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**Journal of
Differential
Equations**

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New existence and symmetry results for least energy positive solutions of Schrödinger systems with mixed competition and cooperation terms

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Received 23 January 2015; revised 7 March 2016

Abstract

In this paper we focus on existence and symmetry properties of solutions to the cubic Schrödinger system

$$-\Delta u_i + \lambda_i u_i = \sum_{j=1}^d \beta_{ij} u_j^2 u_i \quad \text{in } \Omega \subset \mathbb{R}^N, \quad i = 1, \dots, d$$

where $d \geq 2$, $\lambda_i, \beta_{ii} > 0$, $\beta_{ij} = \beta_{ji} \in \mathbb{R}$ for $j \neq i$, $N = 2, 3$. The underlying domain Ω is either bounded or the whole space, and $u_i \in H_0^1(\Omega)$ or $u_i \in H_{\text{rad}}^1(\mathbb{R}^N)$ respectively. We establish new existence and symmetry results for least energy positive solutions in the case of mixed cooperation and competition coefficients, as well as in the purely cooperative case.

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MSC: primary 35J50; secondary 35B06, 35B09, 35J47

Keywords: Competitive and cooperative systems; Foliated Schwarz symmetry; Least energy positive solutions; Nehari manifold; Positive solutions; Schrödinger cubic systems

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<http://dx.doi.org/10.1016/j.jde.2016.03.015>

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1. Introduction

The existence and the qualitative description of *least energy solutions* to the nonlinear elliptic system

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 u^3 + \beta uv^2 \\ -\Delta v + \lambda_2 v = \mu_2 v^3 + \beta u^2 v \\ u, v \in H_0^1(\Omega), \end{cases} \quad \text{with } \Omega \subset \mathbb{R}^N \text{ or } \Omega = \mathbb{R}^N, \text{ and } N = 2, 3, \quad (1.1)$$

have attracted considerable attention in the last ten years, starting from the seminal paper [15] by T.-C. Lin and J. Wei. Collecting all the results contained in several contributions, it is possible to obtain an exhaustive picture of the problem, see the forthcoming Subsection 1.1. In striking contrast, a complete understanding in the case of an arbitrary $d \geq 3$ components system

$$\begin{cases} -\Delta u_i + \lambda_i u_i = \sum_{j=1}^d \beta_{ij} u_j^2 u_i & \text{in } \Omega \\ u_i \not\equiv 0 \\ u_i \in H_0^1(\Omega), \end{cases} \quad i = 1, \dots, d, \quad \beta_{ij} = \beta_{ji} \quad (1.2)$$

is not available, mainly due to the possible coexistence of *cooperation* and *competition*, that is, the existence of two pairs (i_1, j_1) and (i_2, j_2) such that $\beta_{i_1 j_1} > 0$ and $\beta_{i_2 j_2} < 0$. We recall that the sign of the coupling parameter β_{ij} determines the nature of the interaction between the components u_i and u_j : if $\beta_{ij} > 0$, then they cooperate, while if $\beta_{ij} < 0$, then they compete. Very recently, the systematic study of existence of least energy solutions in problems with simultaneous cooperation and competition has been started by the first author in [27] and by Y. Sato and Z.-Q. Wang in [22]. Nevertheless, there are still some gaps to fill in order to obtain a complete picture. In the present paper we give a contribution to fill some of these gaps, and, in the meantime, we analyse the symmetry properties of least energy solutions to (1.2), proving results which are new also in a purely cooperative context ($\beta_{ij} > 0$ for every $i \neq j$), and recover what is known in the purely competitive one ($d = 2$ and $\beta_{12} = \beta < 0$).

In order to motivate our research, in the following we review the results already available in the literature, but before it is worth to observe that thanks to the assumption $\beta_{ij} = \beta_{ji}$, system (1.2) has variational structure, as its solutions are critical points of the functional $J : H_0^1(\Omega, \mathbb{R}^d) \rightarrow \mathbb{R}$ defined by

$$J(\mathbf{u}) := \int_{\Omega} \frac{1}{2} \sum_{i=1}^d (|\nabla u_i|^2 + \lambda_i u_i^2) - \int_{\Omega} \frac{1}{4} \sum_{i,j=1}^d \beta_{ij} u_i^2 u_j^2,$$

where we used the vector notation $\mathbf{u} = (u_1, \dots, u_d)$. Observe that (1.2) admits *semi-trivial solutions*, i.e., solutions $\mathbf{u} \not\equiv 0$ with some zero components. However, we will be only interested in the existence of *positive solutions*: \mathbf{u} solving (1.2) such that $u_i > 0$ for every i . In particular, we will be interested in the existence of *least energy positive solutions*, that is solutions achieving the *least energy positive level*:

$$a := \inf \{ J(\mathbf{u}) : \mathbf{u} \text{ is a solution of (1.2) such that } u_i > 0 \text{ for all } i \},$$

or $a := +\infty$ if (1.2) has no positive solution. Observe that, due to the strong maximum principle, in the definition of the previous energy level, one can replace $u_i > 0$ by $u_i \geq 0$ and $u_i \not\equiv 0$. Observe moreover that, depending on the ranges of β_{ij} , this might not coincide with the least energy level (ground state):

$$\inf\{J(\mathbf{u}) : \mathbf{u} \neq 0 \text{ is a solution of (1.2)}\},$$

(see for instance [2] or [25]), which is an additional difficulty when one looks for least energy positive solutions.

The research of least energy solutions is strictly related to the minimization of the functional J on Nehari type sets such as

$$\mathcal{N}_1 := \left\{ \mathbf{u} \in H_0^1(\Omega, \mathbb{R}^d) : \mathbf{u} \neq \mathbf{0} \text{ and } \sum_{i=1}^d \partial_i J(\mathbf{u}) u_i = 0 \right\}$$

or

$$\mathcal{N}_2 := \left\{ \mathbf{u} \in H_0^1(\Omega, \mathbb{R}^d) : u_i \neq 0 \text{ and } \partial_i J(\mathbf{u}) u_i = 0 \text{ for every } i = 1, \dots, k \right\} \subset \mathcal{N}_1,$$

in the following sense. Let us set

$$c_1 := \inf_{\mathcal{N}_1} J \quad \text{and} \quad c_2 := \inf_{\mathcal{N}_2} J;$$

if a positive solution \mathbf{u} of (1.2) does exist, then $\mathbf{u} \in \mathcal{N}_2$, and hence $a \geq c_2 \geq c_1$. Clearly, this means that if c_2 is attained, then $c_2 = a$ and a positive minimizer for c_2 is a least energy positive solution. In a similar way, if c_1 is attained by a minimizer \mathbf{u} such that $u_i > 0$ for every i , then $c_1 = a$ and \mathbf{u} is a least energy positive solution. We emphasize that in general the inequalities $a \geq c_2 \geq c_1$ could be strict, see the discussion below.

1.1. Known results

Let us first describe the existing results in either the *purely cooperative case* $\beta_{ij} > 0$ for every $i \neq j$, or in the *purely competitive case* $\beta_{ij} < 0$ for every $i \neq j$. Some results deal with Ω bounded, while others with the case $\Omega = \mathbb{R}^N$. An important observation is that, in all the cited contributions dealing with the case \mathbb{R}^N , one is naturally led to work in $H_{\text{rad}}^1(\mathbb{R}^N)$, and with least energy positive *radial* solutions, that is positive radial solutions having minimal energy among all the positive radial solutions. In fact, in the purely cooperative case, each positive solution of (1.2) is radially decreasing, as comes out from [9]; hence a least energy positive level coincides with the radial one. On the other hand, in the purely competitive case, it is proved in [15, Theorem 1] and [2, Proposition 4.1-(i)] that c_2 defined above is not achieved, and that c_1 , which is achieved if $\beta_{12} > -\sqrt{\beta_{11}\beta_{22}}$, is attained by semi-trivial solutions; as a consequence, in order to adopt a Nehari-set approach one is induced to work in spaces of radial functions, taking advantage of the compactness of the Sobolev embedding $H_{\text{rad}}^1(\mathbb{R}^N) \hookrightarrow L^4(\mathbb{R}^N)$ to prove the existence of a minimizer for

$$\inf\{J(\mathbf{u}) : \mathbf{u} \in \mathcal{N}_2 \text{ and is radially symmetric}\}.$$

We notice that such a compactness result makes the research of least energy positive radial solutions in \mathbb{R}^N substantially equivalent to that of least energy positive solutions in bounded domains. For the sake of clarity, in what follows we always refers to a result as it was stated in its original contribution, but the reader has always to keep in mind that, whenever we cite a result for “least energy radial positive solution in \mathbb{R}^N ”, this also yields an existence of “least energy positive solution in a bounded domain”, and vice-versa.

Having this in mind, we focus at first on the 2 components system (1.1) in \mathbb{R}^N ($N = 2, 3$), with $\lambda_1, \lambda_2, \mu_1, \mu_2 \geq 0$. By collecting the main theorems in [2,14,15,20,25], one deduces that there exist $0 < \underline{\beta} \leq \bar{\beta}$ (depending on λ_i and μ_i) such that if either $-\infty < \beta < \underline{\beta}$, or $\beta > \bar{\beta}$, then (1.1) has a least energy radial positive solution (whose level coincides with the least energy positive level for $\beta > 0$). For the expression of the optimal values for $\underline{\beta}$ and $\bar{\beta}$, we refer to [11,21,25].

The results for the 2 components system have been partially extended for systems with an arbitrary number of components. Sufficient conditions for the existence of a least energy radial positive solutions of (1.2) in \mathbb{R}^N are the following.

- *Strong cooperation*: $\lambda_1 = \dots = \lambda_N = \lambda > 0$, $\beta_{ii} > 0$, and $\beta_{ij} = \beta$ for every $i \neq j$ is larger than a positive constant depending on β_{ii} and λ (see Corollary 2.3 and Theorem 2.1 in [19]; see also Theorem 1.6 and Remark 3 in [27]).
- *Weak cooperation*: $\lambda_i > 0$, $\beta_{ii} > 0$, $0 < \beta_{ij}$ smaller than a positive constant depending on λ_i and β_{ii} , and the matrix (β_{ij}) is positive definite (see Theorem 2 in [14]).
- *Competition*: if $\lambda_i > 0$, $\beta_{ii} > 0$, and $\beta_{ij} \leq 0$ for every $i \neq j$, then there exists a least energy radial positive solution (see Theorem 1.1. plus Remark 1.5 in [16] for the case Ω bounded; we refer also to Theorem 3.1 in [19], and to Corollary 1.4 plus Proposition 1.5 in [27]).

Other sufficient conditions in a purely cooperative setting have been given in [25, Section 4], [19, Theorem 2.1] and [10]. It is natural to assume that β_{ij} is either large, or small, with respect to β_{ii} and β_{jj} . Indeed if for instance $\beta_{ii} \leq \beta_{ij} \leq \beta_{jj}$ and $\lambda_i > \lambda_j$, then a positive solution of (1.2) does not exists, see Theorem 1-(ii) in [25] or Theorem 0.2 in [5].

As far as the possible occurrence of simultaneous cooperation and competition is concerned, in [22, Theorem 0.1] Y. Sato and Z.-Q. Wang considered a 3 components system in a bounded domain, showing that a least energy positive solution of (1.2) does exist if $\beta_{13}, \beta_{23} \leq 0$ are fixed, and $\beta_{12} \gg 1$ is very large (depending on β_{13} and β_{23}); we refer to this cases as *competition vs. arbitrarily large cooperation*. In [27, Theorems 1.6, 1.7, 1.9] the author considered an arbitrary d components system, proving the existence of least energy positive solutions whenever the d components are divided into m groups, with $m \leq d$, and

- the relation between components of the same group is purely cooperative, with coupling parameters greater than an explicit positive constant,
- the relation between components of different groups is competitive, and the competition is very strong.

When restricted to a 3 components system, this leads for instance to existence of a least energy solution if $\beta_{12} > \bar{\beta} > 0$, and $\beta_{13}, \beta_{23} \ll -1$ (depending on β_{12}). We refer to this cases as to *strong cooperation vs. arbitrarily large competition*. In the previous two results we would like to stress that the “large” parameters depend on *all* the other interaction terms. Ahead we will give a result which allows to fix *a priori* all the ranges for the parameters, which consists of a novelty when dealing with mixed cooperative and competitive interaction.

Remark 1.1. It is worth to point out the difference between *strong cooperation* and *arbitrarily large cooperation*: in the former case, we mean that some β_{ij} 's are greater than a positive constant which can be large but is fixed and determined as function of β_{ii} and λ_i , more or less explicitly; in the latter one we mean that some β_{ij} have to be thought as very large parameters which are tending to $+\infty$, depending in a non-explicit way on the other parameters. The same discussion holds for the distinction between *competition* and *arbitrarily large competition*.

For further existence results for system (1.2) with mixed cooperative and competitive couplings, which regard solutions not necessarily of least energy, we refer the reader to [15, Theorem 4], [18, Theorem 2.1], [27, Corollary 1.4], [23] and [12].

Concerning the symmetry properties of least energy positive solutions in bounded domains, the main results are contained in [28,30]. We postpone a precise description of them after having introduced some notation. In the next subsections, we describe the main results of this paper.

