



# Global existence for semilinear wave equations with the critical blow-up term in high dimensions <sup>☆</sup>

Hiroyuki Takamura <sup>a,\*</sup>, Kyouhei Wakasa <sup>b,1</sup>

<sup>a</sup> Department of Complex and Intelligent Systems, Faculty of Systems Information Science, Future University Hakodate, 116-2 Kamedanakano-cho, Hakodate, Hokkaido 041-8655, Japan

<sup>b</sup> Department of Mathematics, Hokkaido University, Sapporo 060-0810, Japan

Received 17 June 2015

Available online 15 April 2016

## Abstract

We are interested in almost global existence cases in the general theory for nonlinear wave equations, which are caused by critical exponents of nonlinear terms. Such situations can be found in only three cases in the theory, cubic terms in two space dimensions, quadratic terms in three space dimensions and quadratic terms including a square of unknown functions itself in four space dimensions. Except for the last case, criteria to classify nonlinear terms into the almost global, or global existence case, are well-studied and known to be so-called null condition and non-positive condition.

Our motivation of this work is to find such a kind of the criterion in four space dimensions. In our previous paper, an example of the non-single term for the almost global existence case is introduced. In this paper, we show an example of the global existence case. These two examples have nonlinear integral terms which are closely related to derivative loss due to high dimensions. But it may help us to describe the final form of the criterion.

© 2016 Elsevier Inc. All rights reserved.

MSC: primary 35L71, 35E15; secondary 35A01, 35A09, 35B33, 35B44

<sup>☆</sup> This work is partially supported by the Grant-in-Aid for Scientific Research (C) (No. 24540183), Japan Society for the Promotion of Science.

\* Corresponding author.

E-mail addresses: [takamura@fun.ac.jp](mailto:takamura@fun.ac.jp) (H. Takamura), [wakasa@math.sci.hokudai.ac.jp](mailto:wakasa@math.sci.hokudai.ac.jp), [wakasa@mmm.muroran-it.ac.jp](mailto:wakasa@mmm.muroran-it.ac.jp) (K. Wakasa).

<sup>1</sup> The 2nd year of the doctor course. Current affiliation: College of Liberal Arts, Mathematical Science Research Unit, Muroran Institute of Technology, 27-1, Mizumoto-cho, Muroran, Hokkaido 050-8585, Japan.

**Keywords:** Semilinear wave equation; High dimensions; Critical exponent; Lifespan

## 1. Introduction

First we shall outline the general theory on the initial value problem for fully nonlinear wave equations,

$$\begin{cases} u_{tt} - \Delta u = H(u, Du, D_x Du) & \text{in } \mathbf{R}^n \times [0, \infty), \\ u(x, 0) = \varepsilon f(x), \quad u_t(x, 0) = \varepsilon g(x), \end{cases} \quad (1.1)$$

where  $u = u(x, t)$  is a scalar unknown function of space–time variables,

$$\begin{aligned} Du &= (u_{x_0}, u_{x_1}, \dots, u_{x_n}), \quad x_0 = t, \\ D_x Du &= (u_{x_i x_j}), \quad i, j = 0, 1, \dots, n, \quad i + j \geq 1, \end{aligned}$$

$f, g \in C_0^\infty(\mathbf{R}^n)$  and  $\varepsilon > 0$  is a “small” parameter. We note that it is impossible to construct a general theory for “large”  $\varepsilon$  due to blow-up results. For example, see Glassey [6], Levine [16], or Sideris [20]. Let

$$\widehat{\lambda} = (\lambda; (\lambda_i), i = 0, 1, \dots, n; (\lambda_{ij}), i, j = 0, 1, \dots, n, i + j \geq 1).$$

Suppose that the nonlinear term  $H = H(\widehat{\lambda})$  is a sufficiently smooth function with

$$H(\widehat{\lambda}) = O(|\widehat{\lambda}|^{1+\alpha})$$

in a neighborhood of  $\widehat{\lambda} = 0$ , where  $\alpha \geq 1$  is an integer. Let us define the lifespan  $\widetilde{T}(\varepsilon)$  of classical solutions of (1.1) by

$$\widetilde{T}(\varepsilon) = \sup\{t > 0 : \exists \text{ a classical solution } u(x, t) \text{ of (1.1) for arbitrarily fixed data, } (f, g)\}.$$

When  $\widetilde{T}(\varepsilon) = \infty$ , the problem (1.1) admits a global solution, while we only have a local solution on  $[0, \widetilde{T}(\varepsilon))$  when  $\widetilde{T}(\varepsilon) < \infty$ . For local solutions, one can measure the long time stability of a zero solution by orders of  $\varepsilon$ . Because the uniqueness of the solution of (1.1) may yield that  $\lim_{\varepsilon \rightarrow +0} \widetilde{T}(\varepsilon) = \infty$ . Such an uniqueness theorem can be found in Appendix of John [11] for example.

In Chapter 2 of Li and Chen [18], we have long histories on the estimate for  $\widetilde{T}(\varepsilon)$ . The lower bounds of  $\widetilde{T}(\varepsilon)$  are summarized in the following table. Let  $a = a(\varepsilon)$  satisfy

$$a^2 \varepsilon^2 \log(a + 1) = 1 \quad (1.2)$$

and  $c$  stand for a positive constant independent of  $\varepsilon$ . Then, due to the fact that it is impossible to obtain an  $L^2$  estimate for  $u$  itself by standard energy methods, we have

$\tilde{T}(\varepsilon) \geq$	$\alpha = 1$	$\alpha = 2$	$\alpha \geq 3$
$n = 2$	$c\varepsilon$ in general case, $c\varepsilon^{-1}$ if $\int_{\mathbf{R}^2} g(x)dx = 0$ , $c\varepsilon^{-2}$ if $\partial_u^2 H(0) = 0$	$c\varepsilon^{-6}$ in general case, $c\varepsilon^{-18}$ if $\partial_u^3 H(0) = 0$ , $\exp(c\varepsilon^{-2})$ if $\partial_u^3 H(0) = \partial_u^4 H(0) = 0$	$\infty$
$n = 3$	$c\varepsilon^{-2}$ in general case, $\exp(c\varepsilon^{-1})$ if $\partial_u^2 H(0) = 0$	$\infty$	$\infty$
$n = 4$	$\exp(c\varepsilon^{-2})$ in general case, $\infty$ if $\partial_u^2 H(0) = 0$	$\infty$	$\infty$
$n \geq 5$	$\infty$	$\infty$	$\infty$

The result for  $n = 1$  is that

$$\tilde{T}(\varepsilon) \geq \begin{cases} \varepsilon^{-\alpha/2} & \text{in general case,} \\ c\varepsilon^{-\alpha(1+\alpha)/(2+\alpha)} & \text{if } \int_{\mathbf{R}} g(x)dx = 0, \\ c\varepsilon^{-\alpha} & \text{if } \partial_u^\beta H(0) = 0 \text{ for } 1 + \alpha \leq \forall \beta \leq 2\alpha. \end{cases} \quad (1.3)$$

For references on these results, see Li and Chen [18]. We shall skip to refer them here. But we note that two parts in this table are different from the one in Li and Chen [18]. One is the general case in  $(n, \alpha) = (4, 1)$ . In this part, the lower bound of  $\tilde{T}(\varepsilon)$  is  $\exp(c\varepsilon^{-1})$  in Li and Chen [18]. But later, it has been improved by Li and Zhou [19]. Another is the case for  $\partial_u^3 H(0) = 0$  in  $(n, \alpha) = (2, 2)$ . This part is due to Katayama [13]. But it is missing in Li and Chen [18]. Its reason is closely related to the sharpness of results in the general theory. The sharpness is achieved by the fact that there is no possibility to improve the lower bound of  $\tilde{T}(\varepsilon)$  in sense of order of  $\varepsilon$  by blow-up results for special equations and special data. It is expressed in the upper bound of  $\tilde{T}(\varepsilon)$  with the same order of  $\varepsilon$  as in the lower bound. On this matter, Li and Chen [18] say that all these lower bounds are known to be sharp except for  $(n, \alpha) = (4, 1)$ . But before this article, Li [17] says that  $(n, \alpha) = (2, 2)$  has also open sharpness while the case for  $\partial_u^3 H(0) = 0$  is still missing. Li and Chen [18] might have dropped the open sharpness in  $(n, \alpha) = (2, 2)$  by conjecture that  $\partial_u^4 H(0) = 0$  is a technical condition. No one disagrees with this observation because the model case of  $H = u^4$  has a global solution in two space dimensions,  $n = 2$ . However, Zhou and Han [26] have obtained this final sharpness in  $(n, \alpha) = (2, 2)$  by studying  $H = u_t^2 u + u^4$ . This puts Katayama's result into the table after 20 years from Li and Chen [18]. We note that Godin [7] has showed the sharpness of the case for  $\partial_u^3 H(0) = \partial_u^4 H(0) = 0$  in  $(n, \alpha) = (2, 2)$  by studying  $H = u_t^3$ . This result has been reproved by Zhou and Han [25].

