capacitance or, if no capacitances are present, adjacent to a resistance should be chosen. In this case, if the reference node is adjacent to a capacitance,  $A_C C A_C^T$  contains at least one zero-row and by permutation with  $\Pi_C$  as in (36) this row is removed from  $\tilde{A}_C$ . However, this  
 525 is no guarantee for the success of the  $\Sigma$ -method as Example 2 shows, since the remaining part  $\tilde{A}_C \tilde{C} \tilde{A}_C^T$  can still be singular.

A simple check for failure of the  $\Sigma$ -method in the case that  $n_v = 0$  can be



performed as follows. We permute  $[A_C, A_{\mathcal{R}}]$  as

$$\Pi[A_C, A_{\mathcal{R}}] = \begin{bmatrix} \tilde{A}_C & \tilde{A}_{\mathcal{R},1} & \tilde{n}_1 \\ 0 & \tilde{A}_{\mathcal{R},2} & \tilde{n}_2 \\ 0 & 0 & n_{\eta} - 1 - \tilde{n}_1 - \tilde{n}_2 \end{bmatrix}$$

where  $\Pi$  is a permutation,  $\tilde{A}_C$  is of size  $\tilde{n}_1 \times n_C$ , with  $\tilde{n}_1$  minimal,  $\tilde{A}_{\mathcal{R},1}$  is of size  $\tilde{n}_1 \times n_{\mathcal{R}}$ , and  $\tilde{A}_{\mathcal{R},2}$  is of size  $\tilde{n}_2 \times n_{\mathcal{R}}$  with  $\tilde{n}_2$  minimal. The  $\tilde{n}_1$  nodes corresponding to the first block row are the nodes that are directly connected to a capacitance, the  $\tilde{n}_2$  nodes corresponding to the second block row are the nodes that are directly connected to a resistance, but not connected to a capacitance, and the  $n_{\eta} - 1 - \tilde{n}_1 - \tilde{n}_2$  nodes corresponding to the last block row are the nodes that are neither connected to a capacitance nor connected to a resistance. Then, if  $n_{\nu} = 0$ , the  $\Sigma$ -method will fail if  $\tilde{A}_C C \tilde{A}_C^T$  or  $\tilde{A}_{\mathcal{R},2} \mathcal{G} \tilde{A}_{\mathcal{R},2}^T$  is singular, i.e., if

$$\text{rank } \tilde{A}_C < \tilde{n}_1 \quad \text{or} \quad \text{rank } \tilde{A}_{\mathcal{R},2} < \tilde{n}_2.$$

#### 4.3. The Signature Method for the MLA equations

530 In this section we apply the  $\Sigma$ -method to the MLA equations (7). Again, we start by considering the two examples from Section 4.2.

**Example 4.** We consider again the circuit given in Example 1. The oriented loops contained in the graph are depicted in Figure 7. A reduced loop matrix

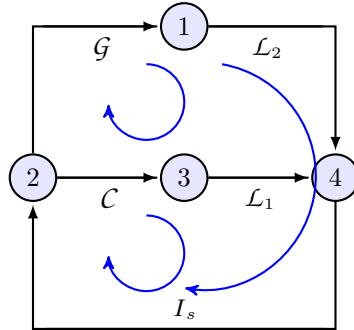


Figure 7: Loops in the  $\mathcal{LRCT}$ -circuit

is given by

$$B = \left[ \begin{array}{c|c|c|c|c} 1 & -1 & -1 & 1 & 0 \\ 0 & 1 & 1 & 0 & 1 \end{array} \right] = \begin{bmatrix} B_{\mathcal{R}} & B_C & B_L & B_I \end{bmatrix}$$

and the MLA equations (7) take the form

$$\begin{aligned} (\mathcal{L}_1 + \mathcal{L}_2) \frac{d}{dt} j_1 - \mathcal{L}_1 \frac{d}{dt} j_2 + \mathcal{R} j_1 - \nu_C &= 0, \\ -\mathcal{L}_1 \frac{d}{dt} j_1 + \mathcal{L}_1 \frac{d}{dt} j_2 + \nu_C + \nu_I &= 0, \\ C \frac{d}{dt} \nu_C + j_1 - j_2 &= 0, \\ j_2 + I_s &= 0. \end{aligned}$$

Thus, the  $\Sigma$ -matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{c|c|c|c} 1 & 1 & 0 & - \\ 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & - \\ - & 0 & - & - \end{array} \right], \quad \mathfrak{J} = \left[ \begin{array}{c|c|c|c} \mathcal{L}_1 + \mathcal{L}_2 & -\mathcal{L}_1 & 0 & 0 \\ -\mathcal{L}_1 & \mathcal{L}_1 & 0 & 1 \\ 0 & 0 & C & 0 \\ 0 & 1 & 0 & 0 \end{array} \right].$$

We see that  $\mathfrak{J}$  is nonsingular and, thus, the  $\Sigma$ -method succeeds with  $\nu_S = 2 = \nu_d$ .

◁

**Example 5.** We consider again the circuit given in Example 2. In this case, the circuit graph contains only one loop and the loop matrix is given by

$$B_0 = B = \begin{bmatrix} 1 & 1 & 1 & 1 \end{bmatrix}.$$

The MLA equations (7) take the form

$$\begin{aligned} \mathcal{L} \frac{d}{dt} j + (\mathcal{R}_1 + \mathcal{R}_2) j + \nu_I &= 0, \\ j + I_s &= 0, \end{aligned}$$

and the  $\Sigma$ -matrix and  $\Sigma$ -Jacobian for this system are given by

$$\Sigma = \left[ \begin{array}{c|c} 1 & 0 \\ 0 & - \end{array} \right], \quad \mathfrak{J} = \begin{bmatrix} \mathcal{L} & 1 \\ 1 & 0 \end{bmatrix}.$$

We see that  $\mathfrak{J}$  is nonsingular such that the  $\Sigma$ -method succeeds with  $\nu_S = 2 = \nu_d$ .

◁

We see that, if we use the MLA equations (7) instead of the MNA equations (6) to describe the behavior of the electrical circuits given in Example 1 and 2, the  $\Sigma$ -method succeeds with  $\nu_S = 2 = \nu_d$ .  
545

A natural question that arises is: Will the  $\Sigma$ -method always succeed for the MLA equations? Unfortunately, this is not the case. For the MLA equations (7) we get a signature matrix of the form

$$\Sigma = \begin{bmatrix} \Sigma_{\mathcal{L}\mathcal{R}} & \Sigma_{B_C} & \Sigma_{B_I} \\ \Sigma_{B_C}^T & \Sigma_C & - \\ \Sigma_{B_I}^T & - & - \end{bmatrix},$$

where  $\Sigma_{\mathcal{L}\mathcal{R}} = \Sigma_{\mathcal{L}\mathcal{R}}^T$  is of size  $n_b - n_\eta + 1 \times n_b - n_\eta + 1$  with entries in  $\{-\infty, 0, 1\}$ ,  $\Sigma_C = \Sigma_C^T$  is of size  $n_c \times n_c$  with entries in  $\{-\infty, 0, 1\}$ ,  $\Sigma_{B_C}$  is of size  $n_b - n_\eta + 1 \times n_c$  with entries in  $\{-\infty, 0\}$ , and  $\Sigma_{B_I}$  is of size  $n_b - n_\eta + 1 \times n_l$  with entries in  $\{-\infty, 0\}$ . Due to Assumption **(A3)** we know that

$$\Sigma_C = \begin{bmatrix} 1 & & \leq 1 \\ & \ddots & \\ \leq 1 & & 1 \end{bmatrix}$$

with ones on the diagonal and entries  $\leq 1$  on the off-diagonal.

