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Leonard pairs, spin models, and distance-regular graphs



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ABSTRACT

A Leonard pair is an ordered pair of diagonalizable linear maps on a finite-dimensional vector space, that each act on an eigenbasis for the other one in an irreducible tridiagonal fashion. In the present paper we consider a type of Leonard pair, said to have spin. The notion of a spin model was introduced by V.F.R. Jones to construct link invariants. A spin model is a symmetric matrix over \mathbb{C} that satisfies two conditions, called the type II and type III conditions. It is known that a spin model W is contained in a certain finite-dimensional algebra $N(W)$, called the Nomura algebra. It often happens that a spin model W satisfies $W \in M \subseteq N(W)$, where M is the Bose-Mesner algebra of a distance-regular graph Γ ; in this case we say that Γ affords W . If Γ affords a spin model, then each irreducible module for every Terwilliger algebra of Γ takes a certain form, recently described by Caughman, Curtin, Nomura, and Wolff. In the present paper we show that the converse is true; if each irreducible module for every Terwilliger algebra of Γ takes this form, then Γ affords a spin model. We explicitly construct this spin model when Γ has q -Racah type. The proof of our main result relies heavily on the theory of spin Leonard pairs.

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1. Introduction

The notion of a Leonard pair was introduced by the second author in [26]. A Leonard pair is roughly described as follows. Let V denote a vector space with finite positive dimension. A Leonard pair on V is an ordered pair of diagonalizable linear maps $A : V \rightarrow V$ and $A^* : V \rightarrow V$ that each act on an eigenbasis for the other one in an irreducible tridiagonal fashion. The Leonard pairs were classified up to isomorphism in [26]; the isomorphism classes correspond to a family of orthogonal polynomials consisting of the q -Racah polynomials and their relatives in the terminating branch of the Askey scheme [28]. See [22,23,26–28] for more information about Leonard pairs.

In the present paper we will consider a type of Leonard pair, said to have spin [7]. The spin condition is described as follows. Let A, A^* denote a Leonard pair on V . A Boltzmann pair for A, A^* is an ordered pair of invertible linear maps $W : V \rightarrow V$ and $W^* : V \rightarrow V$ such that $WA = AW$ and $W^*A^* = A^*W^*$ and $WA^*W^{-1} = (W^*)^{-1}AW^*$. The Leonard pair A, A^* is said to have spin whenever there exists a Boltzmann pair for A, A^* . In [7] the spin Leonard pairs are classified up to isomorphism and their Boltzmann pairs are described.

Next we summarize how Leonard pairs come up in algebraic graph theory. Let Γ denote a distance-regular graph [4, p. 126] with vertex set X and adjacency matrix $A \in \text{Mat}_X(\mathbb{C})$. Let M denote the subalgebra of $\text{Mat}_X(\mathbb{C})$ generated by A . By [4, Theorem 2.6.1] the algebra M has a basis $\{A_i\}_{i=0}^D$ where D is the diameter of Γ [4, p. 432] and A_i is the i th distance-matrix of Γ [4, p. 127] for $0 \leq i \leq D$. The algebra M is called the Bose-Mesner algebra of Γ ; it is a useful tool in the study of Γ [3,4]. To carry out this study in more detail, it is helpful to bring in another algebra called the subconstituent algebra or Terwilliger algebra. This algebra is described as follows. For $x \in X$

and $0 \leq i \leq D$, the diagonal matrix $E_i^* = E_i^*(x) \in \text{Mat}_X(\mathbb{C})$ is the projection onto the i th subconstituent of Γ with respect to x [4, p. 434]. The matrices $\{E_i^*\}_{i=0}^D$ form a basis for a commutative subalgebra $M^* = M^*(x)$ of $\text{Mat}_X(\mathbb{C})$, called the dual Bose-Mesner algebra of Γ with respect to x . The Terwilliger algebra $T = T(x)$ is the subalgebra of $\text{Mat}_X(\mathbb{C})$ generated by M and M^* [25]. The algebra T is semi-simple [25, Lemma 3.4], so it is natural to investigate the irreducible T -modules. For an irreducible T -module U , there is a condition called thin; this means that for the action of M and M^* on U the eigenspaces all have dimension 1. Most of the known distance-regular graphs with large diameter satisfy an algebraic condition called Q -polynomial [4, Section 4.1.E]; if Γ is Q -polynomial then M^* contains a certain matrix A^* called the dual adjacency matrix of Γ with respect to x [25, p. 379]. In this case A, A^* acts on each thin irreducible T -module as a Leonard pair [25, Lemmas 3.9, 3.12].

Motivated by state models in statistical mechanics, V.F.R. Jones [17] introduced the notion of a spin model to construct link invariants. The construction is roughly described as follows. A link is a finite collection of mutually disjoint closed curves in 3-dimensional Euclidean space. A link diagram is a projection of a link onto a plane. For a link diagram L and a square matrix W over \mathbb{C} with all entries nonzero, Jones defined a scalar $Z_W(L) \in \mathbb{C}$. The function Z_W from the set of all link diagrams to \mathbb{C} gives a link invariant, provided that W satisfies two conditions called type II and type III; these are given in [17] and (68), (71) below. The matrix W is called a spin model whenever it is type II and type III.

In [17] Jones gave three examples of spin models; the Potts model, the square model, and the odd cyclic model. In [15] F. Jaeger observed that each of these examples is contained in the Bose-Mesner algebra of a distance-regular graph Γ . Specifically, for the Potts model Γ is a complete graph, for the square model Γ is the 4-cycle, and for the odd cyclic model Γ is an odd cycle. Jaeger then found a new spin model contained in the Bose-Mesner algebra of the Higman-Sims graph [15]. Motivated by this result, other authors found new spin models contained in the Bose-Mesner algebra of the even cycle [2], the Hadamard graph [20], and the double cover of the Higman-Sims graph [19].

For the rest of this section let Γ denote a distance-regular graph with vertex set X and Bose-Mesner algebra M . For a spin model W contained in M , it often happens that W and M are related by the following condition. For $b, c \in X$ let $\mathbf{u}_{b,c}$ denote the column vector with entries indexed by X and y -entry $W(b, y)W(c, y)^{-1}$ for $y \in X$. The condition is that $\mathbf{u}_{b,c}$ is a common eigenvector of M for all $b, c \in X$. In this case we say that Γ affords W . Each of the above mentioned spin models [2,15,17,19,20] is afforded by the associated distance-regular graph; this will be demonstrated in Section 15 below.

Next we describe how spin models are related to spin Leonard pairs. Assume that Γ affords a spin model W . It was shown in [9, Lemma 5.1] that Γ satisfies a condition called formally self-dual, which implies that Γ is Q -polynomial [4, p. 49]. Pick $x \in X$ and write $A^* = A^*(x)$, $M^* = M^*(x)$, $T = T(x)$. Define $W^* \in M^*$ that has (y, y) -entry $|X|^{1/2}W(x, y)^{-1}$ for $y \in X$. Let U denote a thin irreducible T -module. Then A, A^* act on U as a spin Leonard pair, and W, W^* act on U as a Boltzmann pair for this Leonard pair; see Proposition 11.10 and Lemma 11.11 below. Consequently the intersection numbers of

U [28, Definition 11.1] are given by the formulas in [5, Theorem 8.5]. Another consequence for U is that its endpoint [25, Lemma 3.12] is equal to its dual endpoint [25, Lemma 3.9].

Now we describe our main result. Assume that our distance-regular graph Γ is formally self-dual and has q -Racah type [29, Definition 5.1]. Assume that for all $x \in X$ each irreducible $T(x)$ -module U is thin with matching endpoint and dual-endpoint, and the intersection numbers of U are given by the formulas in [5, Theorem 8.5]. Then there exists a spin model afforded by Γ ; this spin model is explicitly constructed. Our main result is Theorem 14.7. We comment on the nature of this result. We have not discovered any new spin model to date. What we have shown, is that a new spin model would result from the discovery of a distance-regular graph with the right sort of irreducible T -modules.

The proof of our main result relies heavily on the theory of spin Leonard pairs. In the first half of the paper we develop this theory, and obtain several new results that may be of independent interest; for instance Theorem 5.25 and Lemmas 6.11, 6.18.

2. Leonard pairs and Leonard systems

We now begin our formal argument. In this section we recall the notion of a Leonard pair. We use the following terms. A square matrix is said to be *tridiagonal* whenever each nonzero entry lies on either the diagonal, the subdiagonal, or the superdiagonal. A tridiagonal matrix is said to be *irreducible* whenever each entry on the subdiagonal is nonzero and each entry on the superdiagonal is nonzero. Let \mathbb{F} denote a field.

Definition 2.1. [26, Definition 1.1] Let V denote a vector space over \mathbb{F} with finite positive dimension. By a *Leonard pair* on V we mean an ordered pair of \mathbb{F} -linear maps $A : V \rightarrow V$ and $A^* : V \rightarrow V$ that satisfy the following (i), (ii).

- (i) There exists a basis for V with respect to which the matrix representing A is irreducible tridiagonal and the matrix representing A^* is diagonal.
- (ii) There exists a basis for V with respect to which the matrix representing A^* is irreducible tridiagonal and the matrix representing A is diagonal.

Note 2.2. According to a common notational convention, for a matrix A its conjugate-transpose is denoted by A^* . We are not using this convention. In a Leonard pair A, A^* the linear maps A, A^* are arbitrary subject to (i) and (ii) above.

Note 2.3. Referring to Definition 2.1, let A, A^* denote a Leonard pair on V . Then A^*, A is a Leonard pair on V .

We refer the reader to [28] for background on Leonard pairs.

From now until the end of Section 6, fix an integer $d \geq 0$ and let V denote a vector space over \mathbb{F} with dimension $d + 1$. Let $\text{End}(V)$ denote the \mathbb{F} -algebra consisting of the \mathbb{F} -linear maps $V \rightarrow V$. We denote by I the identity element of $\text{End}(V)$. For $A \in \text{End}(V)$

let $\langle A \rangle$ denote the subalgebra of $\text{End}(V)$ generated by A . By [24, Corollary 2.133] the center of $\text{End}(V)$ consists of the elements αI ($\alpha \in \mathbb{F}$).

Lemma 2.4. [28, Corollary 5.6] *Let A, A^* denote a Leonard pair on V . Then A, A^* together generate the algebra $\text{End}(V)$.*

For an \mathbb{F} -algebra \mathcal{A} , by an *automorphism* of \mathcal{A} we mean an \mathbb{F} -algebra isomorphism $\mathcal{A} \rightarrow \mathcal{A}$. By an *antiautomorphism* of \mathcal{A} we mean an \mathbb{F} -linear bijection $\tau : \mathcal{A} \rightarrow \mathcal{A}$ such that $(YZ)^\tau = Z^\tau Y^\tau$ for all $Y, Z \in \mathcal{A}$.

Lemma 2.5. [24, Corollary 7.126] (Skolem-Noether) *A map $\sigma : \text{End}(V) \rightarrow \text{End}(V)$ is an automorphism of $\text{End}(V)$ if and only if there exists an invertible $K \in \text{End}(V)$ such that $Y^\sigma = KYK^{-1}$ for all $Y \in \text{End}(V)$.*

Lemma 2.6. [28, Theorem 6.1] *Let A, A^* denote a Leonard pair on V . Then there exists a unique antiautomorphism \dagger of $\text{End}(V)$ such that $A^\dagger = A$ and $A^{*\dagger} = A^*$. Moreover $(Y^\dagger)^\dagger = Y$ for all $Y \in \text{End}(V)$. The map \dagger fixes every element in $\langle A \rangle$ and every element in $\langle A^* \rangle$.*

We recall the notion of an isomorphism for Leonard pairs. Let A, A^* denote a Leonard pair on V . Let V' denote a vector space over \mathbb{F} with dimension $d+1$, and let A', A'^* denote a Leonard pair on V' . By an *isomorphism of Leonard pairs* from A, A^* to A', A'^* we mean an \mathbb{F} -algebra isomorphism $\text{End}(V) \rightarrow \text{End}(V')$ that sends $A \mapsto A'$ and $A^* \mapsto A'^*$. The Leonard pairs A, A^* and A', A'^* are said to be *isomorphic* whenever there exists an isomorphism of Leonard pairs from A, A^* to A', A'^* . In this case the isomorphism is unique by Lemma 2.4.

We recall the self-dual Leonard pairs. A Leonard pair A, A^* is said to be *self-dual* whenever A, A^* is isomorphic to A^*, A . In this case, the isomorphism of Leonard pairs from A, A^* to A^*, A is called the *duality* of A, A^* .

Lemma 2.7. *Assume that A, A^* is self-dual with duality σ . Then $\sigma^2 = 1$.*

Proof. By construction σ^2 fixes each of A, A^* . By this and Lemma 2.4, σ^2 fixes every element of $\text{End}(V)$. The result follows. \square

When working with a Leonard pair, it is convenient to consider a closely related object called a Leonard system [26]. Before we define a Leonard system, we recall a few concepts from linear algebra.

Let $\text{Mat}_{d+1}(\mathbb{F})$ denote the \mathbb{F} -algebra consisting of the $d+1$ by $d+1$ matrices that have all entries in \mathbb{F} . We index the rows and columns by $0, 1, \dots, d$. Let $\{v_i\}_{i=0}^d$ denote a basis for V . For $A \in \text{End}(V)$ and $M \in \text{Mat}_{d+1}(\mathbb{F})$, we say that M *represents A with respect to $\{v_i\}_{i=0}^d$* whenever $Av_j = \sum_{i=0}^d M_{i,j}v_i$ for $0 \leq j \leq d$. Let $A \in \text{End}(V)$. For $\theta \in \mathbb{F}$ define $V(\theta) = \{v \in V \mid Av = \theta v\}$. Observe that $V(\theta)$ is a subspace of V . The scalar θ is

called an *eigenvalue* of A whenever $V(\theta) \neq 0$. In this case, $V(\theta)$ is called the *eigenspace* of A corresponding to θ . We say that A is *diagonalizable* whenever V is spanned by the eigenspaces of A . We say that A is *multiplicity-free* whenever A is diagonalizable, and each eigenspace of A has dimension one. Assume that A is multiplicity-free, and let $\{V_i\}_{i=0}^d$ denote an ordering of the eigenspaces of A . Then $V = \sum_{i=0}^d V_i$ (direct sum). For $0 \leq i \leq d$ let θ_i denote the eigenvalue of A corresponding to V_i . For $0 \leq i \leq d$ define $E_i \in \text{End}(V)$ such that $(E_i - I)V_i = 0$ and $E_i V_j = 0$ if $j \neq i$ ($0 \leq j \leq d$). Thus E_i is the projection onto V_i . Observe that (i) $V_i = E_i V$ ($0 \leq i \leq d$); (ii) $E_i E_j = \delta_{i,j} E_i$ ($0 \leq i, j \leq d$); (iii) $I = \sum_{i=0}^d E_i$; (iv) $A = \sum_{i=0}^d \theta_i E_i$. Also

$$E_i = \prod_{\substack{0 \leq j \leq d \\ j \neq i}} \frac{A - \theta_j I}{\theta_i - \theta_j} \quad (0 \leq i \leq d). \tag{1}$$

We call E_i the *primitive idempotent* of A for θ_i ($0 \leq i \leq d$). Observe that $\{A^i\}_{i=0}^d$ is a basis for the \mathbb{F} -vector space $\langle A \rangle$, and $\prod_{i=0}^d (A - \theta_i I) = 0$. Also observe that $\{E_i\}_{i=0}^d$ is a basis for the \mathbb{F} -vector space $\langle A \rangle$.

Let A, A^* denote a Leonard pair on V . By [26, Lemma 1.3] each of A, A^* is multiplicity-free. Let $\{E_i\}_{i=0}^d$ denote an ordering of the primitive idempotents of A . For $0 \leq i \leq d$ pick $0 \neq v_i \in E_i V$. Then $\{v_i\}_{i=0}^d$ is a basis for V . The ordering $\{E_i\}_{i=0}^d$ is said to be *standard* whenever the basis $\{v_i\}_{i=0}^d$ satisfies Definition 2.1(ii). A standard ordering of the primitive idempotents of A^* is similarly defined.

Definition 2.8. [26, Definition 1.4] By a *Leonard system* on V we mean a sequence

$$\Phi = (A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d) \tag{2}$$

of elements in $\text{End}(V)$ that satisfy the following (i)–(iii):

- (i) A, A^* is a Leonard pair on V ;
- (ii) $\{E_i\}_{i=0}^d$ is a standard ordering of the primitive idempotents of A ;
- (iii) $\{E_i^*\}_{i=0}^d$ is a standard ordering of the primitive idempotents of A^* .

For the rest of this section, we fix a Leonard system Φ as in (2).

Lemma 2.9. *The antiautomorphism \dagger from Lemma 2.6 fixes each of E_i, E_i^* for $0 \leq i \leq d$.*

Proof. By Lemma 2.6. \square

We recall the notion of isomorphism for Leonard systems. Let V' denote a vector space over \mathbb{F} with dimension $d + 1$, and let $\Phi' = (A'; \{E'_i\}_{i=0}^d; A'^*; \{E'^*_i\}_{i=0}^d)$ denote a Leonard system on V' . By an *isomorphism of Leonard systems* from Φ to Φ' we mean an \mathbb{F} -algebra isomorphism $\text{End}(V) \rightarrow \text{End}(V')$ that sends $A \mapsto A', A^* \mapsto A'^*$ and $E_i \mapsto E'_i,$

$E_i^* \mapsto E_i^{*'}$ for $0 \leq i \leq d$. The Leonard systems Φ and Φ' are said to be *isomorphic* whenever there exists an isomorphism of Leonard systems from Φ to Φ' . In this case the isomorphism is unique.

Each of the following is a Leonard system on V :

$$\begin{aligned} \Phi^* &= (A^*; \{E_i^*\}_{i=0}^d; A; \{E_i\}_{i=0}^d), \\ \Phi^\downarrow &= (A; \{E_i\}_{i=0}^d; A^*; \{E_{d-i}^*\}_{i=0}^d), \\ \Phi^\updownarrow &= (A; \{E_{d-i}\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d). \end{aligned}$$

Moreover, for $\alpha, \beta, \alpha^*, \beta^* \in \mathbb{F}$ with $\alpha\alpha^* \neq 0$, the sequence

$$(\alpha A + \beta I; \{E_i\}_{i=0}^d; \alpha^* A^* + \beta^* I; \{E_i^*\}_{i=0}^d)$$

is a Leonard system on V . We call this Leonard system an *affine transformation of Φ* .

For $0 \leq i \leq d$ let θ_i (resp. θ_i^*) denote the eigenvalue of A (resp. A^*) corresponding to E_i (resp. E_i^*). We call $\{\theta_i\}_{i=0}^d$ (resp. $\{\theta_i^*\}_{i=0}^d$) the *eigenvalue sequence* (resp. *dual eigenvalue sequence*) of Φ . Throughout this paper we often discuss the Leonard system Φ^* as well as Φ . The following notational convention will simplify this discussion.

Definition 2.10. Referring to our Leonard system Φ , for any object ω associated with Φ , let ω^* denote the corresponding object for Φ^* .

We recall the Φ -standard basis. Pick $0 \neq u \in E_0V$. By [28, Lemma 10.2] the vectors $\{E_i^*u\}_{i=0}^d$ form a basis for V , said to be Φ -*standard*. With respect to this basis the matrix representing A is irreducible tridiagonal, and the matrix representing A^* is diagonal with (i, i) -entry θ_i^* for $0 \leq i \leq d$. So with respect to the basis $\{E_i^*u\}_{i=0}^d$ the matrices representing A, A^* are

$$A : \begin{pmatrix} a_0 & b_0 & & & & \mathbf{0} \\ c_1 & a_1 & b_1 & & & \\ & c_2 & \cdot & \cdot & & \\ & & \cdot & \cdot & \cdot & \\ & & & \cdot & b_{d-1} & \\ \mathbf{0} & & & & c_d & a_d \end{pmatrix}, \quad A^* : \begin{pmatrix} \theta_0^* & & & & & \mathbf{0} \\ & \theta_1^* & & & & \\ & & \theta_2^* & & & \\ & & & \cdot & & \\ & & & & \cdot & \\ \mathbf{0} & & & & & \theta_d^* \end{pmatrix}, \quad (3)$$

where $\{c_i\}_{i=1}^d, \{a_i\}_{i=0}^d, \{b_i\}_{i=0}^{d-1}$ are scalars in \mathbb{F} such that $b_{i-1}c_i \neq 0$ for $1 \leq i \leq d$. The scalars $\{c_i\}_{i=1}^d, \{a_i\}_{i=0}^d, \{b_i\}_{i=0}^{d-1}$ are uniquely determined by Φ , and called the *intersection numbers of Φ* . The sum of the vectors in a Φ -standard basis is contained in E_0V , and therefore an eigenvector for A with eigenvalue θ_0 . Consequently

$$\theta_0 = c_i + a_i + b_i \quad (0 \leq i \leq d),$$

where $c_0 = 0$ and $b_d = 0$.

