



Dirac operators on quasi-Hamiltonian G -spaces



Yanli Song

University of Toronto, Canada

ARTICLE INFO

Article history:

Received 22 September 2015

Received in revised form 23 January 2016

Accepted 29 January 2016

Available online 12 February 2016

Keywords:

Dirac operators

Geometric quantization

Quasi-Hamiltonian G -space

Loop group

ABSTRACT

We construct twisted spinor bundles as well as twisted pre-quantum bundles on quasi-Hamiltonian G -spaces, using the spin representation of loop group and the Hilbert space of Wess–Zumino–Witten model. We then define a Hilbert space together with a Dirac operator acting on it. The main result of this paper is that we show the Dirac operator has a well-defined index given by positive energy representation of the loop group. This generalizes the geometric quantization of Hamiltonian G -spaces to quasi-Hamiltonian G -spaces.

Crown Copyright © 2016 Published by Elsevier B.V. All rights reserved.

1. Introduction

Let G be a compact, connected Lie group, and (M, ω) a compact symplectic manifold with a Hamiltonian G -action. By choosing a G -invariant ω -compatible almost complex structure on M , we can define a G -equivariant \mathbb{Z}_2 -graded spinor bundle S_M^\pm . If the Hamiltonian G -space M is pre-quantizable and has a G -equivariant pre-quantum line bundle L , we define a \mathbb{Z}_2 -graded Hilbert space by

$$\mathcal{H}^\pm = L^2(M, S_M^\pm \otimes L)$$

and a G -equivariant Spin^c -Dirac operator

$$D^\pm : \mathcal{H}^\pm \rightarrow \mathcal{H}^\mp.$$

Attributed to Bott, the *quantization* of (M, ω) can be defined as the equivariant index

$$Q(M, \omega) = \text{Ind}(D) = [\ker(D^+)] - [\ker(D^-)] \in R(G).$$

The goal of this paper is to generalize the quantization process to the *quasi-Hamiltonian G -space* introduced by Alekseev–Malkin–Meinrenken [1]. The q -Hamiltonian G -space, arising from infinite-dimensional Hamiltonian loop group space, differs in many respects from Hamiltonian G -space. In particular, the moment map takes values in the group G and the 2-form ω does not have to be closed or non-degenerate. Consequently, the two key ingredients in defining $Q(M, \omega)$: the spinor bundle S_M and pre-quantum line bundle L might not exist in general.

Given a q -Hamiltonian G -space (M, ω) , we use the spin representation of loop group to construct twisted spinor bundles S^{spin} on M , and the Hilbert space of Wess–Zumino–Witten model to construct twisted pre-quantum bundles S^{pre} . Both of them are bundles of Hilbert space and play the same roles as the spinor bundle and pre-quantum line bundle for Hamiltonian

E-mail address: songyanl@math.utoronto.ca.

<http://dx.doi.org/10.1016/j.geomphys.2016.01.012>

0393-0440/Crown Copyright © 2016 Published by Elsevier B.V. All rights reserved.

G-space. We analogously define a Hilbert space

$$\mathcal{H} := [L^2(M, S^{\text{spin}} \otimes S^{\text{pre}})]^G.$$

One key in the construction of Dirac operators on \mathcal{H} is the algebraically defined cubic Dirac operator. It was introduced by Kostant for finite-dimensional Lie group, and extensively studied for infinite-dimensional loop group by various people. Our strategy is to construct a Dirac operator as a combination of algebraic cubic Dirac operators and geometric Spin^c -Dirac operators. To be more precise, we choose an open cover of M using the symplectic cross-section theorem for q -Hamiltonian G -space, so that every open subset U has the geometric structure:

$$U \cong G \times_H V,$$

where H is a compact subgroup of G and the slice V is a Hamiltonian H -space. Accordingly, the tangent bundle TU splits equivariantly into “vertical direction” and “horizontal direction”. We define a suitable Dirac operator on U so that it acts as the Spin^c -Dirac operator on the vertical part V and the cubic Dirac operator for loop group on the horizontal part G/H . Using partition of unity, we obtain a global Dirac operator \mathcal{D} on \mathcal{H} by patching together Dirac operators on the open sets U . The main result of this paper is that we show the Dirac operator \mathcal{D} has a well-defined index given by positive energy representations of loop group.

2. Loop group and positive energy representation

We first give a brief review on loop groups and their representations. We use [2] as our primary reference.

2.1. Loop group and central extension

Let G be a compact, simple and simply connected Lie group, and fix a “Sobolev level” $s > 1$. We define LG the loop group as the Banach Lie group consisting of maps $S^1 \rightarrow G$ of Sobolev class $s + \frac{1}{2}$ with the group structure given by pointwise multiplication. The Lie algebra $L\mathfrak{g} = \Omega^0(S^1, \mathfrak{g})$ is given by the space Lie algebra \mathfrak{g} -valued 0-forms of Sobolev class $s + \frac{1}{2}$ and $L\mathfrak{g}^* = \Omega^1(S^1, \mathfrak{g})$ the space of \mathfrak{g} -valued 1-forms of Sobolev class $s - \frac{1}{2}$. Integration over S^1 gives a natural non-degenerate pairing between $L\mathfrak{g}$ with $L\mathfrak{g}^*$.

Note that $L\mathfrak{g}^*$ can be identified with the affine space of connections on the trivial principle G -bundle over S^1 . The loop group LG acts on $L\mathfrak{g}^*$ by gauge transformation

$$g \cdot \xi = \text{Ad}_g(\xi) - dg \cdot g^{-1}, \quad g \in LG, \xi \in L\mathfrak{g}^*, \tag{2.1}$$

where $dg \cdot g^{-1}$ is the pull-back of the right-invariant Maurer–Cartan form on G .

Let LG be the basic central extension of LG , defined in [2, Section 4.4]. The coadjoint action of LG on

$$\widehat{L\mathfrak{g}^*} = L\mathfrak{g}^* \oplus \mathbb{R}$$

is given by the formula

$$g \cdot (\xi, k) = (\text{Ad}_g(\xi) - k \cdot g^{-1}dg, k).$$

One can view the action (2.1) as the coadjoint action on the affine hyperplane $L\mathfrak{g}^* \times \{1\} \subset \widehat{L\mathfrak{g}^*}$.

Fixing a maximal torus T , the choice of a set of positive roots \mathfrak{R}_+ for G determines a positive Weyl chamber \mathfrak{t}_+^* . It is well-known that the orbits of coadjoint G -action on \mathfrak{g}^* are parameterized by points in \mathfrak{t}_+^* . The set of coadjoint LG -orbits can be described as follows. Denote by α_0 the highest root and

$$\rho_G = \frac{1}{2} \sum_{\alpha \in \mathfrak{R}_+} \alpha.$$

There is a unique ad-invariant inner product $\langle \cdot, \cdot \rangle_{\mathfrak{g}}$ on \mathfrak{g} , rescaled so that the highest root of \mathfrak{g} has norm $\sqrt{2}$. The dual Coxeter number of G is defined by

$$h^\vee = 1 + \langle \rho_G, \alpha_0 \rangle_{\mathfrak{g}},$$

and the fundamental Weyl alcove for G is the simplex

$$\mathfrak{A} = \{\xi \in \mathfrak{t}_+^* \mid \langle \alpha_0, \xi \rangle_{\mathfrak{g}} \leq 1\} \subset \mathfrak{t} \subset \mathfrak{g}.$$

Every coadjoint orbit of LG -action on $L\mathfrak{g}^*$ contains a unique point in \mathfrak{A} .

For any $\xi \in L\mathfrak{g}^*$, we define the holonomy map

$$\text{Hol} : L\mathfrak{g}^* \rightarrow G$$

the smooth map that sends ξ to the holonomy of ξ around S^1 . This map satisfies the equivariance property

$$\text{Hol}(g \cdot \xi) = g(0) \cdot \text{Hol}(\xi) \cdot g(0)^{-1}.$$

It follows that the based loop group

$$\Omega G = \{g \in LG \mid g(0) = e\}$$

acts freely on $L\mathfrak{g}^*$.

The isotropy group $(LG)_\xi$ is isomorphic to $G_{\text{Hol}(\xi)}$, and thus compact. It sets up a 1–1 correspondence between the set of coadjoint LG -orbits and conjugacy classes. Moreover $(LG)_\xi$ with $\xi \in \mathfrak{t}_+^*$ depends only on the open face σ of \mathfrak{A} containing ξ and will be denoted by $(LG)_\sigma$. If we introduce a partial order on open faces by setting $\tau \preceq \sigma$ if $\tau \subseteq \bar{\sigma}$, then one has that

$$\sigma \preceq \tau \Rightarrow (LG)_\tau \subseteq (LG)_\sigma.$$

In particular, $(LG)_0 = G$ and $(LG)_{\text{int}\mathfrak{A}} = T$.

2.2. Positive energy representation

Let S_{rot}^1 be the rotation group on S^1 and ∂ its infinitesimal generator. Consider a unitary representation of $S_{\text{rot}}^1 \times \widehat{LG}$ on a Hilbert space V , on which the central circle acts by scalar multiplication.

Definition 2.1. We say that V is a *positive energy representation* if the energy operator ∂ is self-adjoint with spectrum bounded below. Moreover, we say that V has *level k* if the central circle of \widehat{LG} acts with weight k .