We now turn back to another open sharpness of the general case in  $(n, \alpha) = (4, 1)$ . It has been obtained by our previous work, Takamura and Wakasa [23], by studying model case of  $H = u^2$ . This part had been open more than 20 years in the analysis on the critical case for model equations,  $u_{tt} - \Delta u = |u|^p$  ( $p > 1$ ). In this way, the general theory and its optimality have been completed.

After the completion of the general theory, we are interested in the almost global existence, namely, the case where  $\tilde{T}(\varepsilon)$  has an lower bound of the exponential function of  $\varepsilon$  with a negative power. Such a case appears in  $(n, \alpha) = (2, 2), (3, 1), (4, 1)$  in the table of the general theory.

It is remarkable that Klainerman [14] and Christodoulou [4] have independently found a special structure on  $H = H(Du, D_x Du)$  in  $(n, \alpha) = (3, 1)$  which guarantees the global existence. This algebraic condition on nonlinear terms of derivatives of the unknown function is so-called “null condition”. It has been also established independently by Godin [7] for  $H = H(Du)$  and Katayama [12] for  $H = H(Du, D_x Du)$  in  $(n, \alpha) = (2, 2)$ . The null condition had been supposed to be not sufficient for the global existence in  $(n, \alpha) = (2, 2)$ . For this direction, Agemi [1] proposed “non-positive condition” in this case for  $H = H(Du)$ . This conjecture has been verified by Hoshiga [8] and Kubo [15] independently. It might be necessary and sufficient condition to the global existence. On the other hand, the situation in  $(n, \alpha) = (4, 1)$  is completely different from  $(n, \alpha) = (2, 2), (3, 1)$  because  $H$  has to include  $u^2$ .

In our previous paper [24], we get the first attempt to clarify a criterion on  $H$  guaranteeing the global existence by showing different blow-up example of  $H$  from  $u^2$  only. More precisely, we have an almost global existence and its optimality for an equation of the form

$$u_{tt} - \Delta u = u^2 - \frac{1}{\pi^2} \int_0^t d\tau \int_{|\xi| \leq 1} \frac{(u_t u)(x + (t - \tau)\xi, \tau)}{\sqrt{1 - |\xi|^2}} d\xi - \frac{\varepsilon^2}{2\pi^2} \int_{|\xi| \leq 1} \frac{f(x + t\xi)^2}{\sqrt{1 - |\xi|^2}} d\xi \quad (1.4)$$

in  $\mathbf{R}^4 \times [0, \infty)$ . We note that the third term in the right-hand side of (1.4) can be neglected by simple modification. One can say that this result is the first example of the blowing-up of a classical solution to nonlinear wave equation with non-single and indefinitely signed term in high dimensions. We note that (1.4) arises from a neglect of derivative loss factors in Duhamel’s term for positive and single nonlinear term,  $u^2$ . Therefore one can conclude that derivative loss factors in Duhamel’s term due to high dimensions do not contribute to any order of  $\varepsilon$  in the estimate of the lifespan.

In this paper, we show that, in contrast with (1.4), another equation of the form

$$u_{tt} - \Delta u = u^2 - \frac{1}{2\pi^2} \int_0^t d\tau \int_{|\omega|=1} (u_t u)(x + (t - \tau)\omega, \tau) dS_\omega - \frac{\varepsilon}{4\pi^2} \int_{|\omega|=1} (\varepsilon f^2 + \Delta f + 2\omega \cdot \nabla g)(x + t\omega) dS_\omega \quad (1.5)$$

admits a global classical solution in  $\mathbf{R}^4 \times [0, \infty)$ . Both the first integral terms in (1.4) and (1.5) look similar to each others. The essential difference is that the second integral term in (1.5) has linear terms of the initial data. This part mainly comes from a neglect of derivative loss factors in the linear part. Therefore one may say that derivative loss factors in the linear part due to high dimensions contribute to estimates of the lifespan.

This paper is organized as follows. In the next section, our main theorems are stated in more general situation on space dimensions and nonlinear terms as well as our motivation of this work by some integral equation. In section 3, we investigate a relation between such an integral equation and (1.5). The decay estimate of the linear part is studied in section 4. The proof of the

local existence appears in section 5, for which *a priori* estimate in section 6 is required. In the final section, blow-up result is proved to show the optimality of the local existence.

This work has begun since the second author was in the 2nd year of the master course, Graduate School of Systems Information Science, Future University Hakodate.

## 2. Main results

This work is initiated by Agemi and Takamura [2] which attempts to make a new representation formula of a solution of the following initial value problem for inhomogeneous wave equations.

$$\begin{cases} \partial_t^2 u - \Delta u = F & \text{in } \mathbf{R}^n \times [0, \infty), \\ u(x, 0) = \varepsilon f(x), \quad u_t(x, 0) = \varepsilon g(x), \quad x \in \mathbf{R}^n, \end{cases} \quad (2.1)$$

where  $u = u(x, t)$  is an unknown function,  $f, g$  and  $F = F(x, t)$  are given smooth functions. In [2], it has proved that, for  $n \geq 3$ , a smooth solution of (2.1) has to satisfy the following integral equation.

$$\begin{aligned} (n-2)\omega_n u(x, t) = & \varepsilon \int_{|\omega|=1} \{t\omega \cdot \nabla f + (n-2)f + tg\}(x + t\omega) dS_\omega \\ & + (n-3) \int_0^t d\tau \int_{|\omega|=1} u_t(x + (t-\tau)\omega, \tau) dS_\omega \\ & + \int_0^t (t-\tau) d\tau \int_{|\omega|=1} F(x + (t-\tau)\omega, \tau) dS_\omega, \end{aligned} \quad (2.2)$$

where  $\omega_n$  is a measure of the unit sphere in  $\mathbf{R}^n$ , i.e.

$$\omega_n = \frac{2\pi^{n/2}}{\Gamma(n/2)} = \begin{cases} \frac{2(2\pi)^m}{(2m-1)!!} & \text{for } n = 2m + 1, \\ \frac{2\pi^{m+1}}{m!} & \text{for } n = 2m + 2, \end{cases} \quad (m = 1, 2, 3, \dots).$$

In view of (2.2), neglecting the second term in the right-hand side, we obtain a representation formula of a solution of some wave equation. With a small modification, it may have the same initial data as in (2.1). Our problem arises in this way.

In fact, let us define our integral equation of an unknown function  $u$  by

$$u(x, t) = \varepsilon V(x, t) + N(F)(x, t), \quad (2.3)$$

where

$$V(x, t) = \frac{1}{\omega_n} \int_{|\omega|=1} \left( \frac{t\omega \cdot \nabla f}{n-2} + f + tg \right) (x + t\omega) dS_\omega \quad (2.4)$$

and

$$N(F)(x, t) = \frac{1}{(n-2)\omega_n} \int_0^t (t-\tau) d\tau \int_{|\omega|=1} F(x + (t-\tau)\omega, \tau) dS_\omega. \quad (2.5)$$

Then, we have the following theorem.

