



ELSEVIER

Contents lists available at SciVerse ScienceDirect

# Journal of Differential Equations

[www.elsevier.com/locate/jde](http://www.elsevier.com/locate/jde)



## On the compatibility condition for linear systems and a factorization formula for wave functions

Alexander Sakhnovich

Fakultät für Mathematik, Universität Wien, Nordbergstrasse 15, A-1090 Vienna, Austria

### ARTICLE INFO

*Article history:*

Received 10 July 2011

Revised 1 October 2011

Available online 14 November 2011

*Keywords:*

Compatibility

Initial-boundary value problem

Wave function

Fundamental solution

Factorization

Inverse spectral transform

Integrable equation

### ABSTRACT

The well-known compatibility condition for linear systems  $w_x = Gw$  and  $w_t = Fw$  is considered and new results are obtained. In this way, a factorization formula for wave functions, which is basic in the inverse spectral transform approach to initial-boundary value problems, is proved in greater generality than before. Some applications follow.

© 2011 Elsevier Inc. All rights reserved.

### 1. Introduction

Zero curvature representation of the integrable nonlinear equations is a well-known approach (see [1,11,21,38] and references in [11]), which was developed soon after seminal Lax pairs appeared in [19]. Namely, many integrable nonlinear equations admit the representation (zero curvature representation)

$$G_t(x, t, z) - F_x(x, t, z) + [G(x, t, z), F(x, t, z)] = 0, \tag{1.1}$$

$$G_t := \frac{\partial}{\partial t} G, \quad [G, F] := GF - FG,$$

which is the compatibility condition of the auxiliary linear systems

$$\frac{\partial}{\partial x} w(x, t, z) = G(x, t, z)w(x, t, z), \quad \frac{\partial}{\partial t} w(x, t, z) = F(x, t, z)w(x, t, z). \tag{1.2}$$

E-mail address: [Oleksandr.Sakhnovych@univie.ac.at](mailto:Oleksandr.Sakhnovych@univie.ac.at).

Here  $G$  and  $F$  are  $m \times m$  matrix functions, and  $z$  is the spectral parameter, which will be omitted sometimes in our notations.

Solution of the integrable nonlinear equations is closely related to the Lax pairs and zero curvature representations, which have been mentioned above, and has been a great breakthrough in the second half of the 20th century. An active study of the cases, which are close to integrable in a certain sense, followed (see, e.g., some references in [3,16]). Initial–boundary value problems for the integrable nonlinear equations can be considered as an important example, where integrability is “spoiled” by the boundary conditions. These problems are of great current interest, and the inverse spectral transform (ISpT) method [4,5,14,29,30] is one of the fruitful approaches in this domain. In particular, some further developments of the results and methods from [29,30] are given in [22,23,25,26,31,32].

Below we assume that  $x, t$  belong to a semi-strip

$$\mathcal{D} = \{(x, t): 0 \leq x < \infty, 0 \leq t < a\}. \tag{1.3}$$

We normalize fundamental solutions of the auxiliary systems by the initial conditions

$$\frac{d}{dx}W(x, t, z) = G(x, t, z)W(x, t, z), \quad W(0, t, z) = I_m; \tag{1.4}$$

$$\frac{d}{dt}R(x, t, z) = F(x, t, z)R(x, t, z), \quad R(x, 0, z) = I_m, \tag{1.5}$$

where  $I_m$  is the identity matrix of order  $m$ . If condition (1.1) holds, the fundamental solution of (1.4) admits factorization

$$W(x, t, z) = R(x, t, z)W(x, 0, z)R(t, z)^{-1}, \quad R(t, z) := R(0, t, z). \tag{1.6}$$

Formula (1.6) is one of the basic and actively used formulas in the inverse spectral transform method (see [22,23,25,26,29–32] and references therein). It was derived in [29,30] under some smoothness conditions (continuous differentiability of  $G$  and  $F$ , in particular): see formulas (1.6) in [29, p. 22] and in [30, p. 39].

Here we prove (1.6) under weaker conditions and in much greater detail, which is important for applications. Namely, we prove the following theorem.

**Theorem 1.1.** *Let  $m \times m$  matrix functions  $G$  and  $F$  and their derivatives  $G_t$  and  $F_x$  exist on the semi-strip  $\mathcal{D}$ , let  $G, G_t$ , and  $F$  be continuous with respect to  $x$  and  $t$  on  $\mathcal{D}$ , and let (1.1) hold. Then the equality*

$$W(x, t, z)R(t, z) = R(x, t, z)W(x, 0, z), \quad R(t, z) := R(0, t, z), \tag{1.7}$$

holds.

Constructions similar to (1.6) appear also in the theory of Knizhnik–Zamolodchikov equation (see Theorem 3.1 in [34] and see also [33]).