We define some polynomials. Let λ denote an indeterminate, and let $\mathbb{F}[\lambda]$ denote the \mathbb{F} -algebra consisting of the polynomials in λ with all coefficients in \mathbb{F} . For $0 \leq i \leq d$ define the following polynomials in $\mathbb{F}[\lambda]$:

$$\begin{aligned} \tau_i(\lambda) &= (\lambda - \theta_0)(\lambda - \theta_1) \cdots (\lambda - \theta_{i-1}), \\ \eta_i(\lambda) &= (\lambda - \theta_d)(\lambda - \theta_{d-1}) \cdots (\lambda - \theta_{d-i+1}). \end{aligned}$$

For a nonzero $u \in E_0^*V$, consider the vectors $\{\tau_i(A)u\}_{i=0}^d$. By [28, Section 21] these vectors form a basis for V . This basis is called a Φ -split basis for V .

Lemma 2.11. [26, Theorem 3.2] *There exist scalars $\{\varphi_i\}_{i=1}^d$ in \mathbb{F} such that with respect to a Φ -split basis for V the matrices representing A and A^* are*

$$A : \begin{pmatrix} \theta_0 & & & & \mathbf{0} \\ 1 & \theta_1 & & & \\ & 1 & \theta_2 & & \\ & & \ddots & \ddots & \\ \mathbf{0} & & & 1 & \theta_d \end{pmatrix}, \quad A^* : \begin{pmatrix} \theta_0^* & \varphi_1 & & & \mathbf{0} \\ & \theta_1^* & \varphi_2 & & \\ & & \theta_2^* & \ddots & \\ & & & \ddots & \varphi_d \\ \mathbf{0} & & & & \theta_d^* \end{pmatrix}.$$

The sequence $\{\varphi_i\}_{i=1}^d$ is uniquely determined by Φ . Moreover $\varphi_i \neq 0$ for $1 \leq i \leq d$.

Referring to Lemma 2.11, the sequence $\{\varphi_i\}_{i=1}^d$ is called the *first split sequence* of Φ . Let $\{\phi_i\}_{i=1}^d$ denote the first split sequence of Φ^\downarrow . We call $\{\phi_i\}_{i=1}^d$ the *second split sequence* of Φ . By the *parameter array* of Φ we mean the sequence

$$(\{\theta_i\}_{i=0}^d; \{\theta_i^*\}_{i=0}^d; \{\varphi_i\}_{i=1}^d; \{\phi_i\}_{i=1}^d). \tag{4}$$

Lemma 2.12. [26, Theorem 1.9] *A Leonard system is uniquely determined up to isomorphism by its parameter array.*

Lemma 2.13. [28, Theorem 23.5] *The intersection numbers $\{c_i\}_{i=1}^d, \{b_i\}_{i=0}^{d-1}$ of Φ are determined by the parameter array in the following way:*

$$c_i = \phi_i \frac{\eta_{d-i}^*(\theta_i^*)}{\eta_{d-i+1}^*(\theta_{i-1}^*)} \quad (1 \leq i \leq d), \tag{5}$$

$$b_i = \varphi_{i+1} \frac{\tau_i^*(\theta_i^*)}{\tau_{i+1}^*(\theta_{i+1}^*)} \quad (0 \leq i \leq d-1). \tag{6}$$

Lemma 2.14. *The parameter array of Φ is determined by the eigenvalue sequence, dual eigenvalue sequence, and the $\{c_i\}_{i=1}^d, \{b_i\}_{i=0}^{d-1}$.*

Proof. By Lemma 2.13. \square

We recall the self-dual Leonard systems. The Leonard system Φ is said to be *self-dual* whenever Φ is isomorphic to Φ^* . In this case, the isomorphism of Leonard systems from Φ to Φ^* is called the *duality* of Φ . We remark that if Φ is self-dual then so is the Leonard pair A, A^* , and in this case the duality of Φ is the duality of A, A^* .

Lemma 2.15. [23, Lemma 8.6, Proposition 8.7] *The Leonard system Φ is self-dual if and only if $\theta_i = \theta_i^*$ for $0 \leq i \leq d$. In this case, $\phi_i = \phi_{d-i+1}$ for $1 \leq i \leq d$.*

3. The map π

In this section, we introduce a certain map $\pi : \text{End}(V) \rightarrow \mathbb{F}$ attached to a Leonard system. Before defining this map, we briefly recall some basic facts about the trace function $\text{tr} : \text{End}(V) \rightarrow \mathbb{F}$. We have $\text{tr}(YZ) = \text{tr}(ZY)$ for $Y, Z \in \text{End}(V)$.

Lemma 3.1. *Let σ denote an automorphism of $\text{End}(V)$. Then $\text{tr}(Y^\sigma) = \text{tr}(Y)$ for all $Y \in \text{End}(V)$.*

Proof. By Lemma 2.5 there exists an invertible $K \in \text{End}(V)$ such that $Y^\sigma = KYK^{-1}$ for all $Y \in \text{End}(V)$. We have $\text{tr}(KYK^{-1}) = \text{tr}(K^{-1}KY) = \text{tr}(Y)$. The result follows. \square

Lemma 3.2. *Let τ denote an antiautomorphism of $\text{End}(V)$. Then $\text{tr}(Y^\tau) = \text{tr}(Y)$ for all $Y \in \text{End}(V)$.*

Proof. Without loss we may identify $\text{End}(V)$ with $\text{Mat}_{d+1}(\mathbb{F})$. The transpose map is an antiautomorphism of the algebra $\text{Mat}_{d+1}(\mathbb{F})$. So the composition of τ and the transpose map is an automorphism of $\text{Mat}_{d+1}(\mathbb{F})$. The result follows in view of Lemma 3.1 and the fact that a matrix and its transpose have the same trace. \square

We now define the map π . Let $\Phi = (A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system on V with parameter array as in (4).

Definition 3.3. Define the map

$$\pi : \begin{array}{ccc} \text{End}(V) & \longrightarrow & \mathbb{F} \\ Y & \longmapsto & \text{tr}(YE_0) \end{array}$$

Observe that π is \mathbb{F} -linear.

Lemma 3.4. *For $Y \in \text{End}(V)$,*

$$E_0YE_0 = \pi(Y)E_0.$$

Proof. Abbreviate $\mathcal{A} = \text{End}(V)$. The primitive idempotent E_0 is a basis for the subspace $E_0\mathcal{A}E_0$. This subspace contains E_0YE_0 , so there exists $\alpha \in \mathbb{F}$ such that $E_0YE_0 = \alpha E_0$.

In this equation, take the trace of each side to get $\text{tr}(E_0 Y E_0) = \alpha \text{tr}(E_0)$. We have $\text{tr}(E_0 Y E_0) = \text{tr}(Y E_0) = \pi(Y)$. Also $\text{tr}(E_0) = 1$ since E_0 is a primitive idempotent. By these comments, $\alpha = \pi(Y)$. The result follows. \square

Lemma 3.5. *For $0 \leq i \leq d$,*

$$\pi(E_i) = \delta_{i,0}. \tag{7}$$

Proof. We have $\text{tr}(E_0) = 1$. For $1 \leq i \leq d$ we have $E_i E_0 = 0$. By these comments and Definition 3.3 we obtain (7). \square

Lemma 3.6. *We have*

$$\pi(I) = 1, \quad \pi(A) = \theta_0.$$

Proof. We have $\pi(I) = 1$ since $\text{tr}(E_0) = 1$. In $A = \sum_{i=0}^d \theta_i E_i$, apply π to each side and use (7) to get $\pi(A) = \theta_0$. \square

Lemma 3.7. *For $Y \in \langle A \rangle$ and $Z \in \text{End}(V)$,*

$$\pi(YZ) = \pi(ZY) = \pi(Y)\pi(Z).$$

Proof. Write $Y = \sum_{i=0}^d \alpha_i E_i$. We have

$$\pi(YZ) = \text{tr}(YZE_0) = \text{tr}(ZE_0Y) = \alpha_0 \text{tr}(ZE_0) = \alpha_0 \pi(Z).$$

Similarly $\pi(ZY) = \alpha_0 \pi(Z)$. The result follows since $\alpha_0 = \pi(Y)$ by Lemma 3.5. \square

Lemma 3.8. *The restriction of π to $\langle A \rangle$ is an \mathbb{F} -algebra homomorphism $\langle A \rangle \rightarrow \mathbb{F}$.*

Proof. By Lemma 3.7. \square

Recall the antiautomorphism \dagger from Lemma 2.6.

Lemma 3.9. *The map π makes the following diagram commute.*

$$\begin{array}{ccc} \text{End}(V) & \xrightarrow{\dagger} & \text{End}(V) \\ \pi \downarrow & & \downarrow \pi \\ \mathbb{F} & \xrightarrow{\text{id}} & \mathbb{F} \end{array}$$

Proof. Pick any $Y \in \text{End}(V)$. Using Lemma 3.2 we argue

$$\pi(Y^\dagger) = \text{tr}(Y^\dagger E_0) = \text{tr}(Y^\dagger E_0^\dagger) = \text{tr}((E_0 Y)^\dagger) = \text{tr}(E_0 Y) = \text{tr}(Y E_0) = \pi(Y).$$

The result follows. \square

Lemma 3.10. *Assume that Φ is self-dual with duality σ . Then the following diagram commutes.*

$$\begin{array}{ccc}
 \text{End}(V) & \xrightarrow{\sigma} & \text{End}(V) \\
 \pi^* \downarrow & & \downarrow \pi \\
 \mathbb{F} & \xrightarrow{\text{id}} & \mathbb{F}
 \end{array}$$

Proof. For $Z \in \text{End}(V)$ we have $\text{tr}(Z^\sigma) = \text{tr}(Z)$. For $Y \in \text{End}(V)$ we have $Y^\sigma E_0 = Y^\sigma (E_0^*)^\sigma = (Y E_0^*)^\sigma$. By these comments $\text{tr}(Y^\sigma E_0) = \text{tr}(Y E_0^*)$. Thus $\pi(Y^\sigma) = \pi^*(Y)$. The result follows. \square

Lemma 3.11. [28, Lemma 9.4] *There exists $0 \neq \nu \in \mathbb{F}$ such that*

$$\nu E_0 E_0^* E_0 = E_0, \quad \nu E_0^* E_0 E_0^* = E_0^*. \tag{8}$$

Note that $\nu = \nu^*$. By (8) and Lemma 3.4,

$$\nu^{-1} = \pi(E_0^*). \tag{9}$$

Lemma 3.12. [28, Theorem 23.8] *We have*

$$\nu = \frac{\eta_d(\theta_0) \eta_d^*(\theta_0^*)}{\phi_1 \phi_2 \cdots \phi_d}. \tag{10}$$

Definition 3.13. For $0 \leq i \leq d$ define

$$k_i = \nu \pi(E_i^*). \tag{11}$$

Lemma 3.14. *We have*

$$k_0 = 1, \quad \sum_{i=0}^d k_i = \nu. \tag{12}$$

Proof. Use (9) and (11) to get $k_0 = 1$. In $\sum_{i=0}^d E_i^* = I$, apply π to each side, and simplify the result using (11) and Lemma 3.6. This gives the equation on the right in (12). \square

In Lemma 4.30 the scalars k_i are expressed in terms of the parameter array.

Lemma 3.15. *For $0 \leq i \leq d$,*

$$E_0 E_i^* E_0 = \nu^{-1} k_i E_0.$$

Proof. Use Lemma 3.4 (with $Y = E_i^*$) and (11). \square

4. The map ρ

Let $\Phi = (A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system on V . In this section we introduce a certain map $\rho : \langle A \rangle \rightarrow \langle A^* \rangle$. We abbreviate $\mathcal{A} = \text{End}(V)$.

Lemma 4.1. *The map $\langle A \rangle \rightarrow \mathcal{A}E_0^*$, $Y \mapsto YE_0^*$ is an \mathbb{F} -linear bijection.*

Proof. The map is clearly \mathbb{F} -linear. By [28, Corollary 5.8] the elements $\{E_i E_0^*\}_{i=0}^d$ form a basis for $\mathcal{A}E_0^*$. The map $Y \mapsto YE_0^*$ sends the basis $\{E_i\}_{i=0}^d$ for $\langle A \rangle$ to the basis $\{E_i E_0^*\}_{i=0}^d$ for $\mathcal{A}E_0^*$, so it is bijective. \square

Lemma 4.2. *There exists a unique \mathbb{F} -linear map $\rho : \langle A \rangle \rightarrow \langle A^* \rangle$ such that for $Y \in \langle A \rangle$,*

$$YE_0^*E_0 = Y^\rho E_0. \tag{13}$$

Proof. Concerning existence, consider the \mathbb{F} -linear map $g : \langle A \rangle \rightarrow \mathcal{A}E_0$, $Y \mapsto YE_0^*E_0$. By Lemma 4.1 applied to Φ^* , the map $\mu : \langle A^* \rangle \rightarrow \mathcal{A}E_0$, $Y \mapsto YE_0$ is an \mathbb{F} -linear bijection. The composition

$$\rho : \langle A \rangle \xrightarrow{g} \mathcal{A}E_0 \xrightarrow{\mu^{-1}} \langle A^* \rangle$$

satisfies (13). We have shown that ρ exists. The map ρ is unique by Lemma 4.1. \square

Recall the scalar ν from (8).

Lemma 4.3. *The maps ρ and $\nu\rho^*$ are inverses. In particular, the maps ρ, ρ^* are bijective.*

Proof. Pick $Y \in \langle A \rangle$. Using in order (13) for Φ^* , (13), (8) we obtain

$$(Y^\rho)^\rho E_0^* = Y^\rho E_0 E_0^* = YE_0^* E_0 E_0^* = \nu^{-1} Y E_0^*.$$

By this and Lemma 4.1 we get $(Y^\rho)^\rho = \nu^{-1} Y$. Similarly, for $Z \in \langle A^* \rangle$ we get $(Z^\rho)^{\rho^*} = \nu^{-1} Z$. Thus the maps ρ and $\nu\rho^*$ are inverses. \square

Definition 4.4. For $0 \leq i \leq d$ define

$$A_i = \nu(E_i^*)^{\rho^*}. \tag{14}$$

Lemma 4.5. *For $0 \leq i \leq d$, ρ sends $A_i \mapsto E_i^*$ and $E_i \mapsto \nu^{-1} A_i^*$.*

Proof. By Lemma 4.3 and Definition 4.4. \square

Lemma 4.6. For $0 \leq i \leq d$,

$$A_i E_0^* E_0 = E_i^* E_0, \quad E_i E_0^* E_0 = \nu^{-1} A_i^* E_0, \tag{15}$$

$$A_i^* E_0 E_0^* = E_i E_0^*, \quad E_i^* E_0 E_0^* = \nu^{-1} A_i E_0^*. \tag{16}$$

Proof. To get (15), use Lemmas 4.2, 4.5. Applying (15) to Φ^* we obtain (16). \square

Lemma 4.7. Assume that Φ is self-dual with duality σ . Then σ sends $A_i \leftrightarrow A_i^*$ ($0 \leq i \leq d$).

Proof. By Lemmas 4.2 and 4.6, A_i is the unique element in $\langle A \rangle$ such that

$$A_i E_0^* E_0 = E_i^* E_0. \tag{17}$$

Applying this to Φ^* , we find that for $0 \leq i \leq d$, A_i^* is the unique element in $\langle A^* \rangle$ such that

$$A_i^* E_0 E_0^* = E_i E_0^*. \tag{18}$$

Apply σ to (17), and compare the result with (18) to find that σ sends $A_i \mapsto A_i^*$. Applying this fact to Φ^* , we find that σ sends $A_i^* \mapsto A_i$. \square

Lemma 4.8. We have $A_0 = I$.

Proof. For the equation on the right in (16), set $i = 0$ and apply (8) to obtain $E_0^* = A_0 E_0^*$. By this and Lemma 4.1 we get $A_0 = I$. \square

Lemma 4.9. The map ρ sends $I \mapsto E_0^*$ and $E_0 \mapsto \nu^{-1} I$.

Proof. Set $i = 0$ in Lemma 4.5 and use Lemma 4.8. \square

Lemma 4.10. We have

$$\sum_{i=0}^d A_i = \nu E_0.$$

Proof. Apply ρ^* to each side of the equation $\sum_{i=0}^d E_i^* = I$, and evaluate the result using Lemmas 4.5, 4.9 applied to Φ^* . \square

Lemma 4.11. The elements $\{A_i\}_{i=0}^d$ form a basis for the \mathbb{F} -vector space $\langle A \rangle$.

Proof. By Lemmas 4.3, 4.5, and since $\{E_i^*\}_{i=0}^d$ form a basis for the \mathbb{F} -vector space $\langle A^* \rangle$. \square

Recall the antiautomorphism \dagger from Lemma 2.6.

Lemma 4.12. *The antiautomorphism \dagger fixes each of A_i, A_i^* for $0 \leq i \leq d$.*

Proof. By Lemma 2.6 and since $A_i \in \langle A \rangle, A_i^* \in \langle A^* \rangle$ for $0 \leq i \leq d$. \square

Lemma 4.13. *The map ρ from Lemma 4.2 makes the following diagram commute.*

$$\begin{array}{ccc}
 \langle A \rangle & \xrightarrow{\rho} & \langle A^* \rangle \\
 \dagger \downarrow & & \downarrow \dagger \\
 \langle A \rangle & \xrightarrow{\rho} & \langle A^* \rangle
 \end{array}$$

Proof. By Lemma 2.9. \square

Recall the \mathbb{F} -linear map π from Definition 3.3.

Lemma 4.14. *For $Y \in \langle A \rangle$,*

$$\pi(Y^\rho) = \nu^{-1}\pi(Y). \tag{19}$$

Proof. By Definition 3.3, $\pi(Y^\rho) = \text{tr}(Y^\rho E_0)$. By Lemma 4.2, $Y^\rho E_0 = Y E_0^* E_0$. By these comments and Definition 3.3, $\pi(Y^\rho) = \pi(Y E_0^*)$. By this and Lemma 3.7, $\pi(Y^\rho) = \pi(Y)\pi(E_0^*)$. By (9), $\pi(E_0^*) = \nu^{-1}$. Thus (19) holds. \square

Lemma 4.15. *For $0 \leq i \leq d$,*

$$A_i E_0 = k_i E_0. \tag{20}$$

Proof. Since $A_i \in \langle A \rangle$ there exists $\alpha_i \in \mathbb{F}$ such that $A_i E_0 = \alpha_i E_0$. In the first equation of Lemma 4.6, take the trace of each side and use (9), (11) to obtain $\alpha_i = k_i$ after a routine computation. The result follows. \square

Lemma 4.16. *For $0 \leq i \leq d$,*

$$\pi(A_i) = k_i, \quad \pi(A_i^*) = \delta_{i,0}. \tag{21}$$

Proof. In (20), take the trace of each side and use $\text{tr}(E_0) = 1$ to get the equation on the left in (21). Next we obtain the equation on the right in (21). By (14), $A_i^* = \nu E_i^\rho$. By Lemma 4.14, $\pi(E_i^\rho) = \nu^{-1}\pi(E_i)$. By these comments, $\pi(A_i^*) = \pi(E_i)$. By this and Lemma 3.5 we obtain the equation on the right in (21). \square

Recall from (3) the intersection numbers $\{c_i\}_{i=1}^d, \{a_i\}_{i=0}^d, \{b_i\}_{i=0}^{d-1}$ of Φ .

