



New results on the Stieltjes constants: Asymptotic and exact evaluation

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Received 27 September 2004

Available online 22 July 2005

Submitted by H.M. Srivastava

Abstract

The Stieltjes constants $\gamma_k(a)$ are the expansion coefficients in the Laurent series for the Hurwitz zeta function about $s = 1$. We present new asymptotic, summatory, and other exact expressions for these and related constants.

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Keywords: Stieltjes constants; Riemann zeta function; Hurwitz zeta function; Laurent expansion; Integrals of periodic Bernoulli polynomials; Functional equation; Kreminski conjecture

1. Introduction

The Stieltjes (or generalized Euler) constants [2,6,10,16,20,21,23,27,29,30] $\gamma_k(a)$ appear as expansion coefficients in the Laurent series about $s = 1$ of the Hurwitz zeta function $\zeta(s, a)$, a generalization of the Riemann zeta function $\zeta(s)$ [3,11,16,18,24,28]. We present both new asymptotic and exact expressions for these and other fundamental mathematical constants.

In the following, we let $s(n, m)$ denote the Stirling numbers of the first kind and $C_k(a) \equiv \gamma_k(a) - (\ln^k a)/a$. By convention, γ_k represents $\gamma_k(1)$. With $B_n(x)$ the Bernoulli

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polynomials, their periodic extension is denoted $P_n(x) \equiv B_n(x - [x])$. As is customary, B_j denotes the Bernoulli numbers $B_j(0) = (-1)^j B_j(1)$.

Proposition 1. As $n \rightarrow \infty$, we have $C_n(a + 1/2) \rightarrow -C_n(a)$.

Proposition 2. For integers $m > 0$ and $n \geq 1$ we have

$$\gamma_n = (-1)^{n-1} n! m^{1-n} \sum_{k=0}^{n+1} \frac{s(n+1, n+1-k)}{k!} \int_1^\infty P_n(mx) \frac{\ln^k x}{x^{n+1}} dx - \sum_{r=1}^{m-1} C_n\left(\frac{r}{m}\right).$$

Proposition 3. For integers $q \geq 2$, we have

$$\sum_{r=1}^{q-1} \gamma_k\left(\frac{r}{q}\right) = -\gamma_k + q(-1)^k \frac{\ln^{k+1} q}{(k+1)} + q \sum_{j=0}^k \binom{k}{j} (-1)^j (\ln^j q) \gamma_{k-j}.$$

Proposition 4. For $\sigma = \text{Re } s < 1$, we have

$$\sum_{k=0}^\infty \frac{(-1)^k}{k!} (s-1)^k \int_0^1 \gamma_k(a) da = \frac{1}{1-s}, \quad \sigma < 1.$$

Proposition 5. For large but fixed k , $0 < a \leq 1$, we have

$$\gamma_k(a) \sim m \sin[2\pi(a + \phi)]$$

for real m and $0 \leq \phi < 1$.

Proposition 6. For integers $n \geq 1$ and $k \geq 1$ we have

$$\begin{aligned} & \sum_{j=0}^{n-1} C_k\left(a \pm \frac{j}{n}\right) \\ &= (-1)^{k-1} k! n^{1-k} \sum_{\ell=0}^{k+1} \frac{s(k+1, k+1-\ell)}{\ell!} \int_1^\infty P_k[n(x-a)] \frac{\ln^\ell x}{x^{k+1}} dx. \end{aligned}$$

Proposition 7. For integers $j \geq 1$ we have the set of coupled differential equations

$$\frac{(-1)^j}{j!} \frac{d\gamma_j(a)}{da} = - \sum_{k=j-1}^\infty \frac{(-1)^k}{k!} \binom{k+1}{j} \gamma_k(a), \quad j \geq 1,$$

and

$$-\psi'(a) = -1 - \sum_{k=0}^\infty \frac{(-1)^k}{k!} \gamma_k(a),$$

where ψ is the digamma function.