The positive energy representations of loop groups behave quite analogously to the representation theory of compact Lie groups. For example, every irreducible positive energy representation is uniquely determined by the highest weight. To be more precise, let T be a maximal torus of G and Λ^* the weight lattice. We take $S_{\text{rot}}^1 \times T \times S^1$ as the maximal torus of $S_{\text{rot}}^1 \times \widehat{LG}$, where the second S^1 factor comes from the central extension. The *affine weights* of LG are in the forms of (m, λ, k) , where $m \in \mathbb{Z}$ is the *energy*, $\lambda \in \Lambda^*$ is the weight of G , and k is the *level*.

The *affine Weyl group* $W_{\text{aff}} = W \ltimes \Lambda$ acts on affine weights as follows: the Weyl group W acts as usual on Λ^* and the action of $z \in \Lambda$ is given by

$$z \cdot (m, \lambda, k) = \left(m + \langle \lambda, z \rangle + \frac{k}{2} \cdot \|z\|^2, \lambda + k \cdot z, k \right).$$

The level k is fixed by the affine Weyl group action and the energy is shifted so as to preserve the inner product:

$$(m_1, \lambda_1, k_1) \cdot (m_2, \lambda_2, k_2) = \langle \lambda_1, \lambda_2 \rangle - m_1 k_2 - m_2 k_1. \quad (2.2)$$

For a fixed level k , every irreducible positive energy representation V of LG is uniquely determined by the dominant weight λ at the minimum energy m . We call $\lambda = (m, \lambda, k)$ the *highest weight* of V .

Let $V(n)$ be a subspace of the Hilbert space V , on which the energy operator ∂ acts on $V(n)$ with weight n . The positive energy condition asserts that there is an integer n_{\min} so that $V(n) = 0$ for all $n < n_{\min}$. The algebraic direct sum

$$V^{\text{fin}} := \bigoplus_n V(n)$$

consists of vectors of *finite energy*. It is a dense subspace of V . In addition, we can always normalize so that the lowest energy level n_{\min} equals zero.

It is well-known by the Borel–Weil theorem that all the irreducible G -representations are parameterized by the dominant weights $P_{G,+} = \Lambda^* \cap \mathfrak{t}_+^*$. Similarly there is a 1–1 correspondence between irreducible positive energy representation at level k and weight in

$$P_{k,+} = k\mathfrak{A} \cap \Lambda^* = \left\{ \lambda \in \Lambda^* \mid \frac{\lambda}{k} \in \mathfrak{A} \right\}.$$

For any $\lambda \in P_{k,+}$, we denote by V_λ the irreducible positive energy representation at level k with highest weight λ . In addition, we write V_λ^* the dual of V_λ , which can be considered as a negative energy representation. In particular, the energy operator ∂ acts on V_λ^* with spectrum bounded above.

The abelian group $R_k(LG)$ generated by irreducible positive energy representations at level k has a finite basis and a ring structure known as *fusion product*. But we will not discuss it in this paper.

3. Dirac operators in the algebraic setting

The *cubic Dirac operator* is an algebraically defined operator introduced by Kostant [3] for finite-dimensional Lie algebras. It has now been generalized to infinite-dimensional case and plays an important role in the theory of loop groups: its application in representation theory was first demonstrated in the lecture notes of Wassermann [4]. Later Landweber and

Posthuma generalize it to different homogeneous settings [5,6]. A family version of the cubic Dirac operator was used by Freed–Hopkins–Teleman [7] to construct the isomorphism between the twisted K -theory and fusion ring of loop groups. In addition Meinrenken [8] discusses its application in general Kac–Moody algebra.

3.1. Finite dimensional case

Let G be a compact Lie group and \mathfrak{g} its Lie algebra equipped with an ad-invariant inner product $\langle \cdot, \cdot \rangle_{\mathfrak{g}}$. Let $\text{Cliff}(\mathfrak{g})$ be the \mathbb{Z}_2 -graded complex Clifford algebra of \mathfrak{g} and $S_{\mathfrak{g}}$ an irreducible \mathbb{Z}_2 -graded $\text{Cliff}(\mathfrak{g})$ -module.

Fix an orthonormal basis

$$X_a, \quad a = 1, \dots, \dim_{\mathfrak{g}}.$$

We define a map $\text{ad}^{\mathfrak{g}} : \mathfrak{g} \rightarrow \text{Cliff}(\mathfrak{g})$ by the formula

$$\text{ad}^{\mathfrak{g}}(X) := \frac{1}{4} \cdot \sum_{a=1}^{\dim_{\mathfrak{g}}} [X, X_a] \cdot X_a, \quad X \in \mathfrak{g}.$$

Let $U(\mathfrak{g})$ be the universal enveloping algebra of \mathfrak{g} . We consider the *non-commutative Weil algebra*

$$U(\mathfrak{g}) \otimes \text{Cliff}(\mathfrak{g})$$

introduced by Alekseev–Meinrenken [9].

Definition 3.1. The *cubic Dirac operator* $D_{\mathfrak{g}}$ is an element of the algebra $U(\mathfrak{g}) \otimes \text{Cliff}(\mathfrak{g})$ defined as

$$D_{\mathfrak{g}} = \sum_{a=1}^{\dim_{\mathfrak{g}}} \left(X_a \otimes X_a + \frac{1}{3} \otimes \text{ad}^{\mathfrak{g}}(X_a) \cdot X_a \right).$$

The key property of the cubic Dirac operator is that its square is simple and nice:

$$D_{\mathfrak{g}}^2 = -2\Omega_G + \text{constant},$$

where Ω_G is the quadratic Casimir operator.

More generally, let $H \subset G$ be a closed subgroup of the equal rank. Using the inner product we write

$$\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{p}, \quad \mathfrak{p} = \mathfrak{h}^{\perp}.$$

This decomposition induces isomorphisms:

$$\text{Cliff}(\mathfrak{g}) \cong \text{Cliff}(\mathfrak{h}) \otimes \text{Cliff}(\mathfrak{p}), \quad S_{\mathfrak{g}} \cong S_{\mathfrak{h}} \otimes S_{\mathfrak{p}},$$

where $S_{\mathfrak{h}}, S_{\mathfrak{p}}$ are spinor modules of $\text{Cliff}(\mathfrak{h})$ and $\text{Cliff}(\mathfrak{p})$ respectively.

Definition 3.2. We define the *relative cubic Dirac operator*

$$D_{\mathfrak{g},\mathfrak{h}} \in U(\mathfrak{g}) \otimes \text{Cliff}(\mathfrak{p})$$

by

$$D_{\mathfrak{g},\mathfrak{h}} = \sum^{(\mathfrak{p})} \left(X_a \otimes X_a + \frac{1}{3} \otimes \text{ad}^{\mathfrak{p}}(X_a) \cdot X_a \right). \tag{3.1}$$

Here $\sum^{(\mathfrak{p})}$ indicates the summation over the basis of \mathfrak{p} .

As it stands, the element $D_{\mathfrak{g},\mathfrak{h}}$ gives us an operator on $W \otimes S_{\mathfrak{p}}$ for any \mathfrak{g} -representation W . To exhibit the structure of $D_{\mathfrak{g},\mathfrak{h}}$, we decompose $W \otimes S_{\mathfrak{p}}$ with respect to the \mathfrak{h} -action and denote by $M(\nu)$ the isotypic \mathfrak{h} -summand with highest weight ν .

Theorem 3.3 ([3]). *Suppose that W_{λ} is an irreducible \mathfrak{g} -representation with highest weight λ . The following formula holds*

$$D_{\mathfrak{g},\mathfrak{h}}^2 \Big|_{M(\nu)} = \|\lambda + \rho_G\|^2 - \|\nu + \rho_H\|^2.$$

3.2. Dirac operators on homogeneous spaces

Suppose now that $M = G/H$ is an orbit of the coadjoint G -action on \mathfrak{g}^* . It is known that M has a G -invariant complex structure so that it is compatible with the symplectic structure on M . Such a complex structure determines a H -equivariant

splitting: $\mathfrak{p} = \mathfrak{p}^+ \oplus \mathfrak{p}^-$. One can check that the spinor bundle associated to the complex structure on M is given by

$$S_M = G \times_H (\wedge \mathfrak{p}^-) \cong G \times_H (S_{\mathfrak{p}}^* \otimes \mathbb{C}_{\rho_G - \rho_H}), \quad (3.2)$$

and the canonical line bundle

$$K_M = G \times_H \mathbb{C}_{2 \cdot (\rho_G - \rho_H)}. \quad (3.3)$$

Hence, the Hilbert space $L^2(M, S_M)$ can be identified with

$$[L^2(G) \otimes S_{\mathfrak{p}}^* \otimes \mathbb{C}_{\rho_G - \rho_H}]^H \cong \bigoplus_{\lambda \in P_{G,+}} W_{\lambda} \otimes [W_{\lambda}^* \otimes S_{\mathfrak{p}}^* \otimes \mathbb{C}_{\rho_G - \rho_H}]^H, \quad (3.4)$$

where the isomorphism comes from the Peter–Weyl theorem.

We define Dirac operators on M in two different ways. First of all, the Levi–Civita connection ∇^{TM} lifts to a Hermitian connection ∇^{S_M} on S_M . In particular, the connection ∇^{S_M} is compatible with the Clifford action in the sense that

$$[X, \nabla_Y^{S_M}] = \nabla_Y^{TM} X, \quad X, Y \in TM.$$

We define a geometric Spin^c-Dirac operator by

$$D_{\text{geo}} = \sum_{a=1}^{\dim M} X_a \cdot \nabla_{X_a}^{S_M},$$

where $\{X_a\}_{i=1}^{\dim M}$ is an orthonormal basis of TM .