Lemma 4.17. *Assume that $d \geq 1$. Then*

$$AA_0 = c_1A_1 + a_0A_0, \tag{22}$$

$$AA_i = c_{i+1}A_{i+1} + a_iA_i + b_{i-1}A_{i-1} \quad (1 \leq i \leq d-1), \tag{23}$$

$$AA_d = a_dA_d + b_{d-1}A_{d-1}. \tag{24}$$

Proof. Pick $0 \neq u \in E_0V$ and consider the Φ -standard basis $\{E_i^*u\}_{i=0}^d$ for V . By (3),

$$AE_i^*u = c_{i+1}E_{i+1}^*u + a_iE_i^*u + b_{i-1}E_{i-1}^*u \quad (0 \leq i \leq d),$$

where $b_{-1} = 0$, $c_{d+1} = 0$ and $E_{-1}^* = 0$, $E_{d+1}^* = 0$. By this and since u is basis for E_0V ,

$$AE_i^*E_0 = c_{i+1}E_{i+1}^*E_0 + a_iE_i^*E_0 + b_{i-1}E_{i-1}^*E_0 \quad (0 \leq i \leq d).$$

By this and the first equation in Lemma 4.6,

$$(AA_i - c_{i+1}A_{i+1} - a_iA_i - b_{i-1}A_{i-1})E_0^*E_0 = 0 \quad (0 \leq i \leq d),$$

where $A_{-1} = 0$, $A_{d+1} = 0$. The map $\langle A \rangle \rightarrow \mathcal{A}E_0$, $Y \mapsto YE_0^*E_0$ is an \mathbb{F} -linear bijection, since this map is the composition of ρ and the bijection $\langle A^* \rangle \rightarrow \mathcal{A}E_0$, $Y \mapsto YE_0$ from Lemma 4.1. By these comments we obtain the result. \square

Lemma 4.18. *Assume that $d \geq 1$. Then $A = c_1A_1 + a_0I$.*

Proof. By Lemma 4.8 and (22). \square

Lemma 4.19. *There exist scalars p_{ij}^h ($0 \leq h, i, j \leq d$) in \mathbb{F} such that*

$$A_iA_j = \sum_{h=0}^d p_{ij}^h A_h \quad (0 \leq i, j \leq d). \tag{25}$$

Proof. By Lemma 4.11. \square

For notational convenience, define

$$q_{ij}^h = (p_{ij}^h)^* \quad (0 \leq h, i, j \leq d).$$

Applying (25) to Φ^* ,

$$A_i^*A_j^* = \sum_{h=0}^d q_{ij}^h A_h^* \quad (0 \leq i, j \leq d). \tag{26}$$

Lemma 4.20. *The following hold:*

- (i) $p_{ij}^h = p_{ji}^h \quad (0 \leq h, i, j \leq d)$;
- (ii) $\sum_{r=0}^d p_{hr}^t p_{ij}^r = \sum_{s=0}^d p_{sj}^t p_{hi}^s \quad (0 \leq h, i, j, t \leq d)$.

Proof. (i) Since the algebra $\langle A \rangle$ is commutative.

(ii) Expand $A_h(A_i A_j) = (A_h A_i)A_j$ in two ways using (25), and compare the coefficients. \square

Lemma 4.21. *We have $p_{ij}^0 = \delta_{i,j} k_i$ for $0 \leq i, j \leq d$.*

Proof. In (25), multiply each side on the left by $E_0 E_0^*$ and on the right by $E_0^* E_0$. Simplify the result using the equation on the left in (15) and the equation obtained from it by applying \dagger . This gives

$$E_0 E_i^* E_j^* E_0 = \sum_{h=0}^d p_{ij}^h E_0 E_0^* E_h^* E_0.$$

Thus

$$\delta_{i,j} E_0 E_i^* E_0 = p_{ij}^0 E_0 E_0^* E_0.$$

In this line, take the trace of each side and use Definition 3.3 along with (9), (11) to get the result. \square

Lemma 4.22. *For $0 \leq h, i, j \leq d$,*

$$k_h p_{ij}^h = k_i p_{hj}^i = k_j p_{hi}^j.$$

Proof. In view of Lemma 4.20(i), it suffices to show that $k_h p_{ij}^h = k_j p_{hi}^j$. To obtain this equation, set $t = 0$ in Lemma 4.20(ii), and evaluate the result using Lemma 4.21. \square

Lemma 4.23. *Assume that $d \geq 1$. Then*

$$\begin{aligned} c_i &= c_1 p_{1,i-1}^i & (1 \leq i \leq d), \\ a_i &= c_1 p_{1i}^i + a_0 & (0 \leq i \leq d), \\ b_i &= c_1 p_{1,i+1}^i & (0 \leq i \leq d-1). \end{aligned}$$

Proof. Compare Lemma 4.17 and (25) in light of Lemma 4.18. \square

In view of Lemma 4.17 we consider the following polynomials.

Definition 4.24. [28, Lemma 13.3] Define polynomials $\{v_i\}_{i=0}^d$ in $\mathbb{F}[\lambda]$ such that $v_0 = 1$ and for $d \geq 1$,

$$\begin{aligned} \lambda v_0 &= c_1 v_1 + a_0, \\ \lambda v_i &= c_{i+1} v_{i+1} + a_i v_i + b_{i-1} v_{i-1} \quad (1 \leq i \leq d-1). \end{aligned}$$

Lemma 4.25. *Referring to Definition 4.24, the following hold for $0 \leq i \leq d$:*

- (i) v_i has degree i and leading coefficient $(c_1 c_2 \cdots c_i)^{-1}$;
- (ii) $A_i = v_i(A)$;
- (iii) $A_i = \sum_{j=0}^d v_i(\theta_j) E_j$;
- (vi) $E_i = \nu^{-1} \sum_{j=0}^d v_i^*(\theta_j^*) A_j$.

Proof. (i) By Definition 4.24.

(ii) By Lemma 4.17 and Definition 4.24.

(iii) We have $A_i E_j = v_i(A) E_j = v_i(\theta_j) E_j$ for $0 \leq i, j \leq d$. We have $A_i = A_i I = \sum_{j=0}^d A_i E_j$. The result follows.

(iv) Applying (iii) to Φ^* we obtain $A_i^* = \sum_{j=0}^d v_i^*(\theta_j^*) E_j^*$. In this equation, apply ρ^* to each side. Simplify the result using Lemma 4.5 (applied to Φ^*) to get the result. \square

The following lemma is from [28, Theorem 12.4]; we give a short proof for the sake of completeness.

Lemma 4.26. *For $0 \leq i \leq d$,*

$$k_i = \frac{b_0 b_1 \cdots b_{i-1}}{c_1 c_2 \cdots c_i}. \tag{27}$$

Proof. We have $k_0 = 1$ by (12). By Lemma 4.22 (with $h = 1, j = i - 1$) and Lemma 4.23,

$$c_i k_i = b_{i-1} k_{i-1} \quad (1 \leq i \leq d).$$

The result follows. \square

Lemma 4.27. *We have $k_i \neq 0$ for $0 \leq i \leq d$.*

Proof. By Lemma 4.26. \square

Lemma 4.28. *For $0 \leq h, i, j \leq d$ the scalar p_{ij}^h is zero (resp. nonzero) if one of h, i, j is greater than (resp. equal to) the sum of the other two.*

Proof. By Lemmas 4.22 and 4.27, it suffices to consider the case $h \geq i + j$. For this case the result follows from (25) and the observation that for $0 \leq r \leq d$ the elements I, A, A^2, \dots, A^r and A_0, A_1, \dots, A_r form a basis for the same vector space. \square

Lemma 4.29. *We have*

$$k_i = v_i(\theta_0) \quad (0 \leq i \leq d).$$

Proof. For the equation in Lemma 4.25(iii), multiply each side by E_0 to get $A_i E_0 = v_i(\theta_0) E_0$. The result follows from this and Lemma 4.15. \square

Lemma 4.30. [28, Theorem 23.9] *For $0 \leq i \leq d$,*

$$k_i = \frac{\varphi_1 \varphi_2 \cdots \varphi_i}{\phi_1 \phi_2 \cdots \phi_i} \frac{\eta_d^*(\theta_0^*)}{\tau_i^*(\theta_i^*) \eta_{d-i}^*(\theta_i^*)} \quad (0 \leq i \leq d). \tag{28}$$

Proof. By (27) and Lemma 2.13. \square

The following lemma is from [28, Theorem 16.2]. We give a short proof for the sake of completeness.

Lemma 4.31. *Referring to Definition 4.24,*

$$v_i(\theta_j)/k_i = v_j^*(\theta_i^*)/k_j^* \quad (0 \leq i, j \leq d). \tag{29}$$

Proof. We evaluate $E_0^* A_i A_j^* E_0$ in two ways. By the equation on the right in (15),

$$A_j^* E_0 = \nu E_j E_0^* E_0. \tag{30}$$

By Lemma 4.25(iii),

$$A_i E_j = v_i(\theta_j) E_j. \tag{31}$$

By Lemma 3.15 (applied to Φ^*),

$$E_0^* E_j E_0^* = \nu^{-1} k_j^* E_0^*. \tag{32}$$

Evaluating $E_0^* A_i A_j^* E_0$ using (30)–(32) we obtain

$$E_0^* A_i A_j^* E_0 = v_i(\theta_j) k_j^* E_0^* E_0. \tag{33}$$

Apply this to Φ^* , then interchange i, j and apply \dagger to get

$$E_0^* A_i A_j^* E_0 = v_j^*(\theta_i^*) k_i E_0^* E_0. \tag{34}$$

We have $E_0 E_0^* \neq 0$ by (8). Now comparing (33) and (34) we obtain (29). \square

5. Boltzmann pairs and spin Leonard pairs

The notions of a Boltzmann pair and a spin Leonard pair were introduced by Curtin [7]. In this section we first recall these notions, and then prove a theorem about these topics.

Lemma 5.1. *Let A, A^* denote a Leonard pair on V . For invertible $W \in \langle A \rangle$ and invertible $W^* \in \langle A^* \rangle$ the following are equivalent:*

$$WA^*W^{-1} = (W^*)^{-1}AW^*, \tag{35}$$

$$W^{-1}A^*W = W^*A(W^*)^{-1}, \tag{36}$$

$$W^*WA^* = AW^*W, \tag{37}$$

$$A^*WW^* = WW^*A. \tag{38}$$

Proof. By linear algebra (37) is equivalent to (35), and (38) is equivalent to (36). The equations (37), (38) are equivalent since each is obtained from the other by applying †. □

Definition 5.2. [7, Definition 1.2] Let A, A^* denote a Leonard pair on V . By a *Boltzmann pair* for A, A^* we mean an ordered pair W, W^* such that

- (i) W is an invertible element of $\langle A \rangle$;
- (ii) W^* is an invertible element of $\langle A^* \rangle$;
- (iii) W, W^* satisfy the four equivalent conditions in Lemma 5.1.

The Leonard pair A, A^* is called a *spin Leonard pair* whenever there exists a Boltzmann pair for A, A^* .

For the rest of this section, let A, A^* denote a spin Leonard pair on V and let W, W^* denote a Boltzmann pair for A, A^* .

Lemma 5.3. *We have*

$$AW = WA, \quad A^*W^* = W^*A^*.$$

Proof. Since $W \in \langle A \rangle$ and $W^* \in \langle A^* \rangle$. □

Lemma 5.4. *For $\alpha, \beta \in \mathbb{F}$ with $\alpha \neq 0$, consider the Leonard pair $\alpha A + \beta I, \alpha A^* + \beta I$. Then W, W^* is a Boltzmann pair for this Leonard pair.*

Proof. By Definition 5.2. □

Lemma 5.5. [7, Lemma 3.3] *The following hold.*

- (i) *For nonzero $\alpha, \alpha^* \in \mathbb{F}$ the pair $\alpha W, \alpha^* W^*$ is a Boltzmann pair for A, A^* .*
- (ii) *The pair $W^{-1}, (W^*)^{-1}$ is a Boltzmann pair for A, A^* .*

Proof. (i) By Definition 5.2.

(ii) Compare lines (35) and (36). \square

Lemma 5.6. *The pair W^*, W is a Boltzmann pair for the Leonard pair A^*, A .*

Proof. By Lemma 5.1. \square

Lemma 5.7. *We have*

$$AWW^*W = WW^*WA^*, \quad A^*WW^*W = WW^*WA, \quad (39)$$

$$AW^*WW^* = W^*WW^*A^*, \quad A^*W^*WW^* = W^*WW^*A. \quad (40)$$

Proof. By Lemmas 5.1, 5.3. \square

Lemma 5.8. [7, Lemma 5.1] *The following agree up to a nonzero scalar factor in \mathbb{F} :*

$$WW^*W, \quad W^*WW^*.$$

Proof. Set $Z = WW^*W(W^*)^{-1}W^{-1}(W^*)^{-1}$. By Lemma 5.7, Z commutes with A and A^* . By this and Lemma 2.4, Z commutes with everything in $\text{End}(V)$. So Z is contained in the center of $\text{End}(V)$. Thus there exists $\alpha \in \mathbb{F}$ such that $Z = \alpha I$. The result follows. \square

Using W and W^* we obtain an action of the modular group $\text{PSL}_2(\mathbb{Z})$ on $\text{End}(V)$ as a group of automorphisms. Recall from [1] that $\text{PSL}_2(\mathbb{Z})$ has a presentation by generators ψ, σ and relations $\psi^3 = 1, \sigma^2 = 1$. The next result is a variation on [7, Lemma 5.2].

Lemma 5.9. *The group $\text{PSL}_2(\mathbb{Z})$ acts on $\text{End}(V)$ such that ψ sends $Y \mapsto (WW^*)^{-1}YWW^*$ and σ sends $Y \mapsto (WW^*W)^{-1}YWW^*W$ for $Y \in \text{End}(V)$.*

Proof. The automorphism σ swaps A, A^* by (39). Therefore A, A^* is self-dual with duality σ . Thus $\sigma^2 = 1$ by Lemma 2.7. By Lemma 5.8, there exists $\alpha \in \mathbb{F}$ such that $W^*WW^* = \alpha WW^*W$. So

$$(WW^*)^3 = WW^*WW^*WW^* = \alpha(WW^*W)^2.$$

Thus $\sigma^2 = \psi^3$. The result follows. \square

For the rest of this section, we identify σ with the automorphism of $\text{End}(V)$ from Lemma 5.9. We record a result from the proof of Lemma 5.9.

Lemma 5.10. *The spin Leonard pair A, A^* is self-dual with duality σ .*

The following definition is motivated by Lemma 5.8.

Definition 5.11. The Boltzmann pair W, W^* is said to be *balanced* whenever

$$WW^*W = W^*WW^*. \tag{41}$$

Lemma 5.12. *The following (i)–(iii) hold.*

- (i) *There exists a nonzero $\alpha \in \mathbb{F}$ such that the Boltzmann pair $\alpha W, W^*$ is balanced.*
- (ii) *There exists a nonzero $\alpha^* \in \mathbb{F}$ such that the Boltzmann pair $W, \alpha^* W^*$ is balanced.*
- (iii) *Assume that W, W^* is balanced. Then for any nonzero $\xi \in \mathbb{F}$, the Boltzmann pair $\xi W, \xi W^*$ is balanced.*

Proof. (i) By Lemma 5.8 there exists a nonzero $\alpha \in \mathbb{F}$ such that $\alpha WW^*W = W^*WW^*$. Thus the pair $\alpha W, W^*$ is balanced.

(ii) Similar.

(iii) By Definition 5.11. \square

For the rest of this section, assume that W, W^* is balanced.

Lemma 5.13. *The automorphism σ swaps W and W^* .*

Proof. In (41), multiply each side on the left by W to get

$$WWW^*W = WW^*WW^*.$$

In this line, multiply each side on the left by $(WW^*W)^{-1}$ to find that σ sends $W \mapsto W^*$. By this and Lemma 2.7, σ sends $W^* \mapsto W$. \square

Let $\{E_i\}_{i=0}^d$ denote a standard ordering of the primitive idempotents of A . Define

$$E_i^* = E_i^\sigma \quad (0 \leq i \leq d). \tag{42}$$

Observe that the sequence

$$\Phi = (A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d) \tag{43}$$

is a Leonard system on V .

Lemma 5.14. *The Leonard system Φ is self-dual with duality σ .*

Proof. By Lemma 5.10, σ swaps A and A^* . By (42) and Lemma 2.7, σ swaps E_i and E_i^* for $0 \leq i \leq d$. The result follows. \square

Lemma 5.15. For $0 \leq i \leq d$,

$$W^*WE_i^* = E_iW^*W, \quad WW^*E_i = E_i^*WW^*. \tag{44}$$

Proof. First we obtain the equation on the left in (44). By Lemma 5.9 and (42),

$$WW^*WE_i^* = E_iWW^*W.$$

In this line, multiply each side on the left by W^{-1} . Simplify the result using $E_i \in \langle A \rangle$ and Lemma 5.3 to get the equation on the left in (44). In this equation, apply \dagger to each side to get the equation on the right in (44). \square

Lemma 5.16. There exist nonzero scalars $f, \{\tau_i\}_{i=0}^d$ in \mathbb{F} such that $\tau_0 = 1$ and

$$W = f \sum_{i=0}^d \tau_i E_i, \quad W^* = f \sum_{i=0}^d \tau_i E_i^*. \tag{45}$$

Proof. First we obtain the equation on the left in (45). By $W \in \langle A \rangle$ and since $\{E_i\}_{i=0}^d$ is a basis for $\langle A \rangle$, there exist scalars $\{\alpha_i\}_{i=0}^d$ in \mathbb{F} such that $W = \sum_{i=0}^d \alpha_i E_i$. The scalars $\{\alpha_i\}_{i=0}^d$ are nonzero since W is invertible. Define $f = \alpha_0$ and $\tau_i = \alpha_0^{-1} \alpha_i$ for $0 \leq i \leq d$. Then the equation on the left in (45) holds. By construction the scalars $f, \{\tau_i\}_{i=0}^d$ are all nonzero and $\tau_0 = 1$. In the equation on the left in (45), apply σ to each side, and use Lemma 5.13 and (42) to get the equation on the right in (45). \square

Note 5.17. The scalars $\{\tau_i\}_{i=0}^d$ do not change under the adjustments of Lemma 5.12(iii), but f can be adjusted to have any nonzero value.

Lemma 5.18. We have

$$W^{-1} = f^{-1} \sum_{i=0}^d \tau_i^{-1} E_i, \quad (W^*)^{-1} = f^{-1} \sum_{i=0}^d \tau_i^{-1} E_i^*. \tag{46}$$

Proof. Follows from (45). \square

Recall the map π from Definition 3.3.

Lemma 5.19. We have

$$\pi(W) = f, \quad \pi^*(W^*) = f. \tag{47}$$

Proof. Apply π to the equation on the left in (45) and use Lemma 3.5 to get the equation on the left in (47). The equation on the right in (47) is similarly obtained. \square

Lemma 5.20. *There exists $\gamma \in \mathbb{F}$ such that*

$$\pi(W^*) = f\gamma, \qquad \pi^*(W) = f\gamma. \tag{48}$$

Proof. Define $\gamma = f^{-1}\pi(W^*)$. Then the equation on the left in (48) holds. By Lemma 5.13, $W^\sigma = W^*$. By Lemma 3.10, $\pi(W^\sigma) = \pi^*(W)$. By these comments we get the equation on the right in (48). \square

Note 5.21. The scalar γ does not change under the adjustments of Lemma 5.12(iii).

Lemma 5.22. *We have*

$$\pi(WW^*) = f^2\gamma, \qquad \pi(W^*W) = f^2\gamma, \tag{49}$$

$$\pi^*(WW^*) = f^2\gamma, \qquad \pi^*(W^*W) = f^2\gamma. \tag{50}$$

Proof. By Lemma 3.7(i), $\pi(WW^*) = \pi(W^*W) = \pi(W)\pi(W^*)$. By this and (47), (48), we obtain (49). The equations (50) are similarly obtained. \square

Recall the map ρ from Lemma 4.2.