1.2. Main results: existence

We are concerned with the existence of least energy solutions of system (1.2):

$$\begin{cases} -\Delta u_i + \lambda_i u_i = \sum_{j=1}^d \beta_{ij} u_i u_j^2 & \text{in } \Omega \\ u_i = 0 & \text{on } \partial\Omega, \end{cases} \quad i = 1, \dots, d$$

where either

$$\begin{aligned} \Omega & \text{ is a bounded domain of } \mathbb{R}^N \text{ with } N = 2, 3, \\ \lambda_i & > -\mu_1(\Omega) \text{ and } \beta_{ii} > 0 \text{ for every } i = 1, \dots, d, \\ \beta_{ij} & = \beta_{ji} \in \mathbb{R} \text{ for every } i \neq j, \end{aligned} \quad (1.3)$$

and $\mu_1(\Omega)$ is the first eigenvalue of the Laplace operator with homogeneous Dirichlet boundary conditions on Ω , or

$$\begin{aligned} \Omega & = \mathbb{R}^N, \text{ with } N = 2, 3, \\ \lambda_i & > 0 \text{ and } \beta_{ii} > 0 \text{ for every } i = 1, \dots, d, \\ \beta_{ij} & = \beta_{ji} \in \mathbb{R} \text{ for every } i \neq j. \end{aligned} \quad (1.4)$$

In this last case, the boundary condition $u_i = 0$ on $\partial\Omega$ has to be replaced by $u_i \rightarrow 0$ as $|x| \rightarrow \infty$, and in the following instead of $H_0^1(\Omega)$ we have to write $H_{\text{rad}}^1(\mathbb{R}^N)$, the space of H^1 radially symmetric functions in \mathbb{R}^N .

We refer to the forthcoming [Remarks 2.6 and 4.1](#) for more details on the equivalence between the problem in bounded domains and the radial problem in \mathbb{R}^N .

In the following we recall some notations already introduced in [27].

- $B := (\beta_{ij})_{i,j=1,\dots,d}$, and we refer to it as to the *coupling matrix* of system (1.2).
- We endow the Sobolev space $H_0^1(\Omega)$ – or $H_{\text{rad}}^1(\mathbb{R}^N)$ – with scalar products and norms

$$\langle u, v \rangle_i := \int_{\Omega} (\nabla u \cdot \nabla v + \lambda_i uv) \quad \text{and} \quad \|u\|_i^2 := \langle u, u \rangle_i,$$

for every $i = 1, \dots, d$; in light of the assumptions on λ_i , these norms are equivalent to the standard one.

- For an arbitrary $1 \leq m \leq d$, we say that a vector $\mathbf{a} = (a_0, \dots, a_m) \in \mathbb{N}^{m+1}$ is a *m-decomposition of d* if

$$0 = a_0 < a_1 < \dots < a_{m-1} < a_m = d;$$

given a *m-decomposition* \mathbf{a} of d , we set, for $h = 1, \dots, m$,

$$\begin{aligned} I_h &:= \{i \in \{1, \dots, d\} : a_{h-1} < i \leq a_h\}, \\ \mathcal{K}_1 &:= \left\{ (i, j) \in I_h^2 \text{ for some } h = 1, \dots, m, \text{ with } i \neq j \right\}, \\ \mathcal{K}_2 &:= \{(i, j) \in I_h \times I_k \text{ with } h \neq k\}. \end{aligned} \quad (1.5)$$

This way, we have partitioned the set $\{1, \dots, d\}$ into m groups I_1, \dots, I_m , and have consequently splitted the components into m groups: $\{u_i : i \in I_h\}$.

We point out that the limit cases $m = 1$ or $m = d$ are also included in our terminology. This means that we will be able to recover (and sometimes improve) the known results for the purely cooperative case (taking $m = 1$) and for the purely competitive or weakly cooperative one (taking $m = d$).

Let us continue to introduce more notations.

Given a *m-decomposition* \mathbf{a} of d , we introduce the Nehari-type set induced by \mathbf{a} as

$$\mathcal{N} := \left\{ \mathbf{u} \in H_0^1(\Omega; \mathbb{R}^d) \mid \begin{array}{l} \sum_{i \in I_h} \|u_i\|_i \neq 0 \text{ and } \sum_{i \in I_h} \partial_i J(\mathbf{u}) u_i = 0 \\ \text{for every } h = 1, \dots, m \end{array} \right\}. \quad (1.6)$$

Only to fix our minds, we assume from now on that Ω is a bounded domain, and then we suppose that (1.3) is in force. Unless otherwise specified, the results can be extended for systems in \mathbb{R}^N replacing “least energy positive solution” with “least energy radial positive solution”.

For a given $d \geq 2$, let \mathbf{a} be a *m-decomposition* of d . We set

$$c := \inf_{\mathbf{u} \in \mathcal{N}} J(\mathbf{u}),$$

the infimum of J on the Nehari-type set \mathcal{N} .

Theorem 1.2. *There exists $K > 0$, depending only on β_{ii} , λ_i ($i = 1, \dots, d$), such that, whenever B satisfies*

$$\beta_{ij} \geq 0 \quad \forall (i, j) \in \mathcal{K}_1, \quad -\infty < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2,$$

then c is achieved by a nonnegative $\mathbf{u}_{\min} \in \mathcal{N}$. Furthermore, any minimizer is a nonnegative solution of (1.2).

This is an improvement of Theorem 1.3 in [27], where K depends also on β_{ij} with $(i, j) \in \mathcal{K}_1$, and the minimization is considered in an *open subset of \mathcal{N}* (notice that in such case a minimizer with all positive components needs not to be a least energy solution).

Theorem 1.2 gives existence of nonnegative solutions for systems of d equations where the populations are associated in groups, in such a way that inside each group there is cooperation, while between different groups we have either competition or weak cooperation.

Observe that if we can show that \mathbf{u}_{\min} has all positive components, we can immediately conclude that it is a least energy positive solution of the system. We are able to obtain this conclusion in several situations. First, applying Theorem 1.2 in the particular case $m = d$, which leads to the decomposition $\mathbf{a} = (0, 1, \dots, d)$, we have existence of least energy positive solutions of (1.2) also in regimes of *competition and/or weak cooperation*.

Corollary 1.3. *There exists $K > 0$, depending only on β_{ii}, λ_i ($i = 1, \dots, d$), such that if*

$$-\infty < \beta_{ij} < K \quad \text{for every } i \neq j,$$

then c is achieved by a least energy positive solution of (1.2).

Recall from the previous subsection that, up to now, this result was only known in the pure competitive or in the pure cooperative cases. Focusing on this last situation, note in particular that, unlike in [14, Theorem 2] (where the case $\beta_{ij} > 0$ for every $i \neq j$ is considered), Corollary 1.3 does not require the positive definiteness of the matrix B . Moreover, in the particular case of *pure cooperation*, we will show that

$$K := \frac{\min_{i=1, \dots, d} \{S_i^2\}}{2 \sum_{j=1}^d \frac{S_j^2}{\beta_{jj}}}, \quad \text{where } S_i := \inf_{\int_{\Omega} u^4 = 1} \|u\|_i^2, \tag{1.7}$$

which generalizes [30], where only the case $d = 2$ is considered.

If we consider a general m -decomposition with $m < d$, to find new least energy positive solutions we have to find conditions on the coupling parameters ensuring that the minimizer \mathbf{u}_{\min} in Theorem 1.2 has all non-trivial components. In what follows we shall use this argument to prove new existence results with respect to those in [22,27]. As in the quoted papers, the idea is to find conditions on the coupling parameters β_{ij} which ensure that

$$\inf_{\mathcal{N}} J < \inf \{J(\mathbf{u}) : \mathbf{u} \in \mathcal{N} \text{ and } u_i = 0 \text{ for some } i\}.$$

Theorem 1.4. *Let $d \geq 2$, let \mathbf{a} be a m -decomposition of d for some $1 \leq m \leq d$. Let K be the constant defined in Theorem 1.2. If*

Please cite this article in press as: N. Soave, H. Tavares, New existence and symmetry results for least energy positive solutions of Schrödinger systems with mixed competition and cooperation terms, J. Differential Equations (2016), <http://dx.doi.org/10.1016/j.jde.2016.03.015>

- (i) $\beta_{ij} = \beta_h > \max\{\beta_{ii} : i \in I_h\}$ for every $(i, j) \in I_h^2$ with $i \neq j$, $h = 1, \dots, m$;
- (ii) $\beta_{ij} = b < K$ for every $(i, j) \in \mathcal{K}_2$;
- (iii) $\lambda_i = \lambda_h > -\mu_1(\Omega)$ for every $i \in I_h$, $h = 1, \dots, m$;

then any minimizer of J in \mathcal{N} is positive, and hence system (1.2) has a least energy positive solution.

This statement together with Corollary 2.3 in [19] and our Corollary 1.3 provides the natural extension of what is known for the 2 components system for systems with an arbitrary number of equations. To be more precise, let us give an example in the case of 3 components system with the additional assumptions $\lambda_1 = \lambda_2 = \lambda$ and $\beta_{13} = \beta_{23}$. In this case, from our result and the quoted ones we deduce the existence of $0 < \underline{\beta} < \bar{\beta}$ such that (1.2) admits a least energy positive solution when one of the following conditions is verified:

$$\begin{aligned} \beta_{12} = \beta_{13} = \beta_{23} &> \bar{\beta}; & \lambda_3 &= \lambda \\ \beta_{12} &> \bar{\beta} \text{ and } -\infty < \beta_{13} = \beta_{23} < \underline{\beta}; \\ -\infty < \beta_{12}, \beta_{13}, \beta_{23} &< \underline{\beta}. \end{aligned} \tag{1.8}$$

The downsize of the previous theorem is that the restriction $\beta_{ij} = b$ for every $(i, j) \in \mathcal{K}_2$ is quite strong. For this reason we present also an alternative result, which permits to avoid this assumption but requires that $|\beta_{ij}|$ is not too large for $(i, j) \in \mathcal{K}_2$.

Theorem 1.5. *Let $d \geq 2$, let \mathbf{a} be a m -decomposition of d for some $1 \leq m \leq d$. Let K be the constant defined in Theorem 1.2, and fix $\alpha > 1$. If*

- (i) $\beta_{ij} = \beta_h$ for every $(i, j) \in I_h^2$ with $i \neq j$, $h = 1, \dots, m$, and

$$\beta_h > \frac{\alpha}{\alpha - 1} \max\{\beta_{ii} : i \in I_h\};$$

- (ii) for every $(i, j) \in \mathcal{K}_2$ there holds

$$|\beta_{ij}| \leq \frac{K}{\alpha d^2};$$

- (iii) $\lambda_i = \lambda_h > -\mu_1(\Omega)$ for every $i \in I_h$, $h = 1, \dots, m$;

then any minimizer of J in \mathcal{N} is positive, and hence system (1.2) has a least energy positive solution.

We observe that the previous two results seem to be the first dealing simultaneously with strong and weak cooperation.

With these results in hands, together with those in [19,22,27], we can give quite a complete picture for the problem of the existence of a least energy solution for system (1.2) when $d \geq 1$. We have already recalled in Subsection 1.1 that the existence of a least energy positive solution has been proved in regimes of *strong cooperation*, *competition*, *weak cooperation*, *arbitrarily large cooperation vs. competition*, *strong cooperation vs. arbitrarily large competition*. Thanks

to Corollary 1.3, Theorems 1.4 and 1.5, we have existence results in regimes of *competition or weak cooperation* and *weak cooperation or competition vs. strong competition*. Recalling the non-existence of positive solution when $\beta_{ii} \leq \beta_{ij} \leq \beta_{jj}$ proved in [5,25], at least from a qualitative point of view the existence of a least energy positive solutions is proved in all the admissible cases.

We would like to stress that in all the previous theorems, unlike in [22,27], the bounds only depend on λ_i and β_{ij} , for $i = 1, \dots, d$.

1.3. Main results: non-existence of non-radial minimizers in \mathbb{R}^N

As it was observed in Subsection 1.1, in the purely competitive case ($\beta_{ij} < 0$ for every $i \neq j$) the Nehari-set approach fails without any radial constraint, cf. Theorem 1 in [15]. A case of mixed cooperation and competition is treated in [14, Theorem 3], but only when exactly one state repels all the other, and the remaining ones have small attractive coefficients. The general case is left open in [14,15].

As we pointed out, our Theorem 1.2 in case $\Omega = \mathbb{R}^N$ works only in a radial setting, and allows much more combinations of cooperation with competition coefficients. It seems natural to ask whether the radial restriction, in general, is completely justified or not. We can prove that it is in several situations.

Theorem 1.6. Under (1.4), let $d \geq 2$, and let \mathbf{a} be a m -decomposition of d , with $1 \leq m \leq d$. If

- $\beta_{ij} \geq 0 \quad \forall (i, j) \in \mathcal{K}_1$,
- $\beta_{ij} \leq 0 \quad \forall (i, j) \in \mathcal{K}_2$, and there exist $h_1 \neq h_2$ such that $\beta_{ij} < 0$ for every $(i, j) \in I_{h_1} \times I_{h_2}$;

then

$$l := \inf_{\mathcal{M}} J \tag{1.9}$$

is not achieved, where

$$\mathcal{M} := \left\{ \mathbf{u} \in H^1(\mathbb{R}^N, \mathbb{R}^d) \mid \begin{array}{l} \sum_{i \in I_h} \|u_i\| \neq 0 \text{ and } \sum_{i \in I_h} \partial_i J(\mathbf{u})u_i = 0 \\ \text{for every } h = 1, \dots, m \end{array} \right\}.$$

For the reader's convenience, we recall that I_h , \mathcal{K}_1 and \mathcal{K}_2 have been defined in (1.5). Notice also that \mathcal{M} is the Nehari set associated to the m -decomposition \mathbf{a} without the radial constraint.