We note that the solvability of system (1.2) in the domain  $\mathcal{D}$  is of independent interest as one of the well-posedness and compatibility problems in domains with a boundary. The well-posedness of initial and initial–boundary value problems is a difficult area, which is actively studied (see, e.g., recent works [6,7,18] and references therein). Formula (1.7) implies that the condition (1.1) is, indeed, the compatibility condition and  $w(x, t, z) = W(x, t, z)R(t, z) = R(x, t, z)W(x, 0, z)$  satisfies (1.2).

Theorem 1.1 is proved in Section 2. Section 3 is dedicated to applications to initial–boundary value problems, and Theorem 3.2 on the evolution of the Weyl function for the “focusing” modified Korteweg–de Vries (mKdV) equation is proved there as an example.

As usual, by  $\mathbb{N}$  we denote the set of positive integers, by  $\mathbb{C}$  we denote the complex plane, and by  $\mathbb{C}^m$  is denoted the  $m$ -dimensional coordinate space over  $\mathbb{C}$ . By  $\Im z$  is denoted the imaginary part of  $z \in \mathbb{C}$ , and  $\arg z$  is the argument of  $z$ . By  $C^k(\mathcal{D})$  we denote the functions and matrix functions, which are  $k$  times continuously differentiable on  $\mathcal{D}$ .

**2. Proof of Theorem 1.1**

The spectral parameter  $z$  is non-essential for the formulation of Theorem 1.1 and for its proof and we shall omit it in this section. We shall need the proposition below.

**Proposition 2.1.** *Let the  $m \times m$  matrix function  $W$  be given on the semi-strip  $\mathcal{D}$  by Eq. (1.4), where  $G(x, t)$  and  $G_t(x, t)$  are continuous matrix functions in  $x$  and  $t$ .*

(i) *Then the derivative  $W_t$  exists and matrix functions  $W$  and  $W_t$  are continuous with respect to  $x$  and  $t$  on the semi-strip  $\mathcal{D}$ .*

(ii) *Moreover, the mixed derivative  $W_{tx}$  exists and the equality  $W_{tx} = W_{xt}$  holds on  $\mathcal{D}$ .*

**Proof.** Consider system

$$\frac{d}{dx}y = \widehat{G}(x, y), \quad \widehat{G}(x, y) = \widehat{G}(x, y_{m+1}) := \begin{bmatrix} G(x, y_{m+1}) & 0 \\ 0 & 0 \end{bmatrix}, \tag{2.1}$$

where  $\widehat{G}$  is an  $(m + 1) \times (m + 1)$  matrix function and  $y_{m+1}$  is the last entry of the column vector  $y \in \mathbb{C}^{m+1}$ . Denote by  $W_j$  and  $e_j$  the  $j$ -th columns of  $W$  and  $I_m$ , respectively ( $1 \leq j \leq m$ ). It easily follows from (1.4) that the solution of (2.1) with the initial condition

$$y(0) = g = \begin{bmatrix} e_j \\ t \end{bmatrix} \tag{2.2}$$

has the form

$$y(x, g) = \begin{bmatrix} W_j(x, t) \\ t \end{bmatrix}. \tag{2.3}$$

Putting  $G(x, t) = G(0, t)$  for  $-\varepsilon \leq x \leq 0$  whereas  $t \geq 0$ , and putting  $G(x, t) = G(x, 0) + tG_t(x, 0)$  for  $-\varepsilon \leq t \leq 0$  ( $\varepsilon > 0$ ) we extend  $G$  so that  $G$  and  $G_t$  remain continuous on the rectangles

$$\mathcal{D}(a_1, a_2) = \{(x, t): -\varepsilon \leq x \leq a_1, -\varepsilon \leq t \leq a_2 < a\}, \quad a_1, a_2 \in \mathbb{R}_+. \tag{2.4}$$

Hence, it follows from the definition of  $\widehat{G}$  in (2.1) that  $\widehat{G}(x, y)$  and, as a consequence, the vector function  $\widehat{G}(x, y)y$  are continuous on  $\mathcal{D}(a_1, a_2)$  together with their derivatives with respect to the entries of  $y$ . Thus, according to the classical theory of ordinary differential equations (see, for instance, theorem on pp. 305–306 in [36]) the partial first derivatives of  $y(x, g)$  with respect to the entries of  $g$  exist in the interior  $\mathcal{D}_i(a_1, a_2)$  of  $\mathcal{D}(a_1, a_2)$ . Moreover,  $y$  and its partial derivatives with respect to the entries of  $g$  are continuous. In particular, since by (2.2) we have  $g_{m+1} = t$ , the functions  $y$  and  $y_t$  are continuous in all rectangles  $\mathcal{D}_i(a_1, a_2)$ . Taking into account (2.3), we see that  $W$  and  $W_t$  are continuous in the rectangles  $\mathcal{D}_i(a_1, a_2)$ , and the statement (i) is proved.