Hence, the signature matrix has a similar structure as the signature matrix (30) for the MNA equations and therefore also similar problems will arise. Analogously as in Section 4.2 Assumption **(A1)** means that the matrix  $[B_L \ B_{\mathcal{R}} \ B_C \ B_I]$  has full row rank. In the same way Assumption **(A2)** means that the matrix  $B_I$  has full column rank. Moreover, the following statements are equivalent (see [14, 16]):  
550

- $B_L \mathcal{L} B_L^T$  is nonsingular;
- $\ker B_L^T = \{0\}$ ;
- the graph corresponding to the circuit does not contain  $\mathcal{CRV}$ -loops.  
555

Analogously, the following statements are equivalent (see [14, 16]):

- $B_{\mathcal{R}}\mathcal{R}B_{\mathcal{R}}^T$  is nonsingular;
- $\ker B_{\mathcal{R}}^T = \{0\}$ ;
- the graph corresponding to the circuit does not contain  $\mathcal{CLVI}$ -loops.

560 Note that  $\mathcal{CRVI}$ -loops and  $\mathcal{CLVI}$ -loops include  $\mathcal{CV}$ -loops (but also pure  $\mathcal{C}$ -loops) as special cases. So again, certain topological configurations lead to singular blocks  $B_L\mathcal{L}B_L^T$  or  $B_{\mathcal{R}}\mathcal{R}B_{\mathcal{R}}^T$ .

In Example 4 the  $\Sigma$ -Jacobian is of the form

$$\mathfrak{J} = \begin{bmatrix} B_L\mathcal{L}B_L^T & 0 & B_I \\ 0 & \mathcal{C} & 0 \\ B_I^T & 0 & 0 \end{bmatrix},$$

and in Example 5 it is of the form

$$\mathfrak{J} = \begin{bmatrix} B_L\mathcal{L}B_L^T & B_I \\ B_I^T & 0 \end{bmatrix},$$

565 and both are nonsingular. However, we can easily construct an example where the  $\Sigma$ -method fails for the MLA equations (7).

**Example 6.** Consider the circuit given in Figure 8. The graph together with

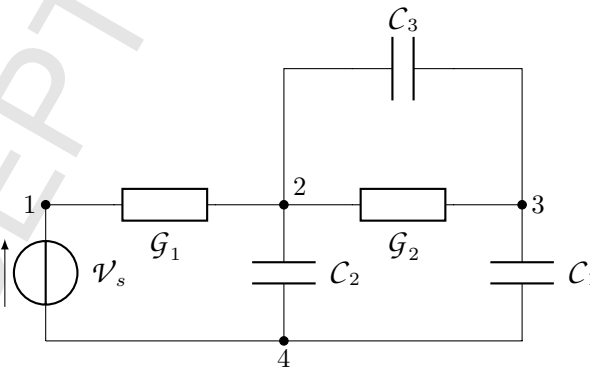


Figure 8:  $\mathcal{RC}$ -circuit

the loops corresponding to the circuit is given in Figure 9. A reduced loop

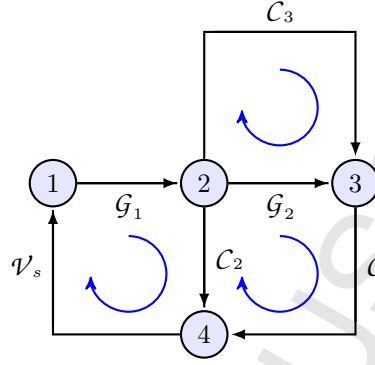


Figure 9: Graph of  $\mathcal{RC}$ -circuit with loops

matrix is given by

$$B = \left[ \begin{array}{cc|cc|c|c} 1 & 0 & 0 & 1 & 0 & 1 \\ 0 & 1 & 1 & -1 & 0 & 0 \\ 0 & -1 & 0 & 0 & 1 & 0 \end{array} \right] = [B_{\mathcal{R}} \quad B_{\mathcal{C}} \quad B_{\mathcal{V}}]$$

and the MLA equations (7) take the form

$$\begin{bmatrix} \mathcal{R}_1 & 0 & 0 \\ 0 & \mathcal{R}_2 & -\mathcal{R}_2 \\ 0 & -\mathcal{R}_2 & \mathcal{R}_2 \end{bmatrix} j + \begin{bmatrix} 0 & 1 & 0 \\ 1 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \nu_{\mathcal{C}} + \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \mathcal{V}_s = 0,$$

$$\begin{bmatrix} \mathcal{C}_1 & 0 & 0 \\ 0 & \mathcal{C}_2 & 0 \\ 0 & 0 & \mathcal{C}_3 \end{bmatrix} \frac{d}{dt} \nu_{\mathcal{C}} - \begin{bmatrix} 0 & 1 & 0 \\ 1 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} j = 0.$$

570 Since the circuit contains a  $\mathcal{C}$ -loop the MLA equations are of d-index  $\nu_d = 2$ .

The corresponding  $\Sigma$ -matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{ccc|ccc} 0 & - & - & - & 0 & - \\ - & 0 & 0 & 0 & 0 & - \\ - & 0 & 0 & - & - & 0 \\ \hline - & 0 & - & 1 & - & - \\ 0 & 0 & - & - & 1 & - \\ - & - & 0 & - & - & 1 \end{array} \right], \quad \mathfrak{J} = \left[ \begin{array}{ccc|ccc} \mathcal{R}_1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \mathcal{R}_2 & -\mathcal{R}_2 & 0 & 0 & 0 \\ 0 & -\mathcal{R}_2 & \mathcal{R}_2 & 0 & 0 & 0 \\ \hline 0 & 1 & 0 & \mathcal{C}_1 & 0 & 0 \\ 1 & -1 & 0 & 0 & \mathcal{C}_2 & 0 \\ 0 & 0 & 1 & 0 & 0 & \mathcal{C}_3 \end{array} \right]$$

and  $\mathfrak{J}$  is singular. ◁

In Example 6 the  $\Sigma$ -Jacobian is of the form

$$\mathfrak{J} = \begin{bmatrix} B_{\mathcal{R}} \mathcal{R} B_{\mathcal{R}}^T & 0 \\ -B_{\mathcal{C}}^T & C \end{bmatrix},$$

and the block  $B_{\mathcal{R}} \mathcal{R} B_{\mathcal{R}}^T$  is singular due to presence of the  $\mathcal{C}$ -loop. Note that, if  
 575 we choose for  $B$  the fundamental loop matrix for the fundamental loops that are  
 defined according to the normal tree  $T = \{C_1, C_2, \mathcal{V}_s\}$  in the graph depicted in  
 Figure 9, then the  $\Sigma$ -method succeeds with nonsingular  $\Sigma$ -Jacobian and  $\nu_S =$   
 $\nu_d = 2$ .

**Remark 17.** If we use the MNA equations (6) to describe the behavior of  
 580 the circuit given in Example 6, then the corresponding DAE system is of d-  
 index  $\nu_d = 1$ , and the  $\Sigma$ -method succeeds for all choices of the reference node.  
 However, for node 1 as reference the  $\Sigma$ -methods determines the structural index  
 $\nu_S = 2 > \nu_d$ , while for all other nodes the  $\Sigma$ -method succeeds with  $\nu_S = \nu_d = 1$ .  
 Both,  $A_C C A_C^T$  and  $A_{\mathcal{R}} \mathcal{G} A_{\mathcal{R}}^T$  are singular (there is a  $\mathcal{C}\mathcal{V}$ -cutset and a  $\mathcal{R}\mathcal{V}$ -  
 585 cutset), however  $\mathfrak{J}$  is regular for all cases.

We get the corresponding result of Theorem 16 for the MLA equations.