**Theorem 2.1.** *Let  $n \geq 3$ . Assume that  $f \in C^3(\mathbf{R}^n)$ ,  $g \in C^2(\mathbf{R}^n)$  and  $F \in C^2(\mathbf{R}^n \times [0, \infty))$ . Then, a solution of the integral equation (2.3) satisfies the following initial value problem for inhomogeneous wave equation.*

$$\begin{cases} \partial_t^2 u(x, t) - \Delta u(x, t) = F(x, t) - H(x, t) & \text{in } \mathbf{R}^n \times [0, \infty), \\ u(x, 0) = \varepsilon f(x), \quad u_t(x, 0) = \varepsilon g(x), & x \in \mathbf{R}^n, \end{cases} \quad (2.6)$$

where  $H$  is defined by

$$\begin{aligned} H(x, t) = & \frac{n-3}{(n-2)\omega_n} \int_0^t d\tau \int_{|\omega|=1} (\partial_t F)(x + (t-\tau)\omega, \tau) dS_\omega \\ & + \frac{n-3}{(n-2)\omega_n} \int_{|\omega|=1} F(x + t\omega, 0) dS_\omega \\ & + \frac{\varepsilon(n-3)}{(n-2)\omega_n} \int_{|\omega|=1} \{\Delta f + (n-2)\omega \cdot \nabla g\}(x + t\omega) dS_\omega. \end{aligned} \quad (2.7)$$

We shall make use of this theorem with  $F(x, t) = u(x, t)^2$  and  $n = 4$ . The proof of this theorem appears in the next section.

**Remark 2.1.** The uniqueness of the solution of (2.6) with  $F(x, t) = |u(x, t)|^p$  ( $p \geq 2$ ) is open. The restricted uniqueness theorem such as in Appendix 1 in John [11] cannot be applicable because (99a) in [11] does not hold for this case.

**Remark 2.2.** It is remarkable that Huygens' principle holds for  $V$  in (2.4) even if the space dimension is even number. See (4.2) below. Moreover, in view of (2.4) and (2.5), we need lower regularities on the data and inhomogeneous term than those from  $H \equiv 0$  to obtain a classical solution.

In order to describe our main theorems, let us define a lifespan  $\widehat{T}(\varepsilon)$  of the integral equation (2.3) by

$$\widehat{T}(\varepsilon) = \sup\{t > 0 : \exists \text{ a solution } u \text{ of (2.3) with } F = F(u) \\ \text{for arbitrarily fixed data, } (f, g)\},$$

where “solution” means a classical solution of (2.6) for  $p \geq 2$ , or the  $C^1$  solution of (2.3) for  $1 < p < 2$ . Our assumption on  $F = F(s)$  is that

$$\begin{cases} \text{there exists a constant } A > 0 \text{ such that } F \in C^1(\mathbf{R}) \text{ satisfies} \\ |F^{(j)}(s)| \leq A|s|^{p-j} \quad (j = 0, 1) \text{ for } s \in \mathbf{R}, \quad 1 < p < 2, \end{cases} \quad (2.8)$$

or

$$\begin{cases} \text{there exists a constant } A > 0 \text{ such that } F \in C^2(\mathbf{R}) \text{ satisfies} \\ |F^{(j)}(s)| \leq A|s|^{p-j} \quad (j = 0, 1, 2) \text{ for } s \in \mathbf{R}, \quad p \geq 2 \end{cases} \quad (2.9)$$

respectively. We also assume on the data that

$$\begin{cases} \text{at least one of } f \in C_0^4(\mathbf{R}^n) \text{ and } g \in C_0^3(\mathbf{R}^n) \text{ does not} \\ \text{vanish identically and have compact support} \\ \text{contained in } \{x \in \mathbf{R}^n : |x| \leq k\} \text{ with some constant } k > 1. \end{cases} \quad (2.10)$$

We now introduce a critical number  $p_1(n)$  as a positive root of the following quadratic equation.

$$\zeta(p, n) \equiv 2 \left( 1 + (n-1)p - (n-2)p^2 \right) = 0. \quad (2.11)$$

This is the analogy to Strauss' number defined by a positive root of  $\gamma(p, n) \equiv 2 + (n+1)p - (n-1)p^2 = 0$ . See Remark 2.3 below.

Then, we have the following lower bounds of the lifespan which mean long time existences of the solution.

**Theorem 2.2.** *Let  $n \geq 3$ . Assume that (2.8), (2.9) and (2.10) are fulfilled. Then there exists a positive constant  $\varepsilon_0 = \varepsilon_0(f, g, n, p, k)$  such that the lifespan  $\widehat{T}(\varepsilon)$  satisfies*

$$\begin{aligned} \widehat{T}(\varepsilon) &= \infty && \text{for } p > p_1(n), \\ \widehat{T}(\varepsilon) &\geq \exp \left( c\varepsilon^{-p(p-1)} \right) && \text{for } p = p_1(n), \\ \widehat{T}(\varepsilon) &\geq c\varepsilon^{-2p(p-1)/\zeta(p, n)} && \text{for } 1 < p < p_1(n) \end{aligned} \quad (2.12)$$

for any  $\varepsilon$  with  $0 < \varepsilon \leq \varepsilon_0$ , where  $c$  is a positive constant independent of  $\varepsilon$ .

**Remark 2.3.** We note that

$$p_1(n) = \frac{n-1 + \sqrt{n^2 + 2n - 7}}{2(n-2)} \leq p_0(n) = \frac{n+1 + \sqrt{n^2 + 10n + 7}}{2(n-1)}$$

and that its equality holds if and only if  $n = 3$ . Here  $p_0(n)$  is Strauss' number on semilinear wave equations,  $u_{tt} - \Delta u = |u|^p$ . See Strauss [21,22] for this number. Also see Takamura and Wakasa [24] for references therein on lifespan estimates for this equation. Therefore the exponent  $(n, p) = (4, 2)$  is in the super critical case for the equation (2.6) with  $F(x, t) = |u(x, t)|^p$ . The key fact is that the linear part  $V$  in (2.4) decays faster than that of a solution of the free wave equation.

For the upper bounds of the lifespan, our assumption on the data is the following.

$$\left\{ \begin{array}{l} \text{Let } f \equiv 0, g(x) = g(|x|) \text{ and } g \in C_0^2([0, \infty)) \text{ satisfy that} \\ \text{(i) } \text{supp } g \subset \{x \in \mathbf{R}^n : |x| \leq k\} \text{ with } k > 0, \\ \text{(ii) there exists } k_0 \text{ such that } g(|x|) > 0 \text{ for } 0 < k_0 < |x| < k. \end{array} \right. \quad (2.13)$$

Then, we have the following theorem.

**Theorem 2.3.** *Let  $n \geq 3$ . Assume that (2.13) is fulfilled. Then there exists a positive constant  $\varepsilon_0 = \varepsilon_0(g, n, p, k)$  such that the lifespan  $\widehat{T}(\varepsilon)$  satisfies*

$$\begin{aligned} \widehat{T}(\varepsilon) &\leq C\varepsilon^{-2p(p-1)/\zeta(p,n)} \quad \text{for } 1 < p < p_1(n), \\ \widehat{T}(\varepsilon) &\leq \exp\left(C\varepsilon^{-p(p-1)}\right) \quad \text{for } p = p_1(n) \end{aligned} \quad (2.14)$$

for any  $\varepsilon$  with  $0 < \varepsilon \leq \varepsilon_0$ , where  $C$  is a positive constant independent of  $\varepsilon$ .

The proofs of both Theorem 2.2 and Theorem 2.3 are similar to those of our previous theorems in [24] which are based on John's iteration argument in a weighted  $L^\infty$  space by John [10]. They are described after the next section.

### 3. Proof of Theorem 2.1

First we shall prove the initial condition in (2.6). It is trivial to get the first condition by setting  $t = 0$  in (2.3). Rewriting

$$\omega \cdot \nabla f(x + t\omega) = \partial_t(f(x + t\omega)),$$

we have that

$$\begin{aligned} u_t(x, t) &= \frac{\varepsilon}{\omega_n} \int_{|\omega|=1} \left\{ \frac{((n-1)\partial_t + t\partial_t^2)f}{n-2} + (1+t\partial_t)g \right\} (x+t\omega) dS_\omega \\ &\quad + \int_0^t d\tau \int_{|\omega|=1} \frac{(1+(t-\tau)\partial_t)F(x+(t-\tau)\omega, \tau)}{(n-2)\omega_n} dS_\omega. \end{aligned} \quad (3.1)$$

Therefore the second condition follows from setting  $t = 0$  in this equation.