In view of (1.4) and the considerations above the derivatives  $W_x$ ,  $W_{xt}$ , and  $W_t$  exist and are continuous in the rectangles  $\mathcal{D}_i(a_1, a_2)$ . Hence, by a stronger formulation (see, e.g., the notes [2,35] or the book [20, p. 201]) of the well-known theorem on mixed derivatives,  $W_{tx}$  exists in  $\mathcal{D}_i(a_1, a_2)$  and  $W_{tx} = W_{xt}$ . Thus, the statement (ii) is valid.  $\square$

Now, we can follow the scheme from [30, Chapter 3] (see also [32, Chapter 12]).

**Proof of Theorem 1.1.** According to statement (i) in Proposition 2.1 the matrix function  $W_t$  exists and is continuous. Introduce  $U(x, t)$  by the equality

$$U := W_t - FW. \tag{2.5}$$

By (1.4), (2.5), and statement (ii) in Proposition 2.1 we have

$$U_x = W_{tx} - F_x W - F W_x = W_{xt} - F_x W - F G W. \quad (2.6)$$

It is immediate also from (1.4) that

$$W_{xt} = (G W)_t = G_t W + G W_t. \quad (2.7)$$

Formulas (2.6) and (2.7) imply

$$U_x = G_t W + G W_t - F_x W - F G W = (G_t - F_x + G F - F G) W + G W_t - G F W. \quad (2.8)$$

From (1.1), (2.8), and definition (2.5) we see that  $U_x = G U$ , that is,  $U$  and  $W$  satisfy the same equation. Taking into account  $W(0, t) = I_m$ , we derive  $W_t(0, t) = 0$ , and so by (2.5) we have  $U(0, t) = -F(0, t)$ . Finally, as

$$U_x = G U, \quad W_x = G W, \quad U(0, t) = -F(0, t), \quad W(0, t) = I_m,$$

we have  $U(x, t) = -W(x, t)F(0, t)$  or, equivalently,

$$W_t(x, t) - F(x, t)W(x, t) = -W(x, t)F(0, t). \quad (2.9)$$

Put

$$Y(x, t) = W(x, t)R(t), \quad Z(x, t) = R(x, t)W(x, 0). \quad (2.10)$$

Recall that  $R(t) = R(0, t)$ . Therefore (1.5), (2.9), and (2.10) imply that

$$\begin{aligned} Y_t(x, t) &= (F(x, t)W(x, t) - W(x, t)F(0, t))R(t) + W(x, t)F(0, t)R(t) \\ &= F(x, t)Y(x, t), \quad Y(x, 0) = W(x, 0). \end{aligned} \quad (2.11)$$

Formulas (1.5) and (2.10) imply that

$$Z_t(x, t) = F(x, t)Z(x, t), \quad Z(x, 0) = W(x, 0). \quad (2.12)$$

By (2.11) and (2.12)  $Y = Z$ , that is, (1.7) holds. The theorem is proved.  $\square$

**Remark 2.2.** Though the case of continuous  $F$  is more convenient for applications, it is immediate from the proof that the statement of Theorem 1.1 is true, when  $F$  is differentiable with respect to  $x$ , and measurable and summable with respect to  $t$  on all finite intervals from  $\mathbb{R}_+$ .

According to the proof of Theorem 1.1 the following remark is also true.

**Remark 2.3.** Theorem 1.1 holds on the domains more general than  $\mathcal{D}$ . In particular, it holds if we consider  $(x, t) \in \mathcal{I}_1 \times \mathcal{I}_2$ , where  $\mathcal{I}_k$  ( $k = 1, 2$ ) is the interval  $[0, b_k)$  ( $0 < b_k \leq \infty$ ).

Another interesting case of matrix factorizations related to boundary value problems is treated in [8,15].

### 3. Some applications

The matrix “focusing” mKdV equation has the form

$$4v_t = v_{xxx} + 3(v_x v^* v + v v^* v_x), \tag{3.1}$$

where  $v(x, t)$  is a  $p \times p$  matrix function. Eq. (3.1) is equivalent (see [9,11,37] and references therein) to zero curvature equation (1.1), where the  $m \times m$  ( $m = 2p$ ) matrix functions  $G(x, t, z)$  and  $F(x, t, z)$  are given by the formulas

$$G = izj + V, \quad j = \begin{bmatrix} I_p & 0 \\ 0 & -I_p \end{bmatrix}, \quad V = \begin{bmatrix} 0 & v \\ -v^* & 0 \end{bmatrix}, \tag{3.2}$$

$$F = -iz^3 j - z^2 V - \frac{iz}{2}(V^2 + V_x j) + \frac{1}{4}(V_{xx} - 2V^3 - V_x V + V V_x). \tag{3.3}$$

At first we omit the variable  $t$  in  $V$  and  $v$ . The Weyl theory of the skew-self-adjoint Dirac system (also called Zakharov–Shabat or AKNS system)

$$\frac{d}{dx} w(x, z) = (izj + V(x)) w(x, z), \quad x \geq 0 \tag{3.4}$$

was treated in [10,12,22,23] (see also preliminaries in [28]).