On the other hand, since the cubic Dirac operator $D_{\mathfrak{g}, \mathfrak{h}}$ is H -equivariant, it restricts to an operator on

$$[W_{\lambda}^* \otimes S_{\mathfrak{p}}^* \otimes \mathbb{C}_{\rho_G - \rho_H}]^H.$$

Tensoring the identity operator on each W_{λ} , and summing over W_{λ} , one obtains an operator D_{alg} on (3.4).

Lemma 3.4. *The difference between D_{alg} and D_{geo} on (3.4) is a bounded operator.*

Proof. We rewrite the geometric connection ∇^{S_M} as

$$\nabla_X^{S_M} = X + \frac{1}{2} \text{ad}^{\mathfrak{p}}(X).$$

It follows that

$$D_{\text{geo}} = \sum_{i=1}^{\dim \mathfrak{p}} X_a \otimes \left(X_a + \frac{1}{2} \text{ad}^{\mathfrak{p}}(X_a) \right).$$

We deduce the lemma by comparing the above with (3.1). An alternative proof can be found in [10, Chapter 9]. \square

3.3. Infinite dimensional case

The definitions of spin representation and cubic Dirac operator can be extended to the infinite-dimensional loop algebra $L\mathfrak{g}$.

Let now G be a compact, simple and simply connected Lie group with Lie algebra \mathfrak{g} . The loop algebra $L\mathfrak{g}$ carries an inner product defined by

$$B(X, Y) = \frac{1}{2\pi} \int_0^{2\pi} \langle X(\theta), Y(\theta) \rangle_{\mathfrak{g}} d\theta, \quad X, Y \in L\mathfrak{g}. \quad (3.5)$$

As in the finite dimensional case, we can define the Clifford algebra $\text{Cliff}(L\mathfrak{g})$, and its spin representation $S_{L\mathfrak{g}}$. Here $S_{L\mathfrak{g}}$ is a \mathbb{Z}_2 -graded complex Hilbert space, and also a positive energy LG -representation with highest weight

$$\rho_G = (0, \rho_G, h^{\vee}).$$

The explicit construction of $S_{L\mathfrak{g}}$ is given in [2]. For the general theory of Clifford algebras and representations on infinite dimensional Hilbert spaces we refer to [11].

Let us fix an orthonormal basis $\{X_a\}$ of \mathfrak{g} . For $n \in \mathbb{Z}$, we write X_a^n for the loop

$$s \mapsto e^{ins} \cdot X_a, \quad s \in \mathbb{R},$$

and $\mathfrak{g}(n)$ the vector space spanned by $\{X_a^n\}_{a=1}^{\dim \mathfrak{g}}$. The algebraic direct sum

$$L\mathfrak{g}^{\text{fin}} := \bigoplus_{n \in \mathbb{Z}} \mathfrak{g}(n)$$

is dense in $L\mathfrak{g}_{\mathbb{C}}$. The dense subspace of the spin representation $S_{L\mathfrak{g}}$ which consists of vectors with finite energy may be realized as

$$S_{L\mathfrak{g}}^{\text{fin}} = S_{\mathfrak{g}} \otimes \bigotimes_{k>0} \Lambda^*(\mathfrak{g}_{\mathbb{C}} z^k).$$

Definition 3.5. We define the *cubic Dirac operator*

$$D_{L\mathfrak{g}} \in U(L\mathfrak{g}) \otimes \text{Cliff}(L\mathfrak{g})$$

by

$$D_{L\mathfrak{g}} = \sum_{n \in \mathbb{Z}} \sum_{a=1}^{\dim \mathfrak{g}} \left(X_a^n \otimes X_a^{-n} + \frac{1}{3} \otimes \text{ad}^{L\mathfrak{g}}(X_a^n) \cdot X_a^{-n} \right). \tag{3.6}$$

Remark 3.6. The map

$$\text{ad}^{L\mathfrak{g}}(X) := \frac{1}{4} \sum_{n \in \mathbb{Z}} \sum_{a=1}^{\dim \mathfrak{g}} [X, X_a^n] \cdot X_a^{-n} \in \text{Cliff}(L\mathfrak{g}), \quad X \in L\mathfrak{g}^{\text{fin}} \tag{3.7}$$

is defined only on vectors with finite energy. To justify the infinite summation in (3.6) and (3.7), we refer to [8,4]. For any positive energy LG -representation V , $D_{L\mathfrak{g}}$ gives us an unbounded operator on $(V \otimes S_{L\mathfrak{g}})^{\text{fin}}$ which is a dense subspace of $V \otimes S_{L\mathfrak{g}}$.

Let \mathfrak{h} be an isotropy Lie algebra \mathfrak{h} of the coadjoint LG -action. We decompose the Lie algebra $L\mathfrak{g} = \mathfrak{p} \oplus \mathfrak{h}$, where $\mathfrak{p} = \mathfrak{h}^{\perp}$. By the multiplicative property of spin representation

$$S_{L\mathfrak{g}} = S_{\mathfrak{p}} \otimes S_{\mathfrak{h}},$$

where $S_{\mathfrak{p}}$ is a \mathbb{Z}_2 -graded irreducible representation of the Clifford algebra $\text{Cliff}(\mathfrak{p})$.

Definition 3.7. We define $D_{L\mathfrak{g},\mathfrak{h}} \in U(L\mathfrak{g}) \otimes \text{Cliff}(\mathfrak{p})$ by the formula

$$D_{L\mathfrak{g},\mathfrak{h}} = \sum_{a,n}^{(\mathfrak{p})} \left(X_a^n \otimes X_a^{-n} + \frac{1}{3} \otimes \text{ad}^{\mathfrak{p}}(X_a^n) \cdot X_a^{-n} \right),$$

where the summation ranges over a basis of \mathfrak{p} . When $\mathfrak{h} \subseteq \mathfrak{g} \subset L\mathfrak{g}$, $D_{L\mathfrak{g},\mathfrak{h}}$ is the same operator in [6].

We denote by \widehat{H} the central extension of H induced by the inclusion $H \hookrightarrow LG$ and the central extension \widehat{LG} . We decompose $V \otimes S_{\mathfrak{p}}$ with respect to the $S_{\text{rot}}^1 \times \widehat{H}$ -action and denote by $M(\mathbf{v})$ the isotypic component labeled by

$$\mathbf{v} = (n, \nu, k + h^{\vee}) \in \mathbb{Z} \times \Lambda^* \times \mathbb{Z}.$$

We have an analog of Theorem 3.3 for the infinite-dimensional case.

Theorem 3.8. Suppose that V_{λ} is an irreducible positive energy representation with highest weight $\lambda = (0, \lambda, k)$. If we restrict to the isotypic component $M(\mathbf{v})$ of $V_{\lambda} \otimes S_{\mathfrak{p}}$, we have that

$$D_{L\mathfrak{g},\mathfrak{h}}^2|_{M(\mathbf{v})} = \|\lambda + \rho_G\|^2 - \|\mathbf{v} + \rho_H\|^2,$$

where $\rho_G = (0, \rho_G, h^{\vee})$ and $\rho_H = (0, \rho_H, 0)$.

Proof. [8, Theorem 7.5]. \square

4. Hamiltonian LG -spaces and q -Hamiltonian G -spaces

The theory of q -Hamiltonian G -spaces was developed in [1]. It provides a finite-dimensional model for Hamiltonian LG -spaces. In this section, we begin by reviewing the basic definitions, and then discuss their cross-section theorems. We assume that G is a compact, simple and simply connected Lie group.

4.1. Basic definitions

Recall that a *Hamiltonian G -space* is a triple (M, ω, μ) , with ω the G -equivariant symplectic 2-form, and $\mu : M \rightarrow \mathfrak{g}^*$ the *moment map* satisfying that

$$\iota_{\xi^M} \omega = d\langle \mu, \xi \rangle, \quad \xi \in \mathfrak{g},$$

where ξ^M is the vector field on M induced by the infinitesimal action of ξ .

The above definition can be extended to the loop group setting. Let \mathcal{M} be an infinite-dimensional Banach manifold. We say that it is *weakly symplectic* if it is equipped with a closed 2-form $\omega \in \Omega^2(\mathcal{M})$ so that the induced map

$$\omega^\flat : T_m\mathcal{M} \rightarrow T_m^*\mathcal{M}$$

is injective.

Definition 4.1. A *Hamiltonian LG-space* is a weakly symplectic Banach manifold (\mathcal{M}, ω) together with a LG-action and a LG-equivariant map $\mu : \mathcal{M} \rightarrow L\mathfrak{g}^*$ so that $\iota_{\xi_M}\omega = d\langle \mu, \xi \rangle$ for all $\xi \in L\mathfrak{g}$.

For example, the coadjoint LG-orbit is a Hamiltonian LG-space, with moment map the inclusion.

Let $(\mathcal{M}, \omega, \mu)$ be a Hamiltonian LG-space. Since the based loop group ΩG acts freely on $L\mathfrak{g}^*$, it acts freely on \mathcal{M} as well by the equivariance of μ . We thus obtain a commuting square

$$\begin{array}{ccc} \mathcal{M} & \xrightarrow{\mu} & L\mathfrak{g}^* \\ \downarrow & & \text{Hol} \downarrow \\ M & \xrightarrow{\phi} & G \end{array}$$

where the quotient $M = \mathcal{M}/\Omega G$ is a finite-dimensional compact smooth manifold provided that μ is proper.

Alekseev–Malkin–Meinrenken [1] give a set of conditions a G-space M must satisfy in order to arise from a Hamiltonian LG-space by such a construction.