For the proof of the equation in (2.6), we shall employ the well-known fact that a function  $M(x, t)$  defined by

$$M(x, t) = \frac{1}{\omega_n} \int_{|\omega|=1} m(x + t\omega) dS_\omega$$

for  $m \in C^2(\mathbf{R}^n)$  satisfies the initial value problem of Darboux equation,

$$\left\{ \begin{array}{l} \left( \partial_t^2 - \Delta + \frac{n-1}{t} \partial_t \right) M = 0 \\ M(x, 0) = m(x), \quad M_t(x, 0) = 0. \end{array} \right. \quad (3.2)$$



Then it follows from (3.1) and (3.2) that

$$u_{tt}(x, t) = \frac{\varepsilon}{\omega_n} \int_{|\omega|=1} \left\{ \frac{(1+t\partial_t)\Delta f}{n-2} + (2\partial_t + t\partial_t^2)g \right\} (x+t\omega) dS_\omega \\ + \int_0^t d\tau \int_{|\omega|=1} \frac{(2\partial_t + (t-\tau)\partial_t^2)F(x+(t-\tau)\omega, \tau)}{(n-2)\omega_n} dS_\omega + \frac{F(x, t)}{n-2}.$$

On the other hand, operating  $\Delta$  to (2.3) yields that

$$\Delta u(x, t) = \frac{\varepsilon}{\omega_n} \int_{|\omega|=1} \left\{ \left( 1 + \frac{t\partial_t}{n-2} \right) \Delta f + t\Delta g \right\} (x+t\omega) dS_\omega \\ + \int_0^t (t-\tau) d\tau \int_{|\omega|=1} \frac{\Delta F(x+(t-\tau)\omega, \tau)}{(n-2)\omega_n} dS_\omega.$$

Therefore, it follows from (3.2) that

$$u_{tt}(x, t) - \Delta u(x, t) = \frac{\varepsilon}{\omega_n} \int_{|\omega|=1} \left\{ \frac{3-n}{n-2} \Delta f + (2\partial_t + t(\partial_t^2 - \Delta))g \right\} (x+t\omega) dS_\omega \\ + \int_0^t d\tau \int_{|\omega|=1} \frac{\{2\partial_t + (t-\tau)(\partial_t^2 - \Delta)\}F(x+(t-\tau)\omega, \tau)}{(n-2)\omega_n} dS_\omega \\ + \frac{F(x, t)}{n-2}.$$

Splitting  $2\partial_t$  into  $(n-1)\partial_t + (3-n)\partial_t$  and making use of (3.2) again, we have that

$$u_{tt}(x, t) - \Delta u(x, t) = \frac{\varepsilon}{\omega_n} \int_{|\omega|=1} \left\{ \frac{3-n}{n-2} \Delta f + (3-n)(\partial_t)g \right\} (x+t\omega) dS_\omega + \frac{F(x, t)}{n-2} \\ + \frac{3-n}{(n-2)\omega_n} \int_0^t d\tau \int_{|\omega|=1} \partial_t(F(x+(t-\tau)\omega, \tau)) dS_\omega.$$

Since

$$\partial_t(F(x+(t-\tau)\omega, \tau)) = (\partial_t F)(x+(t-\tau)\omega, \tau) - \partial_\tau(F(x+(t-\tau)\omega, \tau))$$

and

$$\int_0^t \partial_\tau \left( \int_{|\omega|=1} F(x+(t-\tau)\omega, \tau) dS_\omega \right) d\tau = \omega_n F(x, t) - \int_{|\omega|=1} F(x+t\omega, 0) dS_\omega,$$

we finally obtain that

$$\begin{aligned} & u_{tt}(x, t) - \Delta u(x, t) \\ &= \frac{(3-n)\varepsilon}{(n-2)\omega_n} \int_{|\omega|=1} \{\Delta f + (n-2)(\partial_t)g\}(x + t\omega) dS_\omega \\ & \quad + \frac{3-n}{(n-2)\omega_n} \int_{|\omega|=1} F(x + t\omega, 0) dS_\omega + F(x, t) \\ & \quad + \frac{3-n}{(n-2)\omega_n} \int_0^t d\tau \int_{|\omega|=1} (\partial_t F)(x + (t-\tau)\omega, \tau) dS_\omega. \end{aligned}$$

This ends the proof of [Theorem 2.1](#).  $\square$

#### 4. Decay estimate of the linear part

In this section, we get a space–time decay estimate of  $V$  in [\(2.4\)](#) which plays an essential role to define our weighted  $L^\infty$  space.

**Lemma 4.1.** *Under the same assumption as in [Theorem 2.2](#), there exists a positive constant  $C_{n,k}$  depending only on  $n$  and  $k$  such that  $V$  satisfies*

$$\begin{aligned} & (t + |x| + 2k)^{n-2} |\nabla_x^\alpha V(x, t)| \\ & \leq C_{n,k} \left( \sum_{|\beta| \leq |\alpha|+2} \|\nabla_x^\beta f\|_{L^\infty(\mathbf{R}^n)} + \sum_{|\gamma| \leq |\alpha|+1} \|\nabla_x^\gamma g\|_{L^\infty(\mathbf{R}^n)} \right) \end{aligned} \quad (4.1)$$

for  $|\alpha| \leq 2$ ,  $(x, t) \in \mathbf{R}^n \times [0, \infty)$ , and

$$\text{supp } V \subset \{(x, t) \in \mathbf{R}^n \times [0, \infty) : -k \leq t - |x| \leq k\}. \quad (4.2)$$

**Proof.** First we note that the support property [\(4.2\)](#) immediately follows from the representation of  $V$  in [\(2.4\)](#), and that it is enough to prove the lemma for  $|\alpha| = 0$ . For [\(4.1\)](#) with  $|\alpha| = 0$ , one can employ the standard argument as in Lemma 3.2 in Agemi, Kubota and Takamura [\[3\]](#).

When  $t \geq k$ , taking into account of [\(4.2\)](#), one can make use of

$$t^{n-1} \int_{|\omega|=1} |\varphi(x + t\omega)| dS_\omega \leq \|\nabla_x \varphi\|_{L^1(\mathbf{R}^n)} \quad \text{for } \varphi \in C_0^1(\mathbf{R}^n), \quad t > 0$$

with

$$t \geq \frac{1}{5} (t + |x| + 2k).$$

Hence we obtain that

$$|V(x, t)| \leq \frac{C_{n,k}}{(t + |x| + 2k)^{n-2}} \left( \sum_{1 \leq |\beta| \leq 2} \|\nabla_x^\beta f\|_{L^1(\mathbf{R}^n)} + \sum_{|\gamma|=1} \|\nabla_x^\gamma g\|_{L^1(\mathbf{R}^n)} \right)$$

with some positive constant  $C_{n,k}$  depending only on  $n$  and  $k$ . When  $t \leq k$ , (2.4) yields that

$$|V(x, t)| \leq C_{n,k} \left( \sum_{|\beta| \leq 1} \|\nabla_x^\beta f\|_{L^\infty(\mathbf{R}^n)} + \|g\|_{L^\infty(\mathbf{R}^n)} \right)$$

with a different constant  $C_{n,k} > 0$ . Therefore the proof is completed.  $\square$

## 5. Proof of Theorem 2.2

Following Takamura and Wakasa [24], we prove Theorem 2.2 in this section. We note that its proof is similar to the one of odd dimensional case in [24] because of Huygens' principle for the linear part of the integral equation, (4.2). It is obvious that the theorem follows from the following proposition.