Lemma 5.23. *We have*

$$W^\rho = f^2\gamma(W^*)^{-1}, \qquad (W^*)^{\rho^*} = f^2\gamma W^{-1}. \tag{51}$$

Proof. First we obtain the equation on the left in (51). In the equation on the left in (44) for $i = 0$, multiply each side on the right by E_0 to get

$$W^*WE_0^*E_0 = E_0W^*WE_0.$$

By Lemma 3.4, $E_0W^*WE_0 = \pi(W^*W)E_0$. By these comments and (49),

$$W^*WE_0^*E_0 = f^2\gamma E_0.$$

In this line, multiply each side on the left by $(W^*)^{-1}$ to get

$$WE_0^*E_0 = f^2\gamma(W^*)^{-1}E_0.$$

By this and Lemmas 4.1, 4.2 we get the equation on the left in (51). The equation on the right in (51) is similarly obtained. \square

Lemma 5.24. *We have $\gamma \neq 0$.*

Proof. We have $W^\rho \neq 0$ since $W \neq 0$ and ρ is bijective. By this and (51) we get $\gamma \neq 0$. \square

We now give our main result in this section. Recall the elements $\{A_i\}_{i=0}^d$ from (14) and the scalar ν from (9).

Theorem 5.25. *We have*

$$W = f\gamma \sum_{i=0}^d \tau_i^{-1} A_i, \quad W^{-1} = \nu^{-1} f^{-1} \gamma^{-1} \sum_{i=0}^d \tau_i A_i, \tag{52}$$

$$W^* = f\gamma \sum_{i=0}^d \tau_i^{-1} A_i^*, \quad (W^*)^{-1} = \nu^{-1} f^{-1} \gamma^{-1} \sum_{i=0}^d \tau_i A_i^*. \tag{53}$$

Proof. First we obtain the equation on the left in (52). By (46) and (51),

$$W^\rho = f\gamma \sum_{i=0}^d \tau_i^{-1} E_i^*.$$

In this line, apply ρ^* to each side, and use Lemmas 4.3, 4.5 to get the equation on the left in (52). Next we obtain the equation on the right in (52). By (51), $W^{-1} = f^{-2} \gamma^{-1} (W^*)^{\rho^*}$. In the equation on the right in (45), apply ρ^* to each side and use Lemma 4.5 to get $(W^*)^{\rho^*} = f\nu^{-1} \sum_{i=0}^d \tau_i A_i$. By these comments we get the equation on the right in (52). The equations (53) are similarly obtained. \square

We mention a lemma for later use. Recall the scalars $\{k_i\}_{i=0}^d$ from (11).

Lemma 5.26. *We have*

$$\gamma = \nu^{-1} \sum_{i=0}^d k_i \tau_i. \tag{54}$$

Proof. In the equation on the right in (45), apply π to each side, and use (11), (48). \square

Note 5.27. In [7, Theorem 1.13], the spin Leonard pairs are classified. By that classification, there are five families of spin Leonard pairs, called type I–V; see [7, Lemmas 1.7–1.11]. For each family, the eigenvalue sequence and the intersection numbers are explicitly given. In terms of [28, Section 35], type I is of q -Racah type, type II is of Racah type, type III is of Krawtchouk type, type IV is of Bannai-Ito type with odd diameter, and type V is of Bannai-Ito type with even diameter. In [7, Theorem 1.18], for each spin Leonard pair A, A^* the corresponding Boltzmann pairs are given. We say that two Boltzmann pairs for A, A^* are *equivalent* whenever one is obtained from the other by the adjustments in Lemma 5.5. The Leonard pair A, A^* has precisely one equivalence

class of Boltzmann pairs, unless A, A^* has type V and the intersection number a_i is independent of i for $0 \leq i \leq d$. In this case A, A^* has precisely two equivalence classes of Boltzmann pairs. For one of the equivalence classes we have $\tau_i = (-1)^{\lceil i/2 \rceil}$ for $0 \leq i \leq d$, and for the other equivalence class, we have $\tau_i = (-1)^{\lfloor i/2 \rfloor}$ for $0 \leq i \leq d$.

6. Leonard pairs of q -Racah type

In this section we consider a special case of Leonard pair, said to have q -Racah type [13]. This is the “most general” type of Leonard pair. Throughout this section the following notation is in effect. Fix a nonzero $q \in \mathbb{F}$ such that $q^4 \neq 1$. Let $\Phi = (A; \{E_i\}_{i=0}^d; A^*; \{E_i^*\}_{i=0}^d)$ denote a Leonard system on V with parameter array $(\{\theta_i\}_{i=0}^d; \{\theta_i^*\}_{i=0}^d; \{\varphi_i\}_{i=1}^d; \{\phi_i\}_{i=1}^d)$. To avoid trivialities we assume $d \geq 1$.

Definition 6.1. [13, Definition 5.1] We say that Φ has q -Racah type whenever there exist nonzero $a, b \in \mathbb{F}$ such that

$$\theta_i = aq^{2i-d} + a^{-1}q^{d-2i} \quad (0 \leq i \leq d), \tag{55}$$

$$\theta_i^* = bq^{2i-d} + b^{-1}q^{d-2i} \quad (0 \leq i \leq d). \tag{56}$$

In this case the scalars a, b are unique.

Lemma 6.2. [13, Corollary 6.4] Assume that Φ has q -Racah type. Let $a, b \in \mathbb{F}$ denote nonzero scalars that satisfy (55), (56). Then there exists a nonzero $c \in \mathbb{F}$ such that for $1 \leq i \leq d$,

$$\varphi_i = a^{-1}b^{-1}q^{d+1}(q^i - q^{-i})(q^{i-d-1} - q^{d-i+1})(q^{-i} - abcq^{i-d-1})(q^{-i} - abc^{-1}q^{i-d-1}), \tag{57}$$

$$\phi_i = ab^{-1}q^{d+1}(q^i - q^{-i})(q^{i-d-1} - q^{d-i+1})(q^{-i} - a^{-1}bcq^{i-d-1})(q^{-i} - a^{-1}bc^{-1}q^{i-d-1}). \tag{58}$$

Moreover, c is unique up to inverse.

Definition 6.3. Assume that Φ has q -Racah type. Let $a, b \in \mathbb{F}$ denote nonzero scalars that satisfy (55), (56), and let $c \in \mathbb{F}$ denote a nonzero scalar that satisfies (57), (58). We call the sequence (a, b, c, d) a *Huang data* of Φ .

Note 6.4. Assume that Φ has q -Racah type with Huang data (a, b, c, d) . Then (a, b, c^{-1}, d) is a Huang data for Φ , and Φ has no further Huang data.

For the rest of this section, assume that Φ has q -Racah type with Huang data (a, b, c, d) .

Lemma 6.5. [13, Definition 7.1] The following (i)–(iii) hold:

- (i) $q^{2i} \neq 1$ for $1 \leq i \leq d$;
- (ii) Neither of a^2, b^2 is among $q^{2d-2}, q^{2d-4}, \dots, q^{2-2d}$;
- (iii) None of $abc, a^{-1}bc, ab^{-1}c, abc^{-1}$ is among $q^{d-1}, q^{d-3}, \dots, q^{1-d}$.

Lemma 6.6. *The Leonard system*

$$(-A; \{E_i\}_{i=0}^d; -A^*; \{E_i^*\}_{i=0}^d) \tag{59}$$

has q -Racah type with Huang data $(-a, -b, c, d)$.

Proof. Use Definitions 6.1, 6.3. \square

Recall the intersection numbers $\{b_i\}_{i=0}^{d-1}, \{c_i\}_{i=1}^d$ of Φ from (3).

Lemma 6.7. *We have*

$$b_i = \frac{(q^{i-d} - q^{d-i})(bq^{i-d} - b^{-1}q^{d-i})(abq^i - cq^{d-i-1})(abq^i - c^{-1}q^{d-i-1})}{abq^{d-1}(bq^{2i-d} - b^{-1}q^{d-2i})(bq^{2i-d+1} - b^{-1}q^{d-2i-1})},$$

$$c_i = \frac{(q^i - q^{-i})(bq^i - b^{-1}q^{-i})(bq^i - acq^{d-i+1})(bq^i - ac^{-1}q^{d-i+1})}{abq^{d+1}(bq^{2i-d} - b^{-1}q^{d-2i})(bq^{2i-d-1} - b^{-1}q^{d-2i+1})}$$

for $1 \leq i \leq d - 1$, and

$$b_0 = \frac{(q^{-d} - q^d)(ab - cq^{d-1})(ab - c^{-1}q^{d-1})}{abq^{d-1}(bq^{1-d} - b^{-1}q^{d-1})},$$

$$c_d = \frac{(q^d - q^{-d})(ac - bq^{d-1})(ac^{-1} - bq^{d-1})}{abq^{d-1}(bq^{d-1} - b^{-1}q^{1-d})}.$$

Proof. Evaluate (5), (6) using (55)–(58). \square

We recall some notation. For $\alpha \in \mathbb{F}$,

$$(\alpha; q)_n = (1 - \alpha)(1 - \alpha q) \cdots (1 - \alpha q^{n-1}) \quad n = 0, 1, 2, \dots$$

We interpret $(\alpha; q)_0 = 1$. Recall the scalar ν be from (8).

Lemma 6.8. *We have*

$$\nu = \frac{a^d b^d c^d q^{d(1-d)} (a^{-2}; q^2)_d (b^{-2}; q^2)_d}{(ab^{-1}cq^{1-d}; q^2)_d (a^{-1}bcq^{1-d}; q^2)_d}.$$

Proof. Evaluate (10) using (55)–(58). \square

Recall the scalars $\{k_i\}_{i=0}^d$ from (11).

Lemma 6.9. We have $k_0 = 1$, and

$$k_i = \frac{q^{2id}(1 - b^2q^{4i-2d})(b^2q^{2-2d}; q^2)_i (q^{-2d}; q^2)_i (abcq^{1-d}; q^2)_i (abc^{-1}q^{1-d}; q^2)_i}{a^{2i}(1 - b^2q^{2i-2d})(b^2q^2; q^2)_i (q^2; q^2)_i (a^{-1}bcq^{1-d}; q^2)_i (a^{-1}bc^{-1}q^{1-d}; q^2)_i}$$

for $1 \leq i \leq d - 1$, and

$$k_d = \frac{q^{2d^2}(b^2q^{2-2d}; q^2)_d (q^{-2d}; q^2)_d (abcq^{1-d}; q^2)_d (abc^{-1}q^{1-d}; q^2)_d}{a^{2d}(b^2; q^2)_d (q^2; q^2)_d (a^{-1}bcq^{1-d}; q^2)_d (a^{-1}bc^{-1}q^{1-d}; q^2)_d}.$$

Proof. Evaluate (28) using (55)–(58). \square

We now turn our attention to the self-dual case.

Lemma 6.10. The Leonard system Φ is self-dual if and only if $a = b$.

Proof. By Lemma 2.15 and Definition 6.1. \square

We now turn our attention to the spin case.

Lemma 6.11. [7, Theorem 1.13] The following hold.

- (i) Assume $d \geq 3$, and A, A^* is a spin Leonard pair. Then $b = a$ and $c \in \{a, a^{-1}, -a, -a^{-1}\}$.
- (ii) Assume that $b = a$ and $c \in \{a, a^{-1}, -a, -a^{-1}\}$. Then A, A^* is a spin Leonard pair.

Note 6.12. We have a comment about Lemma 6.11. Assume that $b = a$ and $c \in \{-a, -a^{-1}\}$. By Note 6.4 we may assume that $c = -a$. Replacing Φ by (59) we may assume that $c = a$.

In view of Note 6.12 we focus on the case $a = b = c$.

Lemma 6.13. Assume that $a = b = c$. Then

$$\begin{aligned} b_0 &= \frac{(q^{-d} - q^d)(a^3 - q^{d-1})}{a(a + q^{d-1})}, \\ b_i &= \frac{(q^{i-d} - q^{d-i})(aq^{i-d} - a^{-1}q^{d-i})(a^3 - q^{d-2i-1})}{a(aq^{2i-d} - a^{-1}q^{d-2i})(a + q^{d-2i-1})} \quad (1 \leq i \leq d - 1), \\ c_i &= \frac{a(q^i - q^{-i})(aq^i - a^{-1}q^{-i})(a^{-1} - q^{d-2i+1})}{(aq^{2i-d} - a^{-1}q^{d-2i})(a + q^{d-2i+1})} \quad (1 \leq i \leq d - 1), \\ c_d &= \frac{(q^{-d} - q^d)(a - q^{d-1})}{q^{d-1}(a + q^{1-d})}. \end{aligned}$$

Proof. Set $a = b = c$ in Lemma 6.7. \square

Lemma 6.14. Assume that $a = b = c$. Then

$$\nu = \frac{a^{3d}(a^{-2}; q^2)_d^2}{q^{d(d-1)}(aq^{1-d}; q^2)_d^2}. \tag{60}$$

Proof. Set $a = b = c$ in Lemma 6.8. \square

Lemma 6.15. Assume that $a = b = c$. Then $k_0 = 1$, and

$$k_i = \frac{q^{2id}(1 - a^2q^{4i-2d})(a^2q^{2-2d}; q^2)_i(q^{-2d}; q^2)_i(a^3q^{1-d}; q^2)_i}{a^{2i}(1 - a^2q^{2i-2d})(a^2q^2; q^2)_i(q^2; q^2)_i(a^{-1}q^{1-d}; q^2)_i}$$

for $1 \leq i \leq d - 1$, and

$$k_d = \frac{q^{2d^2}(a^2q^{2-2d}; q^2)_d(q^{-2d}; q^2)_d(a^3q^{1-d}; q^2)_d}{a^{2d}(a^2; q^2)_d(q^2; q^2)_d(a^{-1}q^{1-d}; q^2)_d}.$$

Proof. Set $a = b = c$ in Lemma 6.9. \square

Lemma 6.16. [7, Theorem 1.18] Assume that $a = b = c$. Let $f, \{\tau_i\}_{i=0}^d$ denote nonzero scalars in \mathbb{F} such that $\tau_0 = 1$. Define

$$W = f \sum_{i=0}^d \tau_i E_i, \quad W^* = f \sum_{i=0}^d \tau_i E_i^*. \tag{61}$$

(i) Assume that W, W^* is a Boltzmann pair for A, A^* . Then one of the following holds provided that $d \geq 3$:

$$\tau_i = (-1)^i a^{-i} q^{i(d-i)} \quad (0 \leq i \leq d), \tag{62}$$

$$\tau_i^{-1} = (-1)^i a^{-i} q^{i(d-i)} \quad (0 \leq i \leq d). \tag{63}$$

(ii) Assume that one of (62), (63) holds. Then W, W^* is a balanced Boltzmann pair for A, A^* .

Note 6.17. Referring to Lemma 6.16, if (63) holds then (62) holds for the Boltzmann pair $W^{-1}, (W^*)^{-1}$. For this reason we will focus on (62).

Lemma 6.18. Assume that $b = a$ and define $\{\tau_i\}_{i=0}^d$ by (62). Let $0 \neq f \in \mathbb{F}$ and define W, W^* as in (61). Assume that W, W^* is a Boltzmann pair for A, A^* . Then $c \in \{a, a^{-1}\}$.

Proof. First assume that $d \geq 3$. By way of contradiction assume that $c \notin \{a, a^{-1}\}$. By Lemma 6.11(i) we have $c \in \{-a, -a^{-1}\}$ and $\text{Char}(\mathbb{F}) \neq 2$. By Note 6.4 we may assume that $c = -a$. Thus Φ has Huang data $(a, a, -a, d)$. By Lemma 6.6 the Leonard system (59) has Huang data $(-a, -a, -a, d)$. By Definition 5.11 and Lemma 5.4 the pair W, W^*

is a Boltzmann pair for $-A, -A^*$. Applying Lemma 6.16(i) to the Leonard system (59), one of the following holds:

$$\begin{aligned} \tau_i &= (-1)^i (-a)^{-i} q^{i(d-i)} & (0 \leq i \leq d), \\ \tau_i^{-1} &= (-1)^i (-a)^{-i} q^{i(d-i)} & (0 \leq i \leq d). \end{aligned}$$

In either case, comparing the values of τ_1, τ_3 with (62) we get a contradiction by $\text{Char}(\mathbb{F}) \neq 2$ and Lemma 6.5. We have shown the assertion for the case $d \geq 3$. Next assume that $d \leq 2$. Consider the matrices that represent the elements $A, A^*, \{E_i\}_{i=0}^d, \{E_i^*\}_{i=0}^d, W, W^*$ with respect to the Φ -split basis for V . We represent these matrices in terms of a, c, q using (55)–(58). The matrices representing A, A^* are as in Lemma 2.11. The matrices representing $\{E_i\}_{i=0}^d, \{E_i^*\}_{i=0}^d$ can be obtained using (1). The matrices representing W, W^* are obtained by (61). Now compute each side of (37) in matrix form, and compare the $(d, 0)$ -entries to get $(a - c)(ac - 1) = 0$. Thus $c \in \{a, a^{-1}\}$. \square

The following lemma is a reformulation of a special case of [11, 2.22].

Lemma 6.19. *Let α, β denote commuting indeterminates. Then*

$$\sum_{i=0}^d \frac{(1 - \alpha q^{-4i}) (\alpha^{-1} q^2; q^2)_i (\beta^{-1}; q^2)_i (q^{-2d}; q^2)_i q^{i(2d-i+3)} (-\beta)^i}{(1 - \alpha q^{-2i}) (q^2, q^2)_i (\beta q^2/\alpha; q^2)_i (q^{2d+2}/\alpha; q^2)_i} = \frac{(\alpha q^{-2d}, q^2)_d}{(\alpha q^{-2d}/\beta, q^2)_d}. \tag{64}$$

Proof. Using

$$(x; q^2)_i = (x^{-1}; q^{-2})_i (-x)^i q^{i(i-1)},$$

one routinely checks that the left-hand side of (64) is equal to

$$\sum_{i=0}^d \frac{(\alpha q^{-2}; q^{-2})_i (1 - \alpha q^{-4i}) (\beta; q^{-2})_i (q^{2d}; q^{-2})_i}{(q^{-2}; q^{-2})_i (1 - \alpha q^{-2i}) (\alpha q^{-2}/\beta; q^{-2})_i (\alpha q^{-2d-2}; q^{-2})_i} \left(-\frac{\alpha q^{-2d-2}}{\beta} \right)^i q^{i(1-i)}.$$

By this and [11, 2.22], the left-hand side of (64) is equal to

$$\frac{(\alpha q^{-2}; q^{-2})_d}{(\alpha q^{-2}/\beta; q^{-2})_d}.$$

One checks that

$$(\alpha q^{-2}; q^{-2})_d = (\alpha q^{-2d}; q^2)_d, \quad (\alpha q^{-2}/\beta; q^{-2})_d = (\alpha q^{-2d}/\beta; q^2)_d.$$

The result follows. \square

Lemma 6.20. Referring to Lemma 6.16, assume that the $\{\tau_i\}_{i=0}^d$ satisfy (62). Then

$$\sum_{i=0}^d k_i \tau_i = \frac{(a^{-2}; q^2)_d}{(aq^{1-d}; q^2)_d}. \tag{65}$$

Moreover the scalar γ from Lemma 5.20 satisfies

$$\gamma = \frac{q^{d(d-1)}(aq^{1-d}; q^2)_d}{a^{3d}(a^{-2}; q^2)_d}. \tag{66}$$

Proof. We first show (65). Using the data in Lemma 6.15 we find that the left-hand side of (65) is equal to the sum

$$1 + \sum_{i=1}^{d-1} \frac{(1 - \alpha^{-1}q^{4i})(\alpha^{-1}q^2; q^2)_i (\beta^{-1}; q^2)_i (q^{-2d}; q^2)_i (-\beta)^i q^{i(2d-i+1)}}{(1 - \alpha^{-1}q^{2i})(q^2, q^2)_i (\beta q^2/\alpha; q^2)_i (q^{2d+2}/\alpha; q^2)_i} + \frac{(\alpha^{-1}q^2; q^2)_d (\beta^{-1}; q^2)_d (q^{-2d}; q^2)_d (-\beta)^d q^{d(d+1)}}{(q^2; q^2)_d (\beta q^2/\alpha; q^2)_d (q^{2d}/\alpha; q^2)_d}, \tag{67}$$

where $\alpha = a^{-2}q^{2d}$, $\beta = a^{-3}q^{d-1}$. One routinely checks that the sum (67) is equal to the left-hand side of (64). Thus (65) holds by Lemma 6.19. To get (66), evaluate (54) using (60), (65). \square

7. Type II matrices and spin models

We now turn our attention to type II matrices and spin models. For the rest of this paper, fix a finite nonempty set X . Let $\text{Mat}_X(\mathbb{C})$ denote the \mathbb{C} -algebra consisting of the matrices that have all entries in \mathbb{C} and whose rows and columns are indexed by X . For $R \in \text{Mat}_X(\mathbb{C})$ and $x, y \in X$ the (x, y) -entry of R is denoted by $R(x, y)$. Let I (resp. J) denote the identity matrix (resp. all 1's matrix) in $\text{Mat}_X(\mathbb{C})$. For $R, S \in \text{Mat}_X(\mathbb{C})$ let $R \circ S$ denote their Hadamard (entry-wise) product. Let V denote the vector space over \mathbb{C} consisting of the column vectors whose entries are indexed by X . Note that $\dim V = |X|$. The algebra $\text{Mat}_X(\mathbb{C})$ acts on V by left multiplication. For $y \in X$ define $\hat{y} \in V$ that has y -entry 1 and all other entries 0. Note that $\{\hat{y}\}_{y \in X}$ form a basis for V . For a real number $\alpha > 0$ let $\alpha^{1/2}$ denote the positive square root of α .