Choose an invariant inner product $\langle \cdot, \cdot \rangle_{\mathfrak{g}}$ on \mathfrak{g} and denote by $\theta^L, \theta^R \in \Omega^1(G, \mathfrak{g})$ the left and right invariant Maurer–Cartan forms on G and the *Cartan 3-form*

$$\chi = \frac{1}{12} \langle \theta^L, [\theta^L, \theta^L] \rangle_{\mathfrak{g}} = \frac{1}{12} \langle \theta^R, [\theta^R, \theta^R] \rangle_{\mathfrak{g}} \in \Omega^3(G).$$

Definition 4.2 ([1]). A *q-Hamiltonian G-space* is a compact G-manifold M , together with an equivariant 2-form ω , and an equivariant map $\phi : M \rightarrow G$ satisfying the following properties:

1. $d\omega = \phi^*\chi$;
2. $\iota_{\xi_M}\omega = \frac{1}{2} \langle \phi^*(\theta^L + \theta^R), \xi \rangle_{\mathfrak{g}}$ for all $\xi \in \mathfrak{g}$;
3. $\ker(\omega) \cap \ker(d\phi) = 0$.

We call ϕ the *group-valued moment map*.

According to [1, Theorem 8.3], there is a 1–1 correspondence between Hamiltonian LG-spaces with proper moment map and q-Hamiltonian G-spaces. One can always choose to work with infinite-dimensional Hamiltonian LG-spaces with more conventional definitions or to use finite dimensional q-Hamiltonian G-spaces. The counterparts of coadjoint orbits for q-Hamiltonian G-spaces are conjugacy classes \mathcal{C} in G with group-valued moment map the embedding $\mathcal{C} \hookrightarrow G$.

4.2. Cross-section theorems

The Hamiltonian LG-spaces and their equivalent finite-dimensional models behave in many respects like the usual Hamiltonian G-spaces. This is due to the existence of the cross-section theorem we shall now describe.

Let us first introduce a partial order of open faces of \mathfrak{A} by setting $\tau \preceq \sigma$ if $\tau \subseteq \bar{\sigma}$. The isotropy group $(LG)_{\xi}$ of the coadjoint LG-action on $L\mathfrak{g}^*$ depends only on the open face σ of \mathfrak{A} containing ξ and will be denoted by $(LG)_{\sigma}$ (note however $(LG)_{\sigma}$ will generally contain non-constant loops). One has that

$$\sigma \preceq \tau \Rightarrow (LG)_{\tau} \subseteq (LG)_{\sigma}.$$

In particular, $(LG)_0 = G$ and $(LG)_{\text{int}\mathfrak{A}} = T$.

We define a $(LG)_{\sigma}$ -invariant open subset of $(L\mathfrak{g})_{\sigma}^*$ by

$$\mathcal{A}_{\sigma} = (LG)_{\sigma} \cdot \bigcup_{\sigma \preceq \tau} \tau.$$

Note that \mathcal{A}_{σ} is a slice for all $\xi \in \sigma$ for the action of LG in the sense that

$$LG \times_{(LG)_{\sigma}} \mathcal{A}_{\sigma} \rightarrow LG \cdot \mathcal{A}_{\sigma}$$

is a diffeomorphism of Banach manifolds.

Theorem 4.3. Let $(\mathcal{M}, \omega, \mu)$ be a Hamiltonian LG-space with proper moment map. For every open face σ of \mathfrak{A} , the cross-section

$$\mathcal{V}_{\sigma} = \mu^{-1}(\mathcal{A}_{\sigma})$$

is a finite-dimensional symplectic submanifold with Hamiltonian $(LG)_{\sigma}$ -action. The restriction of $\mu|_{\mathcal{V}_{\sigma}}$ is a moment map of the $(\widehat{LG})_{\sigma}$ -action (the central circle acts trivially on \mathcal{V}_{σ}).

Proof. [12, Theorem 4.8]. □

The symplectic cross-section theorem carries over to q -Hamiltonian G -spaces. The centralizer $G_{\exp(\xi)}$ with $\xi \in \mathfrak{A}$ is isomorphic to $(LG)_{\xi}$ and it depends only on the open face σ of \mathfrak{A} containing ξ . We denoted it by G_{σ} . The subset

$$A_{\sigma} = \text{Ad}(G_{\sigma}) \cdot \exp\left(\bigcup_{\sigma \leq \tau} \tau\right) \subset G_{\sigma} \subset G$$

is smooth and is a slice for the $\text{Ad}(G)$ -action at points in σ .

Theorem 4.4. *Let (M, ω, ϕ) be a q -Hamiltonian G -space. The cross-section*

$$V_{\sigma} = \phi^{-1}(A_{\sigma}) \tag{4.1}$$

is a smooth G_{σ} -invariant submanifold and

$$G \times_{G_{\sigma}} V_{\sigma} \cong G \cdot V_{\sigma}$$

is a G -invariant open subset of M . Moreover, V_{σ} is a q -Hamiltonian G_{σ} -space with the restriction of ϕ as the group-valued moment map.

Proof. [1, Proposition 7.1]. □

Remark 4.5. It is important to point out that if we identify $G_{\sigma} \cong (LG)_{\sigma}$, the two cross-sections

$$\mathcal{V}_{\sigma} \subset \mathcal{M}, \quad V_{\sigma} \subset M$$

are equivariantly diffeomorphic. In particular, every V_{σ} is also a Hamiltonian G_{σ} -space.

5. Twisted spinor bundle and twisted pre-quantum bundle

In this section, we construct twisted spinor bundle and twisted pre-quantum bundles on q -Hamiltonian G -spaces.

5.1. Construction of the twisted spinor bundle

Let G be a compact, simple, and simply connected Lie group. Let (M, ω) be a q -Hamiltonian G -space and \mathcal{M} its corresponding Hamiltonian LG -space.

We first replace the cross-sections V_{σ} in (4.1) with smaller open subsets. To be more precise, for every vertex σ of \mathfrak{A} , let Y_{σ} be a G_{σ} -invariant, open subset of V_{σ} so that $\overline{Y_{\sigma}} \subset V_{\sigma}$, and

$$M/G \subseteq \bigcup_{\sigma, \dim\sigma=0} Y_{\sigma}.$$

Then we form an open cover of M by

$$\{U_{\sigma} = G \times_{G_{\sigma}} Y_{\sigma}\}_{\sigma, \dim\sigma=0}.$$

For all open faces τ of \mathfrak{A} with $\dim\tau > 0$, we define

$$Y_{\tau} = \bigcap_{\sigma \leq \tau, \dim\sigma=0} Y_{\sigma}, \quad U_{\tau} = G \times_{G_{\tau}} Y_{\tau}.$$

Remark at this point that each Y_{τ} is a Hamiltonian G_{τ} -space and admits G_{τ} -invariant almost complex structures.

Lemma 5.1. *There exists a collection of G_{σ} -invariant almost complex structures on the collection of Y_{σ} such that the embedding*

$$Y_{\tau} \hookrightarrow Y_{\sigma}, \quad \sigma < \tau$$

is almost complex. In addition, any two collections of almost complex structures with the required properties are homotopic. We denote by $S_{Y_{\sigma}}$ the spinor bundle on Y_{σ} associated to the almost complex structures.

Proof. [13, Lemma 3.2]. □

Let $\pi : U_{\sigma} \rightarrow G/G_{\sigma}$ be the projection. The tangent bundle decomposes G -equivariantly

$$TU_{\sigma} \cong \pi^*T(G/G_{\sigma}) \oplus G \times_{G_{\sigma}} TY_{\sigma}.$$

The base manifold G/G_{σ} is a conjugacy class and might not have a G -equivariant Spin^c -structure in general. Thus the total space U_{σ} does not have to be Spin^c either.

On the other hand, the coadjoint LG -orbit $\mathcal{O} = LG/(LG)_\sigma$ is a complex manifold. By the discussion in [13], the weight

$$2 \cdot (\rho_G - \rho_\sigma, h^\vee)$$

is fixed by $(\widehat{LG})_\sigma$ and the tensor product

$$S_{\mathfrak{Lg}/(\mathfrak{Lg})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)}$$

is a $(LG)_\sigma$ -space. The associated spinor bundle and canonical line bundle on \mathcal{O} are given by

$$S_{\mathcal{O}} = \widehat{LG} \times_{(\widehat{LG})_\sigma} (S_{\mathfrak{Lg}/(\mathfrak{Lg})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)})$$

and

$$K_{\mathcal{O}} = \widehat{LG} \times_{(\widehat{LG})_\sigma} (\mathbb{C}_{2 \cdot (\rho_G - \rho_\sigma, h^\vee)}).$$

One can compare them with (3.2) and (3.3).

Motivated by the above, we define a bundle of Hilbert space on U_σ by

$$S_{U_\sigma}^{\text{spin}} = G \times_{G_\sigma} (S_{\mathfrak{Lg}/(\mathfrak{Lg})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)} \otimes S_{Y_\sigma}), \tag{5.1}$$

where G_σ acts on

$$S_{\mathfrak{Lg}/(\mathfrak{Lg})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)}$$

factoring through the identification $G_\sigma \cong (LG)_\sigma$. In addition, we equip $S_{U_\sigma}^{\text{spin}}$ with a \mathbb{Z}_2 -grading induced by that on $S_{\mathfrak{Lg}/(\mathfrak{Lg})_\sigma}^*$ and S_{Y_σ} . To sum up, we obtain a collection of G -equivariant bundles of \mathbb{Z}_2 -graded Hilbert space $\{S_{U_\sigma}^{\text{spin}}\}$. We next show that such a collection of bundles can be glued together.

