

Asymptotic stability of the Toda m -soliton

G.N. Benes^{a,b}, A. Hoffman^{c,*}, C.E. Wayne^{a,b}

^a Boston University, Department of Mathematics and Statistics, 111 Cummington Ave., Boston, MA 02215, United States

^b Center for BioDynamics, Boston University, 111 Cummington Ave., Boston, MA 02215, United States

^c Franklin W. Olin College of Engineering, 1000 Olin Way, Needham, MA 02492, United States

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ABSTRACT

We prove that multi-soliton solutions of the Toda lattice are both linearly and nonlinearly stable in exponentially weighted spaces. Our proof uses neither the inverse spectral method nor the Lax pair of the model but instead studies the linearization of the Bäcklund transformation which links the $(m - 1)$ -soliton solution to the m -soliton solution. We use this to construct a conjugation between the Toda flow linearized about an m -soliton solution and the Toda flow linearized about the zero solution, whose stability properties can be determined by explicit calculation.

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1. Introduction

Solitary waves or solitons are a common feature of infinite dimensional dispersive dynamical systems in which the effects of dispersion are exactly off-set by the effects of the nonlinearity. They are also important features of physical systems ranging from water waves to nonlinear optics. As such, questions of stability and interaction of such waves are natural and important. One class of systems which is amenable to study is lattices of coupled, nonlinear oscillators like the Fermi–Pasta–Ulam model. In the present paper we study the stability of multi-soliton (m -soliton) solutions of a particular lattice system, the Toda system, which features an exponential interaction potential between adjacent oscillators.

More precisely, this paper establishes the stability of m -soliton solutions to the Toda lattice equations

$$\ddot{Q} = e^{-(Q-Q_-)} - e^{-(Q_+-Q)}, \tag{1}$$

where $Q = (Q_n)_{n \in \mathbb{Z}}$ is the vector of positions of unit masses in an infinite chain of particles with nearest-neighbor interaction potential $V(x) = e^{-x} - (1 - x)$. Here and throughout the paper X_{\pm} denotes a left or right shift of X , i.e. $(X_{\pm})_n = X_{n \pm 1}$. A solution of (1) of the form $Q_n(t) = \phi(n - ct)$ is called a *traveling wave solution*. In the case that ϕ is monotone and the limits $\phi(\pm\infty)$ are defined, we say that ϕ is a *kink* or *front*. Since (1) is invariant with respect to adding a constant, we can, without loss of generality, restrict attention to the case $\phi(-\infty) = 0$. A special class of front solutions of (1), the solitons, is given explicitly by $Q_n^1(t; \kappa, \gamma) = \log \frac{\cosh(\kappa n - t \sinh \kappa + \gamma)}{\cosh(\kappa(n+1) - t \sinh \kappa + \gamma)} - \kappa$. Note that they form a two-dimensional manifold parameterized by κ and γ , with asymptotic limit $\lim_{n \rightarrow \infty} Q_n^1(t; \kappa, \gamma) = -2\kappa$.

It is well known that the Toda lattice is an integrable system and hence admits multi-soliton solutions [17,1]. To be concrete, in our context an m -soliton solution is a solution Q^m for which there are $3m$ constants $\gamma_1^{\pm}, \dots, \gamma_m^{\pm}$ and $\kappa_1, \dots, \kappa_m$ such that

* Corresponding author.

E-mail address: aaron.hoffman@olin.edu (A. Hoffman).

$$\lim_{t \rightarrow \pm\infty} \left| Q_n^m(t) - \sum_{j=1}^m Q_n^1(t; \kappa_j, \gamma_j^\pm) \right| = 0.$$

In practice $(\gamma_1^+, \dots, \gamma_m^+)$ are determined by $(\gamma_1^-, \dots, \gamma_m^-)$ and $(\kappa_1, \dots, \kappa_m)$ so that the space of m -soliton solutions is $2m$ -dimensional.

The asymptotic behavior of the Toda lattice has previously been considered with the aid of the inverse scattering method [8,7]. Using that technique, Krüger and Teschl showed that given initial data which decays sufficiently rapidly at infinity, the solution of (1) converges pointwise in space, algebraically in time, to an m -soliton plus dispersive modes. Our approach here is quite different. On one hand, we can consider only initial conditions which lie in a neighborhood of some m -soliton, in some weighted function space. However, we obtain a stronger convergence result – namely, the convergence is exponential in time, uniformly in the exponentially weighted space in which the initial data lies. Furthermore, as described above, the method of our proof avoids the inverse spectral method and instead relies on the existence of a Bäcklund transform which connects the m - and $(m - 1)$ -solitons. One advantage of this approach, in our opinion, is that it seems more amenable to extensions to nearby, non-integrable models (see [12,5]).

The main result of this paper is that these m -soliton solutions are asymptotically stable in a sense that we will make precise in Section 2. We do not depend on the existence of a Lax pair or inverse scattering method, but we do use in an essential way the existence of a Bäcklund transformation which relates m - and $(m - 1)$ -solitons [17]:

$$\begin{aligned} P + e^{-(Q' - Q - \kappa_m)} + e^{-(Q - Q' + \kappa_m)} &= 2 \cosh \kappa_m, \\ P' + e^{-(Q' - Q - \kappa_m)} + e^{-(Q + Q' + \kappa_m)} &= 2 \cosh \kappa_m. \end{aligned} \tag{2}$$

Here (Q, P) is an m -soliton solution, (Q', P') is an $(m - 1)$ -soliton solution, and the speed of the additional soliton in (Q, P) is given by $\frac{\sinh \kappa_m}{\kappa_m}$ [18]. Given (Q', P') , Eqs. (2) together with the relations $\dot{Q} = P$ and $\dot{Q}' = P'$ determine (Q, P) up to a constant of integration corresponding to the phase of the new soliton γ_m . The corresponding m -soliton Q has the asymptotic limits $Q_{-\infty} = 0$ and $Q_\infty = -2 \sum_{j=1}^m \kappa_j$.

This paper contains two main results. Theorem 1.1 establishes linear stability for the m -soliton solution. Theorem 1.2 establishes nonlinear stability. The proof that linear stability implies nonlinear stability is based on the now classical theory of modulation equations as developed in [15,3].

The proof of linear stability is based on a more recent idea of using the *linearization* of the Bäcklund transformation to construct a conjugacy between the linearization of the flow of the Toda equations around an m -soliton and the linearization around an $(m - 1)$ -soliton. The existence of such a conjugacy between the linearization about a single soliton and the linearization about the zero solution was implicitly exploited in [10] in the context of the KdV equation and developed in detail for the Toda lattice in [13]. More recently the general case of the linearization around an m -soliton of the KdV equation has been considered in [12].

1.1. Hamiltonian formulation

It is convenient to define $R = Q_+ - Q = (S - I)Q$. Here S is the shift operator and throughout the paper we will use the subscript “+” to denote the action of this operator – i.e. $(Sx) = x_+$. We also define $U = (R, P)$. Then (1) can be rewritten in the Hamiltonian form

$$\dot{U} = JH'(U). \tag{3}$$

Here the symplectic operator J is given by $J = \begin{pmatrix} 0 & S^{-1} \\ I - S^{-1} & 0 \end{pmatrix}$. The Hamiltonian is defined by $H = \sum_{n \in \mathbb{Z}} (\frac{1}{2} P_n^2 + V(R_n))$ where the interaction potential is given by $V(R) = e^{-R} - 1 + R$. Define the weighted norm $\|x\|_a = \|x\|_{\ell_a^2}^2 = \sum_{n \in \mathbb{Z}} e^{2an} x_n^2$ and let ℓ_a^2 denote the associated weighted Hilbert space. Note that we will use $\|x\|$ to denote the standard ℓ^2 norm. We abuse notation and write $\|x\|_{c,a} := e^{-act} \|x\|_a = (\sum_n e^{2a(n-ct)} x_n^2)^{1/2}$ for the weighted norm of x with the weight moving with speed c to the right.

Let $\Phi(t, s)$ be the evolution operator associated with the linearized Toda flow $\dot{u} = JH''(U^m)u$, where U^m is the m -soliton solution written in terms of the (R, P) variables. Define the following subspace of $\ell_a^2 \times \ell_a^2$:

$$X_m(s) := \{u \in \ell_a^2 \times \ell_a^2 \mid \langle u, J^{-1} \partial_{\kappa_i} U^m(s) \rangle = \langle u, J^{-1} \partial_{\gamma_i} U^m(s) \rangle = 0, i = 1, \dots, m\}. \tag{4}$$

Here $\langle \cdot, \cdot \rangle$ is the standard ℓ^2 inner product. The condition $\langle x, J^{-1}y \rangle = 0$ can be thought of as a symplectic orthogonality condition, as discussed in [3].