**Theorem 18.** *Consider the system of MLA equations (7) and assume that  
 (A1)-(A5) hold.*

1. *If the system (7) is of d-index  $\nu_d = 0$ , then the  $\Sigma$ -method succeeds with  
 590  $c = [0, \dots, 0]$ ,  $d = [1, \dots, 1]$  and structural index  $\nu_S = 0$ .*
2. *If  $n_{\mathcal{L}} = 0$  and  $n_I = 0$  and there are no  $\mathcal{C}\mathcal{V}$ -loops (including pure  $\mathcal{C}$ -loops)  
 in the circuit, then the  $\Sigma$ -method for (7) succeeds with structural index  
 $\nu_S = 1$ .*
3. *If  $n_{\mathcal{R}} = 0$  and  $n_I = 0$  and there are no  $\mathcal{C}\mathcal{V}$ -loops (including pure  $\mathcal{C}$ -loops)  
 595 in the circuit, then the  $\Sigma$ -method for (7) succeeds with structural index  
 $\nu_S = 1$ .*

*Proof.* Analogously to the proof of Theorem 16. ◻

#### 4.4. The Signature Method for the Branch-Oriented Model Equations

In the previous sections we have seen that certain circuit configurations lead to failure of the  $\Sigma$ -method for the MNA equations while other configurations lead to failure of the  $\Sigma$ -method for the MLA equations. For more complex circuit examples one could construct configurations where the  $\Sigma$ -method fails for both formulations. In this section we apply the  $\Sigma$ -method to the branch-oriented model equations (13). The key result is formulated in the following theorem.

**Theorem 19.** Consider the branch-oriented model equations (13) and let Assumptions (A1)-(A5) hold. Then the  $\Sigma$ -method applied to system (13) always succeeds with nonsingular  $\Sigma$ -Jacobian and structural index  $\nu_S = \nu_d$ .

*Proof.* If the DAE (13) is of d-index  $\nu_d = 1$ , the normal tree used for the formulation of (13) is actually a proper tree, such that  $\nu_1 = [\nu_C^T, \nu_{\mathcal{V}}^T, \nu_{\mathcal{R}1}^T]^T$  consists of the branch voltages for all capacitors, voltage sources and twig resistors, and  $\nu_2 = [\nu_L^T, \nu_I^T, \nu_{\mathcal{R}2}^T]^T$  consists of the branch voltages for all inductors, current sources and link resistors. Thus, system (13) reduces to

$$\begin{aligned} C(\nu_C) \frac{d}{dt} \nu_C &= \nu_C, \\ \mathcal{L}(\nu_L) \frac{d}{dt} \nu_L &= \nu_L, \\ 0 &= \begin{bmatrix} \nu_{\mathcal{R}1} \\ \nu_{\mathcal{R}2} \end{bmatrix} - g(\nu_{\mathcal{R}1}, \nu_{\mathcal{R}2}), \\ \begin{bmatrix} \nu_L \\ \nu_I \\ \nu_{\mathcal{R}2} \end{bmatrix} &= - \begin{bmatrix} F_{11} & F_{13} & F_{14} \\ F_{31} & F_{33} & F_{34} \\ F_{41} & F_{43} & F_{44} \end{bmatrix} \begin{bmatrix} \nu_C \\ \nu_{\mathcal{V}} \\ \nu_{\mathcal{R}1} \end{bmatrix}, \\ \begin{bmatrix} \nu_C \\ \nu_{\mathcal{V}} \\ \nu_{\mathcal{R}1} \end{bmatrix} &= \begin{bmatrix} F_{11}^T & F_{31}^T & F_{41}^T \\ F_{13}^T & F_{33}^T & F_{43}^T \\ F_{14}^T & F_{34}^T & F_{44}^T \end{bmatrix} \begin{bmatrix} \nu_L \\ \nu_I \\ \nu_{\mathcal{R}2} \end{bmatrix}, \\ 0 &= \nu_{\mathcal{V}} - \mathcal{V}_s(t), \\ 0 &= \nu_I - I_s(t), \end{aligned}$$

and the corresponding  $\Sigma$ -matrix for the variables

$$x = [\nu_C, \nu_{\mathcal{R}1}, \nu_{\mathcal{R}2}, \nu_{\mathcal{V}}, \nu_L, \nu_I, \iota_C, \iota_{\mathcal{R}1}, \iota_{\mathcal{R}2}, \iota_{\mathcal{V}}, \iota_L, \iota_I]$$

takes the form

$$\begin{bmatrix} \Sigma_C & - & - & - & - & - & - & - & - & - & - \\ - & - & - & - & - & - & - & - & - & - & - \\ - & \Sigma_{22} & \Sigma_{23} & - & - & - & - & - & - & - & - \\ - & \Sigma_{32} & \Sigma_{33} & - & - & - & - & - & - & - & - \\ \hline \Sigma_{51} & \Sigma_{52} & - & \Sigma_{54} & \mathcal{O}_L & - & - & - & - & - & - \\ \Sigma_{61} & \Sigma_{62} & - & \Sigma_{64} & - & \mathcal{O}_I & - & - & - & - & - \\ \Sigma_{71} & \Sigma_{72} & \mathcal{O}_{\mathcal{R}2} & \Sigma_{74} & - & - & - & - & - & - & - \\ \hline - & - & - & - & - & - & \mathcal{O}_C & - & \Sigma_{71}^T & - & \Sigma_{51}^T \\ - & - & - & - & - & - & - & - & \Sigma_{74}^T & \mathcal{O}_{\mathcal{V}} & \Sigma_{54}^T \\ - & - & - & - & - & - & - & - & \Sigma_{72}^T & - & \Sigma_{52}^T \\ \hline - & - & - & \mathcal{O}_{\mathcal{V}} & - & - & - & - & - & - & - \\ - & - & - & - & - & - & - & - & - & - & \mathcal{O}_I \end{bmatrix}.$$

Here,  $\mathcal{O}_*$  for  $* \in \{\mathcal{V}, \mathcal{L}, \mathcal{C}, \mathcal{R}_1, \mathcal{R}_2, \mathcal{I}\}$  denotes square blocks of appropriate size with 0-entries on the diagonal and  $-\infty$ -entries elsewhere. Moreover,  $\Sigma_L$  and  $\Sigma_C$  are defined as before. Consequently, the positions for the HVT can be taken from the diagonal entries of the marked blocks, and we get the canonical offsets  $c = [0]$  and

$$d = [\underbrace{1, \dots, 1}_{n_C}, \underbrace{0, \dots, 0}_{2n_{\mathcal{R}}+2n_{\mathcal{V}}+n_L+n_I+n_C}, \underbrace{1, \dots, 1}_{n_L}, \underbrace{0, \dots, 0}_{n_I}].$$



The corresponding  $\Sigma$ -Jacobian is given by

$$\mathfrak{J} = \left[ \begin{array}{cccc|cc|cc|cccc} \mathcal{C} & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 & \mathcal{L} & 0 \\ 0 & \mathcal{G}_{11} & \mathcal{G}_{12} & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 \\ 0 & \mathcal{G}_{22} & \mathcal{G}_{22} & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 \\ \hline 0 & F_{14} & 0 & F_{13} & I & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & F_{34} & 0 & F_{33} & 0 & I & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & F_{44} & I & F_{43} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & I & 0 & -F_{41}^T & 0 & 0 & -F_{31}^T \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -F_{43}^T & I & 0 & -F_{33}^T \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & I & -F_{44}^T & 0 & 0 & -F_{34}^T \\ \hline 0 & 0 & 0 & I & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & I \end{array} \right],$$