Lemma 5.2. *For any $\sigma < \tau$, the normal bundle ν_τ^σ of $Y_\tau \hookrightarrow Y_\sigma$ has a G_τ -equivariant almost complex structure with spinor bundle isomorphic to*

$$S_{(\mathfrak{Lg})_\sigma/(\mathfrak{Lg})_\tau}^* \otimes \mathbb{C}_{(\rho_\sigma - \rho_\tau)}, \tag{5.2}$$

where ρ_τ, ρ_σ are the half-sums of positive roots for G_τ, G_σ respectively.

Proof. By the cross-section theorem, the normal bundle ν_τ^σ is isomorphic to the trivial bundle

$$\mathfrak{g}_\sigma / \mathfrak{g}_\tau \cong (\mathfrak{Lg})_\sigma / (\mathfrak{Lg})_\tau$$

with equivariant almost complex structure. In fact, let $\mathfrak{R}_\sigma, \mathfrak{R}_\tau$ be compatible sets of positive roots for G_σ and G_τ . Then

$$\mathfrak{g}_\sigma / \mathfrak{g}_\tau = \bigoplus_{\alpha \in \mathfrak{R}_\sigma \setminus \mathfrak{R}_\tau} \mathbb{C}_\alpha$$

and

$$\det_{\mathbb{C}}(\nu_\tau^\sigma) = \bigotimes_{\alpha \in \mathfrak{R}_\sigma \setminus \mathfrak{R}_\tau} \mathbb{C}_\alpha = \mathbb{C}_{2(\rho_\sigma - \rho_\tau)}. \quad \square$$

Lemma 5.3. *There are canonical isomorphisms*

$$\Psi_{\tau, \sigma} : S_{U_\tau}^{\text{spin}} \cong S_{U_\sigma}^{\text{spin}}|_{U_\tau}, \quad \sigma < \tau$$

and they automatically satisfy the cocycle condition.

Proof. By the above lemma, we have that

$$S_{Y_\sigma}|_{Y_\tau} \cong S_{(\mathfrak{Lg})_\sigma/(\mathfrak{Lg})_\tau}^* \otimes \mathbb{C}_{(\rho_\sigma - \rho_\tau)} \otimes S_{Y_\tau}.$$

The claim follows from the construction. \square

Definition 5.4. We define the *twisted spinor bundle* S^{spin} to be the G -equivariant bundle of \mathbb{Z}_2 -graded Hilbert space over M with the property that

$$S^{\text{spin}}|_{U_\sigma} \cong S_{U_\sigma}^{\text{spin}},$$

for all vertexes σ of \mathfrak{A} .

The twisted spinor bundle is determined by the choice of almost complex structures on subsets $\{Y_\sigma\}$. By Lemma 5.1, all the choices are homotopic. Hence, the twisted spinor bundle S^{spin} is unique up to homotopy.

5.2. Construction of the twisted pre-quantum bundle

Let us recall that the pre-quantum line bundle of a symplectic manifold is traditionally defined to be a line bundle whose first Chern class is an integral lift of the symplectic 2-form.

Definition 5.5. We say that a Hamiltonian LG -space \mathcal{M} is *pre-quantizable at level k* ($k > 0$) if there exists a \widehat{LG} -equivariant line bundle

$$\mathcal{L} \rightarrow \mathcal{M}$$

such that the central circle acts with weight k and the first Chern class $c_1(\mathcal{L})$ equals the symplectic 2-form on \mathcal{M} .

Because the pre-quantum line bundle \mathcal{L} is \widehat{LG} -equivariant instead of LG -equivariant, it might not descend to an actual line bundle on its corresponding q -Hamiltonian G -space M .

Remark 5.6. The 2-form ω for a q -Hamiltonian G -space M is not closed in general. Instead the condition $d\omega = \phi^*\chi$ and the fact that χ is a closed 3-form imply that the pair (ω, χ) defines a cocycle for the relative de Rham theory (see [14, Appendix B] for a reference). We denote by $[(\omega, \chi)] \in H^3(\phi, \mathbb{R})$ its cohomology class. We say that a q -Hamiltonian G -space (M, ω, ϕ) is *pre-quantizable at level k* if $k \cdot [(\omega, \chi)]$ is integral. By the 1–1 correspondence between q -Hamiltonian G -spaces and Hamiltonian LG -spaces, their pre-quantum conditions are equivalent.

Let Y_σ be the cross-section defined before. We identify it as a subset in \mathcal{M} . If \mathcal{M} has a pre-quantum line bundle \mathcal{L} at level k , then there exists a $(LG)_\sigma$ -equivariant line bundle obtained by restriction

$$L_{Y_\sigma} = \mathcal{L}|_{Y_\sigma} \rightarrow Y_\sigma,$$

on which the central circle acts with weight k . The collection of line bundles $\{L_{Y_\sigma}\}$ satisfy a compatibility condition in the sense that

$$(LG)_\sigma \times_{(LG)_\tau} L_{Y_\tau} \cong L_{Y_\sigma}|_{Y_\sigma^\sigma}, \quad \sigma < \tau \tag{5.3}$$

where

$$Y_\tau^\sigma = (LG)_\sigma \times_{(LG)_\tau} Y_\tau$$

is a $(LG)_\sigma$ -invariant open subset of Y_σ .

For any irreducible positive energy representation V_λ at level k , we denote V_λ^* its dual. Comparing to (5.1), we define a bundle of Hilbert space on

$$U_\sigma = G \times_{G_\sigma} Y_\sigma$$

by

$$S_{\lambda, U_\sigma}^{\text{pre}} = G \times_{G_\sigma} (V_\lambda^* \otimes L_{Y_\sigma}). \tag{5.4}$$

Here V_λ^* is the dual of V_λ . The central circle of $(\widehat{LG})_\sigma$ acts trivially on the tensor product $V_\lambda^* \otimes L_{Y_\sigma}$, and G_σ acts factoring through $G_\sigma \cong (LG)_\sigma$.

By the compatibility condition (5.3), there are canonical isomorphisms

$$\Psi_{\tau, \sigma} : S_{\lambda, U_\tau}^{\text{pre}} \cong S_{\lambda, U_\sigma}^{\text{pre}}|_{U_\tau}, \quad \sigma < \tau,$$

satisfying the cocycle condition. Thus, they can be glued together. We define S_λ^{pre} the unique G -equivariant bundle of Hilbert space over M with the property that

$$S_\lambda^{\text{pre}}|_{U_\sigma} = S_{\lambda, U_\sigma}^{\text{pre}}.$$

Definition 5.7. We define the *twisted pre-quantum bundle* by

$$S^{\text{pre}} := \bigoplus_{\lambda \in P_{k,+}} V_\lambda \otimes S_\lambda^{\text{pre}}. \tag{5.5}$$

It is a $\widehat{LG} \times G$ -equivariant bundle of Hilbert space over the q -Hamiltonian G -space M .

Remark 5.8. There is a global construction of the twisted pre-quantum bundle. Let us introduce a $\widehat{LG} \times \widehat{LG}$ -space:

$$H_{\text{wzw}, k} = \bigoplus_{\lambda \in P_{k,+}} V_\lambda \otimes V_\lambda^*.$$

This is the so-called Hilbert space of the Wess–Zumino–Witten model (see [15] for a reference). The aim of this paper is not to justify this choice of the Hilbert space, but morally one can consider it as the analog of $L^2(G)$ for loop groups, in the spirit

of the Peter–Weyl decomposition of $L^2(G)$. The tensor product

$$(\mathcal{M} \times H_{wz, k}) \otimes \mathcal{L} = \bigoplus_{\lambda \in P_{k,+}} V_\lambda \otimes (V_\lambda^* \otimes \mathcal{L}) \rightarrow \mathcal{M}$$

is a $\widehat{LG} \times LG$ equivariant bundle. We obtain a $\widehat{LG} \times G$ -equivariant bundle of Hilbert space by taking its ΩG -invariant part

$$S^{\text{pre}} = V_\lambda \otimes [V_\lambda^* \otimes \mathcal{L}]^{\Omega G} \rightarrow M = \mathcal{M}/\Omega G.$$

6. Dirac operators on q -Hamiltonian G -spaces

With the twisted spinor bundle and twisted pre-quantum bundle defined in last section, we now proceed to construct Hilbert spaces and Dirac operators. We keep the same notations as in the last section. Let G be a compact, simple and simply connected Lie group, and M a pre-quantizable q -Hamiltonian G -space at level k .

6.1. Dirac operators on cross-sections

The idea of constructing Dirac operators is that we first define Dirac operators on local cross-sections, and then patch them together using partition of unity.