1.2. Main results

We begin with our results on decay of solutions of the linearized Toda equations. These results depend on the fact that dispersive waves in the Toda equation (i.e. solutions of the equations linearized about the zero solution) all propagate

with speed less than or equal to one, while all solitary waves move with speeds strictly greater than one. In particular, for the linearization of the equations about the zero solution an easy and explicit computation shows that solutions decay exponentially in time in any of the weighted Hilbert spaces ℓ_a^2 if the weight translates to either the left or right with a speed greater than one (a detailed proof of this fact is provided in Lemma 3 of [13]). Intuitively, one can imagine that the region in which the norm “sees” the solution just “outruns” the dispersive perturbations. If we now use the fact that the linearized Bäcklund transformation conjugates the semigroup of the linearization about the zero solution to the linearization about the multi-soliton (provided the orthogonality conditions in (4) are satisfied) we see that we also expect exponential decay of solutions of the equations linearized about the multi-solitons, provided the weight in the norm moves faster than the dispersive part of the solution (i.e., faster than one) but slower than the solitons themselves.

Theorem 1.1. *Let $\kappa_i > 0$ and $\gamma_i \in \mathbb{R}$ be given for $i = 1, \dots, m$. Define $\kappa_{\min} = \min\{\kappa_i \mid i = 1, \dots, m\}$. Let $c > 1$, let $a \in (0, 2\kappa_{\min})$, and let $\beta := ca - 2 \sinh(a/2)$.*

Let $\Phi(t, x)$ be the evolution operator for the linearization of the Toda equations about an m -soliton solution with parameters κ_i and $\gamma_i, i = 1, \dots, m$. Then there exists a constant $K > 0$ such that for any $u_0 \in X_m(s)$ and for all $s \geq t$,

$$\|e^{a(n-ct)} \Phi(t, s) u_0\| \leq K e^{-\beta(t-s)} \|e^{a(n-cs)} u_0\|.$$

Remark. We note the perturbation only decays for $\beta > 0$, which corresponds to choices of a and c such that $c > \frac{\sinh a/2}{a/2}$.

Theorem 1.2. *Let $c > 1, \kappa_i, \gamma_i$ and β be as in Theorem 1.1. There is an $a_{\max} \in (0, 2\kappa_{\min})$ such that for each $a \in (0, a_{\max})$ and $\beta' \in (0, \beta)$, there is a $\delta > 0$ with the property that if u_0 satisfies*

$$\|e^{a(\cdot)} u_0\| + \|u_0\| < \delta,$$

then associated to the solution $u(t)$ of Toda equations with initial data $u_0 + U^m(t_0; \kappa_1, \dots, \gamma_m)$, there exist real valued functions of a real variable $\gamma_i(t)$ and $\kappa_i(t)$ for $i = 1, \dots, m$ such that

$$\|e^{a(\cdot-ct)} [u(t) - U^m(t; \kappa(t), \gamma(t))]\| \leq K e^{-\beta'(t-t_0)}$$

and

$$\|u(t) - U^m(t; \kappa(t), \gamma(t))\|^2 + \sum_{i=1}^m |\kappa_i(t) - \kappa_i(t_0)| + |\gamma_i(t) - \gamma_i(t_0)| \leq K \|u(t_0) - U^m(t_0; \kappa(t_0), \gamma(t_0))\|^2.$$

Moreover, the limits $\lim_{t \rightarrow \infty} \gamma_i(t) = \gamma_i(\infty)$ and $\lim_{t \rightarrow \infty} \kappa_i(t) = \kappa_i(\infty)$ exist. Furthermore, the same result holds in backward time upon replacing a with $-a$.

Remark. Note that we restrict our initial perturbation to lie in a weighted space. We believe that a more refined decomposition of the perturbation as a sum of perturbations to the trailing k -solitons as in [12] could lead to orbital stability in the energy space and with constants chosen uniformly in the parameters γ_i and κ_i .

Stability results for multi-soliton solutions in integrable systems date back to Maddocks and Sachs [9] where KdV multi-solitons are shown to be constrained minimizers of a Lyapunov function. While a similar approach was shown to work for NLS-type equations [6], we know of no comparable result for the Toda lattice. An alternate approach, developed (for example) in [11] decomposes the multi-soliton solution into several weakly interacting pieces, each of which is localized about a soliton and then leverages stability theory for a single solitary wave. This approach extends to the FPU lattices [12] and hence to the Toda lattice as well. Its use thus far has been restricted to initial data corresponding to a train of solitons which will not interact in forward time. We are able to show stability for m -soliton solutions even when the interaction between two or more constituent solitons lies in the future, albeit only when the perturbation is sufficiently small in our exponentially weighted space.

One drawback of the Bäcklund transform method that we use here, is that unlike the variational/dispersive estimate approach in [11] it does not extend immediately to the non-integrable case. However, there has been some recent progress in extending this type of idea to near-integrable systems, and in particular to the Fermi–Pasta–Ulam (FPU) lattice in the long-wave, low-amplitude regime. The second two authors have used the linearized Bäcklund transformation to obtain a simplified proof of stability for the case of a single soliton in FPU [5] and to study two counter-propagating solitary waves [4]. Recently in [12], Mizumachi used linearized Bäcklund transformations for the KdV equation linearized about multi-soliton solutions to establish linear (and ultimately nonlinear stability) of well-separated N -soliton solutions in the FPU model.

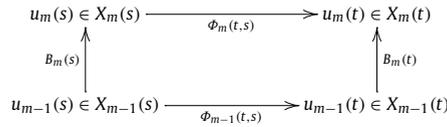


Fig. 1. Commuting diagram used in the induction step.

2. Linear stability of the Toda m -soliton solution

In [13] the authors used the Bäcklund transformation along with its linearizations about the zero- and one-soliton solutions to show that linear stability of the zero solution of the Toda lattice implies the linear stability of the one-soliton solution. In this section we show that this can be regarded as the first step in an induction argument: the stability of the $(m - 1)$ -soliton solution implies that of the m -soliton solution. To proceed we must show (i) that the linearized Bäcklund transformation commutes with the linearized Toda flow, (ii) that the linearized Bäcklund transformation preserves orthogonality with the neutral modes of the linearized Toda system, and (iii) that the linearized Bäcklund transformation is an isomorphism between two spaces (defined below) that depend on these neutral modes. These facts tell us that for each m , the diagram in Fig. 1 commutes.

In Fig. 1, $B_m(\tau)$ is an isomorphism derived from the linearized Bäcklund transform at time τ , $\Phi_k(t, s)$ is the evolution according to the linearization of the Toda equation about the solution U^k , and X_k is given as in (4). We must make X_k a proper subspace of $\ell_a^2 \times \ell_a^2$, in our case by imposing orthogonality conditions, because any perturbation which merely changed the relative amplitude or phase of the different component solitons results in a new multi-soliton solution which cannot converge back to the original one. All other perturbations save those of this type, however, should decay.

Throughout the remainder of this section, (Q', P') and (Q, P) correspond to U_{m-1} and U_m , respectively, for some m . Finally, let (q', p') and (q, p) denote solutions to the linearizations of the Toda differential equations about (Q', P') and (Q, P) , respectively.

There will be several places where we make an argument that uses the fact that an m -soliton solution converges to the sum of m constituent solitons as t tends to either plus or minus infinity. Although this fact is fundamental and well known to experts, we are not aware of a proof that exists in the literature, so we state and prove the following:

Lemma 2.1 (Resolution of m -soliton solutions into one-solitons). *An m -soliton solution converges to a sum of solitons exponentially fast in ℓ^1 as $t \rightarrow \pm\infty$. More concretely: Let $Q^m(t; \kappa_1, \dots, \kappa_m, \gamma_1, \dots, \gamma_m)$ denote the m -soliton solution. Then there exist phase shifts $\zeta_1^\pm, \dots, \zeta_m^\pm$ and $\lambda_0 > 0$ such that*

$$\lim_{t \rightarrow \pm\infty} e^{\lambda_0|t|} \left\| Q^m(t; \kappa_1, \dots, \kappa_m) - \sum_{j=1}^m Q^1(t; \kappa_j, \gamma_j + \zeta_j^\pm) \right\|_{\ell^1} = 0. \tag{5}$$

Proof. See Appendix A. \square

We also show that the derivative of an m -soliton solution with respect to a parameter κ_l or γ_l converges exponentially fast to the same derivative of the phase-shifted l th constituent soliton.