**Proposition 5.1.** *Let  $n \geq 3$ . Suppose that the assumptions (2.8), (2.9) and (2.10) are fulfilled. Then, there exists a positive constant  $\varepsilon_0 = \varepsilon_0(f, g, n, p, k)$  such that (2.3) admits a unique solution  $u \in C^1(\mathbf{R}^n \times [0, T])$  for  $1 < p < 2$ ,  $u \in C^2(\mathbf{R}^n \times [0, T])$  for  $p \geq 2$ , as far as  $T$  satisfies*

$$\begin{aligned} T &\leq c\varepsilon^{-2p(p-1)/\zeta(p,n)} && \text{if } 1 < p < p_1(n), \\ T &\leq \exp\left(c\varepsilon^{-p(p-1)}\right) && \text{if } p = p_1(n), \\ \text{there is no bound} &&& \text{if } p > p_1(n) \end{aligned} \quad (5.1)$$

for  $0 < \varepsilon \leq \varepsilon_0$ , where  $c$  is a positive constant independent of  $\varepsilon$ .

The solution is constructed by almost the same way as in [24]. Actually, we shall set  $U = u - \varepsilon V$  and rewrite (2.3) with  $F = F(u)$  into the following form.

$$U = N(F(U + \varepsilon V)). \quad (5.2)$$

Since  $V$  exists globally in time, we have to consider the lifespan of the solution of (5.2). Let us define the sequence of functions,  $\{U_m\}_{m \in \mathbf{N}}$  by

$$U_m = N(F(U_{m-1} + U_0)) \text{ and } U_0 = \varepsilon V.$$

We also denote a weighted  $L^\infty$  norm of  $U$  by

$$\|U\| = \sup_{(x,t) \in \mathbf{R}^n \times [0,T]} \{w(|x|, t)|U(x, t)|\}$$

with the weighted function

$$w(r, t) = \begin{cases} \tau_+(r, t)^{n-2} \tau_-(r, t)^{\bar{q}} & \text{if } p > \frac{n-1}{n-2}, \\ \tau_+(r, t)^{n-2} \left( \log 4 \frac{\tau_+(r, t)}{\tau_-(r, t)} \right)^{-1} & \text{if } p = \frac{n-1}{n-2}, \\ \tau_+(r, t)^{n-2+\bar{q}} & \text{if } 1 < p < \frac{n-1}{n-2}, \end{cases}$$

where we set

$$\bar{q} = (n-2)p - (n-1)$$

and

$$\tau_+(r, t) = \frac{t+r+2k}{k}, \quad \tau_-(r, t) = \frac{t-r+2k}{k}.$$

**Proof of Proposition 5.1.** In view of Proposition 5.1 in [24], the proof of this proposition follows from the following *a priori* estimate.

**Lemma 5.1.** Let  $n \geq 3$  and  $N$  be a linear integral operator defined in (2.5). Assume that  $U, U_0 \in C^0(\mathbf{R}^n \times [0, T])$  with  $\text{supp } U \subset \{(x, t) \in \mathbf{R}^n \times [0, T] : |x| \leq t+k\}$ ,  $\text{supp } U_0 \subset \{(x, t) \in \mathbf{R}^n \times [0, T] : t-k \leq |x| \leq t+k\}$ , and  $\|U\|, \|\tau_+^{n-2} U_0 w^{-1}\| < \infty$ . Then, there exists a positive constant  $C_{n,v,p}$  depending on  $n, v$  and  $p$  such that

$$\|N(|U_0|^{p-\nu} |U|^\nu)\| \leq C_{n,v,p} k^2 \left\| \frac{\tau_+^{n-2}}{w} U_0 \right\|^{p-\nu} \|U\|^\nu \bar{E}_v(T) \quad (5.3)$$

for  $0 \leq \nu \leq p$ , where  $\bar{E}_v$  is defined by

$$\bar{E}_v(T) = \begin{cases} 1 & \text{if } p > \frac{n-1}{n-2}, \\ \left( \frac{2T+3k}{k} \right)^{\nu\delta} & \text{if } p = \frac{n-1}{n-2}, \\ \left( \frac{2T+3k}{k} \right)^{-\nu\bar{q}} & \text{if } 1 < p < \frac{n-1}{n-2}, \end{cases} \quad (5.4)$$

for  $0 \leq \nu < p$  with any  $\delta > 0$  and

$$\bar{E}_p(T) = \begin{cases} 1 & \text{if } p > p_1(n), \\ \log \frac{2T+3k}{k} & \text{if } p = p_1(n), \\ \left( \frac{2T+3k}{k} \right)^{\zeta(p,n)/2} & \text{if } 1 < p < p_1(n). \end{cases} \quad (5.5)$$

This lemma is proved in the next section.

The construction of the solution in our proposition is completely same as in the proof of lower bounds of the lifespan in odd space dimensions in the section 5 of Takamura and Wakasa [24],

if  $(n-1)/2$  in the exponent of  $\tau_+$ ,  $(n+1)/(n-1)$  in the definition of  $E_\nu(T)$ ,  $q$ ,  $p_0(n)$  and  $\gamma(p, n)$  are substituted by  $(n-2)$ ,  $(n-1)/(n-2)$ ,  $\bar{q}$ ,  $p_1(n)$ , and  $\zeta(p, n)$  in all the questions respectively. Therefore, Proposition 5.1 immediately follows from Lemma 5.1 which is proved in the next section.  $\square$

## 6. A priori estimates

In this section we prove Lemma 5.1 which plays a key role in the proof of Theorem 2.2. The proof follows from the following basic estimate.

**Lemma 6.1** (Basic estimate). *Let  $N$  be the linear integral operator defined by (2.5) and  $a_1 \geq 0$ ,  $a_2 \in \mathbf{R}$  and  $a_3 \geq 0$ . Then, there exists a positive constant  $C_{n,p,a_1,a_2,a_3}$  such that*

$$\begin{aligned} & N \left\{ \tau_+^{-(n-2)p+a_1} \tau_-^{a_2} (\log(4\tau_+/\tau_-))^{a_3} \right\} (x, t) \\ & \leq C_{n,p,a_1,a_2,a_3} k^2 w(r, t)^{-1} \left( \frac{2T+3k}{k} \right)^{a_1} \bar{E}_{a_1,a_2,a_3}(T) \end{aligned} \quad (6.1)$$

for  $|x| \leq t+k$ ,  $t \in [0, T]$ , where  $\bar{E}_{a_1,a_2,a_3}(T)$  is defined by

$$\bar{E}_{a_1,a_2,a_3}(T) = \begin{cases} 1 & \text{if } a_2 < -1 \text{ and } a_3 = 0, \\ \log \frac{2T+3k}{k} & \text{if } a_2 = -1 \text{ and } a_3 = 0, \\ \left( \frac{2T+3k}{k} \right)^{\delta a_3} & \text{if } a_2 \leq -1 \text{ and } a_3 > 0, \\ \left( \frac{2T+3k}{k} \right)^{1+a_2} & \text{if } a_2 > -1, \end{cases} \quad (6.2)$$

where  $\delta$  stands for any positive constant.

To prove this lemma, we shall employ the following lemma which is established by fundamental identity for spherical means by John [9].

**Lemma 6.2.** (See John [9].) *Let  $b \in C([0, \infty))$ . Then, the identity*

$$\int_{|\omega|=1} b(|x + \rho\omega|) dS_\omega = 2^{3-n} \omega_{n-1} (r\rho)^{2-n} \int_{|\rho-r|}^{\rho+r} \lambda b(\lambda) h(\lambda, \rho, r) d\lambda, \quad (6.3)$$

holds for  $x \in \mathbf{R}^n$ ,  $r = |x|$  and  $\rho > 0$ , where

$$h(\lambda, \rho, r) = \{\lambda^2 - (\rho - r)^2\}^{(n-3)/2} \{(\rho + r)^2 - \lambda^2\}^{(n-3)/2}. \quad (6.4)$$

For the proof of this lemma, see Lemma 4.1 in Agemi, Kubota and Takamura [3]. In order to estimate  $h(\lambda, \rho, r)$ , we shall make use of the following four inequalities.