For the case of measurable matrix function  $v$  such that

$$\sup_{0 < x < \infty} \|v(x)\| \leq M, \tag{3.5}$$

the Weyl matrix function  $\varphi$  of system (3.4) is uniquely defined in the semi-plane  $\Im z < -M$  by the inequality

$$\int_0^\infty [\varphi(z)^* \quad I_p] W(x, z)^* W(x, z) \begin{bmatrix} \varphi(z) \\ I_p \end{bmatrix} dx < \infty, \quad \Im z < -M < 0, \tag{3.6}$$

where  $W$  is the fundamental solution of (3.4) normalized by  $W(0, z) = I_m$ . Weyl functions are constructed using pairs of meromorphic  $p \times p$  matrix functions  $P_1(z), P_2(z)$ , which are nonsingular and have property- $j$ , that is,

$$\mathcal{P}(z)^* \mathcal{P}(z) > 0, \quad \mathcal{P}(z)^* j \mathcal{P}(z) \leq 0, \quad \mathcal{P} := \begin{bmatrix} P_1 \\ P_2 \end{bmatrix}. \tag{3.7}$$

**Theorem 3.1.** (See [22].) *The above system (3.4) with  $v$  satisfying (3.5) has a unique Weyl function. This Weyl function is holomorphic in the semi-plane  $\Im z < -M$ . It is given by the equality*

$$\begin{aligned} \varphi(z) &= \lim_{r \rightarrow \infty} (\mathcal{A}_{11}(r, z) P_1(r, z) + \mathcal{A}_{12}(r, z) P_2(r, z)) \\ &\quad \times (\mathcal{A}_{21}(r, z) P_1(r, z) + \mathcal{A}_{22}(r, z) P_2(r, z))^{-1} \quad (\Im z < -M), \end{aligned} \tag{3.8}$$

$$\mathcal{A}(r, z) = \{\mathcal{A}_{kp}(r, z)\}_{k,p=1}^2 := W(r, \bar{z})^*, \tag{3.9}$$

where the pairs  $\{P_1, P_2\}$  are arbitrary pairs satisfying (3.7) and  $W$  is the fundamental solution of (3.4) normalized by  $W(0, z) = I_m$ .

Our next theorem on the evolution of the Weyl function in the case of the focusing mKdV follows from Theorems 1.1 and 3.1. The case of the defocusing mKdV was earlier treated in [29,30,32].

**Theorem 3.2.** *Let a  $p \times p$  matrix function  $v \in C^1(\mathcal{D})$  have a continuous partial second derivative  $v_{xx}$ , and let  $v_{xxx}$  exist. Assume that  $v$  satisfies mKdV (3.1) and that the inequalities*

$$\sup_{(x,t) \in \mathcal{D}} \|v(x,t)\| \leq M, \quad \sup_{(x,t) \in \mathcal{D}} (\|v_x(x,t)\| + \|v_{xx}(x,t)\|) < \infty \tag{3.10}$$

hold.

Then the evolution  $\varphi(t, z)$  of the Weyl function of the skew-self-adjoint Dirac system (1.4), where  $G$  has the form (3.2), is given by the equality

$$\varphi(t, z) = (R_{11}(t, z)\varphi(0, z) + R_{12}(t, z))(R_{21}(t, z)\varphi(0, z) + R_{22}(t, z))^{-1} \tag{3.11}$$

in the semi-plane  $\Im z < -M < 0$ . Here the block matrix function

$$R(t, z) = \{R_{kn}(t, z)\}_{k,n=1}^2 = R(0, t, z) \tag{3.12}$$

is defined by the boundary values  $v(0, t)$ ,  $v_x(0, t)$ , and  $v_{xx}(0, t)$  via formulas (1.5) and (3.3).

**Proof.** As  $V^* = -V$  and  $(V_x j)^* = V_x j$ , it is immediate from (3.3) that  $F(x, t, \bar{z})^* + F(x, t, z) = 0$ . Hence, it follows from (1.5) that

$$\frac{\partial}{\partial t} (R(x, t, \bar{z})^* R(x, t, z)) = 0.$$

Therefore, using equalities  $R(x, 0, z) = I_m$  and (3.12), we get

$$R(x, t, \bar{z})^* R(x, t, z) = I_m, \quad R(t, \bar{z})^* R(t, z) = I_m,$$

or, equivalently,

$$R(x, t, \bar{z})^* = R(x, t, z)^{-1}, \quad R(t, \bar{z})^* = R(t, z)^{-1}. \tag{3.13}$$