Corollary 2.2. *The derivative of an m -soliton solution with respect to one of its parameters γ_j or κ_j converges exponentially fast to the derivative of the l th constituent soliton in an exponentially weighted space moving in the frame of the l th soliton. More concretely: let $Q^m(t; \kappa_1, \dots, \kappa_m, \gamma_1, \dots, \gamma_m)$ denote the m -soliton solution and let $\zeta_1^\pm, \dots, \zeta_m^\pm$ denote the phase shifts described above. There is a $\delta > 0$ and an $a_0 > 0$, both depending upon the soliton parameters, such that if ∂_{\circ_l} describes the derivative with respect to either γ_l or κ_l , then for any $a \in (0, a_0)$ we have*

$$\lim_{t \rightarrow \infty} e^{\delta t} \left\| \partial_{\circ_l} Q^m(t; \kappa_1, \dots, \kappa_m) - \partial_{\circ_l} Q^1(t; \kappa_l, \gamma_l + \zeta_l(\kappa_1, \dots, \kappa_m)) \right\|_{\mathcal{C}_l, \pm a} = 0.$$

Proof. See Appendix A. \square

2.1. Linearized Bäcklund transformation commutes with linearized Toda flow

It is convenient in this subsection to define new variables

$$\alpha := e^{-(Q' - Q - \kappa_m)}, \quad \beta := e^{-(Q - Q' + \kappa_m)}. \tag{6}$$

With these variables we can rewrite the Toda equations as

$$\begin{aligned} \dot{Q} &= P, & \dot{Q}' &= P', \\ \dot{P} &= \alpha_- \beta - \alpha \beta_+, & \dot{P}' &= \alpha \beta - \alpha_+ \beta_+, \end{aligned} \tag{7}$$

their linearization about (Q, P) and (Q', P') as

$$\begin{aligned} \dot{q} &= p, & \dot{q}' &= p', \\ \dot{p} &= (I - S^{-1})\alpha\beta_+(S - I)q, & \dot{p}' &= (I - S^{-1})\alpha_+\beta_+(S - I)q, \end{aligned} \tag{8}$$

and the Bäcklund transformation as

$$\begin{aligned} F_1 &:= P + \alpha + \beta - 2 \cosh \kappa_m = 0, \\ F_2 &:= P' + \alpha + \beta_+ - 2 \cosh \kappa_m = 0. \end{aligned} \tag{9}$$

Linearizing (9) about (Q', P', Q, P) gives us the linearized Bäcklund transformation

$$p + (\alpha - \beta)q + (\beta S^{-1} - \alpha)q' = p' + (\alpha - \beta_+ S)q + (\beta_+ - \alpha)q' = 0 \tag{10}$$

which can also be written as $DF_1(q', p', q, p) = DF_2(q', p', q, p) = 0$.

Proposition 2.3. *Suppose that $(q'(0), p'(0))$ and $(q(0), p(0))$ are related via the linearized Bäcklund transform (10). Then $(q'(t), p'(t))$ and $(q(t), p(t))$ are related via (10) for all $t \in \mathbb{R}$.*

Proof. We claim that there are linear operators A_i, B_i such that

$$\dot{D}F_1 = A_1 DF_1 + B_1 DF_2; \quad \dot{D}F_2 = A_2 DF_1 + B_2 DF_2 \tag{11}$$

holds. Assume the claim. Then since (DF_1, DF_2) satisfies a linear system of equations and $(DF_1(0), DF_2(0)) = (0, 0)$, it follows that $(DF_1(t), DF_2(t)) = (0, 0)$ for all time.

It remains to prove the claim. As the expressions $\dot{\alpha}$ and $\dot{\beta}$ appear in $\dot{D}F_1$ and $\dot{D}F_2$, we differentiate (6) with respect to time and make use of (9) to obtain

$$\dot{\alpha} = \alpha(\beta_+ - \beta); \quad \dot{\beta} = \beta(\alpha - \alpha_-).$$

A calculation now shows that the choice $A_1 = \alpha - \beta, B_1 = \beta S^{-1} - \alpha, A_2 = \alpha - \beta_+ S,$ and $B_2 = \beta_+ - \alpha$ yields equality in (11). \square

2.2. Symplectic orthogonality

The linearization of the Toda equations about an m -soliton solution contains a zero-eigenspace of dimension $2m$. We expect decay of solutions of the linearized equations only if they lie in the complement of this subspace. This condition is incorporated in the definition of our function spaces $X_m(t)$ through the inner products with derivatives of the m -soliton with respect to speed and phase. In order to insure that the sums in these inner products converge we need some information about the spatial asymptotic behavior of these derivatives. This information is contained in the following lemma.

Lemma 2.4 (Spatial asymptotics for neutral modes). *The operator $x \mapsto x_+ - \alpha\beta_+^{-1}x$ is Fredholm with index one and a one-dimensional kernel when regarded as acting on any of the spaces $\ell_a^2, \ell_{-a}^2,$ or $\ell_a^2 \cap \ell_{-a}^2$. Furthermore the norm of its inverse (acting on the orthogonal complement of its kernel) is bounded uniformly in t . Additionally its kernel is spanned by $\partial_{y_m} Q$.*

Moreover $\partial_{y_l} Q$ and $\partial_{y_l} P$ lie in $\ell_{-a}^2 \cap \ell_a^2$ and $\partial_{\kappa_l} Q$ and $\partial_{\kappa_l} P$ live in ℓ_{-a}^2 for $a \in (0, 2\kappa_m)$ and $1 \leq l \leq m$.

Proof. We first examine the Fredholm properties of $x \mapsto x_+ - \alpha\beta_+^{-1}x$. This is a non-autonomous linear first order recurrence, which is moreover asymptotically hyperbolic. At minus infinity the origin for $x_+ = e^{2\kappa_m x}$ is unstable (and since it is one-dimensional this means it has a one-dimensional unstable manifold). At infinity, the origin for $x_+ = e^{-2\kappa_m x}$ is stable (and hence has a zero-dimensional unstable manifold). It follows from Palmer’s theorem and its discrete analogs (e.g. [14]) that the difference operator $x \mapsto x_+ - \alpha\beta_+^{-1}x$ is Fredholm with index one and a one-dimensional kernel when regarded as an operator on any of the spaces $\ell_{-a}^2, \ell_a^2,$ or $\ell_a^2 \cap \ell_{-a}^2$, so long as the weight does not change the stability type of origin at $\pm\infty$, i.e. so long as $a \in [0, 2\kappa_m)$. This establishes the first claim of the lemma.

More concretely, define the semigroup $T(n) = e^{-\sum_{k=0}^{n-1} (2Q'_k - Q_k - Q_{k+1} - 2\kappa_m)}$ so that $x_{n+1} - \frac{\alpha_n}{\beta_{n+1}}x_n = y_n$ has solution $x_n = T(n)x_0 + \sum_{k=0}^{n-1} T(n-k)y_k$. For simplicity let $x_n^* = \sum_{k=0}^n T(n-k)y_k$ and let v be a unit vector which spans the kernel. Then we can write $x = x^* + v\nu$ where $\nu = \frac{\langle x^*, v \rangle}{\|v\|^2}$. Thus $\|x\| \leq (1 + \frac{1}{\|v\|})\|x^*\| = 2\|x^*\|$. It follows from the Hausdorff–Young inequality for convolutions that $\|x^*\|_a \leq \|e^a T\|_{\ell^1} \|y\|_a$ and hence the norm of the inverse restricted to the orthogonal complement of the kernel is bounded above by $2\|e^a T\|_{\ell^1}$, so long as it is finite. Note that T has the asymptotic form

$$T(n) = \begin{cases} \mathcal{O}(e^{-2\kappa_m n}), & n \rightarrow \infty, \\ \mathcal{O}(e^{2\kappa_m n}), & n \rightarrow -\infty, \end{cases} \tag{12}$$

so in particular $\|e^a T\|_{\ell^1}$ is finite. Moreover, in light of Lemma 2.1, as $t \rightarrow \pm\infty$, the semigroup T approaches the semigroup T^\pm corresponding to the single-soliton problem. This semigroup, studied in [4] satisfies $\|e^a T^\pm\|_{\ell^1} < \infty$. Thus we have the bound $\|e^a T\| < K$ uniformly in time. This establishes the second claim of the lemma.

We now establish that the derivatives of Q and P lie in ℓ^2_{-a} . The proof is based on differentiating (2) and solving the resulting first order recurrence. The details are below. Let ∂ denote ∂_{γ_l} or ∂_{κ_l} for some $1 \leq l \leq m$. Differentiate the second line of (2) and solve for ∂Q_+ to obtain

$$\partial Q_+ = \alpha \beta_+^{-1} \partial Q + \beta_+^{-1} \partial P' + (1 - \alpha \beta_+^{-1}) \partial Q' + (\alpha \beta_+^{-1} - 1 - 2 \sinh \kappa_m \beta_+^{-1}) \chi_{\partial=\partial_{\kappa_m}}.$$

This is an inhomogeneous first order recurrence whose homogeneous part: $\partial Q_+ = \alpha \beta_+^{-1} \partial Q$ we have studied above. Note that in the case $\partial = \partial_{\gamma_m}$ the inhomogeneous term is zero, thus $\partial_{\gamma_m} Q$ lies in the kernel as desired. More generally, the inhomogeneous term has the asymptotic form $e^{-\kappa_m} \partial P' + (1 - e^{-2\kappa_m}) \partial Q' - 2(1 - e^{-2\kappa_m}) \chi_{\partial=\partial_{\kappa_m}}$ as $n \rightarrow \infty$ and $e^{\kappa_m} \partial P' + (1 - e^{2\kappa_m}) \partial Q' + e^{-(Q'_n - Q_{n+1})} (e^{2\kappa_m} - 1) (e^{-(Q'_n - Q_n)} - 1) \chi_{\partial=\partial_{\kappa_m}}$ as $n \rightarrow -\infty$. Make the induction hypothesis that $\partial_{\gamma_l} Q'$ and $\partial_{\gamma_l} P'$ lie in $\ell^2_{-a} \cap \ell^2_a$ while $\partial_{\kappa_l} Q'$ and $\partial_{\kappa_l} P'$ lie in ℓ^2_{-a} . Noting that the derivative of κ_m with respect to any of the parameters γ_l or κ_l is either one or zero and $Q'_n - Q_n = \mathcal{O}(e^{2\kappa_m n})$ as $n \rightarrow -\infty$, it follows that the inhomogeneous term lies in ℓ^2_{-a} , and hence that ∂Q lies in ℓ^2_{-a} as well. To see that $\partial P \in \ell^2_{-a}$ simply differentiate the first equation in (2) to obtain

$$\partial P = (\beta - \alpha) \partial Q + \alpha \partial Q' - \beta \partial Q'_- + (\beta - \alpha + 2 \sinh \kappa_m) \chi_{\partial=\partial_{\kappa_m}}.$$

In light of the fact that the coefficient on $\chi_{\partial=\partial_{\kappa_m}}$, which is $e^{\kappa_m} (1 - e^{-(Q' - Q)}) - e^{-\kappa_m} (1 - e^{-(Q - Q'_-)})$, goes to zero rapidly as $n \rightarrow -\infty$ and is bounded as $n \rightarrow \infty$, it follows that ∂P is a linear combination of vectors in ℓ^2_{-a} , hence in ℓ^2_{-a} itself. In the case that $\partial = \partial_{\gamma_l}$ we have that $\partial_{\gamma_l} P$ is a linear combination of vectors in $\ell^2_a \cap \ell^2_{-a}$ hence is itself in that intersection.