1.4. Main results: partial symmetry

We now pass to the statements regarding partial symmetry. First, we recall the following.

Definition 1.7. Let $\Omega \subset \mathbb{R}^N$ be radial with respect to 0. A function $u : \Omega \rightarrow \mathbb{R}$ is called *foliated symmetric Schwarz* with respect to the direction $p \in \mathbb{S}^{N-1}$ if u depends only on $(r, \theta) := (|x|, \arccos(x \cdot p/|x|))$, and is non-increasing in θ .

We write that the vector valued function (u_1, \dots, u_k) is foliated Schwarz symmetric with respect to p if each u_i is foliated Schwarz symmetric with respect to p .

We write that (u_1, \dots, u_h) and (u_{h+1}, \dots, u_k) are foliated Schwarz symmetric with respect to antipodal directions if there exists $p \in \mathbb{S}^{N-1}$ such that u_1, \dots, u_h are foliated Schwarz symmetric with respect to p , while u_{h+1}, \dots, u_k are foliated Schwarz symmetric with respect to $-p$.

We analyse the symmetry properties of \mathbf{u}_{\min} when Ω is a bounded radial domain (since in \mathbb{R}^N we deal with radial solutions, an analogue statement would be trivial in that setting). The following result is stated in the greatest possible generality.

Theorem 1.8. *Let $\Omega \subset \mathbb{R}^N$ be a bounded radially symmetric domain, with $N = 2, 3$. Let $d \geq 2$, let $\mathbf{a} = (a_0, \dots, a_m)$ be a m -decomposition of d for some $1 \leq m \leq d$, assume that (1.3) holds, and take K as in Theorem 1.2. Assume that B satisfies*

$$\beta_{ij} \geq 0 \quad \forall (i, j) \in \mathcal{K}_1, \quad -\infty < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2,$$

and for some $l := a_{\bar{h}}$ with $\bar{h} \in \{1, \dots, m\}$,

$$\begin{aligned} \beta_{ij} &> 0 \quad \text{for every } (i, j) \in \{1, \dots, l\}^2 \cup \{l+1, \dots, d\}^2 \\ \beta_{ij} &< 0 \quad \text{for every } (i, j) \in \{1, \dots, l\} \times \{l+1, \dots, d\}. \end{aligned} \quad (1.10)$$

Then any nonnegative \mathbf{u} achieving $c = \inf_{\mathcal{N}} J$ is such that (u_1, \dots, u_l) and (u_{l+1}, \dots, u_d) are foliated Schwartz symmetric with respect to antipodal directions.

The interpretation of the theorem is the following. A m -decomposition of d induces a separation of the d components into m different groups: $\{u_i : i \in I_h\}$, $h = 1, \dots, m$ (with I_h defined in (1.5)). We join the m groups into two macro groups, the first one collecting the first \bar{h} groups $\{u_i : i \in I_1 \cup \dots \cup I_{\bar{h}}\}$, the second one the remaining components. If we assume that the relation between components in the same macro-group is purely cooperative, while the relation between components in different macro-groups is purely competitive (see assumption (1.10)), then pairs of components of different macro-groups are foliated Schwartz symmetric with respect to antipodal directions.

Up to our knowledge, in the literature all the symmetry results so far were for systems in the pure cooperative or pure competitive cases. We point out also that Theorem 1.8 in the particular case $d = 2$ and $l = 1$ (competition between two components) permits to recover Theorem 1.3 in [28] in the present setting. Moreover, when $m = 1$ and $l = d$ (purely cooperative setting), we recover and significantly extend the results in [30] for $N = 2, 3$. Further remarks and comments are postponed to Section 3, but we would like to remark here that, in general, least energy positive solutions of (1.2) when Ω is radial are *not* radially symmetric, see Remark 5.4 in [28], the results of Section 3 in [30], and Corollary 0.5 in [22]. This break of symmetry can be caused either due to the presence of competition terms, or by the non-convexity of the underlying domain.

Clearly one can now combine Theorem 1.8 with the existence of least energy positive solutions of Subsection 1.2. As particular relevant cases, we would like to highlight that we have existence and symmetry of least energy positive solutions whenever B satisfies the following:

- For a system with $d \geq 2$ equations, if $0 < \beta_{ij} < K$ (where K is given by (1.7)), then (1.1) admits a least energy positive solution, whose components are foliated Schwartz symmetric with respect to the same point.

- In the special case $d = 3$, and assuming moreover that $\lambda_1 = \lambda_2 = \lambda$, we have symmetry for least energy solutions (u_1, u_2, u_3) in all the possible situations described in (1.8):
 - if either $\beta_{12} = \beta_{23} = \beta_{13} > \bar{\beta}$ and $\lambda_3 = \lambda$, or $\beta_{12} > \bar{\beta}$ and $0 < \beta_{13} = \beta_{23} < \bar{\beta}$, then (u_1, u_2, u_3) are foliated Schwartz symmetric with respect to the same point.
 - if $\beta_{12} > \bar{\beta}$ and $\beta_{13} = \beta_{23} < 0$, then (u_1, u_2) and u_3 are foliated Schwartz symmetric with respect to antipodal points.

1.5. Structure of the paper

In Section 2 we prove Theorem 1.2. Although in the introduction we presented the partial symmetry results at last, we point out that Theorem 1.8 regards not only least energy positive solutions, but constrained minimizers found under the assumptions of Theorem 1.2. For this reason, the proof of Theorem 1.8 is the object of Section 3. Section 4 is devoted to the proof of the new results on least energy positive solutions: Theorems 1.4 and 1.5. Finally, in the last section we prove Theorem 1.6.

2. Existence of nonnegative minimizers

This section is devoted to the proof of Theorem 1.2. In the literature, minimization on Nehari type sets as \mathcal{N} is usually addressed by firstly studying properties of minimizing sequences, and then showing that any limit of such a sequence can be projected on \mathcal{N} . The second part of this argument is extremely delicate from a technical point of view when the number of components is arbitrary (see e.g. Lemmas 2.6–2.8 in [27]). In what follows we use a different approach based on the Ekeland’s variational principle for the constrained functional $J|_{\mathcal{N}}$. As we shall see, this permits both to avoid several technicalities, and to obtain an explicit constant K depending only on β_{ii} and λ_i . Fix Ω be a bounded domain, $d \geq 2$, and take $\mathbf{a} = (a_0, a_1, \dots, a_m)$ be a m -decomposition of d for some $1 \leq m \leq d$. We always assume that (1.3) is in force; the case $\Omega = \mathbb{R}^N$ can be treated in the same way as the case Ω bounded, assuming (1.4) instead of (1.3), see Remark 2.6 for more details. We use both the notation introduced in Subsection 1.2, and the following:

- we let

$$S := \inf_{i=1, \dots, d} \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\|u\|_i^2}{|u|_{L^4}^2}.$$

By Sobolev embedding, $S > 0$.

- let $\mathbf{u} \in H_0^1(\Omega, \mathbb{R}^d)$. We set, for $h = 1, \dots, m$,

$$\mathbf{u}_h := (u_{a_{h-1}+1}, \dots, u_{a_h}) \in (H_0^1(\Omega))^{a_h - a_{h-1}}.$$

The space $(H_0^1(\Omega))^{a_h - a_{h-1}}$ is naturally endowed with scalar product and norm

$$\langle \mathbf{v}^1, \mathbf{v}^2 \rangle_h := \sum_{i \in I_h} \langle v_i^1, v_i^2 \rangle_i \quad \text{and} \quad \|\mathbf{v}\|_h^2 := \langle \mathbf{v}, \mathbf{v} \rangle_h.$$

- It will be useful to consider the following set, which contains weak H_0^1 limits of elements of \mathcal{N} :

$$\tilde{\mathcal{N}} = \left\{ \mathbf{u} \in H_0^1(\Omega; \mathbb{R}^d) \mid \begin{array}{l} \|\mathbf{u}_h\|_h \neq 0 \text{ and } \sum_{i \in I_h} \partial_i J(\mathbf{u}) u_i \leq 0 \\ \text{for every } h = 1, \dots, m \end{array} \right\}. \quad (2.1)$$

Finally, we introduce

$$\mathcal{E} := \left\{ \mathbf{u} \in H_0^1(\Omega; \mathbb{R}^d) : M_B(\mathbf{u}) \text{ is strictly diagonally dominant} \right\}, \quad (2.2)$$

where the $m \times m$ matrix $M_B(\mathbf{u})$ is defined by

$$M_B(\mathbf{u}) := \left(\sum_{(i,j) \in I_h \times I_k} \int_{\Omega} \beta_{ij} u_i^2 u_j^2 \right)_{h,k=1,\dots,m}.$$

Recall that a $m \times m$ matrix $A = (a_{ij})_{i,j}$ is strictly diagonally dominant if for every $i = 1, \dots, m$ there holds $|a_{ii}| > \sum_{j \neq i} |a_{ij}|$. Recall that, if a square matrix A is strictly diagonally dominant and has positive diagonal terms, then it is positive definite. Thus, in particular, for each $\mathbf{u} \in \mathcal{E}$, $M_B(\mathbf{u})$ is positive definite. This will be a key property of the matrices $M_B(\mathbf{u})$. Since we deal with $m \times m$ matrices, with m arbitrary, it does not seem easy to check directly that for some \mathbf{u} the matrix $M_B(\mathbf{u})$ is positive definite, and will be always proved by checking that it is strictly diagonally dominant (a condition which involves only the verifications of some inequalities).

Firstly, we recall some basic facts about the geometric structure of $\mathcal{N} \cap \mathcal{E}$, for which we refer to Proposition 1.1 and Remark 9 in [27]. The set \mathcal{E} is an open set in $H_0^1(\Omega, \mathbb{R}^d)$. The set \mathcal{N} is defined by a systems of inequalities ($\|\mathbf{u}_h\|_h > 0$ for $h = 1, \dots, m$) and a system of equations $G_h(\mathbf{u}) = 0$, where

$$G_h(\mathbf{u}) := \|\mathbf{u}_h\|_h^2 - \sum_{k=1}^m \sum_{(i,j) \in I_h \times I_k} \int_{\Omega} \beta_{ij} u_i^2 u_j^2. \quad (2.3)$$

It is not difficult to check that $\mathcal{N} \cap \mathcal{E} \neq \emptyset$, and that if $\mathbf{u} \in \mathcal{N} \cap \mathcal{E}$, then \mathcal{N} is a smooth manifold of codimension m in a neighbourhood of \mathbf{u} . Furthermore, one can show that $\mathcal{N} \cap \mathcal{E}$ is a natural constraint, that is, critical points of J restricted on $\mathcal{N} \cap \mathcal{E}$ are critical points of J in the whole space $H_0^1(\Omega, \mathbb{R}^d)$.

Finally, we recall that with the notation previously introduced the functional J can be written as

$$J(\mathbf{u}) = \frac{1}{2} \sum_{h=1}^m \|\mathbf{u}_h\|_h^2 - \frac{1}{4} M_B(\mathbf{u}) \mathbf{1} \cdot \mathbf{1},$$

where $\mathbf{1} = (1, \dots, 1)$, and \cdot denotes the Euclidean scalar product, and that the constrained functional $J|_{\mathcal{N}}$ reads as

$$J(\mathbf{u}) = \frac{1}{4} \sum_{i=1}^d \|u_i\|_i^2 = \frac{1}{4} \sum_{i,j=1}^d \int_{\Omega} \beta_{ij} u_i^2 u_j^2, \quad (2.4)$$

so that $J|_{\mathcal{N}}$ is coercive and bounded from below. Thus, it makes sense to search for a constrained minimizer for

$$c = \inf_{\mathcal{N}} J \geq 0,$$

where we recall that \mathcal{N} has been defined in (1.6). The following is a refinement of [27, Lemma 2.1], and it is the first step to obtain the constant $K > 0$ in Theorem 1.2.