Let $\{Y_\tau\}$ be the collection of cross-sections defined in last section and open subsets $U_\tau = G \times_{G_\tau} Y_\tau$. Fixing an irreducible positive energy LG -representation V_λ , we define

$$[\Gamma_c^\infty(U_\tau, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G$$

the space of G -invariant, smooth sections of $S^{\text{spin}} \otimes S_\lambda^{\text{pre}}$ with compact support in U_τ , with norm given by

$$\|s\|^2 := \int_{U_\tau} \langle s(m), s(m) \rangle dm.$$

Lemma 6.1. *We have that*

$$[\Gamma_c^\infty(U_\tau, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G \cong [V_\lambda^* \otimes S_{L\mathfrak{g}/(L\mathfrak{g})_\tau}^* \otimes \mathbb{C}_{(\rho_G - \rho_\tau, h^\vee)} \otimes \Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau})]^{G_\tau}. \quad (6.1)$$

Proof. The assertion follows immediately from the isomorphisms (5.1) and (5.4). \square

Let

$$D_{\text{alg}} \in U(L\mathfrak{g}) \otimes \text{Cliff}(L\mathfrak{g}/(L\mathfrak{g})_\tau)$$

be the cubic Dirac operator acting on $V_\lambda^* \otimes S_{L\mathfrak{g}/(L\mathfrak{g})_\tau}^*$ and D_{geo} the equivariant geometric Spin^c -Dirac operator on $\Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau})$. Here we choose D_{geo} so that it is symmetric. Since the sum

$$D_{\text{alg}} \otimes 1 + 1 \otimes D_{\text{geo}}$$

is equivariant, it descends to the G_τ -invariant part. That is, we obtain a collection of operators

$$[D_{\text{alg}} \otimes 1 + 1 \otimes D_{\text{geo}}]^{G_\tau}$$

on

$$[V_\lambda^* \otimes S_p^* \otimes \mathbb{C}_{(\rho_G - \rho_\tau, h^\vee)} \otimes \Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau})]^{G_\tau},$$

and thus an operator on

$$[\Gamma_c^\infty(U_\tau, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G.$$

We denote it by D_{U_τ} . By definition, every D_{U_τ} is an unbounded operator. Since D_{alg} is defined only on vectors with finite energy, the domain of D_{U_τ} is given by

$$[(V_\lambda^* \otimes S_{L\mathfrak{g}/(L\mathfrak{g})_\tau}^*)^{\text{fin}} \otimes \mathbb{C}_{(\rho_G - \rho_\tau, h^\vee)} \otimes \Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau})]^{G_\tau},$$

which is a dense subspace.

Suppose that σ_1, σ_2 are two vertexes of \mathfrak{A} and τ an open face of \mathfrak{A} so that $\sigma_1 \leq \tau$ and $\sigma_2 \leq \tau$. In particular, one has that

$$U_\tau \subseteq U_{\sigma_1}, \quad U_\tau \subseteq U_{\sigma_2}.$$

Proposition 6.2. *If we restrict to*

$$[\Gamma_c^\infty(U_\tau, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G,$$

the difference between $D_{U_{\sigma_1}}$ and $D_{U_{\sigma_2}}$ is a bounded operator.

Proof. Since $U_\tau = G \times_{G_\tau} Y_\tau$, Lemma 6.1 gives us an isomorphism:

$$[\Gamma_c^\infty(U_\tau, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G \cong [V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_\tau}^* \otimes \mathbb{C}_{(\rho_G - \rho_\tau, h^\vee)} \otimes \Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau})]^{G_\tau}. \tag{6.2}$$

The Dirac operator D_{U_τ} is defined as the combination of a cubic Dirac operator on

$$V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_\tau}^* \cong V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_{\sigma_1}}^* \otimes S_{\mathfrak{g}_{\sigma_1}/\mathfrak{g}_\tau}^*$$

and a Spin^c-Dirac operator on

$$\Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau}).$$

On the other hand, we have that

$$U_\tau \cong G \times_{G_{\sigma_1}} Y_\tau^{\sigma_1}, \quad Y_\tau^{\sigma_1} = G_{\sigma_1} \times_{G_\tau} Y_\tau.$$

Applying Lemma 6.1 again,

$$[\Gamma_c^\infty(U_\tau, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G \cong [V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_{\sigma_1}}^* \otimes \mathbb{C}_{(\rho - \rho_{\sigma_1}, h^\vee)} \otimes \Gamma_c^\infty(Y_\tau^{\sigma_1}, S_{Y_{\sigma_1}} \otimes L_{Y_{\sigma_1}})]^{G_{\sigma_1}}. \tag{6.3}$$

Under the isomorphism (6.3), the operator $D_{U_{\sigma_1}}$ decomposes into the sum of a cubic Dirac operator on

$$V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_{\sigma_1}}^*$$

and a Spin^c-Dirac operator on

$$\Gamma_c^\infty(Y_\tau^{\sigma_1}, S_{Y_{\sigma_1}} \otimes L_{Y_{\sigma_1}}).$$

It follows immediately that the two operators D_{U_τ} and $D_{U_{\sigma_1}}$ are identical on

$$V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_{\sigma_1}}^*.$$

Note that the spinor bundle

$$S_{Y_{\sigma_1}}|_{Y_\tau^{\sigma_1}} \cong G_{\sigma_1} \times_{G_\tau} (S_{Y_\tau} \otimes S_{\mathfrak{g}_{\sigma_1}/\mathfrak{g}_\tau}^* \otimes \mathbb{C}_{\rho_{\sigma_1} - \rho_\tau}),$$

and the pre-quantum line bundle

$$L_{Y_{\sigma_1}}|_{Y_\tau^{\sigma_1}} \cong G_{\sigma_1} \times_{G_\tau} L_\tau.$$

Therefore, the space

$$\Gamma_c^\infty(Y_\tau^{\sigma_1}, S_{Y_{\sigma_1}} \otimes L_{Y_{\sigma_1}})$$

decomposes into:

$$[(C^\infty(G_{\sigma_1}) \otimes S_{\mathfrak{g}_{\sigma_1}/\mathfrak{g}_\tau}^*) \otimes \mathbb{C}_{\rho_{\sigma_1} - \rho_\tau} \otimes \Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau})]^{G_\tau}.$$

By definition, both D_{U_τ} and $D_{U_{\sigma_1}}$ act as Spin^c-Dirac operators on

$$\Gamma_c^\infty(Y_\tau, S_{Y_\tau} \otimes L_{Y_\tau});$$

while on the factor

$$C^\infty(G_{\sigma_1}) \otimes S_{\mathfrak{g}_{\sigma_1}/\mathfrak{g}_\tau}^*,$$

D_{U_τ} acts as the cubic Dirac operator, and $D_{U_{\sigma_1}}$ acts as the Spin^c-Dirac operator. By Lemma 3.4, their difference is bounded.

We just show that the difference between $D_{U_{\sigma_1}}$ and D_{U_τ} is a bounded operator on the overlap U_τ . Similarly one can show that the difference between $D_{U_{\sigma_2}}$ and D_{U_τ} is bounded as well. This completes the proof. \square

6.2. Construction of the dirac operator and main theorem

For a fixed irreducible positive energy representation V_λ at level k , we define

$$\mathcal{H}_\lambda = [L^2(M, S^{\text{spin}} \otimes S_\lambda^{\text{pre}})]^G.$$

The \mathbb{Z}_2 -grading on the twisted spinor bundle S^{spin} equips \mathcal{H}_λ with a \mathbb{Z}_2 -grading. That is

$$\mathcal{H}_\lambda = \mathcal{H}_\lambda^+ \oplus \mathcal{H}_\lambda^-.$$

Select a G -invariant, smooth partition of unity $\{f_\sigma^2\}$ which is subordinate to the cover

$$\{U_\sigma = G \times_{G_\sigma} Y_\sigma\}_{\sigma, \dim \sigma = 0}.$$

We define a Dirac operator \mathcal{D} on \mathcal{H} by the formula:

$$\mathcal{D} = \sum_{\sigma, \dim \sigma = 0} f_\sigma \cdot D_{U_\sigma} \cdot f_\sigma.$$

Proposition 6.3. *The Dirac operator \mathcal{D} does not depend on the choice of $\{f_\sigma^2\}$ up to homotopy.*

Proof. Let $\{f'_\sigma\}$ be another partition of unity. We define a Dirac operator \mathcal{D}' on \mathcal{H}_λ by

$$\mathcal{D}' = \sum_{\sigma, \dim \sigma = 0} f'_\sigma \cdot D_{U_\sigma} \cdot f'_\sigma.$$

It suffices to show that

$$\mathcal{D}' - \mathcal{D} \in \mathbb{B}(\mathcal{H}_\lambda).$$

We compute that

$$\begin{aligned} \mathcal{D} &= \sum_{\sigma} f_\sigma \cdot D_{U_\sigma} \cdot f_\sigma = \sum_{\sigma, \tau} f_\sigma \cdot D_{U_\sigma} \cdot f'_\tau \cdot f_\sigma \\ &= \sum_{\sigma, \tau} (f'_\tau \cdot f_\sigma \cdot D_{U_\sigma} \cdot f_\sigma \cdot f'_\tau + f_\sigma \cdot [D_{U_\sigma}, f'_\tau] \cdot f'_\tau \cdot f_\sigma) \\ &= \sum_{\sigma, \tau} (f'_\tau \cdot D_{U_\sigma} \cdot f_\sigma^2 \cdot f'_\tau + f_\sigma \cdot [D_{U_\sigma}, f'_\tau] \cdot f'_\tau \cdot f_\sigma + f'_\tau \cdot [D_{U_\sigma}, f_\sigma] \cdot f'_\tau \cdot f_\sigma). \end{aligned} \quad (6.4)$$

Since the functions f_σ are G -invariant, they commute with the cubic Dirac operator D_{alg} . Their commutators with Spin^c -Dirac operators:

$$c(df_\sigma) = [f_\sigma, D_{\text{geo}}]$$

are all bounded. It shows that

$$[D_{U_\sigma}, f'_\tau], \quad [D_{U_\sigma}, f_\sigma] \in \mathbb{B}(\mathcal{H}_\lambda). \quad (6.5)$$

In addition, by [Proposition 6.2](#),

$$f'_\tau \cdot (D_{U_\sigma} - D_{U_\tau}) \cdot f_\sigma^2 \cdot f'_\tau$$

are bounded for all σ, τ . Therefore,

$$\begin{aligned} \sum_{\sigma, \tau} f'_\tau \cdot D_{U_\sigma} \cdot f_\sigma^2 \cdot f'_\tau &= \sum_{\tau, \sigma} f'_\tau \cdot D_{U_\tau} \cdot f_\sigma^2 \cdot f'_\tau + \sum_{\tau, \sigma} f'_\tau \cdot (D_{U_\sigma} - D_{U_\tau}) \cdot f_\sigma^2 \cdot f'_\tau \\ &= \mathcal{D}' + \sum_{\tau, \sigma} f'_\tau \cdot (D_{U_\sigma} - D_{U_\tau}) \cdot f_\sigma^2 \cdot f'_\tau. \end{aligned} \quad (6.6)$$

The proposition follows from (6.4)–(6.6). \square

Proposition 6.4. *The Dirac operator \mathcal{D} is essentially self-adjoint.*

Proof. It is enough to show that each $f_\sigma \cdot D_{U_\sigma} \cdot f_\sigma$ is essentially self-adjoint. Let us write

$$D_{U_\sigma} = [D_{\text{alg}} \otimes 1 + 1 \otimes D_{\text{geo}}]^{G_\sigma}.$$

The algebraic part D_{alg} is certainly self-adjoint, and the geometric part, D_{geo} is a symmetric Spin^c -Dirac operator on Y_σ . Recall that every symmetric Dirac operator on a complete manifold is essentially self-adjoint. But it does not apply directly to our case since Y_σ is not complete. Nevertheless, we can get around this by the following trick.