Proposition 8. *We have*

$$(a) \quad \gamma_0 = \gamma = \frac{1}{2} \ln 2 - \frac{1}{\ln 2} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \frac{\ln(k+1)}{(k+1)}$$

for the Euler constant,

$$(b) \quad -\gamma_1 = \frac{\ln^2 2}{12} - \frac{1}{2} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \frac{\ln(k+1)}{(k+1)} \\ + \frac{1}{2 \ln 2} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \frac{\ln^2(k+1)}{(k+1)},$$

$$(c) \quad \eta_1 = \frac{\ln^2 2}{12} + \frac{1}{2} \frac{1}{\ln^2 2} \left[\sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \frac{\ln(k+1)}{(k+1)} \right]^2 \\ - \frac{1}{\ln 2} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \frac{\ln^2(k+1)}{(k+1)},$$

$$(d) \quad \ln \pi = \ln 2 - 2 \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \ln(k+1),$$

$$(e) \quad (-1)^{n-1} \frac{2^{2n-1} \pi^{2n}}{(2n)!} B_{2n} = \frac{1}{1-2^{1-2n}} \sum_{\ell=0}^{\infty} \frac{1}{2^{\ell+1}} \sum_{k=0}^{\ell} (-1)^k \binom{\ell}{k} \frac{1}{(k+1)^{2n}} \\ = \frac{1}{(2n-1)} \sum_{\ell=0}^{\infty} \frac{1}{\ell+1} \sum_{k=0}^{\ell} (-1)^k \binom{\ell}{k} \frac{1}{(k+1)^{2n-1}}, \quad n \neq 0,$$

and the summation identity

$$(f) \quad \frac{1}{(2^{2n+1} - 1)} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \ln(k+1)(k+1)^{2n} \\ = (-1)^n \frac{(2n)!}{2(2\pi)^{2n}} \frac{1}{(1-2^{-2n})} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{1}{(k+1)^{2n+1}}.$$

In part (c), the constants η_j are defined by the expansion [7–9] $\ln \zeta(s) = -\ln(s-1) - \sum_{p=1}^{\infty} \frac{\eta_{p-1}}{p} (s-1)^p, s \neq 1$.

These theorems are representative, but not comprehensive, of our results. We do not reproduce relevant background material on properties and functional equations of the special functions involved, but simply mention the useful sources [1,13,14,26,30].

2. Proofs of the propositions

Here we summarize or otherwise simply indicate the method of proof of the propositions above.

Proposition 1. A key observation is that

$$P_{2n}(x - a) = (-1)^{n-1} \frac{2(2n)!}{(2\pi)^{2n}} [\cos 2\pi(x - a) + O(2^{-2n})], \tag{1}$$

and similarly for $P_{2n+1}(x - a)$, based upon the Fourier expansions of these polynomials [1]. So for large n , $P_{2n}(x - a) = -P_{2n}(x - a - 1/2) + O(2^{-2n})$ and we may appeal to the representation [30]

$$C_n(a) = (-1)^{n-1} n! \sum_{k=0}^{n+1} \frac{s(n+1, n+1-k)}{k!} \int_1^\infty P_n(x-a) \frac{\ln^k x}{x^{n+1}} dx, \quad n \geq 1. \tag{2}$$

Due to the boundedness of P_n , $|P_n(x)| \leq [3 + (-1)^n]/(2\pi)^n$ for $n \geq 1$ [3], there is uniform convergence of the integral in this equation. Therefore, we may use expressions like (1) in Eq. (2), interchange the limit and integration, and Proposition 1 follows. We have thereby proved a recent Conjecture II put forth by Kreminski [21].

Proposition 2. The multiplication formula satisfied by the Bernoulli polynomials [1] carries over to P_n . Applying this to the representation (2) yields the proposition.