610 which is nonsingular if the matrix

$$\left[ \begin{array}{cc|cc} \mathcal{G}_{11} & \mathcal{G}_{12} & -I & 0 \\ \mathcal{G}_{21} & \mathcal{G}_{22} & 0 & -I \\ \hline F_{44} & I & 0 & 0 \\ 0 & 0 & I & -F_{44}^T \end{array} \right] \sim \left[ \begin{array}{cc|cc} I & 0 & 0 & -F_{44}^T \\ 0 & I & F_{44} & 0 \\ \hline -\mathcal{G}_{11}^{-1} & \mathcal{G}_{11}^{-1} \mathcal{G}_{12} & I & 0 \\ -\mathcal{G}_{21} \mathcal{G}_{11}^{-1} & \mathcal{G}_{21} \mathcal{G}_{11}^{-1} \mathcal{G}_{12} - \mathcal{G}_{22} & 0 & I \end{array} \right]$$

is nonsingular. Under Assumption **(A5)** and using Lemma 20 and Lemma 21 the nonsingularity of  $\mathfrak{J}$  follows and the  $\Sigma$ -method succeeds with  $\nu_S = \nu_d = 1$ . If the DAE (13) is of d-index  $\nu_d = 2$ , the  $\Sigma$ -matrix for (13) takes the form as in (37) for the variables

$$x = [\nu_{C1}, \nu_{C2}, \nu_{\mathcal{R}1}, \nu_{\mathcal{R}2}, \nu_{\mathcal{V}}, \nu_{L1}, \nu_{L2}, \nu_I, \iota_{C1}, \iota_{C2}, \iota_{\mathcal{R}1}, \iota_{\mathcal{R}2}, \iota_{\mathcal{V}}, \iota_{L1}, \iota_{L2}, \iota_I],$$

where, again, the positions for the HVT can be taken from the diagonal entries of the marked blocks. Since we have chosen a normal tree in the formulation of (13), and  $F_{22} = 0$ ,  $F_{42} = 0$  as well as  $F_{24} = 0$ , this choice for the HVT is unique (up to permutations within each block) and the canonical offset vectors

are given by

$$c = [\underbrace{0 \dots 0}_{n_C + n_{\mathcal{R}}}, \underbrace{\leq 1 \dots \leq 1}_{n_{\mathcal{V}}}, \underbrace{0 \dots 0}_{n_L}, \underbrace{\leq 1 \dots \leq 1}_{n_I}, \underbrace{0 \dots 0}_{n_{L2}}, \underbrace{1 \dots 1}_{n_{C2}}, \underbrace{0 \dots 0}_{n_I + n_{\mathcal{R}2} + n_{C1}}, \underbrace{1 \dots 1}_{n_{L1}}, \underbrace{0 \dots 0}_{n_{\mathcal{V}} + n_{\mathcal{R}1}}],$$

$$d = [\underbrace{1 \dots 1}_{n_C}, \underbrace{0 \dots 0}_{n_{\mathcal{R}}}, \underbrace{\leq 1 \dots \leq 1}_{n_{\mathcal{V}}}, \underbrace{0 \dots 0}_{n_L + n_I + n_C + n_{\mathcal{R}} + n_{\mathcal{V}}}, \underbrace{1 \dots 1}_{n_L}, \underbrace{\leq 1 \dots \leq 1}_{n_I}].$$

The entries with values  $\leq 1$  in the offset vectors depend on the sparsity structure of the blocks  $F_{23}$  and  $F_{32}^T$ . However, these particular values do not influence the result of the  $\Sigma$ -method, since the corresponding  $\Sigma$ -Jacobian always has the form

$$\begin{bmatrix} \mathcal{C}_{11} & \mathcal{C}_{12} & 0 & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \mathcal{C}_{21} & \mathcal{C}_{22} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \mathcal{L}_{11} & \mathcal{L}_{12} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 & 0 & \mathcal{L}_{21} & \mathcal{L}_{22} & 0 \\ 0 & 0 & \mathcal{G}_{11} & \mathcal{G}_{12} & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \mathcal{G}_{21} & \mathcal{G}_{22} & 0 & 0 & 0 & 0 & 0 & 0 & -I & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & F_{14} & 0 & * & F_{12} & I & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ F_{21} & I & 0 & 0 & * & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & F_{34} & 0 & * & F_{32} & 0 & I & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & F_{44} & I & * & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & I & -F_{21}^T & 0 & -F_{41}^T & 0 & 0 & 0 & * \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & I & -F_{12}^T & * & * \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -F_{23}^T & 0 & -F_{43}^T & I & 0 & 0 & * \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & I & -F_{44}^T & 0 & 0 & 0 & * \\ \hline 0 & 0 & 0 & 0 & I & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & I \end{bmatrix},$$

615 which is nonsingular due to Assumptions **(A3)**-(**A5**) no matter which entries occur in the blocks denoted with \*. Thus, the  $\Sigma$ -method succeeds with structural index  $\nu_S = \nu_d = 2$ .  $\square$



**Lemma 20.** Consider a matrix of the form

$$A = \begin{bmatrix} I & \tilde{F} \\ \tilde{G} & I \end{bmatrix},$$

where  $\tilde{F}$  is skew-symmetric and  $\tilde{G}$  is definite. Then  $A$  is nonsingular.

*Proof.* The matrix  $A$  is nonsingular if and only if the Schur complement  $I - \tilde{G}\tilde{F}$  is nonsingular. Assume that

$$(I - \tilde{G}\tilde{F})v = 0 \quad (38)$$

for some vector  $v$ . Then

$$v^T \tilde{F}^T (I - \tilde{G}\tilde{F})v = v^T \tilde{F}^T v - v^T \tilde{F}^T \tilde{G}\tilde{F}v = -v^T \tilde{F}^T \tilde{G}\tilde{F}v = 0$$

due to the skew-symmetry of  $\tilde{F}$ . Since  $\tilde{G}$  is definite we get  $\tilde{F}v = 0$  and from (38) it follows that  $v = 0$ .  $\square$

**Lemma 21.** Let  $G = \begin{bmatrix} G_{11} & G_{12} \\ G_{21} & G_{22} \end{bmatrix}$  be a symmetric positive definite matrix. Then the matrix

$$\tilde{G} = \begin{bmatrix} G_{11}^{-1} & -G_{11}^{-1}G_{12} \\ G_{21}G_{11}^{-1} & G_{22} - G_{21}G_{11}^{-1}G_{12} \end{bmatrix}$$

is again positive definite.

*Proof.* Since  $G$  is symmetric positive definite we have that  $G_{11}$  is symmetric positive definite, and in particular  $G_{11}^{-1}$  exists. Thus, from a relation

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} G_{11} & G_{12} \\ G_{21} & G_{22} \end{bmatrix} \begin{bmatrix} v \\ w \end{bmatrix}$$

we get

$$v = G_{11}^{-1}(x - G_{12}w),$$

$$y = G_{21}G_{11}^{-1}(x - G_{12}w) + G_{22}w.$$

Thus, for arbitrary vectors  $v$  and  $w$  we get

$$\begin{aligned} \begin{bmatrix} v^T & w^T \end{bmatrix} \begin{bmatrix} G_{11} & G_{12} \\ G_{21} & G_{22} \end{bmatrix} \begin{bmatrix} v \\ w \end{bmatrix} &= v^T x + w^T y = x^T v + w^T y \\ &= \begin{bmatrix} x^T & w^T \end{bmatrix} \begin{bmatrix} G_{11}^{-1} & -G_{11}^{-1}G_{12} \\ G_{21}G_{11}^{-1} & G_{22} - G_{21}G_{11}^{-1}G_{12} \end{bmatrix} \begin{bmatrix} x \\ w \end{bmatrix} \end{aligned}$$

and positive definiteness of  $G$  transfers to  $\tilde{G}$ .  $\square$

In contrast to the formulation of the circuit equations using the MNA equations  
 630 (6) or the MLA equations (7), the  $\Sigma$ -method for the branch-oriented model  
 equations (13) will always succeed.