Definition 7.1. [17, Definition 2.1] A matrix $W \in \text{Mat}_X(\mathbb{C})$ is said to be *type II* whenever W is symmetric with all entries nonzero and

$$\sum_{y \in X} \frac{W(a, y)}{W(b, y)} = |X| \delta_{a,b} \quad (a, b \in X). \tag{68}$$

For a symmetric $W \in \text{Mat}_X(\mathbb{C})$ with all entries nonzero, define $W^- \in \text{Mat}_X(\mathbb{C})$ by

$$W^-(y, z) = (W(y, z))^{-1} \quad (y, z \in X). \tag{69}$$

Referring to the equation (68), for $a, b \in X$ the left-hand side is equal to the (a, b) -entry of WW^- .

Lemma 7.2. *Assume $W \in \text{Mat}_X(\mathbb{C})$ is symmetric with all entries nonzero. Then the following (i)–(iii) are equivalent:*

- (i) W is type II;
- (ii) $WW^- = |X|I$;
- (iii) $WW^- = \alpha I$ for some $\alpha \in \mathbb{C}$.

Proof. (i) \Rightarrow (ii) By (68) and (69).
 (ii) \Rightarrow (iii) Clear.
 (iii) \Rightarrow (i) By the comment below (69). \square

Note 7.3. Assume $W \in \text{Mat}_X(\mathbb{C})$ is type II. Then αW is type II for $0 \neq \alpha \in \mathbb{C}$. Moreover, W^{-1} and W^- are type II.

Definition 7.4. [21, Section 1] Assume $W \in \text{Mat}_X(\mathbb{C})$ is type II. For $b, c \in X$ define $\mathbf{u}_{b,c} \in V$ by

$$\mathbf{u}_{b,c} = \sum_{y \in X} \frac{W(b, y)}{W(c, y)} \hat{y}. \tag{70}$$

Define

$$N(W) = \{B \in \text{Mat}_X(\mathbb{C}) \mid B \text{ is symmetric, } B\mathbf{u}_{b,c} \in \mathbb{C}\mathbf{u}_{b,c} \text{ for all } b, c \in X\}.$$

Lemma 7.5. [21, Theorem 6] *Assume $W \in \text{Mat}_X(\mathbb{C})$ is type II. Then $N(W)$ is a commutative subalgebra of $\text{Mat}_X(\mathbb{C})$ that contains J and is closed under the Hadamard product.*

We have a comment.

Lemma 7.6. *Assume $W \in \text{Mat}_X(\mathbb{C})$ is type II. Then $N(\alpha W) = N(W)$ for $0 \neq \alpha \in \mathbb{C}$. Moreover each of $N(W^-)$, $N(W^{-1})$ is equal to $N(W)$.*

Proof. Clearly $N(\alpha W) = N(W)$. For $b, c \in X$ define $\mathbf{u}_{b,c}^- \in V$ by

$$\mathbf{u}_{b,c}^- = \sum_{y \in X} \frac{W^-(b, y)}{W^-(c, y)} \hat{y}.$$

Using (69), (70) we obtain

$$\mathbf{u}_{b,c}^- = \sum_{y \in X} \frac{W(c, y)}{W(b, y)} \hat{y} = \mathbf{u}_{c,b}.$$

By this and Definition 7.4 we find $N(W^-) = N(W)$. We have $W^- = |X|W^{-1}$ by Lemma 7.2, so $N(W^{-1}) = N(W^-)$. The result follows. \square

Definition 7.7. [17, Definition 2.1] A matrix $W \in \text{Mat}_X(\mathbb{C})$ is called a *spin model* whenever W is type II and

$$\sum_{y \in X} \frac{W(a, y)W(b, y)}{W(c, y)} = |X|^{1/2} \frac{W(a, b)}{W(a, c)W(b, c)} \quad (a, b, c \in X). \tag{71}$$

Note 7.8. The equation (71) is often called the type III condition.

Note 7.9. [17, Section 2] (see also [18, Proposition 2.1]). Assume $W \in \text{Mat}_X(\mathbb{C})$ is a spin model. Then $-W$ is a spin model. Moreover W^- is a spin model.

Lemma 7.10. [21, Lemma 8] Assume $W \in \text{Mat}_X(\mathbb{C})$ is a spin model. Then $W \in N(W)$.

8. Hadamard matrices and spin models

In this section we consider a certain class of type II matrices and spin models.

Definition 8.1. A matrix $H \in \text{Mat}_X(\mathbb{C})$ is called *Hadamard* whenever every entry is ± 1 and $HH^t = |X|I$.

Example 8.2. The matrix

$$H = \begin{pmatrix} 1 & -1 & -1 & -1 \\ -1 & 1 & -1 & -1 \\ -1 & -1 & 1 & -1 \\ -1 & -1 & -1 & 1 \end{pmatrix}$$

is Hadamard.

Lemma 8.3. A symmetric Hadamard matrix is type II.

Proof. By Definitions 7.1 and 8.1. \square

Lemma 8.4. For $W \in \text{Mat}_X(\mathbb{C})$ and $0 \neq \alpha \in \mathbb{C}$ the following are equivalent:

- (i) W is type II with all entries $\pm\alpha$;
- (ii) there exists a symmetric Hadamard matrix H such that $W = \alpha H$.

Proof. (i) \Rightarrow (ii) Assume that W is type II with all entries $\pm\alpha$. Define $H = \alpha^{-1}W$. Then the entries of H are ± 1 . Moreover, by Note 7.3, H is type II. Applying Lemma 7.2 to H

and using $H^- = H$, we obtain $H^2 = |X|I$. The matrix H is symmetric since it is type II. By these comments H is a symmetric Hadamard matrix.

(ii) \Rightarrow (i) Assume that there exists a symmetric Hadamard matrix H such that $W = \alpha H$. Then W is symmetric and satisfies (68). So W is type II. Clearly the entries of W are $\pm\alpha$. \square

Definition 8.5. A type II matrix $W \in \text{Mat}_X(\mathbb{C})$ is said to have *Hadamard type* whenever there exists a symmetric Hadamard matrix H and $0 \neq \alpha \in \mathbb{C}$ such that $W = \alpha H$.

We have an example of a spin model having Hadamard type.

Example 8.6. Let the matrix H be from Example 8.2. Then the matrix $W = \sqrt{-1}H$ is a spin model of Hadamard type.

Note 8.7. Spin models of Hadamard type sometimes cause technical problems, so occasionally we will assume that a spin model under discussion does not have Hadamard type.

9. Distance-regular graphs

We now turn our attention to distance-regular graphs [3,4,10]. Let Γ denote an undirected, connected graph, without loops or multiple edges, with vertex set X . For $x, y \in X$ let $\partial(x, y)$ denote the path-length distance between x and y . By the *diameter* of Γ we mean $D = \max\{\partial(x, y) \mid x, y \in X\}$. To avoid trivialities, we always assume $D \geq 1$. For $x \in X$ define

$$\Gamma_i(x) = \{y \in X \mid \partial(x, y) = i\} \quad (0 \leq i \leq D).$$

The graph Γ is said to be *regular* whenever $k = |\Gamma_1(x)|$ is independent of $x \in X$. In this case, k is called the *valency* of Γ . The graph Γ is said to be *distance-regular* whenever for $0 \leq h, i, j \leq D$ and vertices $x, y \in X$ at distance $\partial(x, y) = h$, the number

$$p_{ij}^h = |\Gamma_i(x) \cap \Gamma_j(y)|$$

is independent of x, y . The integers p_{ij}^h are called the *intersection numbers* of Γ .

For the rest of this section, assume that Γ is distance-regular. Define

$$k_i = p_{ii}^0 \quad (0 \leq i \leq D). \tag{72}$$

We have $|\Gamma_i(x)| = k_i$ for $x \in X$ and $0 \leq i \leq D$. Observe that $k_0 = 1$, and that Γ is regular with valency $k = k_1$. Moreover

$$\sum_{i=0}^D k_i = |X|.$$

For $0 \leq h, i, j \leq D$ the scalar p_{ij}^h is zero (resp. nonzero) if one of h, i, j is greater than (resp. equal to) the sum of the other two. Define

$$c_i = p_{1,i-1}^i \quad (1 \leq i \leq D), \tag{73}$$

$$a_i = p_{1i}^i \quad (0 \leq i \leq D), \tag{74}$$

$$b_i = p_{1,i+1}^i \quad (0 \leq i \leq D - 1). \tag{75}$$

For notational convenience define $c_0 = 0$ and $b_D = 0$. We have $c_i b_{i-1} \neq 0$ for $1 \leq i \leq D$. Moreover

$$a_0 = 0, \quad b_0 = k, \quad c_1 = 1, \quad c_i + a_i + b_i = k \quad (0 \leq i \leq D). \tag{76}$$

By [4, Chapter 4, (1c)],

$$k_i = \frac{b_0 b_1 \cdots b_{i-1}}{c_1 c_2 \cdots c_i} \quad (0 \leq i \leq D).$$

For $0 \leq i \leq D$ define $A_i \in \text{Mat}_X(\mathbb{C})$ that has (x, y) -entry 1 if $\partial(x, y) = i$ and 0 if $\partial(x, y) \neq i$ ($x, y \in X$). Note that A_i is symmetric. The matrix A_i is called the *i*th distance-matrix of Γ . Note that $A_0 = I$, and A_1 is the adjacency matrix of Γ . Observe that

$$A_i \circ A_j = \delta_{i,j} A_i \quad (0 \leq i, j \leq D), \quad \sum_{i=0}^D A_i = J, \tag{77}$$

$$A_i A_j = \sum_{h=0}^D p_{ij}^h A_h \quad (0 \leq i, j \leq D). \tag{78}$$

By the equation on the left in (77) the matrices $\{A_i\}_{i=0}^D$ are linearly independent. Let M denote the subspace of $\text{Mat}_X(\mathbb{C})$ with basis $\{A_i\}_{i=0}^D$. We have $p_{ij}^h = p_{ji}^h$ for $0 \leq h, i, j \leq D$. By this and (78), M is a commutative subalgebra of $\text{Mat}_X(\mathbb{C})$. By (77), M is closed under \circ and contains J . We call M the *Bose-Mesner algebra* of Γ .

By [4, Theorem 2.6.1(ii)] there exists a basis $\{E_i\}_{i=0}^D$ for the \mathbb{C} -vector space M such that

$$E_0 = |X|^{-1} J, \quad E_i E_j = \delta_{i,j} E_i \quad (0 \leq i, j \leq D), \quad \sum_{i=0}^D E_i = I. \tag{79}$$

We have

$$V = \sum_{i=0}^D E_i V \quad (\text{direct sum}). \tag{80}$$

In the above sum, the summands are the common eigenspaces for M . For $0 \leq i \leq D$ the matrix E_i is the projection onto the eigenspace E_iV . We call $\{E_i\}_{i=0}^D$ the *primitive idempotents* of Γ . We call E_0 the *trivial primitive idempotent*. The primitive idempotents are unique up to ordering of the nontrivial primitive idempotents.

Since M is closed under \circ , there exist $q_{ij}^h \in \mathbb{C}$ ($0 \leq h, i, j \leq D$) such that

$$E_i \circ E_j = |X|^{-1} \sum_{h=0}^D q_{ij}^h E_h \quad (0 \leq i, j \leq D). \tag{81}$$

The scalars q_{ij}^h are called the *Krein parameters* of Γ [3]. By [3, Theorem 3.8] each Krein parameter is real and nonnegative. Define

$$m_i = q_{ii}^0 \quad (0 \leq i \leq D). \tag{82}$$

We have $m_0 = 1$. By [3, Proposition 3.7], $m_i = \text{rank}(E_i)$ ($0 \leq i \leq D$). So by (80),

$$\sum_{i=0}^D m_i = |X|.$$

Since each of $\{A_i\}_{i=0}^D$ and $\{E_i\}_{i=0}^D$ is a basis for M , there exist matrices $P, Q \in \text{Mat}_{D+1}(\mathbb{C})$ such that

$$A_j = \sum_{i=0}^D P_{ij} E_i, \quad E_j = |X|^{-1} \sum_{i=0}^D Q_{ij} A_i \quad (0 \leq j \leq D). \tag{83}$$

Using (83) we find $PQ = |X|I$. Using $A_0 = I$ and $E_0 = |X|^{-1}J$ we obtain

$$P_{i0} = 1, \quad Q_{i0} = 1 \quad (0 \leq i \leq D).$$

By [4, Lemma 2.2.1(ii)],

$$k_j = P_{0j}, \quad m_j = Q_{0j} \quad (0 \leq j \leq D). \tag{84}$$

Define

$$\theta_i = P_{i1}, \quad \theta_i^* = Q_{i1} \quad (0 \leq i \leq D). \tag{85}$$

By the equation on the left in (83) with $j = 1$,

$$A_1 = \sum_{i=0}^D \theta_i E_i.$$

So θ_i is the eigenvalue of A_1 corresponding to E_i ($0 \leq i \leq D$). We call $\{\theta_i\}_{i=0}^D$ (resp. $\{\theta_i^*\}_{i=0}^D$) the *eigenvalue sequence* (resp. *dual eigenvalue sequence*) of Γ with respect to the ordering $\{E_i\}_{i=0}^D$.

Definition 9.1. Define polynomials $\{v_i\}_{i=0}^D$ in $\mathbb{C}[\lambda]$ such that $v_0 = 1$ and

$$\lambda v_i = c_{i+1}v_{i+1} + a_i v_i + b_{i-1}v_{i-1} \quad (1 \leq i \leq D - 1).$$

Lemma 9.2. [25, Lemma 3.8] *The following (i)–(iv) hold.*

- (i) *For $0 \leq i \leq D$ the polynomial v_i has degree i and leading coefficient $(c_1 c_2 \cdots c_i)^{-1}$.*
- (ii) *For $0 \leq i \leq D$, $A_i = v_i(A_1)$.*
- (iii) *The algebra M is generated by A_1 .*
- (iv) *The scalars $\{\theta_i\}_{i=0}^D$ are mutually distinct.*

Lemma 9.3. *For the polynomials $\{v_i\}_{i=0}^D$ from Definition 9.1,*

$$P_{ij} = v_j(\theta_i) \quad (0 \leq i, j \leq D).$$

Proof. In the equation on the left in (83), multiply each side on the right E_i to get $A_j E_i = P_{ij} E_i$. By Lemma 9.2(ii), $A_j E_i = v_j(A_1) E_i$. Using $A_1 E_i = \theta_i E_i$ one finds that $v_j(A_1) E_i = v_j(\theta_i) E_i$. By these comments we obtain $P_{ij} = v_j(\theta_i)$. \square

We recall the dual Bose-Mesner algebra. For the rest of this section, fix $x \in X$. For $0 \leq i \leq D$ let $E_i^* = E_i^*(x)$ denote the diagonal matrix in $\text{Mat}_X(\mathbb{C})$ that has (y, y) -entry 1 if $\partial(x, y) = i$ and 0 if $\partial(x, y) \neq i$ ($y \in X$). By construction,

$$E_i^* E_j^* = \delta_{i,j} E_i^* \quad (0 \leq i, j \leq D), \quad \sum_{i=0}^D E_i^* = I.$$

Consequently $\{E_i^*\}_{i=0}^D$ form a basis for a commutative subalgebra $M^* = M^*(x)$ of $\text{Mat}_X(\mathbb{C})$. We call M^* the *dual Bose-Mesner algebra of Γ with respect to x* . We call $\{E_i^*\}_{i=0}^D$ the *dual primitive idempotents of Γ with respect to x* . For $0 \leq i \leq D$ we have

$$E_i^* V = \text{Span}\{\hat{y} \mid y \in \Gamma_i(x)\}$$

and $k_i = \text{rank}(E_i^*)$. Moreover

$$V = \sum_{i=0}^D E_i^* V \quad (\text{direct sum}).$$

In the above sum, the summands are the common eigenspaces for M^* . For $0 \leq i \leq D$ the matrix E_i^* is the projection onto the eigenspace $E_i^* V$. For $0 \leq i \leq D$ let $A_i^* = A_i^*(x)$

denote the diagonal matrix in $\text{Mat}_X(\mathbb{C})$ whose (y, y) -entry is the (x, y) -entry of $|X|E_i$ ($y \in X$). By (79) and (81),

$$A_i^* A_j^* = \sum_{h=0}^D q_{ij}^h A_h^* \quad (0 \leq i, j \leq D), \tag{86}$$

$$A_0^* = I, \quad \sum_{i=0}^D A_i^* = |X|E_0^*.$$

By (83) we find that for $0 \leq j \leq D$,

$$E_j^* = |X|^{-1} \sum_{i=0}^D P_{ij} A_i^*, \quad A_j^* = \sum_{i=0}^D Q_{ij} E_i^*. \tag{87}$$

By (87) the matrices $\{A_i^*\}_{i=0}^D$ form a basis for the \mathbb{C} -vector space M^* . We call $\{A_i^*\}_{i=0}^D$ the *dual distance-matrices of Γ with respect to x* . We call A_1^* the *dual adjacency matrix* of Γ with respect to x . By the equation on the right in (87) with $j = 1$,

$$A_1^* = \sum_{i=0}^D \theta_i^* E_i^*.$$

So θ_i^* is the eigenvalue of A_1^* corresponding to E_i^* ($0 \leq i \leq D$).

The algebras M and M^* are related as follows. By [25, Lemma 3.2],

$$E_i^* A_h E_j^* = 0 \quad \text{if and only if} \quad p_{ij}^h = 0 \quad (0 \leq h, i, j \leq D),$$

$$E_i A_h^* E_j = 0 \quad \text{if and only if} \quad q_{ij}^h = 0 \quad (0 \leq h, i, j \leq D).$$

Lemma 9.4. For $0 \leq i \leq D$,

$$A_i E_0^* E_0 = E_i^* E_0, \quad E_i E_0^* E_0 = |X|^{-1} A_i^* E_0, \tag{88}$$

$$A_i^* E_0 E_0^* = E_i E_0^*, \quad E_i^* E_0 E_0^* = |X|^{-1} A_i E_0^*. \tag{89}$$

Proof. For $y, z \in X$, compare the (y, z) -entry of each side using the definition of A_i^* and E_i^* . \square

Lemma 9.5. We have

$$|X| E_0 E_0^* E_0 = E_0, \quad |X| E_0^* E_0 E_0^* = E_0^*.$$

Proof. For the equation on the right in (88), (89) set $i = 0$ and use $A_0 = I, A_0^* = I$. \square

Let $T = T(x)$ denote the subalgebra of $\text{Mat}_X(\mathbb{C})$ generated by M and M^* . We call T the *Terwilliger algebra of Γ with respect to x* [25].