Since the function f_σ has compact support in Y_σ , we can find a smaller subset \tilde{Y}_σ so that

$$\text{Supp} f_\sigma \subset \tilde{Y}_\sigma \subset Y_\sigma$$

and the closure of \tilde{Y}_σ is contained in Y_σ . We denote by g the metric on Y_σ and χ a positive function on Y_σ so that $\chi|_{\tilde{Y}_\sigma} \equiv 1$ and $\chi(m)$ tends to infinity as m tends to the boundary of Y_σ . Under the rescaled metric

$$g^\chi(\cdot, \cdot) = \chi^2 \cdot g(\cdot, \cdot),$$

the manifold Y_σ becomes complete.

Let \tilde{D}_{geo} be an essentially self-adjoint Spin^c -Dirac operator on the complete manifold Y_σ . Consider a new Dirac operator

$$\tilde{D}_{U_\sigma} := [D_{\text{alg}} \otimes 1 + 1 \otimes \tilde{D}_{\text{geo}}]^{G_\sigma},$$

which is essentially self-adjoint. Since the metric remains the same within \tilde{Y}_σ , we have that

$$\tilde{D}_{U_\sigma} s = D_{U_\sigma} s$$

for all s with $\text{Supp} s \subseteq \tilde{Y}_\sigma$. Thus,

$$f_\sigma \cdot \tilde{D}_{U_\sigma} \cdot f_\sigma$$

can be viewed as an operator on \mathcal{H}_λ , and

$$f_\sigma \cdot D_{U_\sigma} \cdot f_\sigma = f_\sigma \cdot \tilde{D}_{U_\sigma} \cdot f_\sigma$$

is essentially self-adjoint. \square

We define an operator by functional calculus,

$$\mathcal{F} = \frac{\mathcal{D}}{\sqrt{1 + \mathcal{D}^2}}.$$

Because \mathcal{D} has a dense domain in \mathcal{H}_λ , \mathcal{F} extends to a self-adjoint bounded operator on \mathcal{H}_λ , which anti-commutes with the \mathbb{Z}_2 -grading on \mathcal{H}_λ .

Theorem 6.5. *The bounded operator \mathcal{F} is Fredholm on \mathcal{H}_λ . We define its Fredholm index by*

$$\text{Ind}(\mathcal{D})_\lambda := \ker(\mathcal{F}) \cap \mathcal{H}_\lambda^+ - \ker(\mathcal{F}) \cap \mathcal{H}_\lambda^- \in \mathbb{Z}.$$

By [Proposition 6.3](#), the index is independent of the choice of partition of unity. Let us define a Hilbert space

$$\mathcal{H} = [L^2(M, S^{\text{spin}} \otimes S^{\text{pre}})]^G \cong \bigoplus_{\lambda \in P_{k,+}} V_\lambda \otimes \mathcal{H}_\lambda,$$

and a Dirac operator

$$\mathcal{D}_M = \bigoplus_{\lambda \in P_{k,+}} 1 \otimes \mathcal{D}|_{\mathcal{H}_\lambda}.$$

Definition 6.6. Let (M, ω) be a q -Hamiltonian G -space. If it is pre-quantizable at level k , we define its *quantization*

$$Q(M) = \text{Ind}(\mathcal{D}_M) = \bigoplus_{\lambda \in P_{k,+}} \text{Ind}(\mathcal{D})_\lambda \cdot V_\lambda \in R_k(LG).$$

This generalizes Bott’s Spin^c -quantization for Hamiltonian G -spaces to q -Hamiltonian G -spaces.

Remark 6.7. In [\[14\]](#) Meinrenken develops a quantization from pre-quantized q -Hamiltonian G -spaces to the equivariant twisted K -homology of G using push-forward maps. By the work of Freed–Hopkins–Teleman [\[7\]](#) the equivariant twisted K -homology of G at level k is isomorphic to fusion ring of loop group $R_k(LG)$.

Example 6.8. Let

$$\mathcal{C} = G \cdot \exp(\xi) \cong G/H, \quad \xi \in \mathfrak{A}$$

be a conjugacy class. We assume that it is pre-quantizable at level k . By [Remark 5.6](#),

$$(\eta, k) = (k \cdot \xi, k) \in \hat{\mathfrak{t}}_+^*$$

is an integral weight. By definition, the twisted spinor bundle and twisted pre-quantum bundle are given by

$$S^{\text{spin}} = G \times_H (S_{L\mathfrak{g}/\mathfrak{h}}^* \otimes \mathbb{C}_{(\rho_G - \rho_H, h^\vee)}), \quad S^{\text{pre}} = G \times_H (H_{\text{wzw}, k} \otimes \mathbb{C}_{(\eta, k)}).$$

The Hilbert space

$$[L^2(\mathcal{C}, S^{\text{spin}} \otimes S^{\text{pre}})]^G \cong \bigoplus_{\lambda \in P_{k,+}} V_\lambda \otimes (V_\lambda^* \otimes S_{L\mathfrak{g}/\mathfrak{h}}^* \otimes \mathbb{C}_{(\rho_G - \rho_H + \eta, k + h^\vee)})^H.$$

By [Theorem 3.8](#), one can calculate that

$$Q(\mathcal{C}) = \text{ind}(\mathcal{D}_M) = V_\eta \in R_k(LG).$$

This is an algebraic version of the Borel–Weil construction for loop groups [\[2\]](#).

6.3. Proof of Theorem 6.5

We will prove the main theorem in this subsection. Let us begin with a lemma.

Lemma 6.9. Let $\mathcal{D}, \mathcal{D}'$ be two self-adjoint unbounded operators on a Hilbert space \mathcal{H} such that $\mathcal{D} - \mathcal{D}' \in \mathbb{B}(\mathcal{H})$. For $\alpha = 0, 1$, if

$$\mathcal{D}^\alpha \cdot (1 + \mathcal{D}^2)^{-1} \in \mathbb{K}(\mathcal{H}),$$

then

$$\mathcal{D}'^\alpha \cdot (1 + \mathcal{D}'^2)^{-1} \in \mathbb{K}(\mathcal{H}).$$

Proof. Suppose that $\mathcal{D}' = \mathcal{D} + B$ with $B \in \mathbb{B}(\mathcal{H})$. By straightforward calculation,

$$\begin{aligned} (1 + \mathcal{D}^2)^{-1} - (1 + \mathcal{D}'^2)^{-1} &= (1 + \mathcal{D}^2)^{-1} \cdot (\mathcal{D}'^2 - \mathcal{D}^2) \cdot (1 + \mathcal{D}'^2)^{-1} \\ &= (1 + \mathcal{D}^2)^{-1} \cdot (\mathcal{D} \cdot B + B \cdot \mathcal{D}') \cdot (1 + \mathcal{D}'^2)^{-1}. \end{aligned} \tag{6.7}$$

Note that the product of a bounded operator and a compact operator is again a compact operator. By the assumption, we deduce that

$$(1 + \mathcal{D}'^2)^{-1} \in \mathbb{K}(\mathcal{H}).$$

We next consider

$$\begin{aligned} \mathcal{D} \cdot (1 + \mathcal{D}^2)^{-1} - \mathcal{D} \cdot (1 + \mathcal{D}'^2)^{-1} &= \mathcal{D} \cdot (1 + \mathcal{D}^2)^{-1} \cdot (\mathcal{D}'^2 - \mathcal{D}^2) \cdot (1 + \mathcal{D}'^2)^{-1} \\ &= \mathcal{D} \cdot (1 + \mathcal{D}^2)^{-1} \cdot (\mathcal{D} \cdot B + B \cdot \mathcal{D}') \cdot (1 + \mathcal{D}'^2)^{-1}. \end{aligned} \tag{6.8}$$

For the same reason as above, it follows that

$$\mathcal{D} \cdot (1 + \mathcal{D}'^2)^{-1} \in \mathbb{K}(\mathcal{H}).$$

Because $\mathcal{D}' - \mathcal{D}$ is bounded, we conclude that

$$\mathcal{D}' \cdot (1 + \mathcal{D}'^2)^{-1} \in \mathbb{K}(\mathcal{H}). \quad \square$$

Fix a vertex σ of \mathfrak{A} . Let us write

$$D_{U_\sigma} = [D_{\text{alg}} \otimes 1 + 1 \otimes D_{\text{geo}}]^{G_\sigma},$$

acting on

$$[V_\lambda^* \otimes S_{L_{\mathfrak{g}}/(L_{\mathfrak{g}})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)} \otimes \Gamma_c^\infty(Y_\sigma, S_{Y_\sigma} \otimes L_{Y_\sigma})]^{G_\sigma}.$$

By the trick used in Proposition 6.4, we can assume that Y_σ is complete and D_{geo} is an essentially self-adjoint Spin^c -Dirac operator on Y_σ .