It remains only to verify the base case for the induction. In the case $m = 1$, we have the vectors $Q' = P' = 0$, hence $\partial Q' = \partial P' = 0 \in \ell^2_{-a}$. This completes the proof. \square

In light of the above together with the Cauchy–Schwartz inequality for the $\ell^2_a - \ell^2_{-a}$ dual pairing it follows that the linear map $(q, p) \mapsto \langle q, \partial P \rangle - \langle p, \partial Q \rangle$ is continuous when regarded as a map from $\ell^2_a \times \ell^2_{-a} \rightarrow \mathbb{R}$. Due to the form of the Bäcklund transformation, it is easier to work with the variables q and p than it is to work with u , so before proceeding we will relate the orthogonality conditions in (4) to q and p .

Lemma 2.5. *Let $U = (R, P) = ((S - I)Q, P)$ be a k -soliton solution of the Toda lattice, and let $u = (r, p) = ((S - I)q, p)$. Then*

$$\langle u, J^{-1} \partial U \rangle = \langle p, \partial Q \rangle - \langle q, \partial P \rangle,$$

where ∂ denotes a partial derivative of U with respect to either κ_j or γ_j , for some $j = 1, \dots, k$.

Proof. We compute

$$\begin{aligned} \langle u, J^{-1} \partial U \rangle &= \left\langle (S - I)q, \sum_{k=-\infty}^0 S^k \partial P \right\rangle + \left\langle p, \sum_{k=-\infty}^{-1} S^k \partial (S - I)Q \right\rangle \\ &= \left\langle q, (S^{-1} - I) \sum_{k=-\infty}^0 S^k \partial P \right\rangle + \left\langle p, \sum_{k=-\infty}^{-1} S^k (S - I) \partial Q \right\rangle \\ &= -\langle q, \partial P \rangle + \langle p, \partial Q \rangle. \end{aligned}$$

Setting both sides to zero proves the lemma. \square

We are now able to redefine the spaces X_k in terms of our preferred (for this subsection) q and p variables:

$$X_k(t) = \{(q, p) \in \ell^2_a \times \ell^2_{-a} \mid \langle q, \partial P(t) \rangle - \langle p, \partial Q(t) \rangle = 0 \text{ for } \partial = \partial_{\gamma_l} \text{ or } \partial_{\kappa_l} \text{ with } 1 \leq l \leq k\}.$$

2.3. Linearized Bäcklund transformation preserves orthogonality conditions

In order to establish the commutativity of the diagram in Fig. 1, we must show that the map B takes X_{m-1} to X_m . In particular, we must show that if (q, p) and (q', p') are related by (10) with $(q', p') \in X_{m-1}$, then $(q, p) \in X_m$.

It is convenient to define the operators

$$\begin{aligned} C &:= \alpha - \beta S^{-1}, & \hat{C} &:= \alpha - S\beta = -\beta_+(S - \alpha\beta_+^{-1}), \\ L &:= \alpha - \beta, & M &:= \alpha - \beta_+, \end{aligned} \quad (13)$$

where α and β are as in the previous subsection. As we shall see, these operators arise naturally both when differentiating and when linearizing the Bäcklund transformation (2). Since we are interested in the orthogonality conditions from (4), we differentiate (2) with respect to γ_i or κ_i , $i = 1, \dots, m-1$, to get

$$\begin{aligned} C\partial Q' &= L\partial Q + \partial P, \\ \partial P' &= \hat{C}\partial Q + M\partial Q', \end{aligned} \quad (14)$$

where $\partial = \partial_{\gamma_i}$ or ∂_{κ_i} , $1 \leq i < m$. These equations help us to prove the following:

Proposition 2.6. *Assume U' and U are related by (2) with u' and u related by (10). Then $\langle u', J^{-1}\partial U' \rangle = \langle u, J^{-1}\partial U \rangle$, where $\partial = \partial_{\kappa_i}$ or ∂_{γ_i} with $1 \leq i < m$.*

Proof. By Lemma 2.5 we need to show that $\langle p', \partial Q' \rangle - \langle q', \partial P' \rangle = \langle p, \partial Q \rangle - \langle q, \partial P \rangle$. Using Eqs. (15)–(14) we compute

$$\begin{aligned} \langle p, \partial Q \rangle &= \langle Cq' - Lq, \partial Q \rangle \\ &= \langle q', \hat{C}\partial Q \rangle - \langle q, L\partial Q \rangle \\ &= \langle q', \partial P' - M\partial Q' \rangle - \langle q, C\partial Q' - \partial P \rangle \\ &= \langle q', \partial P' - M\partial Q' \rangle - \langle \hat{C}q, \partial Q' \rangle + \langle q, \partial P \rangle \\ &= \langle q', \partial P' \rangle - \langle Mq' + \hat{C}q, \partial Q' \rangle + \langle q, \partial P \rangle \\ &= \langle q', \partial P' \rangle - \langle p', \partial Q' \rangle + \langle q, \partial P \rangle, \end{aligned}$$

establishing the desired identity. \square

2.4. Linearized Toda flow preserves orthogonality conditions

In this subsection we consider a solution $(q(t), p(t))$ to (8) which for some $t_0 \in \mathbb{R}$ lies in $X_k(t_0)$. We show that $(q(t), p(t))$ necessarily lies in $X_k(t)$ for all $t \in \mathbb{R}$. This implies that the maps corresponding to horizontal arrows in the commuting diagram in Fig. 1 take $X_k(s)$ to $X_k(t)$ (for $k = m-1, m$).

Proposition 2.7. *Let (Q, P) be a solution to (1) and let $(q, p) \in \ell_a^2 \times \ell_a^2$ be a solution to (8). Then the quantity*

$$\langle p, \partial Q \rangle - \langle q, \partial P \rangle$$

is independent of time.

Proof.

$$\begin{aligned} \frac{d}{dt} \langle p, \partial Q \rangle - \langle q, \partial P \rangle &= \langle \dot{p}, \partial Q \rangle - \langle q, \partial \dot{P} \rangle \\ &= \sum_{n \in \mathbb{Z}} \{ [e^{-(Q_{n+1}-Q_n)}(q_{n+1} - q_n) - e^{-(Q_n-Q_{n-1})}(q_n - q_{n-1})] \partial Q_n \\ &\quad - [e^{-(Q_{n+1}-Q_n)}(\partial Q_{n+1} - \partial Q_n) - e^{-(Q_n-Q_{n-1})}(\partial Q_n - \partial Q_{n-1})] q_n \} \\ &= \sum_{n \in \mathbb{Z}} e^{-(Q_{n+1}-Q_n)} [q_{n+1} \partial Q_n - \partial Q_{n+1} q_n] - e^{-(Q_n-Q_{n-1})} [q_n \partial Q_{n-1} - q_{n-1} \partial Q_n] \\ &= 0. \end{aligned}$$

In the last line we have used the fact that the sum in the third line of the equation is a telescoping series. The boundary terms at $n = \pm\infty$ vanish because $q \in \ell_a^2$ and $\partial Q \in \ell_{-a}^2$. \square

2.5. Linearized Bäcklund transformation is an isomorphism

The linearized Bäcklund transformation gives us a time-invariant relationship between the linearizations about $(m - 1)$ - and m -soliton Toda solutions, and it preserves the orthogonality conditions from X_{m-1} to X_m ; the goal of this subsection is to show that it also implicitly defines an isomorphism $B : (q', p') \mapsto (q, p)$ between subspaces of the form (4) as in Fig. 1.

We can write the linearized Bäcklund transformation in terms of the operators defined in the previous subsection as

$$\begin{aligned} Cq' &= Lq + p, \\ p' &= \hat{C}q + Mq'. \end{aligned} \tag{15}$$

The goal is to define an isomorphism $B : (q', p') \mapsto (q, p)$ from X_{m-1} to X_m . The map is defined as follows: given (q', p') , solve $\hat{C}q = p' - Mq'$ for q and then let $p = Cq' - Lq$. Similarly, given (q, p) , solve $Cq' = Lq + p$ for q' and then let $p' = \hat{C}q + Mq'$. To implement this, we must first understand the solvability conditions, i.e. Fredholm properties, for the maps C and \hat{C} .