Lemma 2.1. *It holds*

$$c := \inf_{\mathcal{N}} J \leq \bar{C},$$

with

$$\bar{C} = \frac{1}{4} \max_{h=1, \dots, m} \min_{i \in I_h} \left\{ \frac{(1 + \lambda_i)^2}{\beta_{ii}} \right\} \cdot \inf_{\substack{\Omega \supset \Omega_1, \dots, \Omega_m \text{ open} \\ \Omega_i \cap \Omega_j = \emptyset, (i \neq j)}} \sum_{h=1}^m \tilde{S}^2(\Omega_h), \tag{2.5}$$

$\tilde{S}(\Omega)$ being the best Sobolev constant for the embedding $H_0^1(\Omega) \hookrightarrow L^4(\Omega)$, characterized as

$$\tilde{S}(\Omega) := \inf_{u \in H^1(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla u|^2 + u^2}{\int_{\Omega} u^4}.$$

Proof. For each h , denote i_h the index achieving $\min_{i \in I_h} \left\{ \frac{(1 + \lambda_i)^2}{\beta_{ii}} \right\}$. Take $\tilde{u}_{i_1}, \dots, \tilde{u}_{i_m} \neq 0$ such that $\tilde{u}_{i_h} \cdot \tilde{u}_{i_k} \equiv 0$ whenever $h \neq k$, and define $t_h = \|\tilde{u}_{i_h}\|_h / (\sqrt{\beta_{i_h i_h}} |\tilde{u}_{i_h}|_{L^4}^2)$. Define $\bar{\mathbf{u}}$ such that $\bar{u}_{i_h} = t_h \tilde{u}_{i_h}$ for $h = 1, \dots, m$, and $\bar{u}_i = 0$ for $i \neq i_1, \dots, i_m$. It is clear that $\bar{\mathbf{u}}_h \neq 0$ for every h , and that $\bar{\mathbf{u}} \in \mathcal{N}$. Thus, by definition,

$$\begin{aligned} c &\leq J(\bar{\mathbf{u}}) = \frac{1}{4} \sum_{i=1}^d \|\bar{u}_i\|_i^2 = \frac{1}{4} \sum_{h=1}^m t_h^2 \|\tilde{u}_{i_h}\|_{i_h}^2 \\ &\leq \frac{1}{4} \sum_{h=1}^m \frac{(1 + \lambda_{i_h})^2}{\beta_{i_h i_h}} \frac{\|\tilde{u}_{i_h}\|_{H^1}^4}{|\tilde{u}_{i_h}|_{L^4}^4} \leq \frac{1}{4} \max_{h=1, \dots, m} \frac{(1 + \lambda_{i_h})^2}{\beta_{i_h i_h}} \sum_{h=1}^m \frac{\|\tilde{u}_{i_h}\|_{H^1}^4}{|\tilde{u}_{i_h}|_{L^4}^4}, \end{aligned}$$

where, as usual, we used the notation

$$\|u\|_{H^1}^2 := \int_{\Omega} |\nabla u|^2 + u^2. \quad \square$$

The following is a key result, both for the existence result of this section, as well as for the symmetry one in the following.

Lemma 2.2. *For $K = S^2/(16\bar{C}) > 0$ (which depends only on $\lambda_i, \beta_{ii}, i = 1, \dots, d$) we have that, whenever*

Please cite this article in press as: N. Soave, H. Tavares, New existence and symmetry results for least energy positive solutions of Schrödinger systems with mixed competition and cooperation terms, J. Differential Equations (2016), <http://dx.doi.org/10.1016/j.jde.2016.03.015>

$$-\infty < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2,$$

the following inclusion holds:

$$\tilde{\mathcal{N}} \cap \left\{ \mathbf{u} : \sum_{i=1}^d \|u_i\|_i^2 \leq 8\bar{C} \right\} \subset \mathcal{E},$$

where we recall that $\tilde{\mathcal{N}}$ and \mathcal{E} have been defined in (2.1) and (2.2) respectively.

Proof. Let us prove that $M_B(\mathbf{u})$ is strictly diagonally dominant, that is, for each $h = 1, \dots, m$,

$$\sum_{(i,j) \in I_h^2 \Omega} \int \beta_{ij} u_i^2 u_j^2 > \sum_{\substack{k=1 \\ k \neq h}}^m \left| \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2 \right| \tag{2.6}$$

Some of the terms inside the absolute value might be positive, while others might be negative. Suppose, without loss of generality, that there exists $\bar{m} \in \{0, \dots, m\}$ such that

$$\sum_{\substack{k=1 \\ k \neq h}}^m \left| \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2 \right| = - \sum_{\substack{k=1 \\ k \neq h}}^{\bar{m}} \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2 + \sum_{\substack{k=\bar{m}+1 \\ k \neq h}}^m \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2.$$

On the other hand, since $\mathbf{u} \in \tilde{\mathcal{N}}$, for each $h = 1, \dots, m$,

$$\sum_{(i,j) \in I_h^2 \Omega} \int \beta_{ij} u_i^2 u_j^2 \geq \sum_{i \in I_h} \|u_i\|_i^2 - \sum_{\substack{k=1 \\ k \neq h}}^m \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2.$$

Thus (2.6) is true if we show that

$$\sum_{i \in I_h} \|u_i\|_i^2 > 2 \sum_{\substack{k=\bar{m}+1 \\ k \neq h}}^m \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2,$$

which holds if, for $(i, j) \in \mathcal{K}_2$, $\beta_{ij} < K := S^2/(16\bar{C})$. Indeed, recalling that by assumption $\sum_i \|u_i\|_i^2 \leq 8\bar{C}$, we have

$$\begin{aligned} 2 \sum_{\substack{k=\bar{m}+1 \\ k \neq h}}^m \sum_{(i,j) \in I_h \times I_k \Omega} \int \beta_{ij} u_i^2 u_j^2 &< \frac{2K}{S^2} \sum_{\substack{k=\bar{m}+1 \\ k \neq h}}^m \sum_{(i,j) \in I_h \times I_k \Omega} \|u_i\|_i^2 \|u_j\|_j^2 \\ &\leq \frac{16K\bar{C}}{S^2} \sum_{i \in I_h} \|u_i\|_i^2 = \sum_{i \in I_h} \|u_i\|_i^2, \end{aligned}$$

thanks to the choice of K . \square

An immediate consequence is the following.

Lemma 2.3. *If*

$$-\infty < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2,$$

then \mathcal{N} is a manifold at each $\mathbf{u} \in \mathcal{N}$ with $J(\mathbf{u}) < 2\bar{C}$. Moreover, constrained critical points of $J|_{\mathcal{N}}$ such that $J(\mathbf{u}) < 2\bar{C}$ are in fact free critical points of J .

Proof. If $\mathbf{u} \in \mathcal{N}$ and $J(\mathbf{u}) < 2\bar{C}$, then $\mathbf{u} \in \tilde{\mathcal{N}}$ and $\sum_i \|u_i\|_i^2 \leq 8\bar{C}$. Thanks to the previous lemma, we deduce that $\mathbf{u} \in \mathcal{N} \cap \mathcal{E}$, thus \mathcal{N} is a manifold at \mathbf{u} (see [27, Remark 9]), while the other conclusion comes from [27, Proposition 1.2]. \square

Now we can prove that minimizing sequences for c are also conveniently bounded from below.

Lemma 2.4. *Take $\delta := S/(2d) > 0$. If*

$$\beta_{ij} \geq 0 \quad \forall (i, j) \in \mathcal{K}_1 \quad \text{and} \quad -\infty < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2,$$

then for every $\mathbf{u} \in \mathcal{N}$ such that $J(\mathbf{u}) \leq 2\bar{C}$ there holds

$$\left(\max_{i,j \in I_h} \beta_{ij} \right) \cdot \sum_{i \in I_h} |u_i|_{L^4}^2 \geq \delta \quad \forall h = 1, \dots, m.$$

Proof. Following the proof of Lemma 2.2 in [27], since $\mathbf{u} \in \mathcal{N}$:

$$\begin{aligned} S \sum_{i \in I_h} |u_i|_{L^4}^2 &\leq \sum_{i \in I_h} \|u_i\|_i^2 \leq \sum_{(i,j) \in I_h^2 \Omega} \int \frac{\beta_{ij}}{2} (u_i^4 + u_j^4) + K \sum_{i \in I_h} \sum_{j \notin I_h} |u_i|_{L^4}^2 |u_j|_{L^4}^2 \\ &< \max_{i,j \in I_h} \{\beta_{ij}\} \sum_{i \in I_h} |u_i|_{L^4}^4 + \frac{8K\bar{C}}{S} \sum_{i \in I_h} |u_i|_{L^4}^2 \\ &< d \max_{i,j \in I_h} \{\beta_{ij}\} \left(\sum_{i \in I_h} |u_i|_{L^4}^2 \right)^2 + \frac{S}{2} \sum_{i \in I_h} |u_i|_{L^4}^2. \quad \square \end{aligned}$$

Having established the basic properties of minimizing sequences, we can proceed with the core of the argument.

Lemma 2.5. *The constrained functional $J|_{\mathcal{N}}$ satisfies the Palais–Smale condition at level c , whenever B is such that*

$$\beta_{ij} \geq 0 \quad \forall (i, j) \in \mathcal{K}_1 \quad \text{and} \quad -\infty < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2.$$

Proof. Take $\{\mathbf{u}_n\} \subset \mathcal{N}$ such that

$$J(\mathbf{u}_n) \rightarrow c, \quad J'(\mathbf{u}_n) = \sum_{h=1}^m \lambda_{h,n} G'_h(\mathbf{u}_n) + o(1), \quad (2.7)$$

where we recall that G_h has been defined in (2.3). We can take n large enough so that $J(\mathbf{u}_n) \leq 2\bar{C}$. By (2.4), up to a subsequence

$$u_{i,n} \rightarrow u_i \quad \text{weakly in } H_0^1(\Omega) \text{ and strongly in } L^4(\Omega),$$

whence

- (i) $\mathbf{u} \in \tilde{\mathcal{N}}$ and $\sum_i \|u_i\|_i^2 \leq 8\bar{C}$;
- (ii) $M_B(\mathbf{u}_n) \rightarrow M_B(\mathbf{u})$ component-wise.

From (i), by Lemma 2.2 we deduce that $\mathbf{u} \in \mathcal{E}$, so that $M_B(\mathbf{u})$ is positive definite. Testing the second equation in (2.7) with $\hat{\mathbf{u}}^h \in H_0^1(\Omega, \mathbb{R}^d)$ defined by

$$\hat{u}_{i,n}^h := \begin{cases} u_{i,n} & \text{if } i \in I_h \\ 0 & \text{if } i \notin I_h, \end{cases}$$

we obtain with (ii) that

$$\mathbf{o}(1) = M_B(\mathbf{u}_n) \begin{pmatrix} \lambda_{1,n} \\ \vdots \\ \lambda_{m,n} \end{pmatrix} = (M_B(\mathbf{u}) + \mathbf{o}(1)) \begin{pmatrix} \lambda_{1,n} \\ \vdots \\ \lambda_{m,n} \end{pmatrix};$$

multiplying by $(\lambda_{1,n}, \dots, \lambda_{m,n})$, and using the fact that $M_B(\mathbf{u})$ is positive definite, we finally infer

$$\mathbf{o}(1)|(\lambda_{m,n}, \dots, \lambda_{m,n})| \geq C|(\lambda_{1,n}, \dots, \lambda_{m,n})|^2 + \mathbf{o}(1)|(\lambda_{1,n}, \dots, \lambda_{m,n})|^2,$$

yielding $\lambda_{i,n} \rightarrow 0$. Moreover, $\{G'_h(\mathbf{u}_n)\}$ is a uniformly bounded family of operators, thanks to the boundedness of $\{\mathbf{u}_n\}$, and hence $J'(\mathbf{u}_n) \rightarrow 0$. This means that \mathbf{u}_n is a standard Palais–Smale sequence, and the result follows easily from now on. \square

We are now ready to prove the main result of this section.

Proof of Theorem 1.2. The proof is a simple consequence of what was established before. In fact, $c \geq 0$, hence we can take a minimizing sequence \mathbf{u}_n , which we can choose, by Ekeland's variational principle, to be a Palais–Smale sequence for $J|_{\mathcal{N}}$ at level c . Note that the Ekeland's principle is applicable, since by Lemma 2.4 the set $\mathcal{N} \cap \{J \leq 2\bar{C}\}$ endowed with the $H_0^1(\Omega)$ topology is a complete metric space. Thus, by the previous lemma, up to a subsequence $\mathbf{u}_n \rightarrow \mathbf{u}$ strongly in $H_0^1(\Omega)$, and by Lemma 2.4 this implies that $\mathbf{u} \in \mathcal{N}$. By convergence, we infer that

$$J(\mathbf{u}) = \lim_{n \rightarrow \infty} J(\mathbf{u}_n) = c,$$

which completes the proof since $\mathcal{N} \cap \{J < 2\bar{C}\}$ is a natural constraint (cf. Lemma 2.3). \square

Remark 2.6. In the proof of Theorem 1.2 we used the validity of the Sobolev inequality (for the definition of S and $\tilde{S}(\Omega)$), the weak lower semi-continuity of $\|\cdot\|_i$ and the compactness of the embedding $H_0^1(\Omega) \hookrightarrow L^4(\Omega)$ (see Lemma 2.5). These tools are still in force if we replace $H_0^1(\Omega)$ with $H_{\text{rad}}^1(\mathbb{R}^N)$, so that, as already observed, Theorem 1.2 holds also for the radial problem in the whole space.

Remark 2.7. For future reference, we observe that the constant K is equal to $S^2/(16\bar{C})$, where

$$S := \inf_{i=1, \dots, d} \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\|u\|_i^2}{|u|_{L^4}^2},$$

and

$$\bar{C} = \frac{1}{4} \max_{h=1, \dots, m} \min_{i \in I_h} \left\{ \frac{(1 + \lambda_i)^2}{\beta_{ii}} \right\} \cdot \inf_{\substack{\Omega \supset \Omega_1, \dots, \Omega_m \text{ open} \\ \Omega_i \cap \Omega_j = \emptyset, (i \neq j)}} \sum_{h=1}^m \tilde{S}^2(\Omega_h)$$

(recall Lemma 2.2 and (2.5)).