**Lemma 6.3.** Let  $h(\lambda, \rho, r)$  be the function defined by (6.4). Suppose that  $|\rho - r| \leq \lambda \leq \rho + r$ , or equivalently  $|\lambda - r| \leq \rho \leq \lambda + r$ , and  $\rho \geq 0$ . Then the following inequalities hold.

$$h(\lambda, \rho, r) \leq 4^{n-3} r^{n-3} \lambda^{n-3}, \quad (6.5)$$

$$h(\lambda, \rho, r) \leq 2^{n-3} \rho^{n-3} r^{(n-3)/2} \lambda^{(n-3)/2}, \quad (6.6)$$

$$h(\lambda, \rho, r) \leq 8^{n-3} \rho^{n-3} r^{n-3}, \quad (6.7)$$

$$h(\lambda, \rho, r) \leq 2^{n-3} \rho^{n-3} \lambda^{n-3}. \quad (6.8)$$

**Proof.** (6.5), (6.6) and (6.7) are due to Lemma 4.2 in Agemi, Kubota and Takamura [3] with elementary computations. (6.8) is due to Lemma 2.2 in Georgiev [5] with geometrical observation. But one may prove (6.8) also by elementary computation as follows.

$$\begin{aligned} & 4\rho^2\lambda^2 - \{\lambda^2 - (\rho - r)^2\}\{(\rho + r)^2 - \lambda^2\} \\ &= \lambda^4 + \{4\rho^2 - (\rho + r)^2 - (\rho - r)^2\}\lambda^2 + (\rho - r)^2(\rho + r)^2 \\ &= (\lambda^2 + \rho^2 - r^2)^2 \geq 0. \quad \square \end{aligned}$$

**Proof of Lemma 6.1.** The proof is almost the same as the one in the estimates for  $I_{\text{odd}}$  in Lemma 4.5 of Takamura and Wakasa [24]. We denote various positive constants depending only on  $n$  and  $p$  by  $C$  which may change at place to place. By virtue of Lemma 6.2, we have that

$$N \left\{ \tau_+^{-(n-2)p+a_1} \tau_-^{a_2} (\log(4\tau_+/\tau_-))^{a_3} \right\} (x, t) = I(r, t),$$

where we set

$$\begin{aligned} I(r, t) &= C r^{2-n} \int_0^t (t - \tau)^{3-n} d\tau \int_{|t-\tau-r|}^{t-\tau+r} \tau_+(\lambda, \tau)^{-(n-2)p+a_1} \tau_-(\lambda, \tau)^{a_2} \times \\ &\quad \times \left( \log 4 \frac{\tau_+(\lambda, \tau)}{\tau_-(\lambda, \tau)} \right)^{a_3} \lambda h(\lambda, t - \tau, r) d\lambda. \end{aligned} \quad (6.9)$$

We shall estimate  $I(r, t)$  on three domains,

$$\begin{aligned} D_1 &= \{(r, t) \mid r \geq t - r > -k \text{ and } r \geq 2k\}, \\ D_2 &= \{(r, t) \mid r \geq t - r > -k \text{ and } r \leq 2k\}, \\ D_3 &= \{(r, t) \mid t - r \geq r\}. \end{aligned}$$

(i) Estimate in  $D_1$ .

Making use of (6.8), we get

$$I(r, t) \leq Cr^{2-n} \int_0^t d\tau \int_{|t-\tau-r|}^{t+r-\tau} \lambda^{n-2} \times \\ \times \tau_+(\lambda, \tau)^{-(n-2)p+a_1} \tau_-(\lambda, \tau)^{a_2} \left( \log 4 \frac{\tau_+(\lambda, \tau)}{\tau_-(\lambda, \tau)} \right)^{a_3} d\lambda.$$

Changing variables in the above integral by

$$\alpha = \tau + \lambda, \quad \beta = \tau - \lambda,$$

we get

$$I(r, t) \leq Cr^{2-n} \int_{-k}^{t-r} \left( \frac{\beta + 2k}{k} \right)^{a_2} d\beta \int_{|t-r|}^{t+r} (\alpha - \beta)^{n-2} \times \\ \times \left( \frac{\alpha + 2k}{k} \right)^{-(n-2)p+a_1} \left( \log 4 \frac{\alpha + 2k}{\beta + 2k} \right)^{a_3} d\alpha.$$

It follows from

$$\frac{r}{k} = \frac{r + 2r + r}{4k} \geq \frac{\tau_+(r, t)}{4}$$

in  $D_1$  that

$$I(r, t) \leq C\tau_+(r, t)^{2-n} \left( \frac{t+r+2k}{k} \right)^{a_1} \int_{-k}^{t-r} \left( \frac{\beta + 2k}{k} \right)^{a_2} d\beta \times \\ \times \int_{t-r}^{t+r} \left( \frac{\alpha + 2k}{k} \right)^{-1-\bar{q}} \left( \log 4 \frac{\alpha + 2k}{\beta + 2k} \right)^{a_3} d\alpha. \quad (6.10)$$

When  $a_3 = 0$ ,  $\alpha$ -integral in (6.10) is dominated by

$$\begin{cases} Ck\tau_-^{-\bar{q}} & \text{if } p > \frac{n-1}{n-2}, \\ k \log \frac{\tau_+}{\tau_-} & \text{if } p = \frac{n-1}{n-2}, \\ Ck\tau_+^{-\bar{q}} & \text{if } 1 < p < \frac{n-1}{n-2} \end{cases}$$

and  $\beta$ -integral in (6.10) is dominated by

$$\begin{cases} \frac{-k}{1+a_2} & \text{if } a_2 < -1, \\ k \log \frac{t-r+2k}{k} & \text{if } a_2 = -1, \\ \frac{k}{1+a_2} \left( \frac{t-r+2k}{k} \right)^{1+a_2} & \text{if } a_2 > -1. \end{cases}$$

(6.1) is now established for  $a_3 = 0$ .

When  $a_3 > 0$ , we employ the following simple lemma.

**Lemma 6.4.** *Let  $\delta > 0$  be any given constant. Then, we have*

$$\log X \leq \frac{X^\delta}{\delta} \text{ for } X \geq 1.$$

The proof of this lemma follows from elementary computation. We shall omit it. Then, it follows from Lemma 6.4 that

$$\begin{aligned} I(r, t) &\leq C(4\delta^{-1})^{a_3} \tau_+(r, t)^{2-n} \left( \frac{t+r+2k}{k} \right)^{a_1+\delta a_3} \times \\ &\quad \times \int_{-k}^{t-r} \left( \frac{\beta+2k}{k} \right)^{a_2-\delta a_3} d\beta \int_{t-r}^{t+r} \left( \frac{\alpha+2k}{k} \right)^{-1-\bar{q}} d\alpha. \end{aligned}$$

The  $\alpha$ -integral above can be estimated by the same manner in the case of  $a_3 = 0$ . The  $\beta$ -integral is dominated by

$$\begin{cases} \frac{-k}{1+a_2-\delta a_3} & \text{if } a_2 \leq -1, \\ \frac{k}{1+a_2-\delta a_3} \left( \frac{t-r+2k}{k} \right)^{1+a_2-\delta a_3} & \text{if } a_2 > -1 \end{cases} \quad (6.11)$$

with  $\delta > 0$  satisfying  $1 + a_2 - \delta a_3 > 0$ . Therefore  $I$  is bounded in  $D_1$  by the quantity in the right-hand side of (6.1) as desired. It is obvious that such a restriction on  $\delta > 0$  is finally removed from the statement.

(ii) Estimate in  $D_2$  or  $D_3$ .