Proposition 3. One procedure is to take derivatives of the special case of the Hurwitz zeta function $\sum_{r=1}^{q-1} \zeta(s, r/q) = (q^s - 1)\zeta(s)$ [14] and to evaluate as $s \rightarrow 1^+$. Another way is to use the defining Laurent expansions for the Stieltjes constants [21,27,30] in this special case, writing

$$\begin{aligned} & \sum_{r=1}^{q-1} \left[\frac{1}{s-1} + \sum_{k=0}^\infty \frac{(-1)^k \gamma_k(r/q)}{k!} (s-1)^k \right] \\ &= [q e^{(s-1) \ln q} - 1] \zeta(s) \\ &= \left[q \sum_{j=0}^\infty \frac{\ln^j q}{j!} (s-1)^j - 1 \right] \left[\frac{1}{s-1} + \sum_{k=0}^\infty \frac{(-1)^k \gamma_k}{k!} (s-1)^k \right]. \end{aligned} \tag{3}$$

We expand the right side of this equation and cancel the term $(q - 1)/(s - 1)$ on both sides. We then reorder the double sum on the right side and equate coefficients of like powers of $s - 1$ on each side of the resulting equation. The proposition follows once again. Proposition 3 takes the same form as Corollary 11 of Dilcher [10], who introduced a generalized digamma function. We explain this more in the Discussion section (Section 3).

Proposition 4. The key here is to note that

$$\int_0^1 \zeta(s, a) da = 0, \quad \sigma < 1, \tag{4}$$

and then to integrate the Laurent expansion of the Hurwitz zeta function.

Remark. Many other related integral-based results for the Stieltjes constants are possible.

Proposition 5. For proof, we give the details for $C_{2n-1}(a)$, with the case for $C_{2n}(a)$ being very similar. From Eq. (2), we have

$$C_{2n-1}(a) = (2n - 1)! \sum_{k=0}^{2n} \frac{s(2n, 2n - k)}{k!} \int_1^\infty P_{2n-1}(x - a) \frac{\ln^k x}{x^{2n}} dx, \quad n \geq 1. \tag{5}$$

We have [1]

$$P_{2n-1}(x - a) = (-1)^n \frac{2(2n - 1)!}{(2\pi)^{2n-1}} [\sin 2\pi(x - a) + O(2^{1-2n})], \tag{6}$$

so that we obtain

$$C_{2n-1}(a) = \kappa_n \sum_{k=0}^{2n} \frac{s(2n, 2n - k)}{k!} \times \int_1^\infty [\cos 2\pi a \sin 2\pi x - \sin 2\pi a \cos 2\pi x + O(2^{1-2n})] \frac{\ln^k x}{x^{2n}} dx, \tag{7}$$

where $\kappa_n \equiv (-1)^n 2[(2n - 1)!]^2 / (2\pi)^{2n-1}$. As a result of the integration in this equation, the leading term in $C_{2n-1}(a)$ may be written as

$$C_{2n-1}(a) \sim r_1(n) \cos 2\pi a - r_2(n) \sin 2\pi a, \tag{8}$$

where

$$r_1(n) = \kappa_n \sum_{k=0}^{2n} \frac{s(2n, 2n - k)}{k!} \int_1^\infty \sin 2\pi x \frac{\ln^k x}{x^{2n}} dx, \tag{9a}$$

and

$$r_2(n) = \kappa_n \sum_{k=0}^{2n} \frac{s(2n, 2n - k)}{k!} \int_1^\infty \cos 2\pi x \frac{\ln^k x}{x^{2n}} dx. \tag{9b}$$

The form (8) can always be written as given in the proposition with $\tan 2\pi \phi = -r_1/r_2$ and $m^2 = r_1^2 + r_2^2$.

Remark. The integrals appearing in Eqs. (9) can be written in several different ways via integration by parts, but we do not pursue this here.

Proposition 6. Proposition follows from Eq. (2), the multiplication formula satisfied by P_n [1], and the interchange of two finite sums.