Remark 2.8. In the particular case of full cooperative systems, we can have a better (and more explicit) constant $K > 0$. Thinking for instance at the d -decomposition $\mathbf{a} = (0, 1, \dots, m)$, one can take:

$$K := \frac{\min_{i=1, \dots, d} \{S_i^2\}}{2 \sum_{j=1}^d \frac{S_j^2}{\beta_{jj}}}, \quad \text{where } S_i := \inf_{\int_{\Omega} u^4 = 1} \|u\|_i^2. \tag{2.8}$$

This constant is similar to the one appearing in assumption (H2) of [30] (for $d = 2$), being ours slightly worse in the framework of Z.-Q. Wang and M. Willem’s paper. This is a price to pay from passing from $d = 2$ to more equations, since we had to prove that $M_B(\mathbf{u})$ is positive definite by proving that actually it is strictly diagonally dominant (while in 2×2 matrices one can perform an explicit computation). The proof of (2.8) is not completely immediate, but since we consider this to be a lateral statement, here we just provide some hints. By taking \tilde{w}_i to be functions achieving S_i , we can take $w_i = \sqrt{S_i/\beta_{ii}}\tilde{w}_i$, which satisfies the equation $-\Delta w_i + \lambda_i w_i = \beta_{ii} w_i^3$, with $\|w_i\|_i^2 = S_i^2/\beta_{ii}$. Moreover, since we are in a full cooperative case, it is straightforward to check that $\mathbf{w} \in \tilde{\mathcal{N}}$. By using (2.8), $\mathbf{w} \in \mathcal{E}$ and there exist $0 < t_i < 1$ such that $(t_1 w_1, \dots, t_d w_d) \in \mathcal{N}$. Thus $c \leq \sum_{j=1}^d S_j^2/(4\beta_{jj})$, and working through the proof of Lemma 2.2 the rest follows.

3. Partial symmetry of nonnegative minimizers

We now turn to the problem of the symmetry of least energy positive solutions when Ω is a radially symmetric bounded domain of \mathbb{R}^N (which we always assume along this section). Observe

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that when Ω is a ball, the classical result by Troy [29] (see also [24]) yields that, in the cooperative case, each positive solution is radially symmetric. However, thanks to [28, Theorem 1.4] (which deal with the competitive case), [22, Corollary 0.5] (mixed cooperation and competition case), and [30, Section 3] (cooperative case, Ω an annulus), it is known that *in general* a least energy positive solution is not radial. On the other hand, it is natural to expect that least energy solutions inherit part of the symmetric structure of the problem. Theorem 1.8 establishes that this is the case when the competition takes place between two groups of cooperative components, or in a purely cooperative setting.

3.1. Comments on Theorem 1.8

Due to its general formulation, Theorem 1.8 might not be easy to read and to understand. In this subsection we present several remarks which should help the reader towards this purpose.

Remark 3.1. The reader could object that only the division in macro-groups is necessary, and that one could drop the original division into m groups considering only a 2 decomposition of d in 2 different groups. This is a particular case of our result, but it is not equivalent, because in our statement we allow to minimize in different Nehari type sets (related to the division in m groups), and so in this way we can deal with a larger class of constrained minimizers.

As an illustrative example, we consider a 3 components system, separating (u_1, u_2, u_3) into 2 macro-groups (u_1, u_2) and u_3 . This can be the result of two different decompositions:

- Let us consider the natural 2-decomposition $(0, 2, 3)$. By Theorem 1.7 in [27], it is known that if $\beta_{12} > C \max\{\beta_{11}, \beta_{22}\}$ for a positive constant $C > 0$ depending on λ_i , and $\beta_{13}, \beta_{23} \ll -1$, then the minimum of J on $\mathcal{N}_{(0,2,3)}$ is achieved by a *positive* solution of (1.2) (here $\mathcal{N}_{(0,2,3)}$ is the Nehari set determined by the 2-decomposition of 3). The same result holds true if $\beta_{13}, \beta_{23} < 0$ and $\beta_{12} \gg 1$, as proved in Theorem 0.1 in [22]. In both cases, Theorem 1.8 applies proving that (u_1, u_2) and u_3 are foliated Schwartz symmetric with respect to antipodal points.
- Let us now consider the 3-decomposition $(0, 1, 2, 3)$, denoting by $\mathcal{N}_{(0,1,2,3)}$ the corresponding Nehari set. By Corollary 1.3, there exists $K > 0$ such that if $\beta_{ij} < K$ for every $i \neq j$, then the minimum of J on $\mathcal{N}_{(0,1,2,3)}$ is achieved by a *positive* solution of (1.2). If we assume further that $0 < \beta_{12} < K$ and $\beta_{13}, \beta_{23} < 0$, then by Theorem 1.8 we obtain that (u_1, u_2) and u_3 are foliated Schwartz symmetric with respect to antipodal points.

The second result would not have been obtained if we had considered only the 2-decomposition $(0, 2, 3)$ in our symmetry result.

Let us now make some comments regarding Theorem 1.8 in the *purely cooperative case*, which correspond to $l = d$ in the assumptions.

Remark 3.2. In [30], Z.-Q. Wang and M. Willem proved partial symmetry results in two situations. For a system with d components, they showed that if the infimum of J on the natural Nehari manifold

$$\left\{ \mathbf{u} \in H_0^1(\Omega; \mathbb{R}^d) \mid \mathbf{u} \neq \mathbf{0} \text{ and } \sum_{i=1}^d \partial_i J(\mathbf{u}) u_i = 0 \right\}$$

is achieved (which is suitable only for large cooperation rates), then any positive minimizer is such that all the components (u_1, \dots, u_d) are foliated Schwartz symmetric with respect to the same point. Moreover, for systems of $d = 2$ components, they showed that if $0 < \beta_{12} < C(\beta_{11}, \beta_{22}, \lambda_1, \lambda_2)$, then the infimum of J on the Nehari set

$$\left\{ \mathbf{u} \in H_0^1(\Omega; \mathbb{R}^2) \mid \begin{array}{l} u_i \neq 0 \text{ and } \partial_i J(\mathbf{u})u_i = 0 \\ \text{for } i = 1, 2 \end{array} \right\}$$

is achieved by a positive solution of (1.2), and any minimizer is such that u_1, u_2 are foliated Schwartz symmetric with respect to the same point. The restriction $d = 2$ is relevant in their proof.

Our result recovers both the ones in [30], extend the second one to systems with an arbitrary number of components, and provide partial symmetry also when m -decompositions of d with $m \neq 1, d$ are considered. For this reason, we think that Theorem 1.8 in the purely cooperative case deserves a statement on its own.

Corollary 3.3. *Let $d \geq 2$, and let \mathbf{a} be a m -decomposition of d . Take $K > 0$ be as in Theorem 1.2. If*

$$\beta_{ij} > 0 \quad \forall (i, j) \in \mathcal{K}_1, \quad 0 < \beta_{ij} < K \quad \forall (i, j) \in \mathcal{K}_2,$$

then any nonnegative minimizer $\tilde{\mathbf{u}}$ of J constrained on \mathcal{N} is such that all the components $\tilde{u}_1, \dots, \tilde{u}_d$ are foliated Schwartz symmetric with respect to the same point.

Remark 3.4. Concerning the proof of Theorem 1.8, in the literature the partial symmetry of solutions of elliptic equations is often obtained through an infsup characterization. Dealing with systems, it turns out to be very complicated to obtain such a variational characterization; for instance, this is why the proof of Theorem 1.2 in [30] works only for systems with 2 components. Thus, we think that it is worth to point out that we will not use any infsup characterization, basing our argument directly on the constrained minimality.

3.2. Proof of Theorem 1.8

In what follows, without loss of generality, we suppose that Ω is radial with respect to 0. We will use polarization techniques, and hence at first we recall some definitions which are by now classic.

Assume H is a closed half-space in \mathbb{R}^N . We denote by $\sigma_H : \mathbb{R}^N \rightarrow \mathbb{R}^N$ the reflection with respect to the boundary ∂H of H . For a measurable function $w : \mathbb{R}^N \rightarrow \mathbb{R}$ we define the polarization w_H of w relative to H by

$$w_H(x) = \begin{cases} \max\{w(x), w(\sigma_H(x))\}, & x \in H, \\ \min\{w(x), w(\sigma_H(x))\}, & x \in \mathbb{R}^N \setminus H. \end{cases}$$

We consider the set \mathcal{H}_0 of all closed half-spaces H in \mathbb{R}^N such that $0 \in \partial H$. Given an unitary vector $p \in \mathbb{S}^{N-1}$, we denote by $\mathcal{H}_0(p)$ the set of all closed half-spaces $H \in \mathcal{H}_0$ such that $p \in \text{int}(H)$.

Recall that a function $f : \mathbb{R}^N \rightarrow \mathbb{R}$ is said to be foliated Schwarz symmetric with respect to a unitary vector p if it is axially symmetric with respect to the axis $\mathbb{R}p$ and non-increasing in the polar angle $\theta = \arccos(p \cdot x/|x|) \in [0, \pi]$. We mention that, up to our knowledge, the link between polarization and foliated Schwarz symmetry appeared firstly in [7,26]. We would like to mention also the precursory works [1,3,8] that brought to light the relation between polarizations and rearrangements in many different settings, and refer to the survey [32] for a detailed history of the subject. The following is a useful alternative characterization of foliated Schwarz symmetry, and we refer to [7, Lemma 4.2] (see also [32, Proposition 2.7]) for the proof.

Lemma 3.5. *Let Ω a radial set centered at the origin, and let $u : \Omega \rightarrow \mathbb{R}$ be a continuous function. Then u is foliated Schwarz symmetric with respect to $p \in \mathbb{S}^{N-1}$ if, and only if, for every $H \in \mathcal{H}_0(p)$ we have $u(x) \geq u(\sigma_H(x))$ whenever $x \in \Omega \cap H$.*

For every $H \in \mathcal{H}$ we denote by $\widehat{H} \in \mathcal{H}_0$ the closure of the complementary half-space $\mathbb{R}^N \setminus H$. In the spirit of [6,28,30], the proof of Theorem 1.8 is based upon a general criterion (cf. for instance Theorem 2.6 in [6] or Theorem 4.3 in [28]).

Proposition 3.6. *Let \mathbf{u} be a nonnegative solution of (1.2). If for every $H \in \mathcal{H}_0$ the function*

$$\mathbf{u}^H := (\mathbf{u}_{1,H}, \dots, \mathbf{u}_{\widehat{H},H}, \mathbf{u}_{\widehat{H}+1,\widehat{H}}, \dots, \mathbf{u}_m, \widehat{H}) = (u_{1,H}, \dots, u_{l,H}, u_{l+1,\widehat{H}}, \dots, u_{d,\widehat{H}})$$

is still a solution of (1.2), then (u_1, \dots, u_l) and (u_{l+1}, \dots, u_d) are foliated Schwartz symmetric with respect to antipodal points.

Proof. Since $\mathbf{u} \in \mathcal{N}$, there exists an index $i \in I_1$ such that $u_i \not\equiv 0$ in Ω . Without loss of generality, we assume $i = 1$. Let $r > 0$ be such that $\partial B_r(0) \subset \Omega$, and let $p \in \mathbb{S}^{N-1}$ be such that $\max_{\partial B_r(0)} u_1 = u_1(rp)$. By assumption, for any $H \in \mathcal{H}_0(p)$,

$$-\Delta u_{1,H} + \lambda_1 u_{1,H} = \sum_{j=1}^l \beta_{1j} u_{1,H} u_{j,H}^2 + \sum_{j=l+1}^d \beta_{1j} u_{1,H} u_{j,\widehat{H}}^2 \quad \text{in } \Omega.$$

Thus, if we let $w := u_{1,H} - u_1$, we obtain

$$\begin{aligned} -\Delta w + \lambda_1 w &= \sum_{j=1}^l \beta_{1j} (u_{1,H} u_{j,H}^2 - u_1 u_j^2) + \sum_{j=l+1}^d \beta_{1j} (u_{1,H} u_{j,\widehat{H}}^2 - u_1 u_j^2) \\ &= \sum_{j=1}^l \beta_{1j} u_{1,H} (u_{j,H}^2 - u_j^2) + \sum_{j=l+1}^d \beta_{1j} u_{1,H} (u_{j,\widehat{H}}^2 - u_j^2) \\ &\quad + \sum_{j=1}^d \beta_{1j} u_j^2 (u_{1,H} - u_1), \end{aligned}$$

in $\Omega \cap H$, that is

$$-\Delta w + \underbrace{\left(\lambda_1 - \sum_{j=1}^d \beta_{1j} u_j^2 \right)}_{=: q(x) \in L_{loc}^\infty(\Omega \cap H)} w = \sum_{j=1}^l \beta_{1j} u_{1,H} (u_{j,H}^2 - u_j^2) + \sum_{j=l+1}^d \beta_{1j} u_{1,H} (u_{j,\hat{H}}^2 - u_j^2). \tag{3.1}$$

By definition $u_{j,H} \geq u_j$ in $\Omega \cap H$ for every $j = 1, \dots, l$, while $u_{j,\hat{H}} \leq u_j$ in $\Omega \cap H$ for every $j = l + 1, \dots, d$. Therefore, recalling that $\beta_{1j} > 0$ for $j = 1, \dots, l$, while $\beta_{1j} < 0$ for $j = l + 1, \dots, d$, we deduce that $-\Delta w + q(x)w \geq 0$ and $w \geq 0$ in $\Omega \cap H$. The strong maximum principle leads to the alternative $w > 0$ in $\Omega \cap H$, or $w \equiv 0$ in $\Omega \cap H$. As $w(rp) = 0$, the latter condition holds true, and coming back to equation (3.1) we deduce that

$$0 = -\Delta w + q(x)w = \sum_{j=1}^l \beta_{1j} u_{1,H} (u_{j,H}^2 - u_j^2) + \sum_{j=l+1}^d \beta_{1j} u_{1,H} (u_{j,\hat{H}}^2 - u_j^2) \geq 0,$$

which in turn implies (since $u_1 \not\equiv 0 \Rightarrow u_1 > 0$ in Ω) that $u_{j,H} \equiv u_j$ in $\Omega \cap H$ for every $j = 1, \dots, l$ and $u_{j,\hat{H}} \equiv u_j$ in $\Omega \cap H$ for every $j = l + 1, \dots, k$. Since $H \in \mathcal{H}_0(p)$ has been arbitrarily chosen, the thesis follows from Lemma 3.5. \square

In the following we will prove that, under the assumptions of Theorem 1.8, it is possible to apply Proposition 3.6. First we need some preliminary results.