In this case, the proof is completely same as the one in the estimates for  $I_{\text{odd}}$  in Lemma 4.5 in Takamura and Wakasa [24], if  $(n-1)/2$  in the exponent of  $\tau_+$  is substituted by  $(n-2)$ . Because the key fact,  $1 - (n-2)p < 0$ , is also trivial. Therefore, the proof of Lemma 6.1 is now completed.  $\square$

**Proof of Lemma 5.1.** Due to Huygens' principle for the linear part  $V$ , (4.2), one can replace  $\tau_-$  by  $\tau_- \chi_{\{-k \leq t-r \leq k\}}$  in (6.1) when  $0 \leq v < p$ . Then, the integral with respect to the variable  $\beta = \tau - \lambda$  is bounded. In order to establish Lemma 5.1, it is sufficient to show

$$\begin{cases} N \left( \tau_+^{-(n-2)(p-v)} w^{-v} \chi_{\{-k \leq t-r \leq k\}} \right) (x, t) \leq C_{n,v,p} k^2 \bar{E}_v(T) & \text{for } 0 \leq v < p, \\ N(w^{-p})(x, t) \leq C_{n,p,p} k^2 \bar{E}_p(T) & \text{for } v = p. \end{cases}$$



To this end, setting

$$\begin{cases} a_1 = a_3 = 0, a_2 = -v\bar{q} & \text{if } p > \frac{n-1}{n-2}, \\ a_1 = a_2 = 0, a_3 = v & \text{if } p = \frac{n-1}{n-2}, \\ a_1 = -v\bar{q}, a_2 = a_3 = 0 & \text{if } 1 < p < \frac{n-1}{n-2} \end{cases}$$

for  $0 \leq v < p$  and

$$\begin{cases} a_1 = a_3 = 0, a_2 = -p\bar{q} & \text{if } p > \frac{n-1}{n-2}, \\ a_1 = a_2 = 0, a_3 = p & \text{if } p = \frac{n-1}{n-2}, \\ a_1 = -p\bar{q}, a_2 = a_3 = 0 & \text{if } 1 < p < \frac{n-1}{n-2} \end{cases}$$

for  $v = p$  in (6.1), we have (5.3).  $\square$

## 7. Proof of Theorem 2.3

In this section, we prove Theorem 2.3 which obviously follows from Proposition 7.1 below. Its proof is almost the same as the one in odd dimensional case of Theorem 2.2 in Takamura and Wakasa [24] once the similar iteration frame is established.

**Proposition 7.1.** *Suppose that the assumptions of Theorem 2.3 are fulfilled. Let  $u$  be a  $C^0$ -solution of (2.3) in  $\mathbf{R}^n \times [0, T]$ . Then, there exists a positive constant  $\varepsilon_0 = \varepsilon_0(g, n, p, k)$  such that  $T$  cannot be taken as*

$$T > \exp\left(c\varepsilon^{-p(p-1)}\right) \quad \text{if } p = p_1(n), \quad (7.1)$$

$$T > c\varepsilon^{-2p(p-1)/\zeta(p,n)} \quad \text{if } 1 < p < p_1(n) \quad (7.2)$$

for  $0 < \varepsilon \leq \varepsilon_0$ , where  $c$  is a positive constant independent of  $\varepsilon$ .

**Proof.** Similarly to the proof of Proposition 7.1 in [24], we may assume that the solution of (2.3) is radially symmetric without loss of the generality. Let  $u = u(r, t)$  be a  $C^0$ -solution of

$$u = \varepsilon V + N(|u|^p) \quad \text{in } (0, \infty) \times [0, T], \quad (7.3)$$

where we set

$$V(r, t) = Cr^{2-n}t^{3-n} \int_{|t-r|}^{t+r} \lambda g(\lambda) h(\lambda, t, r) d\lambda, \quad (7.4)$$

$$N(|u|^p)(r, t) = \overline{C} r^{2-n} \int_0^t (t-\tau)^{3-n} d\tau \int_{|t-\tau-r|}^{t-\tau+r} \lambda h(\lambda, t-\tau, r) |u(\lambda, \tau)|^p d\lambda, \quad (7.5)$$

where  $C$  and  $\overline{C}$  are positive constants depending only on  $n$ .

**[The 1st step] Inequality of  $u$ .**

**Lemma 7.1.** Assume (2.13). Then there exists a positive constant  $C_{n,g,k} > 0$  such that for  $t + k_0 < r < t + k_1$  and  $t \geq k_2$ ,

$$V(r, t) \geq \frac{C_{n,g,k}}{r^{n-2}}, \quad (7.6)$$

where  $k_1 = \frac{k+k_0}{2}$  and  $k_2 = k - k_0$ .

**Proof.** Let  $t + k_0 < r < t + k_1$  and  $t \geq k_2/2$ . Then, (7.4) gives us

$$V(r, t) \geq C r^{2-n} t^{3-n} \int_{k_1}^k \lambda g(\lambda) h(\lambda, t, r) d\lambda.$$

Note that

$$\begin{aligned} r + t + \lambda &\geq r, \quad \lambda + r - t \geq \lambda, \\ r + t - \lambda &\geq r + t - k \geq 2t + k_0 - k \geq t, \quad \lambda + t - r \geq \lambda - k_1 \end{aligned}$$

hold in the domain of the integral above for  $t + k_0 < r < t + k_1$  and  $t \geq k_2$ . Hence, we get

$$\begin{aligned} V(r, t) &\geq C r^{-(n-1)/2} t^{-(n-3)/2} \int_{k_1}^k \lambda^{(n-1)/2} g(\lambda) (\lambda - k_1)^{(n-3)/2} d\lambda \\ &\geq C \left( \frac{k - k_1}{2} \right)^{(n-3)/2} r^{-(n-2)} \int_{(k+k_1)/2}^k \lambda^{(n-1)/2} g(\lambda) d\lambda \end{aligned}$$

for  $t + k_0 < r < t + k_1$  and  $t \geq k_2$ . Therefore we obtain (7.6).  $\square$

Making use of this estimate of  $V$ , we have the following iteration frame.

**Lemma 7.2.** Let  $u$  be a  $C^0$ -solution of (7.3). Assume (2.13). Then  $u$  in  $\Sigma_0 = \{(r, t) : 2k \leq t - r \leq r\}$  satisfies

$$\begin{aligned}
u(r, t) \geq & \frac{\overline{C} 2^{(n-3)/2} (t-r)^{(n-1)/2}}{r^{(3n-7)/2}} \times \\
& \times \int_{R(r,t)} \int \{ (t-r-\tau+\lambda)(t+r-\tau-\lambda) \}^{(n-3)/2} |u(\lambda, \tau)|^p d\lambda d\tau + \\
& + \frac{E_1 (t-r)^{(3n-5)/2-(n-2)p}}{r^{(3n-7)/2}} \varepsilon^p,
\end{aligned} \quad (7.7)$$

where  $\overline{C}$  is the one in (7.5),

$$E_1 = \frac{\overline{C} C_{n,g,k}^p (k_1 - k_0)}{(n-1) 2^{(n-2)p-(3n-11)/2}}$$

and

$$R(r, t) = \{(\lambda, \tau) : t-r \leq \lambda, \tau + \lambda \leq t+r, 2k \leq \tau - \lambda \leq t-r\}.$$

**Proof.** Comparing  $L_{odd}$  in (4.7) of [24] with radially symmetric form of  $N$  in (7.4) of this paper, the difference between the proof of Lemma 7.2 of [24] and the one of this lemma has to appear only in the second term,  $I_2$ , which arises from the estimate of the linear part. In view of the proof of Lemma 7.2 in [24], the desired estimate immediately follows from simple replacement of  $1 - (n-1)p/2$  in the exponent of  $\alpha - \beta$  by  $1 - (n-2)p$ .  $\square$

### [The 2nd Step] Comparison argument.

Let us consider a solution  $w$  of

$$\begin{aligned}
w(t-r) = & \frac{\overline{C} 2^{(n-5)/2} (t-r)^{(n-1)/2}}{r^{(3n-7)/2}} \int_{2k}^{t-r} (t-r-\beta)^{(n-3)/2} d\beta \\
& \times \int_{2(t-r)+\beta}^{t+r} (t+r-\alpha)^{(n-3)/2} |w(\beta)|^p d\alpha \\
& + \frac{E_1 (t-r)^{(3n-5)/2-(n-2)p}}{2r^{(3n-7)/2}} \varepsilon^p.
\end{aligned} \quad (7.8)$$

Then we have the following comparison lemma.