Proposition 7. According to the property $\partial\zeta(s, a)/\partial a = -s\zeta(s + 1, a)$ and the Laurent expansion of $\zeta(s, a)$, we have

$$\sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \frac{\gamma_k(a)}{da} (s-1)^k = -1 - \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \gamma_k(a) s^{k+1}, \tag{10}$$

wherein $d\gamma_0(a)/da = -\psi'(a)$ and ψ' is the trigamma function. If we perform a binomial expansion on the right side of this equation, reorder the sums there, and then equate coefficients of like powers of $s - 1$ on both sides, we arrive at the stated set of equations.

Proposition 8. Valid in the whole complex plane is the form of the Riemann zeta function

$$\zeta(s) = \frac{1}{1 - 2^{1-s}} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{1}{(k+1)^s}, \quad s \neq 1. \tag{11}$$

This expression, due to Hasse [15], can be derived by applying Euler’s series transformation to the alternating zeta function [25].

For purposes of expanding Eq. (11) about $s = 1$ we have

$$1 - 2^{1-s} = - \sum_{j=1}^{\infty} \frac{(-\ln 2)^j}{j!} (s-1)^j, \tag{12a}$$

$$(k+1)^{-s} = \sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{\ell!} \ln^\ell(k+1) \sum_{q=0}^{\ell} \binom{\ell}{q} (s-1)^q, \tag{12b}$$

and

$$[1 - 2^{1-s}]^{-1} = \frac{1}{\ln 2(s-1)} + \frac{1}{2} + \frac{\ln 2}{12} (s-1) - \frac{\ln^3 2}{720} (s-1)^3 + O[(s-1)^5]. \tag{12c}$$

By using the series $\sum_{k=0}^n \binom{n}{k} \frac{(-1)^k}{(k+1)} = \frac{1}{n+1}$ it is easy to see that the $q = 0$ term of Eq. (12b) contributes a $\ln 2$ term in Eq. (11):

$$\sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=0}^n \binom{n}{k} \frac{(-1)^k}{(k+1)} = \ln 2. \tag{13}$$

Therefore, comparing with the Laurent expansion of $\zeta(s)$, we obtain parts (a) and (b).

Similarly, if we write

$$\begin{aligned} \ln \zeta(s) &= -\ln(1 - 2^{1-s}) + \ln \left[\sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{1}{(k+1)^s} \right] \\ &= -\ln(s-1) + \frac{1}{2} \ln 2(s-1) + \frac{1}{24} (s-1)^2 + O[(s-1)^4] \end{aligned}$$

$$+ \ln \left[\sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=0}^n (-1)^k \binom{n}{k} \sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{\ell!} \ln^\ell(k+1) \sum_{q=1}^{\ell} \binom{\ell}{q} (s-1)^q \right], \tag{14}$$

we find $-\eta_0 = \gamma$ and the expression given in part (c).

From Eq. (11), we have

$$\zeta'(s) = -\frac{\ln 2}{(2^{s-1} - 1)} \zeta(s) - \frac{1}{1 - 2^{1-s}} \sum_{n=0}^{\infty} \frac{1}{2^{n+1}} \sum_{k=1}^n (-1)^k \binom{n}{k} \frac{\ln(k+1)}{(k+1)^s}, \quad s \neq 1. \tag{15}$$

Part (d) follows from the well-known value $\zeta'(0)/\zeta(0) = \ln 2\pi$. Similarly, based upon the well-known values of the zeta function at positive even integers, we obtain part (e) for the Bernoulli numbers.

From the functional equation of the zeta function comes the evaluation for positive integers n ,

$$\zeta'(-2n) = (-1)^n \frac{(2n)! \zeta(2n+1)}{2(2\pi)^{2n}}. \tag{16}$$

From the representation (11) and another globally but more slowly convergent series for $\zeta(s)$ due to Hasse [15] we then have the identity of part (f).