Lemma 3.7. *For every $u, v \in H_0^1(\Omega)$ and $H \in \mathcal{H}_0(p)$, then also $u_H, v_H \in H_0^1(\Omega)$, and moreover:*

- (i) $\int_{\Omega} |u|^p dx = \int_{\Omega} |u_H|^p$, for every $p \geq 1$.
- (ii) $\int_{\Omega} |\nabla u_H|^2 = \int_{\Omega} |\nabla u|^2$.
- (iii) $\int_{\Omega} u^2 v^2 \leq \int_{\Omega} u_H^2 v_H^2$ and $\int_{\Omega} u^2 v^2 \geq \int_{\Omega} u_H^2 v_{\hat{H}}^2$.

Proof. For the first two items, check for instance [32, Lemma 3.1]. As for (iii), the first inequality follows by Proposition 31.7 in [33], while the second one is a particular case of Lemma 4.5 in [28]. \square

Assume that B is as in the assumption of Theorem 1.8, let $c = \inf_{\mathcal{N}} J$, and let \mathbf{u} be a nonnegative minimizer of J on \mathcal{N} . In view of Proposition 3.6, we aim at proving that $\mathbf{u}^H := (u_{1,H}, \dots, u_{l,H}, u_{l+1,\hat{H}}, \dots, u_{d,\hat{H}})$ also achieves c for every $H \in \mathcal{H}_0$. In this perspective, the main difficulty consists in showing that $\mathbf{u}^H \in \mathcal{N}$ (in the literature, this is usually the part of the proof which requires an inf sup characterization, see Remark 3.4). To this aim, we study the function of real variables:

$$\Psi : (t_1, \dots, t_m) \in \overline{\mathbb{R}}_+^m \mapsto J(\sqrt{t_1} \mathbf{u}_{1,H}, \dots, \sqrt{t_{\hat{h}}} \mathbf{u}_{\hat{h},H}, \sqrt{t_{\hat{h}+1}} \mathbf{u}_{\hat{h}+1,\hat{H}}, \dots, \sqrt{t_m} \mathbf{u}_{m,\hat{H}}) \in \mathbb{R},$$

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and claim that, under the considered assumption, it has a unique maximum point $\tilde{\mathbf{t}} = (\tilde{t}_1, \dots, \tilde{t}_m)$ such that $\tilde{t}_h > 0$ for every $h = 1, \dots, m$.

By the properties of polarization stated in Lemma 3.7, and the assumptions made on B , we have that

$$\mathbf{u}^H \in \tilde{\mathcal{N}}, \quad \sum_{h=1}^m \|\mathbf{u}_h^H\|_h^2 = \sum_{j=1}^d \|u_j\|_j^2 = 4c \leq 4\bar{C}$$

where \bar{C} is as in (2.5). Hence, thanks to Lemma 2.2, the matrix $M_B(\mathbf{u}^H)$ is positive definite, and observing that

$$\Psi(t_1, \dots, t_m) = \frac{1}{2} \sum_{h=1}^m \|\mathbf{u}_h^H\|_h^2 t_h - \frac{1}{4} M_B(\mathbf{u}^H) \mathbf{t} \cdot \mathbf{t},$$

this implies that there exists $\kappa > 0$ such that

$$\Psi(t_1, \dots, t_m) \leq \frac{1}{2} \sum_{h=1}^m t_h \|\mathbf{u}_h^H\|_h^2 - \kappa \sum_{h=1}^m t_h^2 \rightarrow -\infty \quad \text{as } |\mathbf{t}| \rightarrow \infty.$$

Therefore, Ψ admits a global maximum $\tilde{\mathbf{t}}$ in $\overline{\mathbb{R}_+^m}$. Note also that it is a strictly concave function. Since Ψ is of class C^1 up to the boundary of \mathbb{R}_+^m , the maximality of \tilde{t} entails $\partial_h \Psi(\tilde{\mathbf{t}}) \leq 0$ if $\tilde{t}_h = 0$, and $\partial_h \Psi(\tilde{\mathbf{t}}) = 0$ if $\tilde{t}_h > 0$. In particular, we observe that

$$\|\mathbf{u}_h^H\|_h^2 = \sum_{k=1}^m M_B(\mathbf{u}^H)_{hk} \tilde{t}_k^H \quad \text{whenever } \tilde{t}_h > 0. \tag{3.2}$$

Remark 3.8. Let us consider the function

$$\Phi : (t_1, \dots, t_m) \in \mathbb{R}_+^m \mapsto J(\sqrt{t_1} \mathbf{u}_1, \dots, \sqrt{t_m} \mathbf{u}_m).$$

Under our assumption, by Lemma 2.2 any minimizer for c stays in \mathcal{E} . Therefore, as in the previous discussion, we can check that Φ is strictly concave and has a maximum point in $\overline{\mathbb{R}_+^m}$. Moreover, since $\mathbf{u} \in \mathcal{N}$, the point $\mathbf{1} = (1, \dots, 1)$ is a critical point for Φ . Hence, by strict concavity, $\mathbf{1}$ is the unique critical point of Φ , and is a global maximum.

Lemma 3.9. We have

$$\sum_{h=1}^m \tilde{t}_h \|\mathbf{u}_h^H\|_h^2 \leq 4\bar{C}.$$

Proof. We claim that

$$\|\mathbf{u}_h^H\|_h^2 \tilde{t}_h = \sum_{k=1}^m M_B(\mathbf{u}^H)_{hk} \tilde{t}_h \tilde{t}_k \quad \forall h = 1, \dots, m.$$

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Indeed, if $\tilde{t}_h = 0$ this relation is trivially satisfied. If $\tilde{t}_h > 0$, then it follows by (3.2). Therefore

$$J\left(\sqrt{\tilde{t}_1}\mathbf{u}_1^H, \dots, \sqrt{\tilde{t}_m}\mathbf{u}_m^H\right) = \frac{1}{4} \sum_{h=1}^m \|\mathbf{u}_h^H\|_h^2 \tilde{t}_h.$$

On the other hand, combining Lemma 3.7 with the assumptions on B , we have also

$$\begin{aligned} J\left(\sqrt{\tilde{t}_1}\mathbf{u}_1^H, \dots, \sqrt{\tilde{t}_m}\mathbf{u}_m^H\right) &\leq J\left(\sqrt{\tilde{t}_1}\mathbf{u}_1, \dots, \sqrt{\tilde{t}_m}\mathbf{u}_m\right) \leq \sup_{t_1, \dots, t_m \geq 0} J\left(\sqrt{t_1}\mathbf{u}_1, \dots, \sqrt{t_m}\mathbf{u}_m\right) \\ &= J(\mathbf{u}_1, \dots, \mathbf{u}_m) = c \leq \bar{C}. \end{aligned}$$

Notice that we have used the fact that Φ has the unique maximizer $(1, \dots, 1)$, see the previous remark. \square

Lemma 3.10. *There holds*

$$\tilde{t}_1, \dots, \tilde{t}_m > 0.$$

Proof. Assume, in view of a contradiction, that $\tilde{t}_1 = 0$. We will check that

$$J\left(\sqrt{t_1}\mathbf{u}_1^H, \dots, \sqrt{\tilde{t}_m}\mathbf{u}_m^H\right) > J\left(0, \sqrt{\tilde{t}_2}\mathbf{u}_2^H, \dots, \sqrt{\tilde{t}_m}\mathbf{u}_m^H\right)$$

for $t_1 > 0$ sufficiently close to 0, which contradicts the maximality of $(0, \tilde{t}_2, \dots, \tilde{t}_m)$. The left hand side of the inequality can be rewritten as

$$\frac{1}{2}t_1\|\mathbf{u}_1^H\|_1^2 - \frac{1}{2} \sum_{h=2}^m M_B(\mathbf{u}^H)_{1h}t_1\tilde{t}_h - \frac{1}{4}M_B(\mathbf{u}^H)_{11}t_1^2 + \frac{1}{2} \sum_{h=2}^m \tilde{t}_h\|\mathbf{u}_h^H\|_h^2 - \frac{1}{4} \sum_{h,k=2}^m M_B(\mathbf{u}^H)_{hk}\tilde{t}_h\tilde{t}_k.$$

Observe that, for $\beta_{ij} < K$ in \mathcal{K}_2 (and K as in Remark 2.7), it results

$$\begin{aligned} \sum_{h=2}^m M_B(\mathbf{u}^H)_{1h}\tilde{t}_h &= \sum_{h=2}^m \sum_{(i,j) \in I_1 \times I_h} \int_{\Omega} \beta_{ij}(u_i^H u_j^H)^2 \tilde{t}_h \leq \frac{K}{S^2} \sum_{h=2}^m \sum_{(i,j) \in I_1 \times I_h} \|u_i^H\|_i^2 \|u_j^H\|_j^2 \tilde{t}_h \\ &\leq \frac{K}{S^2} \|\mathbf{u}_1^H\|_1^2 \sum_{h=2}^m \tilde{t}_h \|\mathbf{u}_h^H\|_h^2 \leq \frac{4K\bar{C}}{S^2} \|\mathbf{u}_1^H\|_1^2 \leq \frac{1}{4} \|\mathbf{u}_1^H\|_1^2, \end{aligned}$$

where we used the estimate of the previous lemma. Thus

$$\begin{aligned} &\frac{1}{2}t_1\|\mathbf{u}_1^H\|_1^2 - \frac{1}{2} \sum_{h=2}^m M_B(\mathbf{u}^H)_{1h}t_1\tilde{t}_h - \frac{1}{4}M_B(\mathbf{u}^H)_{11}t_1^2 \\ &= \frac{1}{2}t_1 \left(\|\mathbf{u}_1^H\|_1^2 - \sum_{h=2}^m M_B(\mathbf{u}^H)_{1h}\tilde{t}_h - \frac{1}{2}M_B(\mathbf{u}^H)_{11}t_1 \right) \geq \frac{1}{2}t_1 \left(\frac{3}{4}\|\mathbf{u}_1^H\|_1^2 - t_1\mathbf{M}(\mathbf{u}^H)_{11} \right) > 0 \end{aligned}$$

for sufficiently small $t_1 > 0$, which yields the desired contradiction. \square

End of the proof of Theorem 1.8. Given \mathbf{u} achieving c and $H \in \mathcal{H}_0$, we have concluded that there exists a maximizer \tilde{t} for the function Ψ in $\overline{\mathbb{R}_+^m}$, and $\tilde{t}_1, \dots, \tilde{t}_m > 0$. By (3.2), we infer that $(\sqrt{\tilde{t}_1} \mathbf{u}_1^H, \dots, \sqrt{\tilde{t}_m} \mathbf{u}_m^H) \in \mathcal{N}$; together with Lemma 3.7, this implies that

$$\begin{aligned} c &\leq J\left(\sqrt{\tilde{t}_1} \mathbf{u}_1^H, \dots, \sqrt{\tilde{t}_m} \mathbf{u}_m^H\right) \leq J\left(\sqrt{\tilde{t}_1} \mathbf{u}_1, \dots, \sqrt{\tilde{t}_m} \mathbf{u}_m\right) \\ &\leq \sup_{t_1, \dots, t_m \geq 0} J\left(\sqrt{t_1} \mathbf{u}_1, \dots, \sqrt{t_m} \mathbf{u}_m\right) = J(\mathbf{u}) = c, \end{aligned}$$

which is then a chain of equalities. The uniqueness of the maximum for the function Φ (see Remark 3.8) entails $\tilde{t}_h = 1$ for every h , and thus \mathbf{u}^H also achieves c , being in particular a solution of (1.2) (cf. Lemma 2.3). We can now conclude using the criterion of Proposition 3.6. \square

4. Existence of least energy positive solutions

This section is devoted to the proofs of Theorems 1.4 and 1.5. They are inspired by those of Theorems 1.6 and 1.7 in [27].