Because of the smoothness conditions on  $v$ , we see that  $G$  and  $F$  given by (3.2) and (3.3), respectively, satisfy the requirements of Theorem 1.1, that is, (1.7) holds. In view of (3.13) rewrite (1.7) in the form

$$\mathcal{A}(x, t, z) R(x, t, z) = R(t, z) \mathcal{A}(x, 0, z), \tag{3.14}$$

where  $\mathcal{A}(x, t, z) := W(x, t, \bar{z})^*$  (compare with (3.9)). Let  $\mathcal{P}(x, z)$  satisfy (3.7) and put

$$\tilde{\mathcal{P}}(x, t, z) = \begin{bmatrix} \tilde{P}_1(x, t, z) \\ \tilde{P}_2(x, t, z) \end{bmatrix} := R(x, t, z) \mathcal{P}(x, z). \tag{3.15}$$

By (3.14) and (3.15) we have

$$\mathcal{A}(x, t, z) \tilde{\mathcal{P}}(x, t, z) = R(t, z) \mathcal{A}(x, 0, z) \mathcal{P}(x, z). \tag{3.16}$$

Now, taking into account that  $\mathcal{P}(x, z)$  is a nonsingular pair with property- $j$ , we show that  $\tilde{\mathcal{P}}(x, t, z)$  is a nonsingular pair with property- $j$  too. According to (1.5), (3.3), and (3.10) we get

$$\frac{\partial}{\partial t} (R(x, t, z)^* j R(x, t, z)) = R(x, t, z)^* (i(\bar{z}^3 - z^3) I_m + O(z^2)) R(x, t, z) \tag{3.17}$$

for  $z \rightarrow \infty$ . Formula (3.17) implies that for some

$$M_1 > M > 0 \quad \left( M \geq \sup_{(x,t) \in \mathcal{D}} \|v(x,t)\| \right), \quad (3.18)$$

and for all  $z$  from the domain

$$\mathcal{D}_1 = \{z: z \in \mathbb{C}, \Im z < -M_1, 0 > \arg z > -\pi/4\} \quad (3.19)$$

we have

$$\frac{\partial}{\partial t} (R(x,t,z)^* j R(x,t,z)) \leq 0,$$

and so

$$R(x,t,z)^* j R(x,t,z) \leq j. \quad (3.20)$$

Relations (3.7), (3.15), and (3.20) imply that

$$\tilde{\mathcal{P}}(x,t,z)^* \tilde{\mathcal{P}}(x,t,z) > 0, \quad \tilde{\mathcal{P}}(x,t,z)^* j \tilde{\mathcal{P}}(x,t,z) \leq 0 \quad (z \in \mathcal{D}_1). \quad (3.21)$$

Clearly, it suffices to prove (3.11) for values of  $z$  from  $\mathcal{D}_1$ . (According to (3.18) and (3.19) the domain  $\mathcal{D}_1$  belongs to the semi-plane  $\Im z < -M$ .)

In a way similar to the proofs of (3.13) and (3.20) we derive

$$\mathcal{A}(x,t,z) = W(x,t,z)^{-1}, \quad W(x,t,z)^* j W(x,t,z) \geq j \quad (\Im z < -M). \quad (3.22)$$

It is immediate from (3.22) that

$$\mathcal{A}(x,t,z)^* j \mathcal{A}(x,t,z) \leq j \quad (\Im z < -M). \quad (3.23)$$

Hence, inequalities (3.7) and (3.21) imply

$$\det(\mathcal{A}_{21}(x,0,z)P_1(x,z) + \mathcal{A}_{22}(x,0,z)P_2(x,z)) \neq 0 \quad (\Im z < -M), \quad (3.24)$$

$$\det(\mathcal{A}_{21}(x,t,z)\tilde{\mathcal{P}}_1(x,t,z) + \mathcal{A}_{22}(x,t,z)\tilde{\mathcal{P}}_2(x,t,z)) \neq 0 \quad (z \in \mathcal{D}_1). \quad (3.25)$$

In view of (3.24) rewrite (3.16) as

$$\mathcal{A}(x,t,z)\tilde{\mathcal{P}}(x,t,z) = R(t,z) \begin{bmatrix} \phi(x,0,z) \\ I_p \end{bmatrix} (\mathcal{A}_{21}(x,0,z)P_1(x,z) + \mathcal{A}_{22}(x,0,z)P_2(x,z)), \quad (3.26)$$

$$\begin{aligned} \phi(x,0,z) &:= (\mathcal{A}_{11}(x,0,z)P_1(x,z) + \mathcal{A}_{12}(x,0,z)P_2(x,z)) \\ &\quad \times (\mathcal{A}_{21}(x,0,z)P_1(x,z) + \mathcal{A}_{22}(x,0,z)P_2(x,z))^{-1}. \end{aligned} \quad (3.27)$$