By a T -module we mean a subspace of V that is invariant under the action of T . The T -module V is said to be *standard*. A T -module U is said to be *irreducible* whenever $U \neq 0$ and there is no T -submodule of U other than $0, U$.

Lemma 9.6. [25, Lemma 3.4] *The following hold.*

- (i) *The standard T -module V is a direct sum of irreducible T -modules.*
- (ii) *Each irreducible T -module U is a direct sum of the nonzero subspaces among E_0U, E_1U, \dots, E_DU , and a direct sum of the nonzero subspaces among $E_0^*U, E_1^*U, \dots, E_D^*U$.*

Let U denote an irreducible T -module. We say that U is *thin* (resp. *dual thin*) whenever $\dim E_i^*U \leq 1$ (resp. $\dim E_iU \leq 1$) for $0 \leq i \leq D$. Define the *diameter* $d = d(U)$ and *dual diameter* $d^* = d^*(U)$ by

$$d = |\{i \mid 0 \leq i \leq D, E_i^*U \neq 0\}| - 1, \quad d^* = |\{i \mid 0 \leq i \leq D, E_iU \neq 0\}| - 1.$$

If U is thin (resp. dual thin) then the dimension of U is $d + 1$ (resp. $d^* + 1$). If U is both thin and dual thin, then $d = d^*$.

We recall the primary T -module [25, Lemma 3.6]. Define

$$\mathbf{1} = \sum_{y \in X} \hat{y}.$$

Observe that the entries of $\mathbf{1}$ are all 1.

Lemma 9.7. [25, Lemma 3.6] *For $0 \leq i \leq D$,*

$$A_i \hat{x} = E_i^* \mathbf{1}, \quad A_i^* \mathbf{1} = |X| E_i \hat{x}.$$

Moreover, $M\hat{x} = M^ \mathbf{1}$ is a thin, dual-thin, irreducible T -module of diameter D .*

Definition 9.8. The T -module in Lemma 9.7 is called *primary*. Let U denote the primary T -module.

Lemma 9.9. *The following hold:*

- (i) *the map $M \rightarrow U, B \mapsto B\hat{x}$ is a \mathbb{C} -linear bijection;*
- (ii) *the map $M^* \rightarrow U, B \mapsto B\mathbf{1}$ is a \mathbb{C} -linear bijection.*

Proof. (i) The map is \mathbb{C} -linear by construction, and surjective by Lemma 9.7. It is a bijection since M, U have the same dimension.

(ii) The map is \mathbb{C} -linear by construction, and surjective by Lemma 9.7. It is a bijection since M^*, U have the same dimension. \square

Corollary 9.10. *The following hold.*

- (i) *Let $B \in M$ such that $BU = 0$. Then $B = 0$.*
- (ii) *Let $B \in M^*$ such that $BU = 0$. Then $B = 0$.*

Proof. (i) We have $B\hat{x} = 0$ since $\hat{x} \in U$, so $B = 0$ in view of Lemma 9.9(i).

(ii) We have $B\mathbf{1} = 0$ since $\mathbf{1} \in U$, so $B = 0$ in view of Lemma 9.9(ii). \square

Next we discuss general irreducible T -modules. For an irreducible T -module U , define the *endpoint* $r = r(U)$ by

$$r = \min\{i \mid 0 \leq i \leq D, E_i^*U \neq 0\}.$$

Lemma 9.11. [25, Lemma 3.9] *Let U denote an irreducible T -module with endpoint r and diameter d . Then the following hold.*

- (i) *We have $E_i^*U \neq 0$ if and only if $r \leq i \leq r + d$ ($0 \leq i \leq D$).*
- (ii) *If U is thin then U is dual thin.*

Next we recall the Q -polynomial property. The ordering $\{E_i\}_{i=0}^D$ is said to be Q -polynomial whenever q_{ij}^h is zero (resp. nonzero) if one of h, i, j is greater than (resp. equal to) the sum of the other two ($0 \leq h, i, j \leq D$). The graph Γ is said to be Q -polynomial with respect to $\{E_i\}_{i=0}^D$ whenever the ordering $\{E_i\}_{i=0}^D$ is Q -polynomial. For the rest of this section, assume that Γ is Q -polynomial with respect to $\{E_i\}_{i=0}^D$. Define

$$c_i^* = q_{1,i-1}^i \quad (1 \leq i \leq D), \tag{90}$$

$$a_i^* = q_{1i}^i \quad (0 \leq i \leq D), \tag{91}$$

$$b_i^* = q_{1,i+1}^i \quad (0 \leq i \leq D - 1). \tag{92}$$

For notational convenience define $c_0^* = 0$ and $b_D^* = 0$. We have $c_i^* b_{i-1}^* \neq 0$ for $1 \leq i \leq D$. Abbreviate $m = m_1$. By [3, Proposition 3.7],

$$a_0^* = 0, \quad b_0^* = m, \quad c_1^* = 1, \quad c_i^* + a_i^* + b_i^* = m \quad (0 \leq i \leq D). \tag{93}$$

By [4, Lemma 2.3.1(iv)],

$$m_i = \frac{b_0^* b_1^* \cdots b_{i-1}^*}{c_1^* c_2^* \cdots c_i^*} \quad (0 \leq i \leq D).$$

Definition 9.12. Define polynomials $\{v_i^*\}_{i=0}^D$ in $\mathbb{C}[\lambda]$ such that $v_0^* = 1$ and

$$\lambda v_i^* = c_{i+1}^* v_{i+1}^* + a_i^* v_i^* + b_{i-1}^* v_{i-1}^* \quad (1 \leq i \leq D - 1).$$

Lemma 9.13. [25, Lemma 3.11] *The following (i)–(iv) hold.*

- (i) *For $0 \leq i \leq D$ the polynomial v_i^* has degree i and leading coefficient $(c_1^* c_2^* \cdots c_i^*)^{-1}$.*
- (ii) *For $0 \leq i \leq D$, $A_i^* = v_i^*(A_1^*)$.*
- (iii) *The algebra M^* is generated by A_1^* .*
- (iv) *The scalars $\{\theta_i^*\}_{i=0}^D$ are mutually distinct.*

Lemma 9.14. *For the polynomials $\{v_i^*\}_{i=0}^D$ from Definition 9.12,*

$$Q_{ij} = v_j^*(\theta_i^*) \quad (0 \leq i, j \leq D).$$

Proof. Similar to the proof of Lemma 9.3. \square

Let U denote an irreducible T -module. Define the *dual endpoint* $s = s(U)$ by

$$s = \min\{i \mid 0 \leq i \leq D, E_i U \neq 0\}.$$

Lemma 9.15. [25, Lemma 3.9] *Let U denote an irreducible T -module with dual endpoint s and dual diameter d^* . Then the following hold.*

- (i) *We have $E_i U \neq 0$ if and only if $s \leq i \leq s + d^*$ ($0 \leq i \leq D$).*
- (ii) *If U is dual thin then U is thin.*

Lemma 9.16. [25, Lemmas 3.9, 3.12] *Let U denote a thin irreducible T -module with endpoint r , dual endpoint s , and diameter d . Then the sequence*

$$(A_1; \{E_{s+i}\}_{i=0}^d; A_1^*; \{E_{r+i}^*\}_{i=0}^d) \tag{94}$$

acts on U as a Leonard system, with eigenvalue sequence $\{\theta_{s+i}\}_{i=0}^d$ and dual eigenvalue sequence $\{\theta_{r+i}^\}_{i=0}^d$.*

Consider the Leonard system (94) on U . Any object ω associated with this Leonard system will be denoted by $\omega(U)$.

Recall the primary T -module U . By Lemma 9.16 the sequence

$$(A_1; \{E_i\}_{i=0}^D; A_1^*; \{E_i^*\}_{i=0}^D) \tag{95}$$

acts on U as a Leonard system, with eigenvalue sequence $\{\theta_i\}_{i=0}^D$ and dual eigenvalue sequence $\{\theta_i^*\}_{i=0}^D$.

Lemma 9.17. *We have*

$$\nu(\mathbf{U}) = |X|.$$

Proof. By Lemma 3.11, $\nu(\mathbf{U})E_0E_0^*E_0 = E_0$ holds on \mathbf{U} . Now by Lemma 9.5 and Corollary 9.10 we obtain the result. \square

Lemma 9.18. *For the Leonard system in (95) on \mathbf{U} , consider the elements $\{A_i\}_{i=0}^D$ and $\{A_i^*\}_{i=0}^D$ from Definition 4.4. Then for $0 \leq i \leq D$, A_i acts on \mathbf{U} as A_i , and A_i^* acts on \mathbf{U} as A_i^* .*

Proof. By Lemma 4.6 the following holds on \mathbf{U} :

$$E_i^*E_0E_0^* = \nu(\mathbf{U})^{-1}A_iE_0^*.$$

Comparing this with the equation on the right in (89), and using Lemma 9.17, we find that $A_i = A_i$ on $E_0^*\mathbf{U}$. The subspace $E_0^*\mathbf{U}$ is spanned by \hat{x} , so $A_i = A_i$ on \mathbf{U} in view of Lemma 9.9(i). The result for A_i^* is similarly obtained. \square

Lemma 9.19. *For $0 \leq h, i, j \leq D$,*

$$p_{ij}^h(\mathbf{U}) = p_{ij}^h, \quad q_{ij}^h(\mathbf{U}) = q_{ij}^h.$$

Proof. Compare the equations (25), (26) with (78), (86) and use Lemma 9.18. \square

Lemma 9.20. *We have*

$$c_i(\mathbf{U}) = c_i \quad c_i^*(\mathbf{U}) = c_i^* \quad (1 \leq i \leq D), \tag{96}$$

$$a_i(\mathbf{U}) = a_i \quad a_i^*(\mathbf{U}) = a_i^* \quad (0 \leq i \leq D), \tag{97}$$

$$b_i(\mathbf{U}) = b_i \quad b_i^*(\mathbf{U}) = b_i^* \quad (0 \leq i \leq D - 1). \tag{98}$$

Proof. We first obtain the equations on the left in (96)–(98). Apply Lemma 4.18 to the Leonard system (95) on \mathbf{U} , and use Lemma 9.18 to find that $A_1 = c_1(\mathbf{U})A_1 + a_0(\mathbf{U})I$ holds on \mathbf{U} . By this and Corollary 9.10(i) we obtain $c_1(\mathbf{U}) = 1$ and $a_0(\mathbf{U}) = 0$. Now compare the equations in Lemma 4.23 with (73)–(75), and use Lemma 9.19. This gives the equations on the left in (96)–(98). Next we obtain the equations on the right in (96)–(98). Apply Lemma 4.18 to the dual of the Leonard system (95) on \mathbf{U} , and use Lemma 9.18 to find that $A_1^* = c_1^*(\mathbf{U})A_1^* + a_0^*(\mathbf{U})I$ holds on \mathbf{U} . By this and Corollary 9.10(ii) we obtain $c_1^*(\mathbf{U}) = 1$ and $a_0^*(\mathbf{U}) = 0$. Now apply Lemma 4.23 to the dual of the Leonard system (95) on \mathbf{U} to get

$$\begin{aligned}
 c_i^*(U) &= q_{1,i-1}^i(U) & (1 \leq i \leq D), \\
 a_i^*(U) &= q_{1,i}^i(U) & (0 \leq i \leq D), \\
 b_i^*(U) &= q_{1,i+1}^i(U) & (0 \leq i \leq D-1).
 \end{aligned}$$

Now compare these equations with (90)–(92), and use Lemma 9.19. This gives the equations on the right in (96)–(98). \square

Lemma 9.21. For $0 \leq i \leq D$,

$$k_i(U) = k_i, \quad k_i^*(U) = m_i. \tag{99}$$

Proof. To get the equation on the left in (99), apply Lemma 4.21 to the Leonard system (95) on U , and compare the result with (72) using Lemma 9.19. To get the equation on the right in (99), apply Lemma 4.21 to the dual of the Leonard system (95) on U , and compare the result with (82) using Lemma 9.19. \square

Definition 9.22. Let $0 \neq q \in \mathbb{C}$. The Q -polynomial ordering $\{E_i\}_{i=0}^D$ is said to have q -Racah type whenever there exists an affine transformation of the Leonard system (95) on U , that has q -Racah type.

10. Formally self-dual distance-regular graphs

In this section we discuss a type of distance-regular graph, said to be formally self-dual. Let Γ denote a distance-regular graph with diameter $D \geq 1$. Recall the distance-matrices $\{A_i\}_{i=0}^D$ and primitive idempotents $\{E_i\}_{i=0}^D$ in the Bose-Mesner algebra M of Γ . Recall the matrices $P, Q \in \text{Mat}_{D+1}(\mathbb{C})$ from (83). Fix $x \in X$ and abbreviate $T = T(x)$, $M^* = M^*(x)$, and $E_i^* = E_i^*(x)$, $A_i^* = A_i^*(x)$ for $0 \leq i \leq D$.

Definition 10.1. The graph Γ is said to be *formally self-dual with respect to* $\{E_i\}_{i=0}^D$ whenever $P = Q$.

The following fact is well-known; see [4, Section 2.3]. We give a short proof for completeness.

Lemma 10.2. Assume that Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$. Then $p_{ij}^h = q_{ij}^h$ for $0 \leq h, i, j \leq D$.

Proof. Consider the \mathbb{C} -linear map $\sigma : M \rightarrow M^*$ that sends $E_i \mapsto E_i^*$ for $0 \leq i \leq D$. Clearly σ is bijective. We have $E_i E_j = \delta_{i,j} E_i$ and $E_i^* E_j^* = \delta_{i,j} E_i^*$ for $0 \leq i, j \leq D$. Thus σ is an isomorphism of \mathbb{C} -algebras. By the equation on the left in (83) and the equation on the right in (87) with $P = Q$, we find that σ sends $A_j \mapsto A_j^*$ for $0 \leq j \leq D$. Now we get the result by (78) and (86). \square

By Lemma 10.2, if Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$, then Γ is Q -polynomial with respect to $\{E_i\}_{i=0}^D$.

Lemma 10.3. *Assume that Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$. Let the polynomials $\{v_i\}_{i=0}^D$ (resp. $\{v_i^*\}_{i=0}^D$) be from Definition 9.1 (resp. Definition 9.12). Then $v_i = v_i^*$ for $0 \leq i \leq D$.*

Proof. Recall the intersection numbers $\{c_i\}_{i=1}^D$, $\{a_i\}_{i=0}^D$, $\{b_i\}_{i=0}^{D-1}$ and the Krein parameters $\{c_i^*\}_{i=1}^D$, $\{a_i^*\}_{i=0}^D$, $\{b_i^*\}_{i=0}^{D-1}$ of Γ . By Lemma 10.2, $p_{ij}^h = q_{ij}^h$ for $0 \leq h, i, j \leq D$. By (73)–(75) and (90)–(92) we see that $c_i = c_i^*$ ($1 \leq i \leq D$), $a_i = a_i^*$ ($0 \leq i \leq D$), $b_i = b_i^*$ ($0 \leq i \leq D - 1$). Now $v_i = v_i^*$ ($0 \leq i \leq D$) by Definitions 9.1, 9.12. \square

The formally self-dual condition is characterized as follows. Recall the eigenvalue sequence $\{\theta_i\}_{i=0}^D$ and dual eigenvalue sequence $\{\theta_i^*\}_{i=0}^D$ of Γ with respect to $\{E_i\}_{i=0}^D$.

Proposition 10.4. *The following (i)–(iii) are equivalent:*

- (i) Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$;
- (ii) Γ is Q -polynomial with respect to $\{E_i\}_{i=0}^D$ and $\theta_i = \theta_i^*$ for $0 \leq i \leq D$;
- (iii) the sequence (95) acts on the primary T -module U as a self-dual Leonard system.

Proof. (i) \Rightarrow (ii) As we saw above Lemma 10.3, Γ is Q -polynomial with respect to $\{E_i\}_{i=0}^D$. By Definition 10.1 we have $P = Q$. By this and (85) we obtain $\theta_i = \theta_i^*$ ($0 \leq i \leq D$).

(ii) \Rightarrow (iii) Let Φ denote the Leonard system in (95) on U . Then Φ has eigenvalue sequence $\{\theta_i\}_{i=0}^D$ and dual eigenvalue sequence $\{\theta_i^*\}_{i=0}^D$. We assume $\theta_i = \theta_i^*$ for $0 \leq i \leq D$, so Φ is self-dual by Lemma 2.15.

(iii) \Rightarrow (i) We show that $P = Q$. Let σ denote the duality of the Leonard system in (95). Then, on U , σ swaps E_i and E_i^* for $0 \leq i \leq D$. By Lemmas 4.7 and 9.18, on U , σ swaps A_i and A_i^* for $0 \leq i \leq D$. Now comparing (83) and (87), we obtain $P = Q$. \square

11. From a spin model to spin Leonard pairs

Let Γ denote a distance-regular graph with vertex set X and diameter $D \geq 1$. In this section we recall what it means for a spin model W to be afforded by Γ . For such W and all $x \in X$, we construct a spin Leonard pair on each irreducible $T(x)$ -module.

Recall the Bose-Mesner algebra M of Γ . Let W denote a spin model in $\text{Mat}_X(\mathbb{C})$. Recall the algebra $N(W)$ from Definition 7.4 and Lemma 7.5.

Definition 11.1. We say that Γ affords W whenever $W \in M \subseteq N(W)$.

Lemma 11.2. *Assume that W is afforded by Γ . Then each of the spin models $-W$, W^- is afforded by Γ .*

Proof. By Definition 11.1, $W \in M \subseteq N(W)$. We have $-W \in M$, and $N(-W) = N(W)$ by Lemma 7.6. Thus $-W \in M \subseteq N(-W)$, and so $-W$ is afforded by Γ . We have $W^{-1} \in M$, since for any invertible $B \in \text{Mat}_X(\mathbb{C})$ the element B^{-1} is a polynomial in B . By this and Lemma 7.2(ii), $W^{-1} \in M$. By Lemma 7.6, $N(W^{-1}) = N(W)$. Thus $W^{-1} \in M \subseteq N(W^{-1})$, and so W^{-1} is afforded by Γ . \square

For the rest of this section, assume that W is afforded by Γ .

Definition 11.3. [16, Section 3.1] We define a map $\Psi : N(W) \rightarrow \text{Mat}_X(\mathbb{C})$ as follows. For $B \in N(W)$ and $b, c \in X$ the (b, c) -entry of $\Psi(B)$ is the eigenvalue of B for the eigenvector $\mathbf{u}_{b,c}$ from Definition 7.4. So

$$B\mathbf{u}_{b,c} = \Psi(B)(b, c)\mathbf{u}_{b,c}.$$

Note that the map Ψ is \mathbb{C} -linear.

Recall the distance-matrices $\{A_i\}_{i=0}^D$ of Γ .

Lemma 11.4. [9, Lemma 5.1] *There exists an ordering $\{E_i\}_{i=1}^D$ of the nontrivial primitive idempotents of Γ such that*

$$\Psi(A_i) = |X|E_i \quad (0 \leq i \leq D).$$

Moreover Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$.

For the rest of this section, fix the ordering of the primitive idempotents from Lemma 11.4. Since $\{E_i\}_{i=0}^D$ is a basis for M and W is an invertible element in M , there exist nonzero scalars $f, \{\tau_i\}_{i=0}^D$ in \mathbb{C} such that $\tau_0 = 1$ and

$$W = f \sum_{i=0}^D \tau_i E_i. \tag{100}$$

Lemma 11.5. *We have*

$$W^{-1} = f^{-1} \sum_{i=0}^D \tau_i^{-1} E_i, \quad W^{-1} = |X|f^{-1} \sum_{i=0}^D \tau_i^{-1} E_i.$$

Proof. The first equation follows from (100). The second equation follows from the first equation and Lemma 7.2(ii). \square

For the rest of this section fix $x \in X$. Abbreviate $T = T(x)$ and $A_i^* = A_i^*(x), E_i^* = E_i^*(x)$ for $0 \leq i \leq D$. Define

$$W^* = f \sum_{i=0}^D \tau_i E_i^*. \tag{101}$$

Note that W^* is invertible.