The argument for the proof of both results is the same up to a certain point, thus we proceed in a unified way. Under the considered assumptions, by Theorem 1.2 there exists a nonnegative solution \mathbf{u} of (1.2) which minimizes J in the Nehari set \mathcal{N} . We wish to show that

$$\inf_{\mathcal{N}} J < \inf_{\mathcal{N} \cap \{u_i = 0 \text{ for some } i\}} J.$$

If this is true, then by minimality $u_i \neq 0$ for every i . By contradiction, let us assume that for some index l there holds $u_l = 0$. Let $\bar{h} \in \{1, \dots, m\}$ be such that $l \in I_{\bar{h}}$. By definition of \mathcal{N} , there exists $p \in I_{\bar{h}}$ such that $u_p \neq 0$, and actually, by Lemma 2.4 and since $\beta_{\bar{h}} > \max_{i \in I_{\bar{h}}} \{\beta_{ii}\}$, there exists $p \in I_{\bar{h}}$ such that

$$\beta_{\bar{h}} |u_p|_{L^4}^2 \geq \frac{\delta}{|I_{\bar{h}}|} \geq \frac{\delta}{d}. \tag{4.1}$$

By Lemma 2.2 we know that $\mathbf{u} \in \mathcal{E}$, and hence \mathcal{N} defines, in a neighbourhood of \mathbf{u} , a smooth manifold (actually a C^2 -manifold, as it is immediate to verify) of codimension m in $H_0^1(\Omega; \mathbb{R}^d)$. We claim that

$$d^2 J(\mathbf{u})[\mathbf{v}, \mathbf{v}] \geq 0 \quad \text{for every } \mathbf{v} \in T_{\mathbf{u}}(\mathcal{N}), \tag{4.2}$$

where $T_{\mathbf{u}}(\mathcal{N})$ denotes the tangent space to \mathcal{N} at the point \mathbf{u} . To prove this, we observe that since \mathcal{N} is of class C^2 , for any $\mathbf{v} \in T_{\mathbf{u}}(\mathcal{N})$ there exists a C^2 curve $\gamma : (-\varepsilon, \varepsilon) \rightarrow \mathcal{N}$ for some $\varepsilon > 0$ such that $\gamma(0) = \mathbf{u}$ and $\gamma'(0) = \mathbf{v}$. Now by minimality of \mathbf{u} , and recalling that $dJ(\mathbf{u}) = 0$, we infer that

$$0 \leq \frac{d^2}{dt^2} J(\gamma(t)) \Big|_{t=0} = d^2 J(\gamma(t))[\gamma'(t), \gamma'(t)] \Big|_{t=0} + dJ(\gamma(t))[\gamma''(t)] \Big|_{t=0} = d^2 J(\mathbf{u})[\mathbf{v}, \mathbf{v}],$$

which proves the claim (4.2). By direct computations, one can easily check that

$$d^2 J(\mathbf{u})[\mathbf{v}, \mathbf{v}] = \sum_{i=1}^d \|v_i\|_i^2 - \sum_{i,j=1}^d \int_{\Omega} \beta_{ij} u_i^2 v_j^2 - 2 \sum_{i,j=1}^d \int_{\Omega} \beta_{ij} u_i u_j v_i v_j.$$

We consider the variation \mathbf{v} defined by

$$v_i := \begin{cases} 0 & \text{if } i \neq l \\ u_p & \text{if } i = l. \end{cases}$$

Since $\langle \nabla G_h(\mathbf{u}), \mathbf{v} \rangle = 0$ for every h , we have that $\mathbf{v} \in T_{\mathbf{u}}(\mathcal{N})$ (for the reader’s convenience, we recall that G_h has been defined in (2.3)). Plugging this choice of \mathbf{v} into (4.2), we infer that

$$0 \leq \|u_p\|_l^2 - \sum_{i \in I_{\bar{h}} \Omega} \beta_{il} u_i^2 u_p^2 - \sum_{i \notin I_{\bar{h}} \Omega} \beta_{il} u_i^2 u_p^2. \tag{4.3}$$

On the other hand, testing the equation for u_p against u_p itself, and recalling that $u_l \equiv 0$, we deduce that

$$\|u_p\|_p^2 = \sum_{i \in I_{\bar{h}} \setminus \{l\} \Omega} \beta_{ip} u_i^2 u_p^2 + \sum_{i \notin I_{\bar{h}} \Omega} \beta_{ip} u_i^2 u_p^2 \tag{4.4}$$

Conclusion of the proof of Theorem 1.4. Using assumptions (i)–(iii) in the theorem, (4.3) and (4.4) give

$$0 \leq \|u_p\|_p^2 - \sum_{i \in I_{\bar{h}} \setminus \{l\} \Omega} \beta_{\bar{h}} u_i^2 u_p^2 - \sum_{i \notin I_{\bar{h}} \Omega} b u_i^2 u_p^2,$$

and

$$\|u_p\|_p^2 = \int_{\Omega} \beta_{pp} u_p^4 + \sum_{i \in I_{\bar{h}} \setminus \{l, p\} \Omega} \beta_{\bar{h}} u_i^2 u_p^2 + \sum_{i \notin I_{\bar{h}} \Omega} b u_i^2 u_p^2,$$

respectively. Therefore

$$0 \leq \int_{\Omega} \beta_{pp} u_p^4 + \sum_{i \in I_{\bar{h}} \setminus \{l, p\} \Omega} \beta_{\bar{h}} u_i^2 u_p^2 - \sum_{i \in I_{\bar{h}} \setminus \{l\} \Omega} \beta_{\bar{h}} u_i^2 u_p^2 = \int_{\Omega} (\beta_{pp} - \beta_{\bar{h}}) u_p^4 < 0$$

whenever $\beta_{\bar{h}} > \beta_{pp}$, which is guaranteed by our assumptions. \square

Conclusion of the proof of Theorem 1.5. Recalling the explicit shapes $K = S^2/16\bar{C}$ and $\delta = S/2d$ from Lemma 2.4 and Remark 2.7, (4.1), (4.3) and (4.4) give

$$\begin{aligned}
 0 &\leq \int_{\Omega} \beta_{pp} u_p^4 + \sum_{i \in I_{\bar{h}} \setminus \{l, p\}} \int_{\Omega} \beta_{\bar{h}} u_i^2 u_p^2 - \sum_{i \in I_{\bar{h}} \setminus \{l\}} \int_{\Omega} \beta_{\bar{h}} u_i^2 u_p^2 + \sum_{i \notin I_{\bar{h}}} \int_{\Omega} (\beta_{ip} - \beta_{il}) u_i^2 u_p^2 \\
 &= \int_{\Omega} (\beta_{pp} - \beta_{\bar{h}}) u_p^4 + \sum_{i \notin I_{\bar{h}}} \int_{\Omega} (\beta_{ip} - \beta_{il}) u_i^2 u_p^2 \\
 &\leq \underbrace{(\beta_{pp} - \beta_{\bar{h}})}_{<0} |u_p|_{L^4}^4 + \frac{2K}{\alpha d^2 S} |u_p|_{L^4}^2 \sum_{i \notin I_{\bar{h}}} \|u_i\|_i^2 \\
 &\leq \frac{\beta_{pp} - \beta_{\bar{h}}}{\beta_{\bar{h}}} \frac{\delta}{d} |u_p|_{L^4}^2 + \frac{8K\bar{C}}{\alpha d^2 S} |u_p|_{L^4}^2 \\
 &= \frac{S}{2d^2} \left(\frac{\beta_{pp}}{\beta_{\bar{h}}} - 1 + \frac{1}{\alpha} \right) |u_p|_{L^4}^2 = \frac{S}{2d^2} \left(\frac{\beta_{pp}}{\beta_{\bar{h}}} - \frac{\alpha - 1}{\alpha} \right) |u_p|_{L^4}^2 < 0,
 \end{aligned}$$

where we used the assumption on the coupling parameters, and the estimates of [Lemmas 2.1 and 2.4](#). \square

Remark 4.1. If Ω is replaced by \mathbb{R}^N , the proofs of [Theorems 1.4 and 1.5](#) remain unchanged.

5. Non-existence results in $H^1(\mathbb{R}^N)$

In this section we prove [Theorem 1.6](#), which illustrates that when working in $\Omega = \mathbb{R}^N$ in presence of simultaneous cooperation and competition, in order to find some kind of least energy solution it is often necessary to work in $H_{\text{rad}}^1(\mathbb{R}^N)$ instead that in $H^1(\mathbb{R}^N)$. We choose a m -decomposition \mathbf{a} of d , and we assume that the basic assumption [\(1.4\)](#) holds true. Throughout this section we assume that $\beta_{ij} > 0$ for every $(i, j) \in \mathcal{K}_1$ (recall the definition [\(1.5\)](#) of \mathcal{K}_1) and, for every $h = 1, \dots, m$, we consider the sub-system

$$\begin{cases} -\Delta v_i + \lambda_i v_i = \sum_{j \in I_h} \beta_{ij} v_i v_j^2 & \text{in } \mathbb{R}^N \\ v_i \in H^1(\mathbb{R}^N), \end{cases} \quad \forall i \in I_h. \tag{5.1}$$

We introduce the functional

$$E_h(\mathbf{v}) := \int_{\mathbb{R}^N} \sum_{i \in I_h} \frac{1}{2} (|\nabla v_i|^2 + v_i^2) - \frac{1}{4} \sum_{(i,j) \in I_h^2} \beta_{ij} v_i^2 v_j^2,$$

and the Nehari manifold for the system [\(5.1\)](#), defined by

$$\mathcal{M}_h := \left\{ \mathbf{v} \in (H^1(\mathbb{R}^N))^{a_h - a_{h-1}} : \mathbf{v} \neq \mathbf{0} \text{ and } \langle \nabla E_h(\mathbf{v}), \mathbf{v} \rangle = 0 \right\}.$$

We set

$$l_h := \inf_{\mathcal{M}_h} E_h.$$

The strategy consists in showing that l , defined in (1.9), coincides with the sum of the least energy levels l_h of the uncoupled sub-systems (5.1). This is inspired by Theorem 1 in [15], which is a particular case of our Theorem 1.6 (for $m = d$, $\mathcal{K}_1 = \emptyset$). We point out that our proof present substantial differences with respect to the one in [15], referring to the forthcoming Remark 5.6 for more details.

Before proceeding, we need the following preliminary result. Although it is essentially known by the community, we present a short proof of it here, as we were not able to find any reference.

Lemma 5.1. *Let $(u_i)_{i \in I_h}$ be a nonnegative solution of (5.1) with $u_i \in H^1(\mathbb{R}^N)$. Then for each $0 < \beta < \min_i \{\lambda_i\}$ there exists $\alpha > 0$ such that*

$$|u_i(x)| \leq \alpha e^{-\sqrt{1+\beta}|x|^2}, \quad \forall x \in \mathbb{R}^N, i \in I_h.$$

Proof. By a Brezis–Kato type argument, one has that $u_i \in L^\infty(\mathbb{R}^N)$. Thus, by standard gradient estimates for Poisson’s equation (see [13, eq. (3.15)]), we deduce also that $\nabla u_i \in L^\infty(\mathbb{R}^N)$. Since we also assume that $u_i \in L^2(\mathbb{R}^N)$, this clearly implies that $u_i \rightarrow 0$ as $|x| \rightarrow \infty$.

Defining $z(x) = \alpha \exp(-\sqrt{1 + \beta}|x|^2)$, a straightforward computation gives

$$-\Delta z + \beta z \geq \frac{\alpha\beta}{1 + \beta|x|^2} e^{-\sqrt{1+\beta}|x|^2}.$$

The difference $z - u_i$ satisfies:

$$-\Delta(z - u_i) + \beta(z - u_i) \geq (\lambda_i - \beta)u_i - \sum_{j \in I_h} \beta_{ij} u_i u_j^2 + \frac{\alpha\beta}{1 + \beta|x|^2} e^{-\sqrt{1+\beta}|x|^2},$$

where $k = |I_h|$. Fix any $0 < \beta < \lambda_i$. Since $\sum_{j \in I_h} \beta_{ij} u_i u_j^2 \rightarrow 0$ as $|x| \rightarrow \infty$, we can take $R > 0$ such that

$$\sum_{j \in I_h} \beta_{ij} u_i u_j^2 \leq (\lambda_i - \beta)u_i, \quad \text{for } |x| \geq R,$$

which implies that $-\Delta(z - u_i) + \beta(z - u_i) \geq 0$ for $|x| \geq R$. On the other hand, there exists $C > 0$ and a sufficiently large $\alpha > 0$, such that

$$\sum_{j \in I_h} \beta_{ij} u_i u_j^2 \leq C \leq \frac{\alpha\beta}{1 + \beta R^2} e^{-\sqrt{1+\beta}R^2} \leq \frac{\alpha\beta}{1 + \beta|x|^2} e^{-\sqrt{1+\beta}|x|^2} \quad \text{for } |x| \leq R.$$

To sum up, we show that it is possible to choose $\alpha > 0$ in such a way that

$$-\Delta(z - u_i) + \beta(z - u_i) \geq 0 \quad \text{in } \mathbb{R}^N,$$

and testing the inequality with $(z - u_i)^-$, we deduce that $u_i \leq z$. \square

For $h = 1, \dots, m$, we introduce

$$\begin{aligned} \tilde{E}_h(\mathbf{v}) &:= \frac{1}{4} \sum_{i \in I_h} \int_{\Omega} (|\nabla v_i|^2 + \lambda_i v_i^2) = \frac{1}{4} \|\mathbf{v}\|_h^2 \\ \tilde{\mathcal{M}}_h &:= \left\{ \mathbf{v} : \mathbf{v} \neq \mathbf{0} \text{ and } \|\mathbf{v}\|_h^2 \leq \sum_{(i,j) \in I_h^2 \mathbb{R}^N} \int \beta_{ij} v_i^2 v_j^2 \right\} \\ \tilde{l}_h &:= \inf_{\tilde{\mathcal{M}}_h} \tilde{E}_h. \end{aligned}$$

Lemma 5.2. Both l_h and \tilde{l}_h are achieved, $l_h = \tilde{l}_h$, and any minimizer for \tilde{l}_h is a minimizer for l_h .