According to (3.24)–(3.26) we get

$$\begin{aligned} &(\mathcal{A}_{11}(x,t,z)\tilde{\mathcal{P}}_1(x,t,z) + \mathcal{A}_{12}(x,t,z)\tilde{\mathcal{P}}_2(x,t,z)) \\ &\quad \times (\mathcal{A}_{21}(x,t,z)\tilde{\mathcal{P}}_1(x,t,z) + \mathcal{A}_{22}(x,t,z)\tilde{\mathcal{P}}_2(x,t,z))^{-1} \\ &= (R_{11}(t,z)\phi(x,0,z) + R_{12}(t,z)) (R_{21}(t,z)\phi(x,0,z) + R_{22}(t,z))^{-1}. \end{aligned} \quad (3.28)$$

As  $\tilde{P}(x, t, z)$  satisfies (3.21) for  $z \in \mathcal{D}_1$ , using (3.8) we derive

$$\begin{aligned} \varphi(t, z) &= \lim_{x \rightarrow \infty} (\mathcal{A}_{11}(x, t, z) \tilde{P}_1(x, t, z) + \mathcal{A}_{12}(x, t, z) \tilde{P}_2(x, t, z)) \\ &\quad \times (\mathcal{A}_{21}(x, t, z) \tilde{P}_1(x, t, z) + \mathcal{A}_{22}(x, t, z) \tilde{P}_2(x, t, z))^{-1} \quad (z \in \mathcal{D}_1). \end{aligned} \tag{3.29}$$

In a similar way we derive from (3.8) and (3.27) that

$$\varphi(0, z) = \lim_{x \rightarrow \infty} \phi(x, 0, z) \quad (\Im z < -M). \tag{3.30}$$

Let us show that

$$\det(R_{21}(t, z)\varphi(0, z) + R_{22}(t, z)) \neq 0 \quad (z \in \mathcal{D}_1). \tag{3.31}$$

Indeed, it follows from (3.7), (3.23), and (3.27) that

$$[\phi(x, 0, z)^* \quad I_p] j \begin{bmatrix} \phi(x, 0, z) \\ I_p \end{bmatrix} \leq 0. \tag{3.32}$$

By (3.30) and (3.32) the inequality

$$[\varphi(0, z)^* \quad I_p] j \begin{bmatrix} \varphi(0, z) \\ I_p \end{bmatrix} \leq 0 \tag{3.33}$$

is valid. Finally, inequalities (3.20) and (3.33) imply

$$[\varphi(0, z)^* \quad I_p] R(t, z)^* j R(t, z) \begin{bmatrix} \varphi(0, z) \\ I_p \end{bmatrix} \leq 0 \quad (z \in \mathcal{D}_1). \tag{3.34}$$

It is immediate from (3.34) that (3.31) holds. In view of analyticity of both parts of (3.11), relations (3.28)–(3.31) imply (3.11) in the whole semi-plane  $\Im z < -M$ .  $\square$

In a way similar to [17] and the more general constructions for self-adjoint systems in [30,32] (see also some references therein), one can use structured operators to solve inverse problem for the skew-self-adjoint system (3.4) too. Namely, to recover  $v$ , which satisfies condition (3.5), from the Weyl function  $\varphi$  we use operators  $S_l$  (acting in  $L^2_p(0, l)$ ,  $0 < l < \infty$ ) of the form

$$S_l f = f(x) + \frac{1}{2} \int_0^l \int_{|x-r|}^{x+r} s' \left( \frac{\zeta + x - r}{2} \right) s' \left( \frac{\zeta + r - x}{2} \right)^* d\zeta f(r) dr. \tag{3.35}$$

Here  $s' := \frac{d}{dx}s$ . Below we give the procedure from [22] modified in accordance with [12,23].

First, we recover a  $p \times p$  matrix function  $s(x)$  with the entries from  $L^2(0, l)$  (i.e.,  $s(x) \in L^2_{p \times p}(0, l)$ ) via the Fourier transform. That is, we put

$$s(x) = \frac{i}{2\pi} e^{-\eta x} \text{l.i.m.}_{a \rightarrow \infty} \int_{-a}^a e^{i\xi x} z^{-1} \varphi(z/2) d\xi \quad (z = \xi + i\eta, \eta < -2M), \tag{3.36}$$

the limit l.i.m. being the limit in  $L^2_{p \times p}(0, l)$ . Formula (3.36) has sense for any  $l < \infty$ , and so the matrix function  $s(x)$  is defined on the non-negative real semi-axis  $x \geq 0$ . Moreover,  $s$  is absolutely continuous,

it does not depend on the choice of  $\eta < -2M$ ,  $s'$  is bounded on any finite interval, and  $s(0) = 0$ . To define the operator  $S_l$  we substitute  $s'(x)$  into (3.35).