**Example 7.** We consider again the circuit given in Example 2. The only proper  
 tree for the graph depicted in Figure 4 is given by the set of branches  $\{\mathcal{G}_1, \mathcal{G}_2, \mathcal{L}\}$   
 and the branch  $I_s$  is the only link. Based on this tree the fundamental cutset  
 matrix and fundamental loop matrix are given by

$$Q = \left[ \begin{array}{ccc|c} 1 & 0 & 0 & -1 \\ 0 & 1 & 0 & -1 \\ 0 & 0 & 1 & -1 \end{array} \right], \quad B = \left[ \begin{array}{c|ccc} 1 & 1 & 1 & 1 \end{array} \right],$$

i.e., we have  $F = \begin{bmatrix} 1 & 1 & 1 \end{bmatrix}$ . Thus, the branch-oriented model equations (13)  
 take the form

$$\begin{aligned} \mathcal{L} \frac{d}{dt} \iota_L &= \nu_L, \\ \iota_{\mathcal{R}_1} &= \mathcal{G}_1 \nu_{\mathcal{R}_1}, \\ \iota_{\mathcal{R}_2} &= \mathcal{G}_2 \nu_{\mathcal{R}_2}, \\ \nu_I &= -(\nu_L + \nu_{\mathcal{R}_1} + \nu_{\mathcal{R}_2}), \\ \iota_L &= \iota_I, \\ \iota_{\mathcal{R}_1} &= \iota_I, \\ \iota_{\mathcal{R}_2} &= \iota_I, \\ 0 &= \iota_I - I_s(t), \end{aligned}$$

and the corresponding  $\Sigma$ -matrix and  $\Sigma$ -Jacobian are given by

$$\Sigma = \left[ \begin{array}{ccc|c|ccc|c} - & - & 0 & - & - & - & 1 & - \\ 0 & - & - & - & 0 & - & - & - \\ - & 0 & - & - & - & 0 & - & - \\ \hline 0 & 0 & 0 & 0 & - & - & - & - \\ \hline - & - & - & - & - & - & 0 & 0 \\ - & - & - & - & 0 & - & - & 0 \\ - & - & - & - & - & 0 & - & 0 \\ \hline - & - & - & - & - & - & - & 0 \end{array} \right]$$

and

$$\mathfrak{J} = \left[ \begin{array}{ccc|c|ccc|c} 0 & 0 & -1 & 0 & 0 & 0 & \mathcal{L} & 0 \\ \mathcal{G}_1 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & \mathcal{G}_2 & 0 & 0 & 0 & -1 & 0 & 0 \\ \hline 1 & 1 & 1 & 1 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 1 & -1 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & -1 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{array} \right].$$

Here,  $\mathfrak{J}$  is nonsingular such that the  $\Sigma$ -method succeeds with  $\nu_S = 2$ .

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#### 4.5. The Signature Method for the pHC equations

Finally, we apply the  $\Sigma$ -method to the pHC equations (16). We consider the pHC equations (16) given in the form

$$\mathcal{C}(t) \frac{d}{dt} \nu_C - \imath_C = 0, \quad (39a)$$

$$\imath_{\mathcal{R}} - \mathcal{G}(t) \nu_{\mathcal{R}} = 0, \quad (39b)$$

$$-\nu_C + A_C^T \eta = 0, \quad (39c)$$

$$-\nu_{\mathcal{R}} + A_{\mathcal{R}}^T \eta = 0, \quad (39d)$$

$$A_{\eta}^T \eta - \mathcal{V}_s(t) = 0, \quad (39e)$$

$$\mathcal{L}(t) \frac{d}{dt} \imath_{\mathcal{L}} - A_{\mathcal{L}}^T \eta = 0, \quad (39f)$$

$$-A_C \imath_C - A_{\mathcal{R}} \imath_{\mathcal{R}} - A_I I_s(t) - A_{\mathcal{V}} \imath_{\mathcal{V}} - A_{\mathcal{L}} \imath_{\mathcal{L}} = 0, \quad (39g)$$

with symmetric and pointwise positive definite matrices  $\mathcal{C}(t)$ ,  $\mathcal{G}(t)$  and  $\mathcal{L}(t)$  due to Assumptions (A3)-(A5). For such a system we get a  $\Sigma$ -matrix of the form

$$\Sigma_{pHC} = \left[ \begin{array}{cc|cccc|c} \Gamma_C & - & \mathcal{O}_C & - & - & - & - \\ - & \Gamma_{\mathcal{R}} & - & \mathcal{O}_{\mathcal{R}} & - & - & - \\ \hline \mathcal{O}_C & - & - & - & - & - & \Sigma_{A_C}^T \\ - & \mathcal{O}_{\mathcal{R}} & - & - & - & - & \Sigma_{A_{\mathcal{R}}}^T \\ - & - & - & - & - & - & \Sigma_{A_{\mathcal{V}}}^T \\ - & - & - & - & - & \Sigma_{\mathcal{L}} & \Sigma_{A_{\mathcal{L}}}^T \\ \hline - & - & \Sigma_{A_C} & \Sigma_{A_{\mathcal{R}}} & \Sigma_{A_{\mathcal{V}}} & \Sigma_{A_{\mathcal{L}}} & - \end{array} \right], \quad (40)$$

where the blocks  $\mathcal{O}_C$  and  $\mathcal{O}_{\mathcal{R}}$  of size  $n_C \times n_C$  and  $n_{\mathcal{R}} \times n_{\mathcal{R}}$ , respectively, have 0-diagonal and entries  $-\infty$  on the off-diagonal, the block  $\Gamma_C$  of size  $n_C \times n_C$  has 1-diagonal and entries  $\leq 1$  on the off-diagonal, and the block  $\Gamma_{\mathcal{R}}$  of size  $n_{\mathcal{R}} \times n_{\mathcal{R}}$  has 0-diagonal and entries  $\leq 0$  on the off-diagonal. The other blocks are defined as before. In this case we get the following result.

**Theorem 22.** Consider an electrical circuit that contains neither  $\mathcal{L}I$ -cutsets nor  $\mathcal{CV}$ -loops (including pure  $\mathcal{C}$ -loops) and let Assumptions (A1)-(A5) hold.

Then the  $pHC$  equations (39) are of  $d$ -index  $\nu_d = 1$  and the  $\Sigma$ -method succeeds with nonsingular  $\Sigma$ -Jacobian and structural index  $\nu_S = 1$ .

*Proof.* Due to **(A1)**, **(A2)** we can always rearrange the rows and columns of the matrix  $[A_C, A_{\mathcal{R}}, A_{\mathcal{V}}, A_L]$  such that

$$\begin{bmatrix} A_C & A_{\mathcal{R}} & A_{\mathcal{V}} & A_L \end{bmatrix} \sim \begin{bmatrix} A_{\mathcal{V}_1} & A_{C_1} & A_{\mathcal{R}_1} & A_{L_1} \\ A_{\mathcal{V}_2} & A_{C_2} & A_{\mathcal{R}_2} & A_{L_2} \end{bmatrix} \begin{matrix} n_{\mathcal{V}} \\ n_{\eta} - 1 - n_{\mathcal{V}} \end{matrix}$$

where  $A_{\mathcal{V}_1}$  is regular and  $\begin{bmatrix} A_{C_2} & A_{\mathcal{R}_2} & A_{L_2} \end{bmatrix}$  has full row rank  $n_{\eta} - 1 - n_{\mathcal{V}}$ . Here, the assumption that  $n_{\eta} - 1 \geq n_{\mathcal{V}}$  is reasonable, since otherwise there would be  $\mathcal{V}$ -loops in the circuit contradicting **(A1)**. As a consequence, we can always find  $n_{\eta} - 1$  positions for the HVT in the last block row and block column of (40) and  $n_{\mathcal{V}}$  of these positions can be chosen from the blocks  $\Sigma_{A_{\mathcal{V}}}$  and  $\Sigma_{A_{\mathcal{V}}}^T$  guaranteeing the existence of a HVT. If the circuit contains no  $\mathcal{C}\mathcal{V}$ -loops and also no  $\mathcal{V}$ -loops or  $\mathcal{C}$ -loops we have  $n_c + n_{\mathcal{V}} \leq n_{\eta} - 1$  as well as