Lemma 6.10. For $\alpha = 0, 1$, one has that

$$f_\sigma \cdot D_{U_\sigma}^\alpha \cdot (1 + D_{U_\sigma}^2)^{-1} \cdot f_\sigma \in \mathbb{K}(\mathcal{H}_\lambda).$$

Proof. We decompose

$$L^2(Y_\sigma, S_{Y_\sigma} \otimes L_{Y_\sigma})$$

with respect to the $S_{\text{rot}}^1 \times \widehat{G}_\sigma$ -action and denote by $M(\mathfrak{v})$ the isotypic component labeled by

$$\mathfrak{v} = (n, \nu, k) \in \mathbb{Z} \times \Lambda^* \times \mathbb{Z}.$$

Since f_σ has compact support, it follows from Rellich's lemma that

$$f_\sigma \cdot D_{\text{geo}}^\alpha \cdot (1 + D_{\text{geo}}^2)^{-1} \cdot f_\sigma$$

is a compact operator on

$$L^2(Y_\sigma, S_{Y_\sigma} \otimes L_{Y_\sigma}).$$

It implies that the norm of the restriction of

$$f_\sigma \cdot D_{U_\sigma}^\alpha \cdot (1 + D_{U_\sigma}^2)^{-1} \cdot f_\sigma$$

to

$$[V_\lambda^* \otimes S_{L\mathfrak{g}/(L\mathfrak{g})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)} \otimes M(\mathbf{v})]^{G_\sigma}$$

tends to 0 as $\|\mathbf{v}\|$ tends to infinity. Thus, it suffices to show that the operator

$$f_\sigma \cdot D_{U_\sigma}^\alpha \cdot (1 + D_{U_\sigma}^2)^{-1} \cdot f_\sigma$$

is compact after restricted to a fixed component

$$[V_\lambda^* \otimes S_{L\mathfrak{g}/(L\mathfrak{g})_\sigma}^* \otimes \mathbb{C}_{(\rho_G - \rho_\sigma, h^\vee)} \otimes M(\mathbf{v})]^{G_\sigma}.$$

By [Theorem 3.8](#), the operator $D_{\text{alg}}^2 = D_{L\mathfrak{g}, (L\mathfrak{g})_\sigma}^2$ acts on the above space by

$$\|\lambda + \rho_G\|^2 - \|\mathbf{v} + \rho_G\|^2.$$

By the formula in [\(2.2\)](#),

$$\|\lambda + \rho_G\|^2 - \|\mathbf{v} + \rho_G\|^2 = \|\lambda + \rho_G\|^2 + 2 \cdot n \cdot (k + h^\vee) - \|\mathbf{v} + \rho_G\|^2.$$

Since $\|\lambda + \rho_G\|^2$ and $\|\mathbf{v} + \rho_G\|^2$ are fixed constants, the operator D_{alg}^2 tends to infinity as the energy level n goes to infinity. This proves the assertion. \square

Lemma 6.11. *One has that*

$$f_\sigma \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma, \quad [f_\sigma, (1 + \mathcal{D}^2)^{-1}] \cdot f_\sigma \in \mathbb{K}(\mathcal{H}_\lambda).$$

Proof. As shown in [Proposition 6.3](#), there exists a bounded operator B so that

$$\mathcal{D} \cdot s = D_{U_\sigma} \cdot s + B \cdot s$$

for any $s \in \mathcal{H}_\lambda$ with $\text{Supp} s \subset U_\sigma$. By the choice of function f_σ , the support of $f_\sigma \cdot s$ is automatically contained in U_σ for all $s \in \mathcal{H}_\lambda$. One can verify that

$$f_\sigma \cdot \mathcal{D}^\alpha \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma = f_\sigma \cdot (D_{U_\sigma} + B)^\alpha \cdot (1 + (D_{U_\sigma} + B)^2)^{-1} \cdot f_\sigma.$$

By [Lemmas 6.9](#) and [6.10](#), we conclude that for $\alpha = 0, 1$,

$$f_\sigma \cdot \mathcal{D}^\alpha \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma \in \mathbb{K}(\mathcal{H}_\lambda).$$

For the second half, we calculate that

$$\begin{aligned} [f_\sigma, (1 + \mathcal{D}^2)^{-1}] \cdot f_\sigma &= (1 + \mathcal{D}^2)^{-1} \cdot [f_\sigma, \mathcal{D}^2] \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma \\ &= (1 + \mathcal{D}^2)^{-1} \cdot c(df_\sigma) \cdot \mathcal{D} \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma + (1 + \mathcal{D}^2)^{-1} \cdot \mathcal{D} \cdot c(df_\sigma) \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma. \end{aligned} \tag{6.9}$$

We point out here that $c(df_\sigma)$ is a bounded operator with support in U_σ . As in [Lemma 6.10](#), one can similarly prove that

$$c(df_\sigma) \cdot \mathcal{D}^\alpha \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma \in \mathbb{K}(\mathcal{H}_\lambda).$$

This completes the proof. \square

Because the operator

$$\mathcal{F} = \frac{\mathcal{D}}{\sqrt{1 + \mathcal{D}^2}} \in \mathbb{B}(\mathcal{H}_\lambda),$$

we have that

$$\begin{aligned} 1 - \mathcal{F}^2 &= (1 + \mathcal{D}^2)^{-1} = \sum_{\dim \sigma = 0} (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma^2 \\ &= \sum_{\dim \sigma = 0} (f_\sigma \cdot (1 + \mathcal{D}^2)^{-1} \cdot f_\sigma + [f_\sigma, (1 + \mathcal{D}^2)^{-1}] \cdot f_\sigma). \end{aligned} \tag{6.10}$$

By [Lemma 6.11](#),

$$1 - \mathcal{F}^2 \in \mathbb{K}(\mathcal{H}_\lambda),$$

which implies that \mathcal{F} is a Fredholm operator on \mathcal{H}_λ .

Acknowledgments

The author received many useful suggestions as well as warm help from Eckhard Meinrenken. The author also would like to thank Nigel Higson and Mathai Varghese for inspiring conversations.

References

- [1] Anton Alekseev, Anton Malkin, Eckhard Meinrenken, Lie group valued moment maps, *J. Differential Geom.* 48 (3) (1998) 445–495.
- [2] Andrew Pressley, Graeme Segal, Loop groups, in: *Oxford Mathematical Monographs*, The Clarendon Press Oxford University Press, New York, 1986, Oxford Science Publications.
- [3] Bertram Kostant, A cubic Dirac operator and the emergence of Euler number multiplets of representations for equal rank subgroups, *Duke Math. J.* 100 (3) (1999) 447–501.
- [4] Antony Wassermann, Kac-moody and virasoro algebras, 2010, <http://arxiv.org/abs/1004.1287>.
- [5] Gregory D. Landweber, Multiplets of representations and Kostant's Dirac operator for equal rank loop groups, *Duke Math. J.* 110 (1) (2001) 121–160.
- [6] Hessel Posthuma, Dirac induction for loop groups, *Lett. Math. Phys.* 95 (1) (2011) 89–107.
- [7] Daniel S. Freed, Michael J. Hopkins, Constantin Teleman, Loop groups and twisted K -theory II, *J. Amer. Math. Soc.* 26 (3) (2013) 595–644.
- [8] Eckhard Meinrenken, The cubic Dirac operator for infinite-dimensional Lie algebras, *Canad. J. Math.* 63 (6) (2011) 1364–1387.
- [9] Anton Alekseev, Eckhard Meinrenken, The non-commutative Weil algebra, *Invent. Math.* 139 (1) (2000) 135–172.
- [10] Eckhard Meinrenken, Clifford algebras and Lie theory, in: *Ergebnisse der Mathematik und ihrer Grenzgebiete. 3. Folge. A Series of Modern Surveys in Mathematics [Results in Mathematics and Related Areas. 3rd Series. A Series of Modern Surveys in Mathematics]*, vol. 58, Springer, Heidelberg, 2013.
- [11] R.J. Plymen, P.L. Robinson, Spinors in Hilbert space, in: *Cambridge Tracts in Mathematics*, vol. 114, Cambridge University Press, Cambridge, 1994.
- [12] Eckhard Meinrenken, Chris Woodward, Hamiltonian loop group actions and Verlinde factorization, *J. Differential Geom.* 50 (3) (1998) 417–469.
- [13] Eckhard Meinrenken, Chris Woodward, Canonical bundles for Hamiltonian loop group manifolds, *Pacific J. Math.* 198 (2) (2001) 477–487.
- [14] Eckhard Meinrenken, Twisted K -homology and group-valued moment maps, *Int. Math. Res. Not. IMRN* (20) (2012) 4563–4618.
- [15] Krzysztof Gawedzki, Conformal field theory: a case study, in: *Conformal Field Theory (Istanbul, 1998)*, in: *Front. Phys.*, vol. 102, Adv. Book Program, Perseus Publ., Cambridge, MA, 2000, p. 55.