3. Discussion

We may note the special values of the Stirling numbers of the first kind [1] $s(n, 0) = \delta_{n0}$ and $s(n, n) = 1$. Therefore, the general integral representation of Eq. (2) can be simplified to

$$C_n(a) = (-1)^{n-1} n! \left[\int_1^{\infty} \frac{P_n(x-a)}{x^{n+1}} dx + \sum_{k=1}^n \frac{s(n+1, n+1-k)}{k!} \int_1^{\infty} P_n(x-a) \frac{\ln^k x}{x^{n+1}} dx \right], \quad n \geq 1. \tag{17}$$

When $a = 1$, the first term on the right side of this equation may be evaluated by using integration by parts on the result of Appendix A.

Despite some of the complexity of the representation (11), we have several motivations for examining it:

- (i) This form gives rise to series that play a significant role in renormalization and quantum field theory [5,12].
- (ii) Equation (11) has connections to random variables in an analytic number theory setting [4]. There is a possibility in particular to link this structure with that of Brownian processes. This brings us to a third point of interest.

(iii) We would like to know not only the structure of the Stieltjes constants but also of the constants $\{\eta_j\}$ that enter the Laurent expansion of the logarithmic derivative of $\zeta(s)$.

Such knowledge would be very influential in deciding the nature of a sum denoted $S_2(n)$ [7–9] and thereby the resulting Li/Keiper constants λ_k [19,22].

As a byproduct of this work we obtain interesting infinite series for fundamental constants such as the Euler constant and $\ln \pi$. The rapidity of convergence may make some of these suitable for applications. Indeed, even a naive numerical implementation of part (a) or (b) of Proposition 8 appears to provide these constants to 16 decimal places after summing over n to 51 or 52 terms.

The generalized digamma function used by Dilcher [10] is given by

$$\psi_k(a) \equiv -\gamma_k - \frac{1}{a} \ln^k a - \sum_{v=1}^{\infty} \left[\frac{\ln^k(v+a)}{v+a} - \frac{\ln^k a}{v} \right], \quad k \geq 0, \tag{18}$$

and we may relate it to other functions defined in the literature. We have $R_m(a) = (-1)^{m+1} (\partial^m / \partial s^m) \zeta(0, a)$ [17] and find that $\psi_k(a) = R'_{k+1}(a) / (k + 1)$. We have $R_j(a) = (-1)^j j! - \sum_{\ell=0}^{\infty} \gamma_{j+\ell}(a) / \ell!$ and by using our Proposition 7, we determine that $\psi_j(a) = -\gamma_j(a)$. This relation explains how our Proposition 3 takes exactly the same form as Dilcher’s Corollary 11 and permits the re-expression of several of his other results as well.

In his computational work based upon Newton–Cotes integration for the high accuracy approximation of the Stieltjes constants, Kreminski observed that for large values of k , $\gamma_k(1/2) \approx -\gamma_k$ and that more generally $C_k(a + 1/2) \approx -C_k(a)$. That is, in a sense, the Stieltjes constants for large index are anti-periodic with period 1/2. Our Proposition 1 makes this precise.

Much earlier, Hansen and Patrick [14] showed that a fundamental interval in a for the Hurwitz zeta function need only be of length 1/2. Perhaps solely on that basis one would then suspect that some sort of relationship(s) should exist between $\gamma_k(a)$ and $\gamma_k(a + 1/2)$.

Much of our development relies on the underlying theory of periodized Bernoulli polynomials and corresponding integral representations of the Stieltjes constants [30]. Equations such as (17) and still others that we have obtained help to expose more of the analytic structure of the Stieltjes and η_j constants.