$$\text{rank}[A_C, A_{\mathcal{R}}, A_{\mathcal{V}}] = n_{\eta} - 1 \quad \text{and} \quad \ker[A_C, A_{\mathcal{V}}] = \{0\}$$

due to Theorem 10, such that the  $n_{\eta} - 1$  positions for the HVT in the last block row can be picked from  $[\Sigma_{A_C}, \Sigma_{A_{\mathcal{R}}}, \Sigma_{A_{\mathcal{V}}}]$  only. If  $n_{\mathcal{V}} > 0$ , then  $n_{\mathcal{V}}$  position of the HVT in the last block row are fixed in  $\Sigma_{A_{\mathcal{V}}}$ , and the remaining  $n_{\eta} - 1 - n_{\mathcal{V}}$  positions can be chosen from  $[\Sigma_{A_C}, \Sigma_{A_{\mathcal{R}}}]$ . In particular,  $n_c$  positions can be chosen from  $\Sigma_{A_C}$ . If  $n_{\eta} - 1 > n_c + n_{\mathcal{V}}$ , then the remaining  $n_{\eta} - 1 - n_{\mathcal{V}} - n_c$  positions for the last block row have to be taken from  $\Sigma_{A_{\mathcal{R}}}$ . Due to the symmetry of (40) the transposed entries can be picked as positions for the HVT in the last block column. The remaining positions for the HVT can be picked from the diagonal of  $\Gamma_C$ ,  $\Gamma_{\mathcal{R}}$ ,  $\mathcal{O}_{\mathcal{R}}$  and  $\Sigma_L$ . Concluding, we get a HVT of value  $\text{Val}(HVT) = n_c + n_L$  and the canonical offset vectors are given by  $c = [0]$  and

$$d = [\underbrace{1 \dots 1}_{n_c}, \underbrace{0 \dots 0}_{n_{\mathcal{R}}}, \underbrace{0 \dots 0}_{n_c}, \underbrace{0 \dots 0}_{n_{\mathcal{R}}}, \underbrace{0 \dots 0}_{n_{\mathcal{V}}}, \underbrace{1 \dots 1}_{n_L}, \underbrace{0 \dots 0}_{n_{\eta}-1}]$$



according to the block structure of (40). The resulting  $\Sigma$ -Jacobian is given by

$$\mathfrak{J} = \left[ \begin{array}{cc|cccc|c} \mathcal{C} & 0 & -I & 0 & 0 & 0 & 0 \\ 0 & -\mathcal{G} & 0 & I & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 & 0 & A_c^T \\ 0 & -I & 0 & 0 & 0 & 0 & A_{\mathcal{R}}^T \\ 0 & 0 & 0 & 0 & 0 & 0 & A_{\nu}^T \\ 0 & 0 & 0 & 0 & 0 & \mathcal{L} & -A_{\mathcal{L}}^T \\ \hline 0 & 0 & -A_c & -A_{\mathcal{R}} & -A_{\nu} & 0 & 0 \end{array} \right].$$

Under the d-index-1 conditions for the pHC-equations (39) this matrix is non-singular (see also the proof of Theorem 10) and the  $\Sigma$ -method succeeds with  $\nu_S = 1 = \nu_d$ .

□

**Remark 23.** If the circuit contains  $\mathcal{LI}$ -cutsets or  $\mathcal{CV}$ -loops, i.e., the pHC equations (39) are of d-index  $\nu_d = 2$ , then the  $\Sigma$ -method for the pHC equations (39) might still work successfully in many cases. However, in this case the results (and the success) can depend on the selection of the reference node as can be seen in Example 8. If we assume that the branches are ordered in such a way that  $A$  is split into

$$A = \left[ \begin{array}{cccc|cccc} A_{C_1} & A_{\mathcal{L}_1} & A_{\mathcal{R}_1} & A_{\nu_1} & A_{C_2} & A_{\mathcal{L}_2} & A_{\mathcal{R}_2} & A_I \end{array} \right] =: \left[ \begin{array}{c|c} A_T & A_{coT} \end{array} \right],$$

where  $A_T$  contains the twigs that belong to a normal tree and  $A_{coT}$  contains the links of the corresponding cotree, then due to Lemma 28 we know that  $A_T$  is regular and w.l.o.g we can assume that the nodes are ordered such that

$$\left[ \begin{array}{c|c} A_T & A_{coT} \end{array} \right] = \left[ \begin{array}{cccc|cccc} A_{C_{11}} & A_{\mathcal{L}_{11}} & A_{\mathcal{R}_{11}} & A_{\nu_1} & A_{C_{21}} & A_{\mathcal{L}_{21}} & A_{\mathcal{R}_{21}} & A_{I_1} \\ A_{C_{12}} & A_{\mathcal{L}_{12}} & A_{\mathcal{R}_{12}} & A_{\nu_2} & A_{C_{22}} & A_{\mathcal{L}_{22}} & A_{\mathcal{R}_{22}} & A_{I_2} \\ A_{C_{13}} & A_{\mathcal{L}_{13}} & A_{\mathcal{R}_{13}} & A_{\nu_3} & A_{C_{23}} & A_{\mathcal{L}_{23}} & A_{\mathcal{R}_{23}} & A_{I_3} \\ A_{C_{14}} & A_{\mathcal{L}_{14}} & A_{\mathcal{R}_{14}} & A_{\nu_4} & A_{C_{24}} & A_{\mathcal{L}_{24}} & A_{\mathcal{R}_{24}} & A_{I_4} \end{array} \right]$$

with regular blocks  $A_{C_{11}}$ ,  $A_{\mathcal{L}_{12}}$ ,  $A_{\mathcal{R}_{13}}$ , and  $A_{\nu_4}$ . In this case, we can show that the  $\Sigma$ -method succeeds with structural index  $\nu_S = \nu_d = 2$  if  $A_{C_{12}} = 0$ ,  $A_{\mathcal{R}_{12}} = 0$ ,

660  $A_{q_2} = 0$ ,  $A_{\mathcal{R}_{22}} = 0$  and  $A_{C_2} = []$  is an empty matrix, i.e., there are no link capacitors. In particular, this means that we have to demand that nodes that are connected to twig inductors have no connections to tree capacitors, resistors or voltage sources, and that there are no CV-loops. Thus, in general we can say that the  $\Sigma$ -method might fail if there exist nodes in the circuit graph that  
 665 are incident to a twig inductor as well as to a link capacitor for an (arbitrarily chosen) normal tree. If such a node exists it should be chosen as reference node in order to prevent the failure of the  $\Sigma$ -method.