Next, we denote the  $p \times 2p$  block rows of  $W$  by  $\omega_1$  and  $\omega_2$ :

$$\omega_1(x) = [I_p \quad 0] W(x, 0), \quad \omega_2(x) = [0 \quad I_p] W(x, 0). \quad (3.37)$$

It follows from (3.2) and (3.4) that  $W(x, 0)^* W(x, 0) = I_m$ . Hence, by (3.2), (3.4), and (3.37) we have

$$v(x) = \omega_1'(x) \omega_2(x)^*, \quad (3.38)$$

and  $\omega_1, \omega_2$  satisfy the equalities

$$\omega_1(0) = [I_p \quad 0], \quad \omega_1 \omega_1^* \equiv I_p, \quad \omega_1' \omega_1^* \equiv 0, \quad \omega_1 \omega_2^* \equiv 0. \quad (3.39)$$

It is immediate that  $\omega_1$  is uniquely recovered from  $\omega_2$  using (3.39).

Finally, we obtain  $\omega_2$  via the formula

$$\omega_2(l) = [0 \quad I_p] - \int_0^l (S_l^{-1} s')(x)^* [I_p \quad s(x)] dx \quad (0 < l < \infty), \quad (3.40)$$

where  $S_l^{-1}$  is applied to  $s'$  columnwise.

**Theorem 3.3.** Assume that  $\varphi$  is the Weyl function of system (3.4), where  $j$  and  $V$  have the form (3.2) and  $v$  satisfies (3.5). Then  $v$  is recovered from  $\varphi$  via formulas (3.38)–(3.40), where  $s$  and  $S_l$  are given by equalities (3.35) and (3.36). All the mentioned above relations are well defined and the inequalities  $S_l \geq I$  hold.

Another inverse problem, where condition (3.5) on  $v$  is substituted by a condition on  $\varphi$ , is also solved in [12,23,28] using the same procedure.

**Remark 3.4.** One can apply Theorems 3.2 and 3.3 to recover solutions of mKdV. Theorems on the evolution of the Weyl functions constitute also the first step in proofs of uniqueness and existence of the solutions of nonlinear equations via ISpT method (see, for instance, [28]).

The application of Theorem 1.1 to the GBDT version (see [13,24,27] and references therein) of the Bäcklund–Darboux transformation is of interest and shall be discussed in the next paper.

## Acknowledgments

The work of A.L. Sakhnovich was supported by the Austrian Science Fund (FWF) under Grant No. Y330. The author is grateful to the referee for detailed and useful remarks.

## References

- [1] M.J. Ablowitz, D.J. Kaup, A.C. Newell, H. Segur, The inverse scattering transform – Fourier analysis for nonlinear problems, *Stud. Appl. Math.* 53 (4) (1974) 249–315.
- [2] A. Aksoy, M. Martelli, Mixed partial derivatives and Fubini's theorem, *College Math. J.* 33 (2) (2002) 126–130.
- [3] H. Berestycki, et al. (Eds.), *Perspectives in Nonlinear Partial Differential Equations*, *Contemp. Math.*, vol. 446, Amer. Math. Soc., Providence, RI, 2007.
- [4] Yu.M. Berezanskii, Integration of non-linear difference equations by means of inverse problem technique, *Dokl. Akad. Nauk SSSR* 281 (1) (1985) 16–19.
- [5] Yu.M. Berezanskii, M.I. Gekhtman, Inverse problem of spectral analysis and nonabelian chains of nonlinear equations, *Ukrainian Math. J.* 42 (6) (1990) 645–658.