Appendix A. Integrals over periodic Bernoulli polynomials

Herein we evaluate integrals over the periodic Bernoulli polynomial P_1 in terms of polygamma functions $\psi^{(j)}$. We demonstrate

Proposition. For positive integers n and m ,

$$\begin{aligned} & \frac{(-1)^n}{(n+1)!} \psi^{(n)}(s) + \frac{1}{(n+1)} \sum_{k=0}^m \frac{1}{(x+k)^{n+1}} \\ &= \int_m^{\infty} \frac{P_1(x)}{(x+s+1)^{n+2}} dx + \frac{1}{(n+1)} \frac{1}{(s+m)^{n+1}} - \frac{1}{n(n+1)} \frac{1}{(s+m+1)^n} \end{aligned}$$

$$-\frac{1}{2} \frac{1}{(n+1)} \frac{1}{(s+m+1)^{n+1}}. \tag{A.1}$$

Our starting point is [11]

$$\ln \Gamma(s+1) = (s+1/2) \ln s - s + \frac{1}{2} \ln 2\pi - \int_0^\infty \frac{P_1(x)}{(x+s)} dx, \tag{A.2}$$

so that

$$\psi(s+1) = \ln s + \frac{1}{2s} + \int_0^\infty \frac{P_1(x)}{(x+s)^2} dx, \tag{A.3}$$

where $\psi(s) = \psi(s+1) - 1/s$ is the digamma function. We now take n derivatives of relation (A.3), obtaining

$$\frac{(-1)^n}{(n+1)!} \psi^{(n)}(s+1) = -\frac{1}{n(n+1)} \frac{1}{s^n} + \frac{1}{2} \frac{1}{(n+1)} \frac{1}{s^{n+1}} + \int_0^\infty \frac{P_1(x)}{(x+s)^{n+2}} dx. \tag{A.4}$$

Finding that

$$\begin{aligned} \int_0^\infty \frac{P_1(x)}{(x+s)^{n+2}} dx &= \int_1^\infty \frac{P_1(x)}{(x+s)^{n+2}} dx + \frac{1}{n(n+1)} \left[\frac{1}{s^n} - \frac{1}{(s+1)^n} \right] \\ &\quad + \frac{1}{2(n+1)} \left[\frac{1}{s^{n+1}} - \frac{1}{(s+1)^{n+1}} \right], \end{aligned} \tag{A.5}$$

we obtain

$$\begin{aligned} \frac{(-1)^n}{(n+1)!} \psi^{(n)}(s+1) &= \frac{1}{(n+1)} \frac{1}{s^{n+1}} - \frac{1}{n(n+1)} \frac{1}{(s+1)^n} - \frac{1}{2} \frac{1}{(n+1)} \frac{1}{(s+1)^{n+1}} \\ &\quad + \int_1^\infty \frac{P_1(x)}{(x+s)^{n+2}} dx. \end{aligned} \tag{A.6}$$

By making a simple change of variable in the integral of Eq. (A.6) and using the periodicity of P_1 , we have

$$\begin{aligned} &\frac{(-1)^n}{(n+1)!} \psi^{(n)}(s+m+1) \\ &= \frac{1}{(n+1)} \frac{1}{(s+m)^{n+1}} - \frac{1}{n(n+1)} \frac{1}{(s+m+1)^n} - \frac{1}{2} \frac{1}{(n+1)} \frac{1}{(s+m+1)^{n+1}} \\ &\quad + \int_m^\infty \frac{P_1(x)}{(x+s+1)^{n+2}} dx. \end{aligned} \tag{A.7}$$

By applying the functional equation of the polygamma function [1], we obtain the proposition.

When applying Eq. (A.1), it is useful to keep in mind a relation [13,26] between the polygamma function and the Hurwitz zeta function: $\psi^{(n)}(x) = (-1)^{n+1} n! \zeta(n+1, x)$. In particular, when $x = 1/2$ or $x = 1$ in the latter relation, values of the Riemann zeta function appear. In fact, another starting point for evaluating integrals similar to (A.1) is to use integral representations for $\zeta(s)$ [1,16,26,28] in terms of $P_1(x) + 1/2$.

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