Note that it follows from Lemma 2.4 together with the fact that the multiplication operator β_+ is bounded with bounded inverse that \hat{C} is Fredholm with index one and a one-dimensional kernel when regarded as acting on either ℓ_a^2 or ℓ_{-a}^2 . Together with the Fredholm alternative and the fact that ℓ_a^2 is dual to ℓ_{-a}^2 this implies that C is Fredholm with index -1 , injective, and has range given by the orthogonal complement to $\partial_{\gamma_m} Q$ when regarded as acting on ℓ_a^2 . Since M and L are multiplication operators and the sequences α and β are uniformly bounded, it follows that M and L are bounded as well.

Since the range of C is orthogonal to $\partial_{\gamma_m} Q$, it follows that in order to implement the map B we must impose the orthogonality condition $\langle Lq + p, \partial_{\gamma_m} Q \rangle = 0$. The following lemma relates this orthogonality condition back to the familiar symplectic orthogonality condition that is used in the definition of X_m .

Lemma 2.8. For $q, p \in \ell_a^2$, $(Lq + p) \in \text{ran}(C)$ if and only if $\langle u, J^{-1} \partial_{\gamma_m} U \rangle = 0$.

Proof. Since C is Fredholm with $\text{coker}(C) = \text{span}\{\partial_{\gamma_m} Q\}$, $(Lq + p) \in \text{ran}(C)$ if and only if

$$\begin{aligned} 0 &= \langle Lq + p, \partial_{\gamma_m} Q \rangle \\ &= \langle q, L \partial_{\gamma_m} Q \rangle + \langle p, \partial_{\gamma_m} Q \rangle \\ &= -\langle q, \partial_{\gamma_m} P \rangle + \langle p, \partial_{\gamma_m} Q \rangle. \end{aligned}$$

This combined with Lemma 2.5 completes the proof. \square

We can also differentiate (2) with respect to κ_m , using the fact that $\partial_{\kappa_m} Q' = \partial_{\kappa_m} P' = 0$, to obtain

$$L \partial_{\kappa_m} Q + \partial_{\kappa_m} P = \hat{C} \partial_{\kappa_m} Q = 2 \sinh \kappa_m. \tag{16}$$

These equations help us to show that the map $B^{-1} : (q, p) \mapsto (q', p')$ is injective, provided our second orthogonality condition is satisfied.

Lemma 2.9. Suppose (15) is satisfied with $q' = p' = 0$, and that $\langle u, J^{-1} \partial_{\kappa_m} U \rangle = 0$. Then $q = p = 0$.

Proof. It is a consequence of Proposition 2.3 that the solution $(q(t), p(t))$ to (8) with initial condition (q, p) is related via (15) to $(q', p') = (0, 0)$. From (15) we have $\hat{C}q = p' - Mq' = 0$ so

$$q_n(t) = \mu(t) \partial_{\gamma_m} Q_n(t) \tag{17}$$

for some $\mu : \mathbb{R} \rightarrow \mathbb{R}$. Similarly

$$p = Cq' - Lq = -\mu L \partial_{\gamma_m} Q. \tag{18}$$

Differentiate (17) with respect to time to obtain $p = \dot{\mu} \partial_{\gamma_m} Q + \mu \partial_{\gamma_m} P$. Use (18) to obtain

$$\dot{\mu} \partial_{\gamma_m} Q = -\mu [L \partial_{\gamma_m} Q + \partial_{\gamma_m} P] = 0.$$

The last equality follows from (14) together with the fact that $\partial_{\gamma_m} Q' = 0$. Since $\partial_{\gamma_m} Q(t)$ is non-vanishing as an element of ℓ^∞ it follows that $\dot{\mu} \equiv 0$. We have now established that $(q, p) = \mu(\partial_{\gamma_m} Q, -K \partial_{\gamma_m} Q)$ with μ independent of t . Using the second orthogonality condition we compute

$$\begin{aligned}
 0 &= \lim_{t \rightarrow \infty} \langle p, \partial_{\kappa_m} Q \rangle - \langle q, \partial_{\kappa_m} P \rangle \\
 &= -\mu \lim_{t \rightarrow \infty} \langle L \partial_{\gamma_m} Q, \partial_{\kappa_m} Q \rangle - \mu \langle \partial_{\gamma_m} Q, \partial_{\kappa_m} P \rangle \\
 &= -\mu \lim_{t \rightarrow \infty} \langle \partial_{\gamma_m} Q, (L \partial_{\kappa_m} Q + \partial_{\kappa_m} P) \rangle \\
 &= -\mu (2 \sinh \kappa_m) \lim_{t \rightarrow \infty} \sum_{n \in \mathbb{Z}} \partial_{\gamma_m} Q_n \\
 &= -\mu (2 \sinh \kappa_m) \lim_{t \rightarrow \infty} \sum_{n \in \mathbb{Z}} \sum_{l=1}^m \partial_{\gamma_m} Q_n^1(t; \gamma_l, \kappa_l) \\
 &= -\mu (2 \sinh \kappa_m) \lim_{t \rightarrow \infty} \sum_{n \in \mathbb{Z}} \partial_{\gamma_m} Q_n^1(t; \gamma_m, \kappa_m).
 \end{aligned}$$

In the fourth line we have used (16). In the fifth line we have used Lemma 2.1. Note that $Q_n^1(t; \gamma_m, \kappa_m)$ is a monotone function of $n - ct$ and $\partial_{\gamma_m} Q_n^1$ is the derivative of this function, hence is of one sign. Thus

$$\sum_{n \in \mathbb{Z}} \partial_{\gamma_m} Q_n^1(t; \gamma_m, \kappa_m) \neq 0.$$

This implies that $\mu = 0$ as desired. \square

Theorem 2.10. Fix $t \in \mathbb{R}$ and define the space

$$X(t) := \{ (q, p) \in \ell_a^2 \times \ell_a^2 : \langle u, J^{-1} \partial_{\gamma_m} U \rangle = \langle u, J^{-1} \partial_{\kappa_m} U \rangle = 0 \}.$$

Let

$$\begin{aligned}
 B(t) &: \ell_a^2 \times \ell_a^2 \rightarrow X(t) \\
 &\quad (q'(t), p'(t)) \mapsto (q(t), p(t)), \\
 \tilde{B}(t) &: X(t) \rightarrow \ell_a^2 \times \ell_a^2 \\
 &\quad (q(t), p(t)) \mapsto (q'(t), p'(t))
 \end{aligned}$$

be defined implicitly by Eqs. (15). Then $B(t)$ is an isomorphism with inverse $B^{-1}(t) = \tilde{B}(t)$.

Proof. We will suppress dependence on t throughout this proof. Suppose $(q, p) \in X$ is given. By Lemma 2.8, $Lq + p$ is in the range of C and so there exists $q' \in \ell_a^2$ satisfying the first equation of (15). Since $\ker(C) = \{0\}$, q' is uniquely defined, as is $p' := (Mq' + \hat{C}q) \in \ell_a^2$. By Lemma 2.9, \tilde{B} is also injective.

Given $q', p' \in \ell_a^2$, since $\text{ran}(\hat{C}) = \ell_a^2$ there exists $q \in \ell_a^2$ satisfying the second equation of (15). Then $p \in \ell_a^2$ is defined by the first equation of (15). Moreover $Lq + p \in \text{ran}(C)$, and since $\langle \partial_{\kappa_m} U, \partial_{\gamma_m} U \rangle \neq 0$ for $m \geq 1$ we can shift u by a multiple of $J^{-1} \partial_{\gamma_m} U$ to get $u \perp J^{-1} \partial_{\kappa_m} U$. Thus, \tilde{B} is surjective and therefore an isomorphism with inverse B . \square

Corollary 2.11. $B(t) : X_{m-1}(t) \rightarrow X_m(t)$ is an isomorphism.

Proof. This follows from Proposition 2.6 and the fact that $X_{m-1}(t)$ and $X_m(t)$ are codimension $2(m - 1)$ subspaces of $\ell_a^2 \times \ell_a^2$ and $X(t)$, respectively. \square

In the following we will need to show that $B(t)$ and $B(t)^{-1}$ are bounded uniformly in time.

Lemma 2.12. The maps $B(t)$ and $B^{-1}(t)$ are uniformly bounded for $t \in \mathbb{R}$.

Proof. This follows from the uniform boundedness of the operators L, M, C and \hat{C} , as well as the inverses of C and \hat{C} , appropriately restricted. L and M are multiplication operators by sequences which are uniformly bounded, hence their operator norms are uniformly bounded. C and \hat{C} are the sum of a multiplication operator by a uniformly bounded sequence and a shift composed with a multiplication operator by a uniformly bounded sequence, hence their operator norms are bounded uniformly too. The uniform boundedness of \hat{C}^{-1} was addressed in Lemma 2.4. The uniform boundedness of C^{-1} is similar. \square

2.6. Linear stability of the Toda flow

To prove Theorem 1.1 it only remains to show that the decay rates of $\Phi_{m-1}(t, s)$ are inherited by $\Phi_m(t, s)$ in Fig. 1.