**Example 8.** Consider the circuit given in Figure 10 containing a  $\mathcal{CV}$ -loop and, thus, having d-index  $\nu_d = 2$ . The graph corresponding to the circuit is given in

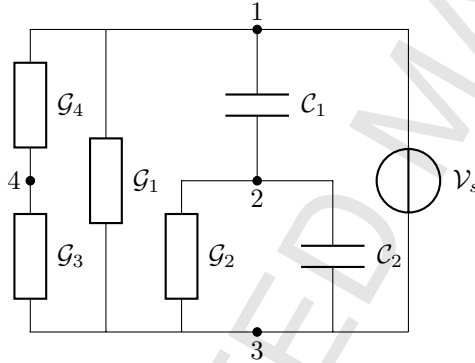


Figure 10:  $\mathcal{RCV}$ -circuit

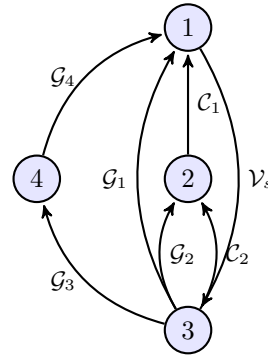


Figure 11: Graph for the  $\mathcal{RCV}$ -circuit

670 Figure 11. If we select node 4 as reference node, then the  $\Sigma$ -method fails due to a singular  $\Sigma$ -Jacobian. Note that node 4 is the only node that is not adjacent to a (link) capacitor. If we select node 1, 2 or 3 as reference node, the  $\Sigma$ -method succeeds with nonsingular  $\Sigma$ -Jacobian and  $\nu_S = 2$ .  $\triangleleft$

## 5. Conclusion

675 The structural analysis is a commonly used and powerful tool in the automatized numerical treatment of DAEs. In this paper we have used the  $\Sigma$ -method for the structural analysis of DAEs that arise in the modeling and simulation of

electrical circuits. The presented examples show that success or failure of the  $\Sigma$ -method for the commonly used MNA equations (6) can depend on the topology of the circuit. In the case of failure of the  $\Sigma$ -method the computed offset vectors will not give the required information on the index and on the system structure that can otherwise be used for the regularization of the system (as proposed e.g. in [24]). As a consequence, a robust and stable numerical integration cannot be guaranteed. Moreover, the result of the  $\Sigma$ -method for the MNA equations (6) can also depend on the choice of the reference node. This choice does not influence the analytical properties of the system (as e.g. the index) and typically a node that has a large number of adjacent edges is chosen as reference. Our investigations suggest that a node that is adjacent to a capacitance or, if no capacitances are present, adjacent to a resistance should be chosen as reference node. However, this is no guarantee for the success of the  $\Sigma$ -method for the MNA equations (6) as Example 2 shows. Similar observations can be made if the  $\Sigma$ -method is applied to the MLA equations (7). We have also seen that the  $\Sigma$ -method can give different results for the MNA and MLA equations depending on the topology of the circuit. The main consequence of the results obtained in Section 4.2 and 4.3 is that Signature Method cannot be guaranteed to work successfully for circuit equations of index  $\geq 1$  if they are modeled using the MNA or MLA equations. In contrast, for the branch-oriented model equations (10) we have shown in Theorem 19 that the  $\Sigma$ -method always succeeds with a structural index that corresponds to the d-index of the system.

The MNA equations are widely used in circuit simulation because of their compact form and structural properties (e.g. symmetries) that allow for an efficient regularization and numerical integration. However, we see that from the view point of structural analysis the MNA equations are not well-suited and more sparser representations as the branch-oriented model equations are beneficial for a structural analysis. This knowledge is of particular importance if one is not only interested in circuit simulation (for which sophisticated simulation packages exist), but rather looks at a circuit as a component of a multi-physical applications in the context of automatized modeling and simulation.

We have also considered a port-Hamiltonian formulation of the circuit equations. In this case, the  $\Sigma$ -method applied to the pH equations (39) will always succeed if the circuit contains neither  $\mathcal{LI}$ -cutsets nor  $\mathcal{CV}$ -loops (including pure  $\mathcal{C}$ -loops). If the circuit contains  $\mathcal{LI}$ -cutsets or  $\mathcal{CV}$ -loops, i.e., the pH equations (39) are of d-index  $\nu_d = 2$ , the  $\Sigma$ -method might still work successfully in many cases, but the results (and the success) can depend on the selection of the reference node as examples show. In particular, the  $\Sigma$ -method can fail for the pH equations (39) if there exists nodes in the circuit graph that are incident to a twig inductor as well as to a link capacitor for an arbitrarily chosen normal tree. If this is the case, such a node should be chosen as reference node in order to prevent the failure of the  $\Sigma$ -method in the pH approach. However, still the  $\Sigma$ -method cannot be guaranteed to work successfully for circuit equations of index 2 if they are modeled using the pH approach.

## References

- [1] K. Brenan, S. Campbell, L. Petzold, Numerical Solution of Initial-Value Problems in Differential Algebraic Equations, Vol. 14 of Classics in Applied Mathematics, SIAM, Philadelphia, PA, 1996.
- [2] E. Hairer, G. Wanner, Solving Ordinary Differential Equations II: Stiff and Differential-Algebraic Problems, 2nd Edition, Springer, Berlin, 1996.
- [3] P. Kunkel, V. Mehrmann, Differential-Algebraic Equations — Analysis and Numerical Solution, EMS Publishing House, Zürich, Switzerland, 2006.
- [4] L. Scholz, A. Steinbrecher, DAEs in applications, in: Numerical Algebra, Matrix Theory, Differential-Algebraic Equations and Control Theory: Festschrift in Honor of Volker Mehrmann, Springer International Publishing, Cham, 2015, pp. 463–501. doi:10.1007/978-3-319-15260-8\_17.
- [5] C. C. Pantelides, The consistent initialization of differential-algebraic systems, SIAM J. Sci. Statist. Comput. 9 (1988) 213–231. doi:10.1137/0909014.

- [6] J. Pryce, A simple structural analysis method for DAEs, BIT Numerical Mathematics 41 (2001) 364–394. doi:10.1023/A:1021998624799.
- [7] S. Mattsson, G. Söderlind, Index reduction in differential-algebraic equations using dummy derivatives, SIAM J. Sci. Statist. Comput. 14 (1993) 677–692. doi:10.1137/0914043.
- [8] G. Reißig, W. S. Martinson, P. I. Barton, Differential-algebraic equations of index 1 may have an arbitrarily high structural index, SIAM Journal on Scientific Computing 21 (2000) 1987–1990. doi:10.1137/S1064827599353853.
- [9] L. Scholz, A. Steinbrecher, Structural-algebraic regularization for coupled systems of DAEs, BIT Numerical Mathematics 56 (2) (2016) 777–804. doi:10.1007/s10543-015-0572-y.
- [10] S. Campbell, C. Gear, The index of general nonlinear DAEs, Numerische Mathematik 72 (2) (1995) 173–196. doi:10.1007/s002110050165.
- [11] N. Deo, Graph Theory with Applications to Engineering and Computer Science, Prentice-Hall, Englewood Cliffs, NJ, 1974.
- [12] R. Riaza, Differential-algebraic Systems: Analytical Aspects and Circuit Applications, World Scientific, 2008.
- [13] J. Vlach, K. Singhal, Computer Methods for Circuit Analysis and Design, Van Nostrand Reinhold, New York, NY, 1994.
- [14] D. Estévez Schwarz, C. Tischendorf, Structural analysis of electric circuits and consequences for MNA, International Journal of Circuit Theory and Applications 28 (2) (2000) 131–162. doi:10.1002/(SICI)1097-007X(200003/04)28:2<131::AID-CTA100>3.0.CO;2-W.
- [15] M. Günther, U. Feldmann, CAD-based electric-circuit modeling in industry, I. Mathematical structure and index of network equations, Surveys on Mathematics for Industry 8 (1999) 97–129.