Lemma 11.6. [5, Theorem 5.3, Corollary 5.4] *We have*

$$A_1 W^* W = W^* W A_1^*, \tag{102}$$

$$W W^* W = W^* W W^*. \tag{103}$$

Lemma 11.7. [6, Lemma 2.7] *We have*

$$W = |X|^{1/2} f^{-1} \sum_{i=0}^D \tau_i^{-1} A_i, \quad W^{-1} = |X|^{-3/2} f \sum_{i=0}^D \tau_i A_i, \tag{104}$$

$$W^* = |X|^{1/2} f^{-1} \sum_{i=0}^D \tau_i^{-1} A_i^*, \quad (W^*)^{-1} = |X|^{-3/2} f \sum_{i=0}^D \tau_i A_i^*. \tag{105}$$

Lemma 11.8. *For $y \in X$,*

$$W^*(y, y) = \frac{|X|^{1/2}}{W(x, y)}.$$

Proof. Let $i = \partial(x, y)$. By (101) we obtain $W^*(y, y) = f\tau_i$. By the equation on the left in (104) we obtain $W(x, y) = |X|^{1/2} f^{-1} \tau_i^{-1}$. The result follows. \square

Lemma 11.9. *The scalar f satisfies*

$$f^{-2} = |X|^{-3/2} \sum_{i=0}^D k_i \tau_i. \tag{106}$$

Proof. Recall the primary T -module U . By Lemma 9.16 the sequence

$$(A_1; \{E_i\}_{i=0}^D; A_1^*; \{E_i^*\}_{i=0}^D)$$

acts on U as a Leonard system. By Theorem 5.25 and Lemmas 5.26, 9.17,

$$W = f\gamma \sum_{i=0}^D \tau_i^{-1} A_i, \quad \gamma = |X|^{-1} \sum_{i=0}^D k_i \tau_i.$$

Comparing this with the equation on the left in (104), we get the result. \square

We now describe the irreducible T -modules. The following result is essentially due to Curtin [6, Theorem 8.3].

Proposition 11.10. *Every irreducible T-module is thin, provided that W does not have Hadamard type.*

Proof. We refer to the equation on the left in (104). We claim that

$$\tau_{i-1}, \tau_i, \tau_{i+1} \text{ are mutually distinct for } 1 \leq i \leq D - 1. \tag{107}$$

By way of contradiction, assume that there exists an integer j ($1 \leq j \leq D - 1$) such that $\tau_{j-1}, \tau_j, \tau_{j+1}$ are not mutually distinct. Then either (i) $\tau_j \in \{\tau_{j-1}, \tau_{j+1}\}$; or (ii) $\tau_{j-1} = \tau_{j+1}$. For the moment assume (i). Then by [8, Corollary 4.6] we get $\tau_i \in \{\tau_0, -\tau_0\}$ for $0 \leq i \leq D$, contradicting Lemma 8.4. Next assume (ii). Then $\tau_1 \in \{\tau_0, -\tau_0\}$ by [8, Corollary 4.7]. By this and [8, Corollary 4.6] we get $\tau_i \in \{\tau_0, -\tau_0\}$ for $0 \leq i \leq D$, contradicting Lemma 8.4. The claim is proved. In [6, Theorem 3.8] Curtin showed that every irreducible T-module is thin under the condition that

$$\tau_i \notin \{\tau_0, -\tau_0\} \text{ for } 1 \leq i \leq D. \tag{108}$$

Analyzing the proof of this theorem, we find that this theorem is implied by Lemmas 3.4 and 3.7 in [6]. Analyzing the proof of these lemmas, we find that these lemmas still hold if we replace the condition (108) with (107). Thus the conclusion of [6, Theorem 3.8] holds under the assumption (107). The result follows. \square

Let U denote a thin irreducible T-module with endpoint r , dual endpoint s , and diameter d . By Lemmas 9.16 the sequence

$$(A_1; \{E_{s+i}\}_{i=0}^d; A_1^*; \{E_{r+i}\}_{i=0}^d) \tag{109}$$

acts on U as a Leonard system. By [9, Lemma 4.4, Theorems 4.1, 5.5] we have $r = s$, and the Leonard system (109) on U is self-dual.

Lemma 11.11. *The pair A_1, A_1^* acts on U as a spin Leonard pair, and W, W^* acts on U as a balanced Boltzmann pair for this Leonard pair.*

Proof. By Definition 5.2 and (102) we see that the pair A_1, A_1^* acts on U as a spin Leonard pair, and the pair W, W^* acts on U as a Boltzmann pair for this Leonard pair. This Boltzmann pair is balanced by Definition 5.11 and (103). \square

12. From a spin model to spin Leonard pairs; the q -Racah case

Let Γ denote a distance-regular graph with vertex set X , diameter $D \geq 3$, and valency at least 3. Recall the distance-matrices $\{A_i\}_{i=0}^D$ in the Bose-Mesner algebra M of Γ . Let W denote a spin model afforded by Γ . Fix an ordering $\{E_i\}_{i=0}^D$ of the primitive idempotents of Γ from Lemma 11.4. Note that Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$. Let

$\{\theta_i\}_{i=0}^D$ denote the eigenvalue sequence of Γ with respect to the ordering $\{E_i\}_{i=0}^D$. As we saw below Lemma 11.4, there exist nonzero scalars $f, \{\tau_i\}_{i=0}^D$ in \mathbb{C} such that $\tau_0 = 1$ and

$$W = f \sum_{i=0}^D \tau_i E_i.$$

Fix $x \in X$. Abbreviate $T = T(x)$ and $A_i^* = A_i^*(x), E_i^* = E_i^*(x)$ for $0 \leq i \leq D$. Define

$$W^* = f \sum_{i=0}^D \tau_i E_i^*.$$

Recall the primary T -module U . By Proposition 10.4 the sequence

$$(A_1; \{E_i\}_{i=0}^D; A_1^*; \{E_i^*\}_{i=0}^D)$$

acts on U as a self-dual Leonard system. This Leonard system has eigenvalue sequence $\{\theta_i\}_{i=0}^D$. Fix $0 \neq q \in \mathbb{C}$, and assume that the ordering $\{E_i\}_{i=0}^D$ has q -Racah type.

Lemma 12.1. *There exist scalars a, α, β in \mathbb{C} with a, α nonzero that satisfy the following (i), (ii).*

(i) *The sequence*

$$(A; \{E_i\}_{i=0}^D; A^*; \{E_i^*\}_{i=0}^D) \tag{110}$$

acts on U as a Leonard system of q -Racah type, where

$$A_1 = \alpha A + \beta I, \quad A_1^* = \alpha A^* + \beta I. \tag{111}$$

(ii) *The Leonard system (110) on U has Huang data (a, a, a, D) .*

Proof. By Definition 9.22 there exists $\alpha, \beta \in \mathbb{C}$ with $\alpha \neq 0$ such that the sequence (110) acts on U as a Leonard system of q -Racah type, where A, A^* are from (111). This Leonard system is self-dual. By Definition 6.1 there exists $0 \neq a \in \mathbb{C}$ such that

$$\theta_i = \alpha(aq^{2i-D} + a^{-1}q^{D-2i}) + \beta \quad (0 \leq i \leq D). \tag{112}$$

By (111) and (112) this Leonard system has eigenvalue sequence $\{aq^{2i-D} + a^{-1}q^{D-2i}\}_{i=0}^D$. By Lemmas 5.4, 11.11 the pair A, A^* acts on U as a spin Leonard pair. By these comments and Lemmas 6.10, 6.11(i) the Leonard system (110) has Huang data (a, a, c, D) with $c \in \{a, a^{-1}, -a, -a^{-1}\}$. By Note 6.4 we may assume $c \in \{a, -a\}$. Replacing A, A^* with $-A, -A^*$ and replacing α with $-\alpha$, we may assume by Lemma 6.6 that $c = a$. The result follows. \square

For the rest of this section, let the scalars a, α, β be as in Lemma 12.1.

Lemma 12.2. [5, Corollaries 6.5, 6.7] *We have*

$$q^{2i} \neq 1 \quad (1 \leq i \leq D), \tag{113}$$

$$a^2 q^{2i} \neq 1 \quad (1 - D \leq i \leq D - 1), \tag{114}$$

$$a^3 q^{2i-D-1} \neq 1 \quad (1 \leq i \leq D). \tag{115}$$

Proof. Apply Lemma 6.5 to the Leonard system (110) on U . \square

Lemma 12.3. *One of the following holds:*

$$\tau_i = (-1)^i a^{-i} q^{i(D-i)} \quad (0 \leq i \leq D), \tag{116}$$

$$\tau_i^{-1} = (-1)^i a^{-i} q^{i(D-i)} \quad (0 \leq i \leq D). \tag{117}$$

Proof. Apply Lemma 6.16(i) to the Leonard system (110) on U , and use Lemma 12.1. \square

Note 12.4. Replacing W by W^- if necessary, and using Lemmas 11.2, 11.5, we may assume that (116) holds. We adopt this assumption for the rest of the section.

We mention a lemma for later use. For the moment, let D, d, r denote arbitrary nonnegative integers such that $r + d \leq D$, and let a, q denote arbitrary nonzero scalars in \mathbb{C} .

Lemma 12.5. *With the above notation, define $\tilde{a} = aq^{2r+d-D}$. Then the following hold.*

(i) *Define*

$$\theta_i = aq^{2i-D} + a^{-1}q^{D-2i} \quad (0 \leq i \leq D).$$

Then

$$\theta_{r+i} = \tilde{a}q^{2i-d} + (\tilde{a})^{-1}q^{d-2i} \quad (0 \leq i \leq d).$$

(ii) *Define*

$$\tau_i = (-1)^i a^{-i} q^{i(D-i)} \quad (0 \leq i \leq D).$$

Then

$$\tau_{r+i} = \tau_r \tilde{\tau}_i \quad (0 \leq i \leq d), \tag{118}$$

where

$$\tilde{\tau}_i = (-1)^i (\tilde{a})^{-i} q^{i(d-i)} \quad (0 \leq i \leq d). \tag{119}$$

Proof. Routine verification. \square

We return our attention to Γ . Let U denote a thin irreducible \mathbb{T} -module, with endpoint r and diameter d . Note that U has dual endpoint r by the sentence below (109). Consider the sequence

$$\Phi = (A_1; \{E_{r+i}\}_{i=0}^d; A_1^*; \{E_{r+i}^*\}_{i=0}^d).$$

Lemma 12.6. *The sequence Φ acts on U as a self-dual Leonard system, with eigenvalue sequence $\{\theta_{r+i}\}_{i=0}^d$.*

Proof. Apply Lemma 9.16 to the Leonard system Φ on U . \square

Lemma 12.7. *On U we have*

$$W = f\tau_r \sum_{i=0}^d \tilde{\tau}_i E_{r+i}, \quad W^* = f\tau_r \sum_{i=0}^d \tilde{\tau}_i E_{r+i}^*, \tag{120}$$

where $\{\tilde{\tau}_i\}_{i=0}^d$ are from (119).

Proof. By construction, on U we have

$$W = f \sum_{i=0}^d \tau_{r+i} E_{r+i}, \quad W^* = f \sum_{i=0}^d \tau_{r+i} E_{r+i}^*.$$

By this and (118) we get the result. \square

Our next goal is to obtain the intersection numbers of the Leonard system Φ on U .

Lemma 12.8. *The sequence*

$$(A; \{E_{r+i}\}_{i=0}^d; A^*; \{E_{r+i}^*\}_{i=0}^d) \tag{121}$$

acts on U as a Leonard system of q -Racah type with Huang data

$$(aq^{2r+d-D}, aq^{2r+d-D}, aq^{2r+d-D}, d).$$

Proof. By Lemma 12.6 and (111) the sequence (121) acts on U as a self-dual Leonard system. For $0 \leq i \leq D$ define

$$\tilde{\theta}_i = aq^{2i-D} + a^{-1}q^{D-2i},$$

and observe by (112) that $\theta_i = \alpha\tilde{\theta}_i + \beta$. Now by (111) and Lemma 12.6 the Leonard system (121) on U has eigenvalue sequence and dual eigenvalue sequence $\{\tilde{\theta}_{r+i}\}_{i=0}^d$. By Lemma 12.5(i) the Leonard system (121) on U has q -Racah type with Huang data $(\tilde{a}, \tilde{a}, \tilde{c}, d)$ for some nonzero $\tilde{c} \in \mathbb{C}$, where \tilde{a} is from Lemma 12.5. By Lemmas 5.4, 11.11 the pair A, A^* acts on U as a spin Leonard pair, and W, W^* acts on U as a Boltzmann pair for this Leonard pair. By this and Lemmas 6.18, 12.7 we obtain $\tilde{c} \in \{\tilde{a}, (\tilde{a})^{-1}\}$. By this and Note 6.4 the sequence $(\tilde{a}, \tilde{a}, \tilde{a}, d)$ is a Huang data for Φ on U . The result follows. \square

Recall the intersection numbers $\{c_i(U)\}_{i=1}^d, \{b_i(U)\}_{i=0}^{d-1}$ for the Leonard system Φ on U .

Lemma 12.9. [5, Theorem 8.5] *For the Leonard system Φ on U the intersection numbers satisfy*

$$b_i(U) = \frac{\alpha(q^{i-d} - q^{d-i})(aq^{2r+i-D} - a^{-1}q^{D-2r-i})(a^3 - q^{3D-2d-6r-2i-1})}{aq^{D-d-2r}(aq^{2r+2i-D} - a^{-1}q^{D-2r-2i})(a + q^{D-2r-2i-1})}, \tag{122}$$

$$c_i(U) = \frac{\alpha a(q^i - q^{-i})(aq^{d+2r+i-D} - a^{-1}q^{D-d-2r-i})(a^{-1} - q^{2d-D+2r-2i+1})}{q^{d-D+2r}(aq^{2r+2i-D} - a^{-1}q^{D-2r-2i})(a + q^{D-2r-2i+1})} \tag{123}$$

for $1 \leq i \leq d - 1$ and

$$b_0(U) = \frac{\alpha(q^{-d} - q^d)(a^3 - q^{3D-2d-6r-1})}{aq^{D-d-2r}(a + q^{D-2r-1})}, \tag{124}$$

$$c_d(U) = \frac{\alpha(q^{-d} - q^d)(a - q^{D-2r-1})}{q^{d-1}(a + q^{D-2d-2r+1})}. \tag{125}$$

Proof. For the Leonard system (121) on U , compute the intersection numbers using Lemmas 6.13 and 12.8. Adjust these intersection numbers using (111) to get the result. \square

Recall the intersection numbers $\{c_i\}_{i=1}^D, \{b_i\}_{i=0}^{D-1}$ of Γ .

Lemma 12.10. [8, Theorem 1.1] *The intersection numbers of Γ satisfy*

$$b_0 = \frac{\alpha(q^{-D} - q^D)(a^3 - q^{D-1})}{a(a + q^{D-1})}, \tag{126}$$

$$b_i = \frac{\alpha(q^{i-D} - q^{D-i})(aq^{i-D} - a^{-1}q^{D-i})(a^3 - q^{D-2i-1})}{a(aq^{2i-D} - a^{-1}q^{D-2i})(a + q^{D-2i-1})} \tag{1 \leq i \leq D - 1}, \tag{127}$$

$$c_i = \frac{\alpha a(q^i - q^{-i})(aq^i - a^{-1}q^{-i})(a^{-1} - q^{D-2i+1})}{(aq^{2i-D} - a^{-1}q^{D-2i})(a + q^{D-2i+1})} \tag{1 \leq i \leq D - 1}, \tag{128}$$

$$c_D = \frac{\alpha(q^{-D} - q^D)(a - q^{D-1})}{q^{D-1}(a + q^{1-D})}. \tag{129}$$

Proof. Apply Lemma 12.9 with $U = U$ and use Lemma 9.20. \square

Lemma 12.11. [5, Theorem 6.4, 6.8] *We have*

$$\alpha = \frac{(aq^{2-D} - a^{-1}q^{D-2})(a + q^{D-1})}{q^{D-1}(q^{-1} - q)(aq - a^{-1}q^{-1})(a - q^{1-D})}, \tag{130}$$

$$\beta = \frac{q(a + a^{-1})(a + q^{-D-1})(aq^{2-D} - a^{-1}q^{D-2})}{(q - q^{-1})(a - q^{1-D})(aq - a^{-1}q^{-1})}. \tag{131}$$

Proof. By (76), $c_1 = 1$. By this and (128) we obtain (130). By (76), (84), (85) we have $b_0 = \theta_0$. By this and (112), (126), (130) we obtain (131). \square

Lemma 12.12. *We have*

$$|X| = \frac{a^{3D}(a^{-2}; q^2)_D^2}{q^{D(D-1)}(aq^{1-D}; q^2)_D^2}.$$

Proof. By Lemmas 6.14 and 9.17. \square

Recall the scalars $\{k_i\}_{i=0}^D$ from (72).

Lemma 12.13. *We have $k_0 = 1$, and*

$$k_i = \frac{q^{2iD}(1 - a^2q^{4i-2D})(a^2q^{2-2D}; q^2)_i (q^{-2D}; q^2)_i (a^3q^{1-D}; q^2)_i}{a^{2i}(1 - a^2q^{2i-2D})(a^2q^2; q^2)_i (q^2; q^2)_i (a^{-1}q^{1-D}; q^2)_i}$$

for $1 \leq i \leq D - 1$, and

$$k_D = \frac{q^{2D^2}(a^2q^{2-2D}; q^2)_D (q^{-2D}; q^2)_D (a^3q^{1-D}; q^2)_D}{a^{2D}(a^2; q^2)_D (q^2; q^2)_D (a^{-1}q^{1-D}; q^2)_D}.$$

Proof. By Lemmas 6.15 and 9.21. \square

Lemma 12.14. *We have*

$$\sum_{i=0}^D k_i \tau_i = \frac{(a^{-2}; q^2)_D}{(aq^{1-D}; q^2)_D}.$$

Proof. By Lemmas 6.20 and 9.21. \square

We have been discussing the q -Racah case. For a discussion of the non q -Racah cases, see [8].

13. From spin Leonard pairs to a spin model

In Sections 11, 12 we considered a distance-regular graph Γ that affords a spin model W . We used W to construct a family of spin Leonard pairs. We now reverse the logical direction, and obtain a spin model afforded by Γ from some spin Leonard pairs. In this section we work with general Γ , and in the next section we consider the q -Racah case. For this section our main results are Theorems 13.7 and 13.9.

Let Γ denote a distance-regular graph with vertex set X and diameter $D \geq 1$. Recall the distance-matrices $\{A_i\}_{i=0}^D$ and the primitive idempotents $\{E_i\}_{i=0}^D$ in the Bose-Mesner algebra M of Γ . Assume that Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$. For $x \in X$, consider the dual Bose-Mesner algebra $M^*(x)$ and the Terwilliger algebra $T(x)$. For $0 \leq i \leq D$ we have the elements $E_i^*(x)$ and $A_i^*(x)$ from the paragraph below Lemma 9.3. Let $f, \{\tau_i\}_{i=0}^D$ denote nonzero scalars in \mathbb{C} such that $\tau_0 = 1$. Define

$$W = f \sum_{i=0}^D \tau_i E_i.$$

For $x \in X$ define

$$W^*(x) = f \sum_{i=0}^D \tau_i E_i^*(x). \tag{132}$$

Note that W and $W^*(x)$ are invertible. For the rest of this section, we make the following assumption.