Proposition 2.13. *Let $c > 1$, $a \in (0, 2\kappa)$ with $\kappa = \min_{1 \leq i \leq m} \kappa_i$, and $\beta = ca - 2 \sinh(a/2)$. Suppose $u'(t) \in X \subset \ell_a^2 \times \ell_a^2$ satisfies*

$$u'_t = JH''(U')u',$$

where as above $U' = ((S - I)Q', P')$ is an $(m - 1)$ -soliton solution with solitons moving at speeds $\kappa_1, \dots, \kappa_{m-1}$. Suppose further that for some constant $\tilde{K} > 0$ we have

$$\|e^{a(n-ct-T)}u'(t)\| \leq \tilde{K}e^{-\beta(t-s)}\|e^{a(n-cs-T)}u'(s)\|.$$

Then for any $u \in Y := \{u \in X \mid \langle u, J^{-1}\partial_{\gamma_m}U \rangle = \langle u, J^{-1}\partial_{\kappa_m}U \rangle = 0\}$, where U is an m -soliton solution with solitons moving at speeds $\kappa_1, \dots, \kappa_m$ and is related to U' by (2), there exists a constant $K > 0$ such that

$$\|e^{a(n-ct-T)}u(t)\| \leq Ke^{-\beta(t-s)}\|e^{a(n-cs-T)}u(s)\|,$$

where $u(t)$ is a solution of the Toda equations, linearized about the m -soliton solution $U(t)$.

Proof. This follows from the fact that $u(t) = B(t)u'(t)$ (see Fig. 1), the bound on u' in the hypothesis of the prop, and Lemma 2.12. \square

The proof of Theorem 1.1 now follows from induction on m . The initial step is provided by [13, Lemma 10] and the above lemma proves the inductive step.

3. Nonlinear stability

The purpose of this section is to leverage the linear stability result, Theorem 1.1 proved in the last section, in order to prove the nonlinear stability result Theorem 1.2. The strategy of the proof is a natural generalization of [15,3] to the setting of multi-soliton solutions. Roughly speaking, in Section 3.1 we show that so long as we are willing to let the parameters κ_i and γ_i modulate, then we can impose an orthogonality condition on any (sufficiently small) perturbation of an m -soliton solution; in Section 3.2 we derive coupled ODEs for the modulating parameters κ and γ and the perturbation v ; in Section 3.3 we obtain estimates which allow us to determine the long-time behavior of solutions to these ODEs with small initial data via the method of bootstrapping.

3.1. (v, ξ) coordinates

Let $\kappa = (\kappa_1, \dots, \kappa_m)$ and $\gamma = (\gamma_1, \dots, \gamma_m)$ be functions of time which we will specify later. Let $\xi = (\gamma_1, \kappa_1, \gamma_2, \kappa_2, \dots, \gamma_m, \kappa_m)$. Let v be defined by the equation

$$u = U(t; \kappa(t), \gamma(t)) + v(t) \equiv U(t; \xi(t)) + v(t), \tag{19}$$

where as in the preceding subsection, U is an m -soliton solution of the Toda system.

We compute

$$[\partial_t - JH''(U(t; \kappa(t_0), \gamma(t_0)))]v = JR - \sum_{j=1}^{2m} \partial_{\xi_j} U \dot{\xi}_j. \tag{20}$$

Here $R = R_1 + R_2$ with $R_1 = H'(U + v) - H'(U) - H''(U)v$, and $R_2 = [H''(U(t; \xi(t_0))) - H''(U(t; \xi(t)))]v$. We work in the space X_m as defined in (4).

Lemma 3.1 (Symplectic form restricted to the tangent space). *Let*

$$\mathcal{A}_{ij}(t) := \langle J^{-1}\partial_{\xi_i}U(t, \xi_*), \partial_{\xi_j}U(t, \xi_*) \rangle.$$

Then \mathcal{A} is independent of t . Moreover, \mathcal{A} is block diagonal with 2×2 blocks given by

$$\begin{pmatrix} 0 & \alpha_{0,k} \\ -\alpha_{0,k} & \alpha_{1,k} \end{pmatrix}$$

where $\alpha_{0,k} = \frac{1}{c_k} \frac{d}{dc} H(u_{c_k}) \frac{dc}{dk}$ and $\alpha_{1,k} = (\frac{d}{dc} \int_{\mathbb{R}} r_{c_k}) \frac{d}{dc} (c \int_{\mathbb{R}} r_{c_k}) (\frac{dc}{dk})^2$.

Proof. Observe

$$\begin{aligned} \frac{d}{dt} \mathcal{A}_{ij} &= \langle J^{-1} \partial_{\xi_i} \partial_t U, \partial_{\xi_i} U \rangle + \langle J^{-1} \partial_{\xi_i} U, \partial_{\xi_i} \partial_t U \rangle \\ &= \langle J^{-1} J H''(U) \partial_{\xi_i} U, \partial_{\xi_j} U \rangle + \langle J^{-1} U_{\xi_i}, J H''(U) \partial_{\xi_i} U \rangle \\ &= \langle H''(U) \partial_{\xi_i} U - H''(U) \partial_{\xi_i} U, \partial_{\xi_j} U \rangle = 0. \end{aligned}$$

Here we have used the fact that $H''(U)$ is self-adjoint and also that the form $\langle \cdot, J^{-1} \partial_{\xi_i} U \rangle$ is skew-symmetric when acting on zero mean sequences. Evaluate \mathcal{A} as $t \rightarrow \infty$ and use Lemma 2.1 in [3] to see that \mathcal{A} is block diagonal with 2×2 blocks given as in the statement of the lemma. \square

Corollary 3.2. $\mathcal{A}(t)$ is invertible with uniformly bounded inverse.

Proof. This follows from the fact that $\alpha_{0,k} \neq 0$ for all k , a fact that is proved in [2]. \square

Proposition 3.3 (Tubular coordinates). *There is a $\delta_* > 0$ such that (v, ξ) coordinates correspond to a unique solution u of the Toda lattice equations in a neighborhood $\|v\| + |\xi - \xi_*| < \delta_*$.*

Proof. We proceed as in [2]. Define $F : \mathbb{R}^{2m} \times \mathbb{R} \times \ell_a^2 \rightarrow \mathbb{R}^{2m}$ by $F_j(\xi, t, u) = \langle J^{-1} \partial_{\xi_j} U(t, \xi), u - U(t, \xi) \rangle$ so that the zero-level set of F is $\mathbb{R}^{2m} \times \bigcup_{t \in \mathbb{R}} X_m(t)$. Let $\xi_* \in \mathbb{R}^{2m}$ and $t_* \in \mathbb{R}$ be given. Observe that $F(\xi_*, t, U(t, \xi_*)) = 0$. Compute $\partial_{\xi_i} F_j(\xi_*, t, U(t, \xi_*)) = -\mathcal{A}_{ij}$ with \mathcal{A} given as in Lemma 3.1. In particular $D_\xi F(\xi_*, t, U(t, \xi_*))$ is invertible. It now follows from the implicit function theorem that there is a smooth function $(u, t) \mapsto \xi(u, t)$ mapping a neighborhood of $U(0, \xi_*) \times \{t_*\}$ in $\ell_a^2 \times \mathbb{R}$ to a neighborhood of ξ_* in \mathbb{R}^{2m} such that $u - U(t, \xi(u, t)) \in X_m(t)$. Since the derivative $\partial_{\xi_i} F_j(\xi_*, t, U(t, \xi_*))$ is independent of time, so is the neighborhood of ξ_* on which the function ξ is defined. Moreover, the map $(u, t) \mapsto (\xi, v)$ given by $\xi = \xi(u, t)$ and $v = u - U(t, \xi(u, t))$ is, for each fixed t , an isomorphism from ℓ_a^2 to $X_m(t) \times \mathbb{R}^{2m}$. \square

3.2. Modulation equations

Take the inner product with $J^{-1} \partial_{\xi_j} U$ in (20) to obtain

$$\langle (\partial_t - J H''(U))v, J^{-1} \partial_{\xi_j} U \rangle - \langle JR, J^{-1} \partial_{\xi_j} U \rangle = \sum_{i=1}^{2m} \langle \partial_{\xi_i} U, J^{-1} \partial_{\xi_j} U \rangle \dot{\xi}_i. \tag{21}$$

Differentiate the identity $\langle v(t), J^{-1} \partial_{\xi_j} U(t; \xi(t_0)) \rangle \equiv 0$ with respect to time to obtain

$$\langle [\partial_t - J H'(U)]v, J^{-1} \partial_{\xi_j} U \rangle \equiv 0$$

thus the modulation equation (21) reduces to

$$\mathcal{A} \dot{\xi} = b \tag{22}$$

where $b_j = -\langle JR, J^{-1} \partial_{\xi_j} U \rangle = \langle R, \partial_{\xi_j} U \rangle$ and \mathcal{A} is given as in Lemma 3.1. In light of the fact that \mathcal{A} is invertible with \mathcal{A}^{-1} block diagonal we have

$$|\dot{\kappa}_i| + |\dot{\gamma}_i| \leq K \|R\|_{c,a} (\|\partial_{\kappa_i} U\|_{c,-a} + \|\partial_{\gamma_i} U\|_{c,-a}).$$

3.3. Stability estimates

Proposition 3.4. *Let $M > 0$ be given. Then there exists a constant K such that*

$$\|R\|_{a,c} \leq K (\|v\|_{\ell^\infty} + |\xi(t) - \xi(t_0)|) \|v\|_{a,c} \tag{23}$$

holds so long as $\|v\|_{a,c} + \|v\|_{\ell^\infty} \leq M$.