Proof. Since $\mathcal{M}_h \subset \tilde{\mathcal{M}}_h$ and

$$E_h(\mathbf{v}) = \frac{1}{4} \|\mathbf{v}\|_h^2 = \tilde{E}_h(\mathbf{v}) \quad \forall \mathbf{v} \in \mathcal{M}_h,$$

we have $\tilde{l}_h \leq l_h$. As far as \tilde{l}_h is concerned, we start by observing that, if $\mathbf{v} \in \tilde{\mathcal{M}}_h$, then also its Schwarz symmetrization $\mathbf{v}^* \in \tilde{\mathcal{M}}_h$, and by the Polya–Szegő inequality $\|\mathbf{v}^*\|_h^2 \leq \|\mathbf{v}\|_h^2$. Therefore,

$$\tilde{l}_h = \inf \{ \tilde{E}_h(\mathbf{v}) : \mathbf{v} \in \tilde{\mathcal{M}}_h, \mathbf{v} \text{ is radial} \}.$$

The functional \tilde{E}_h is coercive in $H_{\text{rad}}^1(\mathbb{R}^N)$, so that any minimizing sequence is bounded from above. Reasoning exactly as in Lemma 2.4, any such sequence is also bounded from below. This permits immediately to obtain the existence of a minimizer for \tilde{l}_h (in this step it is used the fact that $H_{\text{rad}}^1(\mathbb{R}^N)$ compactly embeds into $L^4(\mathbb{R}^N)$ for $N = 2, 3$). To complete the proof, we show that if $\tilde{\mathbf{v}}$ is a minimizer for \tilde{l}_h , then $\tilde{\mathbf{v}} \in \mathcal{M}_h$ and $\tilde{E}_h(\tilde{\mathbf{v}}) = l_h$. Let $\Psi(t) := E_h(\sqrt{t}\tilde{\mathbf{v}})$. By definition we have that

$$t > 0 \quad \text{and} \quad \Psi'(t) = 0 \quad \iff \quad \sqrt{t}\tilde{\mathbf{v}} \in \tilde{\mathcal{M}}_h.$$

By direct computations, it is easy to check that the unique positive critical point of Ψ is given by

$$\tilde{t} = \frac{\|\tilde{\mathbf{v}}\|_h^2}{\sum_{(i,j) \in I_h^2} \int_{\mathbb{R}^N} \tilde{v}_i^2 \tilde{v}_j^2} \leq 1,$$

where the last estimate follows by the fact that $\tilde{\mathbf{v}} \in \tilde{\mathcal{M}}_h$. Thus, we have

$$l_h \leq E_h(\sqrt{\tilde{t}}\tilde{\mathbf{v}}) = \frac{1}{4} \|\tilde{\mathbf{v}}\|_h^2 \tilde{t} \leq \frac{1}{4} \|\tilde{\mathbf{v}}\|_h^2 = \tilde{l}_h,$$

which implies $l_h = \tilde{l}_h$, and in turn forces $\tilde{t} = 1$, that is $\tilde{\mathbf{v}} \in \mathcal{M}_h$. \square

From now on, for each $h = 1, \dots, m$, we fix a minimizer \mathbf{v}^h for l_h , hence a non-trivial solution of (5.1). We have the following decay estimate.

Lemma 5.3. *Let $e_1 \neq e_2 \in \mathbb{S}^{N-1}$. Then, whenever $h_1 \neq h_2$,*

$$\lim_{R \rightarrow +\infty} \int_{\mathbb{R}^N} \sum_{(i,j) \in I_{h_1} \times I_{h_2}} \left(v_i^{h_1}(x - Re_1) v_j^{h_2}(x - Re_2) \right)^2 dx = 0.$$

Proof. Since $e_1 \neq e_2$, $R|e_1 - e_2| \rightarrow +\infty$ as $R \rightarrow +\infty$. Recalling from Lemma 5.1 that each v_i^h is exponentially decaying as $|x| \rightarrow +\infty$, the thesis follows easily. \square

In the next lemma we show that the least energy level of the complete d system (1.2), which is denoted by l , can be controlled by the sum of the least energy levels l_h of the sub-systems (5.1).

Lemma 5.4. *In the previous notation, we have $l \leq \sum_{h=1}^m l_h$.*

Proof. First of all, we observe that

$$J(\mathbf{u}) = \sum_{h=1}^m E_h(\mathbf{u}_h) + \sum_{(i,j) \in \mathcal{K}_{2\mathbb{R}^N}} \int \beta_{ij} u_i^2 u_j^2.$$

Let $e_1, e_2, \dots, e_m \in \mathbb{S}^{N-1}$ with $e_i \neq e_j$ for $i \neq j$ be m different directions in \mathbb{R}^N . For $R > 0$, we define \mathbf{u}^R by means of

$$u_i^R(x) := v_i^h(x - Re_h) \quad \text{for } i \in I_h, \text{ for } h = 1, \dots, m.$$

Clearly, by a change of variables we have that

$$E_h(\mathbf{u}_h^R) = l_h \quad \forall R > 0.$$

We aim at solving the linear system in t_1, \dots, t_m

$$\frac{\partial}{\partial t_h} J \left(\sqrt{t_1} \mathbf{u}_1^R, \dots, \sqrt{t_m} \mathbf{u}_m^R \right) = 0 \iff \sum_{k=1}^m M_B(\mathbf{u}^R)_{hk} t_k = \|\mathbf{u}_h^R\|_h^2.$$

We claim that, for every $R \gg 1$ sufficiently large, this system has a solution (t_1^R, \dots, t_m^R) with $0 < t_h^R \rightarrow 1$ as $R \rightarrow +\infty$. Once that this is proved, we deduce that

$$\left(\sqrt{t_1^R} \mathbf{u}_1^R, \dots, \sqrt{t_m^R} \mathbf{u}_m^R \right) \in \mathcal{M}. \tag{5.2}$$

To prove the claim, we observe that by Lemma 5.3 $M_B(\mathbf{u}^R)_{hk} \rightarrow 0$ as $R \rightarrow +\infty$ for every $h \neq k$. Since on the contrary $M_B(\mathbf{u}^R)_{hh}$ is positive and constant in R , we deduce that $M_B(\mathbf{u}^R)$ is strictly

Please cite this article in press as: N. Soave, H. Tavares, New existence and symmetry results for least energy positive solutions of Schrödinger systems with mixed competition and cooperation terms, J. Differential Equations (2016), <http://dx.doi.org/10.1016/j.jde.2016.03.015>

diagonally dominant, and hence invertible, for every R sufficiently large. As a consequence, for any such R we can compute

$$\begin{pmatrix} t_1^R \\ \vdots \\ t_m^R \end{pmatrix} = M_B(\mathbf{u}^R)^{-1} \begin{pmatrix} \|\mathbf{u}_1^R\|_1^2 \\ \vdots \\ \|\mathbf{u}_m^R\|_m^2 \end{pmatrix} = M_B(\mathbf{u}^R)^{-1} \begin{pmatrix} \|\mathbf{v}^1\|_1^2 \\ \vdots \\ \|\mathbf{v}^m\|_m^2 \end{pmatrix}.$$

But, as already observed, $M_B(\mathbf{u}^R)$ is converging to a diagonal matrix, whose diagonal entry in the h -th row is equal to

$$\sum_{(i,j) \in I_h^2 \mathbb{R}^N} \int (v_i^h(x - Re_h)v_j^h(x - Re_h))^2 dx = \sum_{(i,j) \in I_h^2 \mathbb{R}^N} \int (v_i^h v_j^h)^2,$$

so that

$$\lim_{R \rightarrow +\infty} t_h^R = \frac{\|\mathbf{v}_h\|_h^2}{\sum_{(i,j) \in I_h^2 \mathbb{R}^N} \int (v_i^h v_j^h)^2} = 1,$$

where we used the fact that by assumption $\mathbf{v}^h \in \mathcal{M}_h$.

To sum up, we have just showed that for any large R there exists $\mathbf{t}^R \simeq \mathbf{1}$ such that (5.2) holds. Therefore

$$\begin{aligned} l &\leq \lim_{R \rightarrow +\infty} J\left(\sqrt{t_1^R} \mathbf{u}_1^R, \dots, \sqrt{t_m^R} \mathbf{u}_m^R\right) \\ &= \lim_{R \rightarrow +\infty} \sum_{h=1}^m E_h\left(\sqrt{t_h^R} \mathbf{u}_h^R\right) + \lim_{R \rightarrow +\infty} \sum_{h \neq k} M_B\left(\sqrt{t_1^R} \mathbf{u}_1^R, \dots, \sqrt{t_m^R} \mathbf{u}_m^R\right) \\ &= \sum_{h=1}^m E_h\left(\mathbf{v}^h\right) = \sum_{h=1}^m l_h. \quad \square \end{aligned}$$

Now we have to show that the opposite inequality holds.

Lemma 5.5. *There holds $l \geq \sum_{h=1}^m l_h$.*

Proof. Let $\mathbf{u} \in \mathcal{M}$. Thanks to the assumption $\beta_{ij} \leq 0$ for every $(i, j) \in \mathcal{K}_2$, we have

$$0 < \|\mathbf{u}_h\|_h^2 = \sum_{k=1}^m M_B(\mathbf{u})_{hk} \leq M_B(\mathbf{u})_{hh},$$

so that $\mathbf{u}_h \in \widetilde{\mathcal{M}}_h$ for every $h = 1, \dots, m$. As a consequence

$$\frac{1}{4} \|\mathbf{u}_h\|_h^2 \geq \inf_{\mathbf{v} \in \widetilde{\mathcal{M}}_h} \frac{1}{4} \|\mathbf{v}\|_h^2 = \tilde{l}_h = l_h, \quad \text{and} \quad J(\mathbf{u}) = \sum_{h=1}^m \frac{1}{4} \|\mathbf{u}_h\|_h^2 \geq \sum_{h=1}^m l_h. \quad \square \quad (5.3)$$

Conclusion of the proof of Theorem 1.6. By contradiction, we suppose that there exists $\mathbf{u} \in \mathcal{M}$ such that $J(\mathbf{u}) = l$. Thanks to the fact that \mathcal{M} is a natural constraint (cf. Lemma 2.3, which holds also in $H^1(\mathbb{R}^N)$), \mathbf{u} is a solution of (1.2), which we can assume to be nonnegative. By the strong maximum principle and the definition of \mathcal{M} , we deduce that for every h there exists i_h such that $u_{i_h} > 0$ in \mathbb{R}^N . To reach a contradiction, we observe that the first equation in (5.3), together with the fact that $\sum_h l_h = l$, imply that necessarily

$$\frac{1}{4} \|\mathbf{u}_h\|_h^2 = l_h = \tilde{l}_h.$$

Since $\beta_{ij} \leq 0$ for every $(i, j) \in \mathcal{K}_2$, we have $\mathbf{u}_h \in \widetilde{\mathcal{M}}_h$, so that \mathbf{u}_h minimizes $\tilde{l}_h = l_h$. Thus, by the last statement of Lemma 5.2, also $\mathbf{u}_h \in \mathcal{M}_h$, and in particular

$$\|\mathbf{u}_h\|_h^2 = M_B(\mathbf{u})_{hh} \quad \forall h = 1, \dots, m.$$

This is in contradiction with the fact that, since \mathbf{u} was supposed to be in \mathcal{M} , we have

$$\|\mathbf{u}_{h_1}\|_{h_1}^2 = \sum_{k=1}^m M_B(\mathbf{u})_{h_1 k} \leq M_B(\mathbf{u})_{h_1 h_1} + \int_{\mathbb{R}^N} \beta_{i_{h_1} i_{h_2}} u_{i_{h_1}}^2 u_{i_{h_2}}^2 < M_B(\mathbf{u})_{h_1 h_1},$$

(for the reader's convenience we recall that h_1 and h_2 have been introduced in the assumptions of Theorem 1.6). \square

Remark 5.6. A remarkable fact, which marks a significant difference in our proof with respect to that of Theorem 1 in [15], is that we do not assume that each l_h admits a unique minimizer, nor that it is achieved by a unique positive solution, of the form $u_i(x) = \alpha_i w(\sigma_i x)$ (with w the unique positive radially decreasing solution of the related single equation problem). This was used in [15]. Instead, our argument is only based upon the decay estimate provided by Lemma 5.1.

Concerning the uniqueness of ground states, although sufficient conditions that imply uniqueness are already known in the literature, it is still an open problem to completely determine the range of parameters for which completely cooperative systems as (5.1) have a unique solution, corresponding to the least energy positive level. This is known in the 2 component case with $\lambda_1 = \lambda_2$ and $\beta_{12} > \max\{\beta_{11}, \beta_{22}\}$, see [31]. For systems of more than 2 components, we refer to Theorem 1.1 in [17], where it is shown in particular that if $\lambda_i \equiv \lambda$ and β_{ij} , $i \neq j$ are large and satisfy additional technical assumptions, then any sub-system (5.1) has a unique least energy positive solution. The results in [19, Section 2], in [25, Section 4], and [4, Proposition 2.1] suggest that uniqueness should hold also in more general situations.

Acknowledgments

H. Tavares is supported by Fundação para a Ciência e Tecnologia through the program *Investigador FCT* and the project PEst-OE/EEI/LA0009/2013. N. Soave and H. Tavares are partially supported through the project ERC Advanced Grant 2013 n. 339958 “Complex Patterns for Strongly Interacting Dynamical Systems – COMPAT”.

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