**Lemma 7.3.** *Let  $u$  be a solution of (7.3) and  $w$  be a solution of (7.8). Then,  $u$  and  $w$  satisfy*

$$u > w \quad \text{in } \Sigma_0.$$

**Proof.** Comparing the relation between  $u$  in Lemma 7.3 of [24] and  $w$  in (7.6) of [24] with the one between  $u$  in Lemma 7.2 and  $w$  in (7.8), one can find no difference in the structure of proofs of both Lemma 7.4 of [24] and this lemma.  $\square$

By definition of  $w$  in (7.8), we have

$$w(\xi) \geq \frac{\overline{C}\xi^{3-n}}{2^{n-1}} \int_{2k}^{\xi} (\xi - \beta)^{(n-3)/2} |w(\beta)|^p d\beta \\ \times \int_{2\xi+\beta}^{3\xi} (3\xi - \alpha)^{(n-3)/2} d\alpha + \frac{E_1}{2^{(3n-5)/2}} \xi^{-\bar{q}-(n-2)} \varepsilon^p$$

in  $\Gamma_0$ , where we set

$$\xi = \frac{r}{2}, \quad \Gamma_0 = \{t - r = \xi, r \geq 4k\}.$$

Hence we obtain that

$$w(\xi) \geq \frac{\overline{C}\xi^{3-n}}{2^{n-2}(n-1)} \int_{2k}^{\xi} (\xi - \beta)^{n-2} |w(\beta)|^p d\beta + \frac{E_1 \xi^{-\bar{q}-(n-2)}}{2^{(3n-5)/2}} \varepsilon^p$$

for  $\xi \geq 2k$ . Then, it follows from the setting

$$W(\xi) = \xi^{\bar{q}+n-2} w(\xi)$$

that

$$W(\xi) \geq D_n \xi^{\bar{q}+1} \int_{2k}^{\xi} \frac{(\xi - \beta)^{n-2} |W(\beta)|^p d\beta}{\beta^{(n-2)p+p\bar{q}}} + E_2 \varepsilon^p \quad \text{for } \xi \geq 2k, \quad (7.9)$$

where we set

$$D_n = \frac{\overline{C}}{2^{n-2}(n-1)}, \quad E_2 = \frac{E_1}{2^{(3n-5)/2}}.$$

### Iteration frame in the case of $p = p_1(n)$ .

By virtue of (7.9), we get

$$W(\xi) \geq D_n \int_{2k}^{\xi} \left( \frac{\xi - \beta}{\xi} \right)^{n-2} \frac{|W(\beta)|^p}{\beta^{p\bar{q}}} d\beta + E_2 \varepsilon^p \quad \text{for } \xi \geq 2k. \quad (7.10)$$

The above inequality is the iteration frame for the critical case. This inequality is the same as the one in (7.8) in [24], if  $q$  is substituted by  $\bar{q}$ .

### Iteration frame in the case of $1 < p < p_1(n)$ .

Because of the fact that  $-(n-2)p - p\bar{q} < 0$  for  $n \geq 3$ , (7.9) yields

$$W(\xi) \geq D_n \xi^{-(n-2)-p\bar{q}} \int_{2k}^{\xi} (\xi - \beta)^{n-2} |W(\beta)|^p d\beta + E_2 \varepsilon^p \quad \text{for } \xi \geq 2k. \quad (7.11)$$

The above inequality is the iteration frame for the subcritical case. This inequality is the same as the one in (8.2) in [24], if  $q$  is substituted by  $\bar{q}$ .

Making use of (7.10) and (7.11), one can obtain Proposition 7.1 immediately by the same argument in [24]. Therefore the proof of Theorem 2.3 is now completed.  $\square$

## References

- [1] R. Agemi, Oral communication.
- [2] R. Agemi, H. Takamura, Remarks on representations of solutions to the wave equations, Mathematical research note 94-004, Institute of Mathematics, University of Tsukuba, June 1994.
- [3] R. Agemi, K. Kubota, H. Takamura, On certain integral equations related to nonlinear wave equations, Hokkaido Math. J. 23 (1994) 241–276.
- [4] D. Christodoulou, Global solutions of nonlinear hyperbolic equations for small initial data, Comm. Pure Appl. Math. 39 (1986) 267–282.
- [5] V. Georgiev, Existence of global solutions to supercritical semilinear wave equations, Serdica Math. J. 22 (1996) 125–164.
- [6] R.T. Glassey, Blow-up theorems for nonlinear wave equations, Math. Z. 132 (1973) 183–203.
- [7] P. Godin, Lifespan of solutions of semilinear wave equations in two space dimensions, Comm. Partial Differential Equations 18 (1993) 895–916.
- [8] A. Hoshiga, The existence of the global solutions to semilinear wave equations with a class of cubic nonlinearities in 2-dimensional space, Hokkaido Math. J. 37 (2008) 669–688.
- [9] F. John, Plane Waves and Spherical Means, Applied to Partial Differential Equations, Interscience Publishers, Inc., New York, 1955.
- [10] F. John, Blow-up of solutions of nonlinear wave equations in three space dimensions, Manuscripta Math. 28 (1979) 235–268.
- [11] F. John, Nonlinear Wave Equations, Formation of Singularities, Pitcher Lectures in Mathematical Sciences, Lehigh University, AMS, 1990.
- [12] S. Katayama, Global existence for systems of nonlinear wave equations in two space dimensions. II, Publ. Res. Inst. Math. Sci. 31 (1995) 645–665.
- [13] S. Katayama, Lifespan of solutions for two space dimensional wave equations with cubic nonlinearity, Comm. Partial Differential Equations 26 (2001) 205–232.
- [14] S. Klainerman, The null condition and global existence to nonlinear wave equations, in: Nonlinear Systems of Partial Differential Equations in Applied Mathematics, Part 1, Santa Fe, N.M., 1984, in: Lectures in Appl. Math., vol. 23, Amer. Math. Soc., Providence, RI, 1986, pp. 293–326.
- [15] H. Kubo, Asymptotic behavior of solutions of semilinear wave equations with dissipative structure, Discrete Contin. Dyn. Syst. (Suppl.) (2007) 602–613, Proceedings of the 6th AIMS International Conference; in: Dynamical Systems and Differential Equations, 2007.
- [16] H.A. Levine, Instability and nonexistence of global solutions to nonlinear wave equations of the form  $Pu_{tt} = -Au + \mathcal{F}(u)$ , Trans. Amer. Math. Soc. 192 (1974) 1–21.
- [17] T-T. Li, Lower bounds of the life-span of small classical solutions for nonlinear wave equations, in: M. Beals, R.B. Melrose, J. Rauch (Eds.), Microlocal Analysis and Nonlinear Waves, Minneapolis, MN, 1988–1989, in: The IMA Volumes in Mathematics and Its Applications, vol. 30, Springer-Verlag New York Inc., 1991, pp. 125–136.
- [18] T-T. Li, Y. Chen, Global Classical Solutions for Nonlinear Evolution Equations, Pitman Monographs and Surveys in Pure and Applied Mathematics, vol. 45, Longman Scientific & Technical, 1992.
- [19] T-T. Li, Y. Zhou, A note on the life-span of classical solutions to nonlinear wave equations in four space dimensions, Indiana Univ. Math. J. 44 (1995) 1207–1248.
- [20] T.C. Sideris, Formation of singularities in solutions to nonlinear hyperbolic equations, Arch. Ration. Mech. Anal. 86 (1984) 369–381.
- [21] W.A. Strauss, Nonlinear scattering theory at low energy, J. Funct. Anal. 41 (1981) 110–133.

- [22] W.A. Strauss, *Nonlinear Wave Equations*, CBMS Regional Conference Series in Mathematics, vol. 73, American Mathematical Society, Providence, RI, 1989, Published for the Conference Board of the Mathematical Sciences, Washington, DC.
- [23] H. Takamura, K. Wakasa, The sharp upper bound of the lifespan of solutions to critical semilinear wave equations in high dimensions, *J. Differential Equations* 251 (2011) 1157–1171.
- [24] H. Takamura, K. Wakasa, Almost global solutions of semilinear wave equations with the critical exponent in high dimensions, *Nonlinear Anal.* 109 (2014) 187–229.
- [25] Y. Zhou, W. Han, Sharpness on the lower bound of the lifespan of solutions to nonlinear wave equations, *Chin. Ann. Math. Ser. B* 4 (2011) 521–526.
- [26] Y. Zhou, W. Han, Blow up for some semilinear wave equations in multi-space dimensions, *Comm. Partial Differential Equations* 39 (2014) 651–665.