Proof. We estimate $\|R_1\|_{a,c} \leq K \|v\|_a \|v\|_{\ell^\infty}$ using Taylor's theorem with remainder. We estimate $\|R_2\|_{a,c} \leq K \|v\|_{a,c} (|\xi(t) - \xi(t_0)|)$ using the fact that H'' is locally Lipschitz. Note that for H'' to be locally Lipschitz with respect to the operator norm topology for operators acting on weighted spaces it suffices that U be Lipschitz with respect to ξ with U regarded as taking values in unweighted ℓ^∞ . \square

In order to proceed we must check that $\|v\|_{\ell^\infty}$ (which is controlled by $\|v\|$) remains well behaved for long times. As a preliminary step we prove

Lemma 3.5.

$$\langle H'(U(t; \kappa(t_0), \gamma(t_0))), v(t) \rangle \equiv 0.$$

Proof. The proof follows from the following identity satisfied by the m -soliton profile.

Lemma 3.6. *If U is an m -soliton solution of the Toda equations, then*

$$\partial_t U = - \sum_i 2 \sinh \kappa_i \partial_{\gamma_i} U. \tag{24}$$

This lemma is proved in Appendix A – assuming that it holds, and recalling the orthogonality condition $\langle J^{-1} \partial_{\gamma_i} U, v \rangle \equiv 0$ it follows that $\langle J^{-1} \partial_t U, v \rangle \equiv 0$, thus $\langle H'(U), v \rangle \equiv 0$, as desired. \square

We now derive our first estimate on $v(t)$, the perturbation of the m -soliton.

Proposition 3.7. *There is a constant K such that if $v(t)$ is the function defined in (19), then*

$$\|v(t)\|^2 \leq K (\|v(t_0)\|^2 + |\xi(t) - \xi(t_0)|) \tag{25}$$

holds for $t \geq t_0$ so long as $\|v(t)\| \leq \delta_*$.

Proof. Let $u(t) := U(t; \kappa(t), \gamma(t)) + v(t)$. From the form of the Hamiltonian function we know

$$K_- \|v\|^2 \leq H(u) - H(U) - \langle H'(U), v \rangle \leq K_+ \|v\|^2.$$

Thus

$$\begin{aligned} K_- \|v(t)\|^2 &\leq H(u(t)) - H(U(t; \xi(t))) - \langle H'(U(t; \xi(t))), v(t) \rangle && \text{(i)} \\ &= H(u(t)) - H(u(t_0)) && \text{(ii)} \\ &\quad - H(U(t; \xi(t))) + H(U(t; \xi(t_0))) && \text{(iii)} \\ &\quad - H(U(t; \xi(t_0))) + H(U(t_0; \xi(t_0))) && \text{(iv)} \\ &\quad + \langle H'(U(t; \xi(t_0))), v(t) \rangle - \langle H'(U(t; \xi(t))), v(t) \rangle && \text{(v)} \\ &\quad + \langle H'(U(t_0; \xi(t_0))), v(t_0) \rangle - \langle H'(U(t; \xi(t_0))), v(t) \rangle && \text{(vi)} \\ &\quad + H(u(t_0)) - H(U(t_0; \xi(t_0))) - \langle H'(U(t_0; \xi(t_0))), v(t_0) \rangle. \end{aligned}$$

We estimate (i) = 0 by conservation of the Hamiltonian, (ii) $\leq C|\xi(t) - \xi(t_0)|$ because H and U are locally Lipschitz in their arguments. (iii) = 0 by conservation of the Hamiltonian (iv) $\leq K|\xi(t) - \xi(t_0)|\|v(t)\| \leq K|\xi(t) - \xi(t_0)|$, (v) = 0 by Lemma 3.5, (vi) $\leq K\|v(t_0)\|^2$. Summing these estimates yields (25). \square

We now write the evolution equation (20) in Duhamel form:

$$v(t) = \Phi(t, 0)v(0) + \int_0^t \Phi(t, s) \left[JR(s) + \sum_{j=1}^{2m} \partial_{\xi_j} U(s) \dot{\xi}_j(s) \right] ds.$$

Thus

$$\begin{aligned} \|v(t)\|_{a,c} &\leq e^{-\beta t} \|v(t_0)\|_{a,c} + \int_0^t e^{-\beta(t-s)} \left[\|JR(s)\|_{a,c} + \sum_{j=1}^{2m} \|\partial_{\xi_j} U(s)\|_{a,c} |\dot{\xi}_j(s)| \right] ds \\ &\leq e^{-\beta t} \|v(t_0)\|_{a,c} + \int_0^t e^{-\beta(t-s)} \left[\|JR(s)\|_{a,c} + K \sum_{j=1}^{2m} \|\partial_{\xi_j} U(s)\|_{a,c} \|\partial_{\xi_j} U(s)\|_{-a,c} \|R(s)\|_{a,c} \right] ds \\ &\leq e^{-\beta t} \|v(t_0)\|_{a,c} + K \int_0^t e^{-\beta(t-s)} \|R(s)\|_{a,c} ds. \end{aligned}$$

In the last line we have used that $\|f\|_{a,c}\|g\|_{-a,c}$ is independent of c for any functions f and g together with the fact that $\|\partial_{\xi_j} U\|_{\pm a,c_j}$ is bounded uniformly in t . This second fact follows from the fact that $\partial_{\xi_j} U$ converges exponentially fast to $\partial_{\xi_j} U^1(t; \gamma_j + \zeta_j, \kappa_j)$. For this function, we have explicit formulas for both components of U^1 which live in $\ell^2_{-a} \cap \ell^2_a$ and evolves by rigidly translating to the right at speed c_j and hence is uniformly bounded in the $\|\cdot\|_{a,c_j}$ norm.

Lemma 3.8. Let δ_0, δ_1 and δ_2 be sufficiently small positive numbers. Suppose that the following estimates hold for $t \in [t_0, t_1]$:

$$|\dot{\xi}(t) - \dot{\xi}(t_0)| < \delta_0; \quad \|v(t)\|_{a,c} < \delta_1 e^{-\beta'(t-t_0)}; \quad \|v(t)\| < \delta_2. \tag{26}$$

Then in fact, the estimates $|\dot{\xi}(t) - \dot{\xi}(t_0)| < \frac{K(\delta_2 + \delta_0)\delta_1}{\beta}$, $\|v(t)\| \leq K(\|v(t_0)\| + \sqrt{\delta_0})$, and $\|v(t)\|_{a,c} \leq \delta_1 e^{-\beta'(t-t_0)} \left(\frac{K(\delta_2 + \delta_0)}{\beta - \beta'} + \frac{K}{\delta_1} e^{-(\beta - \beta')(t-t_0)}\|v(t_0)\|_a\right)$ hold for all $t \in [t_0, t_1]$.

Proof. Use (23) to see that $\|R\|_{a,c} \leq K\delta_1(\delta_0 + \delta_2)e^{-\beta'(t-t_0)}$ and hence that

$$\begin{aligned} \|v(t)\|_{a,c} &\leq e^{-\beta(t-t_0)}\|v(t_0)\|_{c,a} + K\delta_1(\delta_0 + \delta_2) \int_0^t e^{-\beta(t-s)}e^{-\beta'(s-t_0)} ds \\ &\leq e^{-\beta(t-t_0)}\|v(t_0)\|_{c,a} + \frac{K\delta_1(\delta_0 + \delta_2)}{(\beta - \beta')}e^{-\beta'(t-t_0)} \\ &< \delta_1 e^{-\beta'(t-t_0)} \left(\frac{K(\delta_2 + \delta_0)}{\beta - \beta'} + \frac{K}{\delta_1} e^{-(\beta - \beta')(t-t_0)}\|v(t_0)\|_a \right) \end{aligned}$$

so long as δ_0, δ_2 and $\|v(t_0)\|_{c,a}$ are chosen small enough.

Substitute (26) into (25) to obtain $\|v(t)\|^2 \leq K(\|v(t_0)\|^2 + \delta_0)$. Make use of the estimate on $\|R\|_{a,c}$ to obtain $|\dot{\xi}| \leq K\delta_1(\delta_2 + \delta_0)e^{-\beta'(t-t_0)}$. Integrating yields $|\dot{\xi}(t) - \dot{\xi}(t_0)| \leq \frac{K\delta_1(\delta_2 + \delta_0)}{\beta}$. \square

Proof of Theorem 2. Let $\delta_2 > 0$ be given. Let $\delta_0 < \frac{\delta_2^2}{2K}$ and let $\delta_1 < \frac{\delta_0\beta'}{K(\delta_2 + \delta_0)}$. Finally, restrict attention to $v(t_0)$ so small so that $\|v(t_0)\| < \frac{\delta_2}{\sqrt{2K}}$ and $\|v(t_0)\|_a < \frac{\delta_1}{2K}$. Apply Lemma 4.5 to conclude that we can take $t_1 = \infty$ without loss of generality. To establish the fact that κ_i and γ_i converge to limits at $\pm\infty$ note that from what we have already proved $\|v(t)\|_a$, and hence $\|R\|_a$ decay exponentially fast. It now follows from (22) that $\dot{\xi}_j$ decays exponentially as well and hence that $\xi_j(t)$ has a limit as $t \rightarrow \infty$. \square

Remark. In fact, although we will not need this, one can prove from (22) and the estimates on $\|v\|_{a,c}$ derived above, that the modulation parameters $\gamma_i(t)$ and $\kappa_i(t)$ approach their asymptotic limits as $t \rightarrow \pm\infty$ exponentially rapidly.