- [6] J.L. Bona, A.S. Fokas, Initial-boundary value problems for linear and integrable nonlinear dispersive equations, *Nonlinearity* 21 (2008) T195–T203.
- [7] A. Bressan, G. Crasta, B. Piccoli, Well-posedness of the Cauchy problem for  $n \times n$  systems of conservation laws, *Mem. Amer. Math. Soc.* 694 (2000).
- [8] I. Brunner, M. Herbst, W. Lerche, J. Walcher, Matrix factorizations and mirror symmetry: the cubic curve, *J. High Energy Phys.* 11 (2006) 006, 19 pp.
- [9] F. Calogero, A. Degasperis, *Spectral Transform and Solitons*, North-Holland, Amsterdam, New York, 1982.
- [10] F. Clark, F. Gesztesy, On self-adjoint and  $J$ -self-adjoint Dirac-type operators: a case study, *Contemp. Math.* 412 (2006) 103–140.
- [11] L.D. Faddeev, L.A. Takhtajan, *Hamiltonian Methods in the Theory of Solitons*, Springer, New York, 1986.
- [12] B. Fritzsche, B. Kirstein, A.L. Sakhnovich, Semiseparable integral operators and explicit solution of an inverse problem for the skew-self-adjoint Dirac-type system, *Integral Equations Operator Theory* 66 (2010) 231–251.
- [13] I. Gohberg, M.A. Kaashoek, A.L. Sakhnovich, Pseudocanonical systems with rational Weyl functions: explicit formulas and applications, *J. Differential Equations* 146 (1998) 375–398.
- [14] M. Kac, P. van Moerbeke, A complete solution of the periodic Toda problem, *Proc. Natl. Acad. Sci. USA* 72 (1975) 2879–2880.
- [15] A. Kapustin, Y. Li, Topological correlators in Landau–Ginzburg models with boundaries, *Adv. Theor. Math. Phys.* 7 (2004) 727–749.
- [16] D.J. Kaup, The forced Toda lattice: an example of an almost integrable system, *J. Math. Phys.* 25 (2) (1984) 277–281.
- [17] M.G. Krein, Continuous analogues of propositions on polynomials orthogonal on the unit circle, *Dokl. Akad. Nauk SSSR* 105 (1955) 637–640 (in Russian).
- [18] I. Kukavica, R. Temam, V.C. Vicol, M. Ziane, Local existence and uniqueness for the hydrostatic Euler equations on a bounded domain, *J. Differential Equations* 250 (2011) 1719–1746.
- [19] P.D. Lax, Integrals of nonlinear equations of evolution and solitary waves, *Comm. Pure Appl. Math.* 21 (1968) 467–490.
- [20] J.E. Marsden, *Elementary Classical Analysis*, W.H. Freeman and Co., San Francisco, 1974.
- [21] S.P. Novikov, A periodic problem for the Korteweg–de Vries equation, *Funct. Anal. Appl.* 8 (3) (1974) 236–246.
- [22] A.L. Sakhnovich, Nonlinear Schrödinger equation on a semi-axis and an inverse problem associated with it, *Ukrainian Math. J.* 42 (3) (1990) 316–323.
- [23] A.L. Sakhnovich, The Goursat problem for the sine-Gordon equation and the inverse spectral problem, *Russian Math. (Iz. VUZ)* 36 (11) (1992) 42–52.
- [24] A.L. Sakhnovich, Dressing procedure for solutions of nonlinear equations and the method of operator identities, *Inverse Problems* 10 (1994) 699–710.
- [25] A.L. Sakhnovich, Second harmonic generation: Goursat problem on the semi-strip, Weyl functions and explicit solutions, *Inverse Problems* 21 (2005) 703–716.
- [26] A.L. Sakhnovich, Weyl functions, inverse problem and special solutions for the system auxiliary to the nonlinear optics equation, *Inverse Problems* 24 (2008) 025026, 23 pp.
- [27] A.L. Sakhnovich, On the GBDT version of the Bäcklund–Darboux transformation and its applications to the linear and nonlinear equations and spectral theory, *Math. Model. Nat. Phenom.* 5 (4) (2010) 340–389.
- [28] A.L. Sakhnovich, Sine-Gordon theory in a semi-strip, *Nonlinear Anal.* 75 (2012) 964–974.
- [29] L.A. Sakhnovich, Nonlinear equations and inverse problems on the half-axis, Preprint 30 Inst. Mat. AN Ukr. SSR, Izd-vo Inst. Matem. AN Ukr. SSR, Kiev, 1987.
- [30] L.A. Sakhnovich, The method of operator identities and problems in analysis, *St. Petersburg Math. J.* 5 (1) (1994) 1–69.
- [31] L.A. Sakhnovich, *Interpolation Theory and Its Applications*, Dordrecht, Kluwer, 1997.
- [32] L.A. Sakhnovich, *Spectral Theory of Canonical Differential Systems. Method of Operator Identities*, Oper. Theory Adv. Appl., vol. 107, Birkhäuser, Basel, 1999.
- [33] L.A. Sakhnovich, On the solutions of Knizhnik–Zamolodchikov system, *Cent. Eur. J. Math.* 7 (1) (2009) 145–162.
- [34] L.A. Sakhnovich, Rationality of the Knizhnik–Zamolodchikov equation solution, *Theoret. and Math. Phys.* 163 (1) (2010) 472–478.
- [35] R.T. Seeley, Fubini implies Leibniz implies  $F_{yx} = F_{xy}$ , *Amer. Math. Monthly* 68 (1961) 56–57.
- [36] V.I. Smirnov, *A Course of Higher Mathematics*, IV, Pergamon Press, Oxford, New York, 1964.
- [37] M. Wadati, The exact solution of the modified Korteweg–de Vries equation, *J. Phys. Soc. Japan* 32 (1972) 1681.
- [38] V.E. Zakharov, A.V. Mikhailov, Relativistically invariant two-dimensional models of field theory which are integrable by means of the inverse scattering problem method, *Soviet Phys. JETP* 74 (6) (1978) 1953–1973.