- [16] T. Reis, Mathematical modeling and analysis of nonlinear time-invariant  
765 RLC circuits, in: P. Benner, R. Findeisen, D. Flockerzi, U. Reichl,  
K. Sundmacher (Eds.), Large Scale Networks in Engineering and Life  
Sciences, Springer International Publishing, Cham, 2014, pp. 125–198.  
doi:10.1007/978-3-319-08437-4\_2.
- [17] C. Tischendorf, Topological index calculation of differential-algebraic equa-  
770 tions in circuit simulation, Surveys on Mathematics for Industry 8 (3-4)  
(1999) 187–199.
- [18] L. Weinbreg, A. Ruehili, Combined modified loop analysis, modified nodal  
analysis for large-scale circuits, IBM Research Report RC 10407 (1984).
- [19] V. Mehrmann, Index concepts for differential-algebraic equations, in:  
775 B. Engquist (Ed.), Encyclopedia of Applied and Computational Mathe-  
matics, Springer, Berlin, Heidelberg, 2015, pp. 676–681. doi:10.1007/  
978-3-540-70529-1\_120.
- [20] H. Paynter, Analysis and Design of Engineering Systems, MIT Press, Cam-  
bridge, 1961.
- [21] A. J. van der Schaft, Port-Hamiltonian Differential-Algebraic Systems,  
780 in: A. Ilchmann, T. Reis (Eds.), Surveys in Differential-Algebraic Equa-  
tions I, Springer, Berlin, Heidelberg, 2013, pp. 173–226. doi:10.1007/  
978-3-642-34928-7\_5.
- [22] C. Beattie, V. Mehrmann, H. Xu, H. Zwart, Port-hamiltonian descriptor  
785 systems, Preprint 06-2017, Institut für Mathematik, TU Berlin (2017).
- [23] A. J. van der Schaft, Port-Hamiltonian systems: network modeling and  
control of nonlinear physical systems, Advanced Dynamics and Control of  
Structures and Machines, CISM Courses and Lectures 444. doi:10.1007/  
978-3-7091-2774-2\_9.

- [24] L. Scholz, A. Steinbrecher, Regularization of DAEs based on the Signature method, BIT Numerical Mathematics 56 (2016) 319–340. doi:10.1007/s10543-015-0565-x.

### Appendix A. Graph theoretical results

The following graph theoretical results can be found e.g. in [11, 12].

**Theorem 24.** Let  $G$  be a directed graph consisting of  $n_\eta$  nodes,  $n_b$  branches and containing  $n_\ell$  loops and  $n_q$  cutsets. Moreover, let  $A_0 \in \mathbb{R}^{n_\eta, n_b}$  be the all-node incidence matrix,  $B_0 \in \mathbb{R}^{n_\ell, n_b}$  the loop matrix and  $Q_0 \in \mathbb{R}^{n_q, n_b}$  the cutset matrix with columns arranged according to the same order of branches. Furthermore, let  $k$  be the number of connected components of  $G$ . Then

- (i)  $\text{rank } A_0 = n_\eta - k$ ,
- (ii)  $\text{rank } B_0 = n_b - n_\eta + k$ ,
- (iii)  $\text{rank } Q_0 = n_\eta - k$ ,
- (iv)  $\text{im } B_0^T = \ker A_0 = \ker Q_0$ ,
- (v)  $A_0 B_0^T = 0$  and  $B_0 Q_0^T = 0$ ,
- (vi)  $G$  is a tree if and only if  $A_0 \in \mathbb{R}^{n_\eta, n_\eta-1}$  and  $\ker A_0 = \{0\}$ .

Under the reasonable assumption that the electrical circuit graph is connected we can always select  $n_\eta - 1$  linearly independent rows of  $A_0$  to obtain the reduced matrix  $A$ , and  $n_b - n_\eta + 1$  linearly independent rows of  $B_0$  to obtain the reduced matrix  $B$ .

**Remark 25.** The all-node incidence matrix  $A_0$  is a submatrix of the cutset matrix  $Q_0$  since for each node of the graph the set of all branches incident to this node forms a cutset of the graph.

A subgraph  $K = (V', B', \Psi|_{B'})$  of a connected graph  $G$  with  $V' = V$ ,  $B' \subset B$  is called a *spanning subgraph*. For a spanning subgraph  $K$  of a directed graph  $G$  let  $A_K$  ( $A_{G-K}$ ) denote the submatrix of the incidence matrix  $A$  that is formed by the columns corresponding to branches in  $K$  (respectively the complementary

graph  $G - K$ ). Analogously, let  $B_K$  and  $B_{G-K}$  denote the corresponding loop matrices. By suitable reordering we can always get  $A = [A_K \ A_{G-K}]$  and  $B = [B_K \ B_{G-K}]$ .

820 **Lemma 26.** *Let  $G$  be a connected directed graph with reduced incidence and loop matrices  $A \in \mathbb{R}^{n_\eta-1, n_b}$  and  $B \in \mathbb{R}^{n_b-n_\eta+1, n_b}$ . Furthermore, let  $K$  be a spanning subgraph of  $G$  and assume that the branches of  $G$  are sorted such that*

$$A = [A_K \ A_{G-K}], \quad B = [B_K \ B_{G-K}].$$

1. *The following assertions are equivalent:*

- $G$  does not contain  $K$ -cutsets;
- 825 •  $\ker A_{G-K}^T = \{0\}$ , i.e.,  $A_{G-K}$  has full row rank;
- $\ker B_K = \{0\}$ , i.e.,  $B_K$  has full column rank.

2. *The following assertions are equivalent:*

- $G$  does not contain  $K$ -loops;
- $\ker A_K = \{0\}$ , i.e.,  $A_K$  has full column rank;
- 830 •  $\ker B_{G-K}^T = \{0\}$ , i.e.,  $B_{G-K}$  has full row rank.

From Lemma 26 we get that an electrical circuit contains no  $\mathcal{CV}$ -loops if and only if the matrix  $[A_C \ A_{\mathcal{V}}]$  has full column rank. Similar the circuit contains no  $\mathcal{LI}$ -cutsets if and only if the matrix  $[A_{\mathcal{R}} \ A_C \ A_{\mathcal{V}}]$  has full row rank.

835 **Lemma 27.** *Let  $G = (V, B)$  be a connected graph and let  $J, K$  be disjoint subsets of  $B$ . Then there exist a tree which contains all branches from  $J$  and no branches from  $K$  if and only if  $J$  has no loops and  $K$  has no cutsets.*

**Lemma 28.** *Let  $K$  be a set of  $n_\eta - 1$  branches of a connected directed graph. Then  $A_K$  is nonsingular if and only if  $K$  defines a tree. In this case  $\det(A_K) = \pm 1$ .*

840 Thus, if we partition  $A$  as  $A = [A_T \ A_{coT}]$  where the columns of  $A_T$  corresponds to the twigs of a tree, then  $A_T$  is regular.



Let  $T$  be a tree of the connected and directed graph  $G$ , and let  $L$  be the set of all branches that do not belong to the tree (i.e., the set of all links). Then, for every  $z \in L$ , the set  $T \cup \{z\}$  forms a loop. These are the so-called *fundamental loops* with orientations defined as the orientation of the corresponding link  $z$ .  
 845 Since each tree contains  $n_\eta - 1$  twigs there are exactly  $n_b - n_\eta + 1$  fundamental loops. The fundamental loop matrix that only contains the fundamental loops has full row rank and by a suitable ordering of the branches takes the form

$$B = \begin{bmatrix} I & B_T \end{bmatrix},$$

where the columns of  $B_T$  corresponds to the twigs of the tree  $T$ .

850 On the other hand, let  $b \in T$  be a twig of the tree. If we remove the branch  $b$ , then  $T$  decomposes into two separated but connected subtrees  $T_1$  and  $T_2$ . If we denote the set of all nodes in  $T_i$  as  $N_i$ , for  $i = 1, 2$ , then the set of all branches of  $G$  that connect nodes from  $N_1$  with nodes from  $N_2$  forms a cutset of  $G$ . This cutset can be uniquely identified with the corresponding twig  $b$  of  $T$ . These  
 855 cutsets are the so-called *fundamental cutsets* with orientations defined as the orientation of the corresponding twig  $b$ . Since the tree  $T$  contains  $n_\eta - 1$  twigs there are  $n_\eta - 1$  fundamental cutsets. The fundamental cutset matrix that only contains the fundamental cutsets can be represented by a suitable ordering of the branches as

$$Q = \begin{bmatrix} Q_L & I \end{bmatrix},$$

860 where the branches corresponding to the last  $n_\eta - 1$  columns belong to the tree  $T$ .