Assumption 13.1. Assume that for all $x \in X$ and all irreducible $T(x)$ -modules U ,

- (i) U is thin;
- (ii) U has the same endpoint and dual-endpoint;
- (iii) the pair $A_1, A_1^*(x)$ acts on U as a spin Leonard pair, and $W, W^*(x)$ acts on U as a balanced Boltzmann pair for this spin Leonard pair;
- (iv) f satisfies (106).

Under Assumption 13.1 we show that W is a spin model afforded by Γ . From now until the beginning of Theorem 13.9, fix $x \in X$ and abbreviate $W^* = W^*(x)$, $T = T(x)$ and $E_i^* = E_i^*(x)$, $A_i^* = A_i^*(x)$ for $0 \leq i \leq D$.

Lemma 13.2. *We have*

$$W = |X|^{1/2} f^{-1} \sum_{i=0}^D \tau_i^{-1} A_i, \quad W^{-1} = |X|^{-3/2} f \sum_{i=0}^D \tau_i A_i. \tag{133}$$

Proof. We first show that (133) holds on the primary T -module U . By Proposition 10.4(i), (iii) the sequence

$$(A_1; \{E_i\}_{i=0}^D; A_1^*; \{E_i^*\}_{i=0}^D) \tag{134}$$

acts on U as a self-dual Leonard system. Let $\sigma : \text{End}(U) \rightarrow \text{End}(U)$ denote the duality of this Leonard system. So σ swaps the action of A_1, A_1^* on U , and swaps the action of E_i, E_i^* on U for $0 \leq i \leq D$. By Assumption 13.1(iii), A_1, A_1^* acts on U as a spin Leonard pair, and W, W^* acts on U as a balanced Boltzmann pair for this Leonard pair. In Lemmas 5.9, 5.10 we used this Boltzmann pair to obtain a duality for this Leonard pair, which must coincide with the above duality σ by the uniqueness of the duality. By this and (42), (43) we find that the results below (43) in Section 5 apply to the Leonard system (134) on U . Apply Theorem 5.25, Lemma 5.26 and (106) to this Leonard system, and use Lemmas 9.17, 9.18, 9.21 to show that (133) holds on U . The result follows in view of Corollary 9.10. \square

Lemma 13.3. *The matrix W is type II.*

Proof. We invoke Lemma 7.2. By construction W is symmetric. Each entry of W is nonzero by the equation on the left in (133). By that equation and (69) we obtain

$$W^- = |X|^{-1/2} f \sum_{i=0}^D \tau_i A_i. \tag{135}$$

Comparing (135) with the equation on the right in (133), we obtain $W^- = |X|W^{-1}$. Therefore $WW^- = |X|I$. The result follows by Lemma 7.2. \square

Lemma 13.4. *We have*

$$WW^*W = W^*WW^*. \tag{136}$$

Proof. By Assumption 13.1(iii) the pair W, W^* acts on each irreducible T -module as a balanced Boltzmann pair. The result follows in view of Lemma 9.6(i). \square

Lemma 13.5. *For $y \in X$,*

$$W^*(y, y) = \frac{|X|^{1/2}}{W(x, y)}.$$

Proof. Let $i = \partial(x, y)$. By (132) we obtain $W^*(y, y) = f\tau_i$. By the equation on the left in (133) we obtain $W(x, y) = |X|^{1/2} f^{-1}\tau_i^{-1}$. The result follows. \square

Lemma 13.6. For $a, b \in X$,

$$\sum_{y \in X} \frac{W(a, y)W(b, y)}{W(x, y)} = |X|^{1/2} \frac{W(a, b)}{W(a, x)W(b, x)}.$$

Proof. For each side of (136), compute the (a, b) -entry and evaluate the result using Lemma 13.5. Here are some details:

$$\begin{aligned} (WW^*W)(a, b) &= \sum_{y \in X} W(a, y)W^*(y, y)W(y, b) \\ &= |X|^{1/2} \sum_{y \in X} \frac{W(a, y)W(b, y)}{W(x, y)}, \end{aligned}$$

and

$$\begin{aligned} (W^*WW^*)(a, b) &= W^*(a, a)W(a, b)W^*(b, b) \\ &= |X| \frac{W(a, b)}{W(a, x)W(b, x)}. \quad \square \end{aligned}$$

Theorem 13.7. The matrix W is a spin model.

Proof. The matrix W satisfies the requirements of Definition 7.7 by Lemmas 13.3 and 13.6. \square

Next we show that the spin model W is afforded by Γ .

Lemma 13.8. For $0 \leq i \leq D$,

$$WW^*E_i = E_i^*WW^*. \tag{137}$$

Proof. Let U denote an irreducible T -module with endpoint r and diameter d . By Assumption 13.1(i), (ii) we see that U is thin with dual endpoint r . By these comments and Lemma 9.16 the sequence

$$(A_1; \{E_{r+j}\}_{j=0}^d; A_1^*; \{E_{r+j}^*\}_{j=0}^d)$$

acts on U as a Leonard system. By Assumption 13.1(iii) the pair W, W^* acts U as a balanced Boltzmann pair for the Leonard pair A_1, A_1^* on U . By these comments and Lemma 5.15,

$$WW^*E_{r+j} = E_{r+j}^*WW^*$$

holds on U for $0 \leq j \leq d$. Thus (137) holds on each irreducible T -module. Now (137) holds by Lemma 9.6(i). \square

Theorem 13.9. *The spin model W is afforded by Γ .*

Proof. By Definition 11.1 we must show that $W \in M \subseteq N(W)$. By construction $W \in M$. To obtain $M \subseteq N(W)$, we show that for all $b, c \in X$ the vector $\mathbf{u}_{b,c}$ is contained in E_iV , where $i = \partial(b, c)$. Using Definition 7.4 and Lemma 13.5,

$$\mathbf{u}_{b,c} = (W^*(b))^{-1}W^*(c)\mathbf{1} = (WW^*(b))^{-1}WW^*(c)\mathbf{1}.$$

Using Lemma 13.8 with $i = 0$,

$$WW^*(c)\mathbf{1} \in WW^*(c)E_0V = E_0^*(c)WW^*(c)V = E_0^*(c)V = \text{Span}\{\widehat{c}\} \subseteq E_i^*(b)V.$$

By the above comments and Lemma 13.8,

$$\mathbf{u}_{b,c} \in (WW^*(b))^{-1}E_i^*(b)V = E_i(WW^*(b))^{-1}V = E_iV.$$

We have shown that $\mathbf{u}_{b,c}$ is contained in E_iV and is therefore a common eigenvector for M . So $M \subseteq N(W)$ in view of Definition 7.4. The result follows. \square

14. From spin Leonard pairs to a spin model; the q -Racah case

In this section we start with a suitable distance-regular graph Γ of q -Racah type, and obtain a spin model afforded by Γ . The result is given in Theorem 14.7.

Let Γ denote a distance-regular graph with vertex set X , diameter $D \geq 1$, and intersection numbers $\{c_i\}_{i=1}^D, \{a_i\}_{i=0}^D, \{b_i\}_{i=0}^{D-1}$. Recall the distance-matrices $\{A_i\}_{i=0}^D$ and primitive idempotents $\{E_i\}_{i=0}^D$ in the Bose-Mesner algebra M of Γ . Throughout this section we make the following assumptions.

Assumption 14.1. Assume that Γ is formally self-dual with respect to $\{E_i\}_{i=0}^D$, with eigenvalue sequence $\{\theta_i\}_{i=0}^D$. Assume that there exist nonzero scalars $a, q \in \mathbb{C}$ that satisfy (112)–(115) and (126)–(129) with α, β from (130), (131). Assume that for all $x \in X$ and all irreducible $T(x)$ -modules U ,

- (i) U is thin;
- (ii) U has the same endpoint and dual-endpoint;
- (iii) the intersection numbers $\{c_i(U)\}_{i=1}^d, \{b_i(U)\}_{i=0}^{d-1}$ satisfy (122)–(125).

Under Assumption 14.1 we construct a spin model W that is afforded by Γ .

Definition 14.2. Define scalars $\{\tau_i\}_{i=0}^D$ in \mathbb{C} by

$$\tau_i = (-1)^i a^{-i} q^{i(D-i)} \quad (0 \leq i \leq D). \tag{138}$$

Define $f \in \mathbb{C}$ such that

$$f^2 = \frac{|X|^{3/2}(aq^{1-D}; q^2)_D}{(a^{-2}; q^2)_D}. \tag{139}$$

Note that $f, \{\tau_i\}_{i=0}^D$ are nonzero. Define

$$W = f \sum_{i=0}^D \tau_i E_i.$$

We are going to show that W is a spin model afforded by Γ . Fix $x \in X$ and abbreviate $T = T(x)$ and $E_i^* = E_i^*(x), A_i^* = A_i^*(x)$ for $0 \leq i \leq D$. Recall the scalars α, β from (130), (131). Define $A \in M$ and $A^* \in M^*$ such that

$$A_1 = \alpha A + \beta I, \quad A_1^* = \alpha A^* + \beta I. \tag{140}$$

Let U denote an irreducible T -module with endpoint r and diameter d . By Lemma 9.16 and since U is thin, each of the sequences

$$\begin{aligned} \Phi &= (A_1; \{E_{r+i}\}_{i=0}^d; A_1^*; \{E_{r+i}^*\}_{i=0}^d), \\ \tilde{\Phi} &= (A; \{E_{r+i}\}_{i=0}^d; A^*; \{E_{r+i}^*\}_{i=0}^d) \end{aligned}$$

acts on U as a Leonard system. By construction these Leonard systems are self-dual.

Lemma 14.3. *The Leonard system $\tilde{\Phi}$ on U has q -Racah type with Huang data*

$$(aq^{2r+d-D}, aq^{2r+d-D}, aq^{2r+d-D}, d). \tag{141}$$

Proof. By [13, Lemma 7.3] there exists a Leonard system Ψ of q -Racah type with Huang data (141). By Lemma 6.10 this Leonard system is self-dual. The eigenvalue sequence and intersection numbers of Ψ are obtained using Definition 6.3 and Lemma 6.13. The eigenvalue sequence and intersection numbers of Φ on U are given in Assumption 14.1. Adjusting these using (140) we obtain the eigenvalue sequence and intersection numbers of $\tilde{\Phi}$ on U . These match the eigenvalue sequence and intersection numbers of Ψ . By Lemmas 2.12, 2.14 the Leonard system Ψ is isomorphic to the Leonard system $\tilde{\Phi}$ on U . The result follows. \square

Define

$$W^* = f \sum_{i=0}^D \tau_i E_i^*.$$

Lemma 14.4. *On U we have*

$$W = f\tau_r \sum_{i=0}^d \tilde{\tau}_i E_{r+i}, \quad W^* = f\tau_r \sum_{i=0}^d \tilde{\tau}_i E_{r+i}^*, \tag{142}$$

where $\{\tilde{\tau}_i\}_{i=0}^d$ are from (119).

Proof. Similar to the proof of Lemma 12.7. \square

Lemma 14.5. *The pair A_1, A_1^* acts on U as a spin Leonard pair, and the pair W, W^* acts on U as a balanced Boltzmann pair of this Leonard pair.*

Proof. By Lemmas 6.11(ii), 14.3 the pair A, A^* acts on U as a spin Leonard pair. By (119) and Lemmas 6.16(ii), 14.4 the pair W, W^* acts on U as a balanced Boltzmann pair for the Leonard pair A, A^* on U . The result follows in view of Lemma 5.4. \square

Recall the integers $\{k_i\}_{i=0}^D$ from (72).

Lemma 14.6. *The scalar f from Definition 14.2 satisfies*

$$f^{-2} = |X|^{-3/2} \sum_{i=0}^D k_i \tau_i.$$

Proof. By Lemma 14.3 the sequence

$$(A; \{E_i\}_{i=0}^D; A^*; \{E_i^*\}_{i=0}^D)$$

acts on the primary T -module U as a Leonard system of q -Racah type with Huang data (a, a, a, D) . Applying Lemma 6.20 to this Leonard system we obtain

$$\sum_{i=0}^D k_i \tau_i = \frac{(a^{-2}; q^2)_D}{(aq^{1-D}; q^2)_D}. \tag{143}$$

The result follows from (139) and (143). \square

Theorem 14.7. *The matrix W in Definition 14.2 is a spin model afforded by Γ .*

Proof. For $x \in X$ let U denote an irreducible $T(x)$ -module. By Lemmas 5.4 and 14.5 the pair $A_1, A_1^*(x)$ acts on U as a spin Leonard pair, and $W, W^*(x)$ acts on U as a balanced Boltzmann pair for this Leonard pair. By Lemma 14.6 the scalar f satisfies (106). Thus Assumption 13.1 is satisfied. Now W is a spin model by Theorem 13.7. This spin model is afforded by Γ by Theorem 13.9. \square

Note 14.8. The spin model W in Theorem 14.7 does not have Hadamard type, since $\tau_1 \neq \pm\tau_0$ by (114) and (138).

15. Examples

The article [8, Section 9] includes a list of the known spin models that are contained in the Bose-Mesner algebra of a distance-regular graph Γ . In this section we display the listed spin models for which Γ has q -Racah type. Each displayed example satisfies the requirements of Assumption 14.1. For each displayed example we describe Γ , and give the parameters q, a that satisfy (112)–(115), (126)–(129). We now display the examples. The presentation of the display is explained in Example 0.

- (0) Name of the graph Γ [Reference to definition].
 - diameter D , intersection numbers $\{\mathbf{b}_0, \dots, \mathbf{b}_{D-1}; \mathbf{c}_1, \dots, \mathbf{c}_D\}$, number of vertices $|X|$.
 - parameters q, a .
- (1) Complete graph K_n ($n \geq 2$) [4, Appendix A1].
 - $D = 1, \{n - 1; 1\}, |X| = n$.
 - q is any nonzero complex number such that $q^4 \neq 1$, and a is any complex number such that $a \neq 1$ and $a + a^{-1} = n - 2$.
- (2) 4-cycle C_4 .
 - $D = 2, \{2, 1; 1, 2\}, |X| = 4$.
 - q is any nonzero complex number such that $q^4 \neq 1$, and a is any complex number such that $a^2 = -1$.
- (3) Higman-Sims graph [12].
 - $D = 2, \{22, 21; 1, 6\}, |X| = 100$.
 - q is any complex number such that $q^2 + q^{-2} = -3$, and $a = -q^{-3}$.
- (4) Hadamard graph H_m ($m \geq 2$) [14].
 - $D = 4, \{4m, 4m - 1, 2m, 1; 1, 2m, 4m - 1, 4m\}, |X| = 16m$.
 - q is any complex number such that $q^4 + q^{-4} = 4m - 2$, and a is any complex number such that $a^2 = -1$.
- (5) Double cover of the Higman-Sims graph [4, Appendix A.5].
 - $D = 5, \{22, 21, 16, 6, 1; 1, 6, 16, 21, 22\}, |X| = 200$.
 - q is any complex number such that $q^2 + q^{-2} = 3$, and a is any complex number such that $a^2 = -1$.
- (6) Odd cycle C_{2m+1} ($m \geq 2$).
 - $D = m, \{2, 1, \dots, 1; 1, 1, \dots, 1\}, |X| = 2m + 1$.
 - $q \in \mathbb{C}$ is any primitive $2m + 1$ root of unity, and $a = -q^m$.
- (7) Even cycle C_{2m} ($m \geq 3$).
 - $D = m, \{2, 1, \dots, 1; 1, \dots, 1, 2\}, |X| = 2m$.
 - $q \in \mathbb{C}$ is any primitive $4m$ root of unity, and a is any complex number such that $a^2 = -1$.

By Theorem 14.7, for each of the above Examples (1)–(7), the matrix W from Definition 14.2 is a spin model afforded by Γ .

References

- [1] R.C. Alperin, $\mathrm{PSL}_2(\mathbb{Z}) = \mathbb{Z}_2 \star \mathbb{Z}_3$, *Am. Math. Mon.* 100 (1993) 385–386.
- [2] E. Bannai, Et. Bannai, Spin models on finite cyclic groups, *J. Algebraic Comb.* 3 (1994) 243–259.
- [3] E. Bannai, T. Ito, *Algebraic Combinatorics I, Association Schemes*, Benjamin/Cummings Publishing Co. Inc., Menlo Park, CA, 1984.
- [4] A.E. Brower, A.M. Cohen, A. Neumaier, *Distance-Regular Graphs*, Springer-Verlag, Berlin, 1989.
- [5] J.S. Caughman IV, N. Wolff, The Terwilliger algebra of a distance-regular graph that supports a spin model, *J. Algebraic Comb.* 21 (2005) 289–310.
- [6] B. Curtin, Distance-regular graphs which support a spin model are thin, *Discrete Math.* 197/198 (1999) 205–216.
- [7] B. Curtin, Spin Leonard pairs, *Ramanujan J.* 13 (2007) 319–332.
- [8] B. Curtin, K. Nomura, Some formulas for spin models on distance-regular graphs, *J. Comb. Theory, Ser. B* 75 (1999) 206–236.
- [9] B. Curtin, K. Nomura, Spin models and strongly hyper-self-dual Bose-Mesner algebras, *J. Algebraic Comb.* 13 (2001) 173–186.
- [10] E.R. van Dam, J.H. Koolen, H. Tanaka, Distance-regular graphs, *Electron. J. Comb.* DS22 (Apr 2016) 15.
- [11] G. Gasper, M. Rahman, *Basic Hypergeometric Series*, *Encyclopedia of Mathematics and Its Applications*, vol. 96, Cambridge University Press, 2004.
- [12] D.G. Higman, C.C. Sims, A simple group of order 44,352,000, *Math. Z.* 105 (1968) 110–113.
- [13] H. Huang, The classification of Leonard triples of QRacah type, *Linear Algebra Appl.* 436 (2012) 1442–1472.
- [14] N. Ito, Hadamard graphs, I, *Graphs Comb.* 1 (1985) 57–64.
- [15] F. Jaeger, Strongly regular graphs and spin models for the Kauffman polynomial, *Geom. Dedic.* 44 (1992) 23–52.
- [16] F. Jaeger, M. Matsumoto, K. Nomura, Bose-Mesner algebras related with type II matrices and spin models, *J. Algebraic Comb.* 8 (1998) 39–72.
- [17] V.F.R. Jones, On knot invariants related to some statistical mechanical models, *Pac. J. Math.* 137 (1989) 311–336.
- [18] K. Kawagoe, A. Munemasa, Y. Watatani, Generalized spin models, *J. Knot Theory Ramif.* 3 (1994) 465–475.
- [19] A. Munemasa, private communication.
- [20] K. Nomura, Spin models constructed from Hadamard matrices, *J. Comb. Theory, Ser. A* 68 (1994) 251–261.
- [21] K. Nomura, An algebra associated with a spin model, *J. Algebraic Comb.* 6 (1997) 53–58.
- [22] K. Nomura, P. Terwilliger, Linear transformations that are tridiagonal with respect to both eigenbases of a Leonard pair, *Linear Algebra Appl.* 420 (2007) 198–207.
- [23] K. Nomura, P. Terwilliger, Affine transformations of a Leonard pair, *Electron. J. Linear Algebra* 16 (2007) 389–418.
- [24] J. Rotman, *Advanced Modern Algebra*, 2nd edition, AMS, Providence, RI, 2010.
- [25] P. Terwilliger, The subconstituent algebra of an association scheme I, *J. Algebraic Comb.* 1 (1992) 363–388.
- [26] P. Terwilliger, Two linear transformations each tridiagonal with respect to an eigenbasis of the other, *Linear Algebra Appl.* 330 (2001) 149–203.
- [27] P. Terwilliger, Leonard pairs and the q -Racah polynomials, *Linear Algebra Appl.* 387 (2004) 235–276.
- [28] P. Terwilliger, *An Algebraic Approach to the Askey Scheme of Orthogonal Polynomials, Orthogonal Polynomials and Special Functions*, *Lecture Notes in Math.*, vol. 1883, Springer, Berlin, 2006, pp. 255–330, arXiv:math/0408390.
- [29] P. Terwilliger, A. Žitnik, Distance-regular graphs of q -Racah type and the universal Askey-Wilson algebra, *J. Comb. Theory, Ser. A* 125 (2014) 98–112.