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Appendix A. Form of the m -soliton solutions of the Toda model

In the present appendix we prove various facts about the m -soliton solutions of the Toda model. While these facts will probably not be surprising to experts, we were unable to find their proofs in the literature and thus collect them here for future reference.

The proof of Lemma 2.1 (resolution of m -soliton solutions into one-solitons):

Proof. While the asymptotic resolution into one-solitons is one of the defining characteristics of an m -soliton solution this lemma shows that the decomposition takes place exponentially rapidly in the space ℓ^1 (and hence in all other ℓ^p spaces). We begin by recalling the explicit expression for the m -soliton [16, Section 3.6] namely

$$Q^m = Q^m(\gamma_1, \dots, \gamma_m, \kappa_1, \dots, \kappa_m) = (I - S^{-1}) \log d \tag{27}$$

where $d = \det(I + C)$ and

$$C_{ij} = C_{ij}(n, t) = \frac{e^{-(\eta_i(n,t) + \gamma_i + \eta_j(n,t) + \gamma_j)}}{1 - e^{-(\kappa_i + \kappa_j)}}, \tag{28}$$

with $\eta_j(n, t) = \kappa_j n - \sinh(\kappa_j)t$. Note that this notation also encompasses the formula for one-soliton solutions which can be written as $Q^1(\gamma, \kappa) = (I - S^{-1}) \log(1 + \frac{e^{-2(\eta+\gamma)}}{1-e^{-2\kappa}})$, with $\eta(n, t) = \kappa n - \sinh(\kappa)t$.

We begin by noting that there is a simple, explicit formula for the phase shifts, ζ_j^\pm , undergone by the constituent solitons in the m -soliton as they interact with one another. Let $\alpha_{ij} = \frac{1}{1-e^{-(\kappa_i+\kappa_j)}}$, let $\alpha^{k,+}$ denote the $(m-k+1) \times (m-k+1)$ matrix with entries α_{ij} for $k \leq i, j \leq m$, and let $\alpha^{k,-}$ denote the $k \times k$ matrix with entries α_{ij} for $1 \leq i, j \leq k$. Since $\kappa_i > 0$ for each i , it follows that each α_{ij} is positive, hence that each matrix $\alpha^{k,\pm}$ is positive-definite, hence that $\det \alpha^{k,\pm}$ is positive.

Thus we may define $\zeta_j^+ = \frac{1}{2} \log \frac{\det \alpha^{j+1,+}}{\det \alpha^{j,+}}$ for $1 \leq j < m$ with $\zeta_m^+ = -\frac{1}{2} \log \det \alpha^{m,+}$ and $\zeta_j^- = \frac{1}{2} \log \frac{\det \alpha^{j-1,-}}{\det \alpha^{j,-}}$ for $1 < j \leq m$ with $\zeta_1^- = -\frac{1}{2} \log \det \alpha^{1,-}$.

We now proceed to the proof of (5). We write $Q^m = Q^m(\gamma_1, \dots, \gamma_m, \kappa_1, \dots, \kappa_m) = (I - S^{-1}) \log d$ where $d = \det(I + C)$ and $C_{ij} = \frac{e^{-(\eta_i+\gamma_i+\eta_j+\gamma_j)}}{1-e^{-(\kappa_i+\kappa_j)}}$. Similarly we write $Q^1(\gamma, \kappa) = (I - S^{-1}) \log(1 + \frac{e^{-2(\eta+\gamma)}}{1-e^{-2\kappa}})$,

$$Q^m - \sum_{i=1}^m Q^1(\gamma_i + \zeta_i^\pm, \kappa_i) = (I - S^{-1}) \log \left(\sum_{\sigma \in S^m} \text{sgn}(\sigma) \prod_{j=1}^n \frac{\delta_{j,\sigma(j)} + C_{j,\sigma(j)}}{1 + e^{-2\zeta_j^\pm} C_{jj}} \right) = (I - S^{-1}) \log(1 + \Delta^\pm)$$

where

$$\Delta^\pm = \frac{[\sum_{\sigma \in S^m} \text{sgn}(\sigma) \prod_{j=1}^m \delta_{j,\sigma(j)} + C_{j,\sigma(j)}] - [\prod_{j=1}^m (1 + e^{-2\zeta_j^\pm} C_{jj})]}{\prod_{j=1}^m (1 + e^{-2\zeta_j^\pm} C_{j,j})}$$

In light of the fact that $(I - S^{-1}) \log(1 + x) = \log(1 + \frac{x-S^{-1}x}{1+S^{-1}x})$ it suffices to show that $\|\Delta^\pm\|_{\ell^1} \rightarrow 0$ exponentially fast as $t \rightarrow \pm\infty$. Let $\xi_j = e^{-\eta_j - \gamma_j}$ be so that $C_{ij} = \xi_i \xi_j \alpha_{ij}$. Then we can write

$$\Delta^\pm = \frac{\sum_{J \subset \{1, \dots, m\}} (\prod_{k \in J} \xi_k^2) [\sum_{\{\sigma \in S^m | J \subset \text{Fix} \sigma\}} \text{sgn}(\sigma) \prod_{j \in J} \alpha_{j,\sigma(j)}] - \prod_{k \in J} e^{-2\zeta_k^\pm}}{\prod_{j=1}^m (1 + e^{-2\zeta_j^\pm} \alpha_{j,j} \xi_j^2)} = \frac{\sum_{J \subset \{1, \dots, m\}} (\prod_{k \in J} \xi_k^2) [\det \alpha^J - e^{-2\sum_{k \in J} \zeta_k^\pm}]}{1 + \sum_{J \subset \{1, \dots, m\}} \det \alpha^J e^{-2\sum_{k \in J} \zeta_k^\pm} \prod_{k \in J} \xi_k^2}$$

where the sum runs over all non-empty subsets of the set of integers $\{1, 2, \dots, m\}$, and α^J is the $|J| \times |J|$ matrix with entries α_{ij} for $i, j \in J$.

Note that for any fixed $t > 0$ and n , the polynomials $\prod_{k \in J} \xi_k^2 = \prod_{k \in J} \xi_k(n, t)^2$ can be ordered in size and the largest will correspond to a subset J in which $\xi_k > 1$ if and only if $k \in J$. Moreover, such a J will be of the form $\{j, j+1, \dots, m\}$, i.e. when $t > 0$ if $\xi_j(n, t) > 1$, then $\xi_k(n, t) > 1$ for $k = j+1, \dots, m$ and if $\xi_j(n, t) < 1$, then $\xi_k(n, t) < 1$ for $k = 1, \dots, j-1$. If $t < 0$, a similar argument shows that the dominant polynomial corresponds to $J = \{1, 2, \dots, k\}$ for some k (as well as the empty set). For the remainder of the proof, we concentrate on the case $t > 0$. The case when $t \rightarrow -\infty$ is handled in a similar fashion.

The key observation is that in the numerator of the expression for $\Delta^+(n, t)$, the definition the asymptotic phase shifts, ζ^\pm , insures that the coefficient of the largest polynomial is zero. This is sufficient to prove the lemma.

In more detail, set $n_j(t) = \lfloor \frac{\sinh(\kappa_j)t}{\kappa_j} \rfloor$, where $\lfloor x \rfloor$ denotes the integer part of x . Also, set $\lambda(t) = \min_{j=1, \dots, m-1} (n_{j+1}(t) - n_j(t))$. Note that there exists $\lambda_0 > 0$ such that $\lambda(t) \geq \lambda_0 t$ for t sufficiently large. Now let $J_{max} = \{k^*, k^* + 1, \dots, m\}$ denote the index set for the dominant terms in Δ^\pm and consider any quotient of the form

$$\frac{\prod_{j \in J} \xi_j^2(n, t)}{\prod_{j \in J_{max}} \xi_j^2(n, t)} \tag{29}$$

for which the coefficient in the numerator of Δ^\pm is nonzero. Note that $J \neq \{k^* + 1, k^* + 2, \dots, m\}$ or $\{k^* - 1, k^*, \dots, m\}$ since such terms also have zero coefficients due to the definition of ζ^\pm . Hence, for any nonzero term, there is either an ‘‘extra’’ factor of the form $\xi_j^2(n, t)$, with $j < k^* - 1$ in the numerator of (29), or a ‘‘missing’’ factor of $\xi_j^2(n, t)$, with $j > k^* + 1$. If we define $\Gamma = \max |\gamma_j|$, then from the fact that $\kappa_1 \leq \kappa_j \leq \kappa_m$, and the definition of $\xi_j(n, t)$, we see that the quotient in (29) can be bounded by

$$\left| \frac{\prod_{j \in J} \xi_j^2(n, t)}{\prod_{j \in J_{max}} \xi_j^2(n, t)} \right| \leq e^{2\kappa_m} e^{2\Gamma} e^{-\kappa_1 |n - n_{k^*}(t)|} e^{-\kappa_1 \lambda(t)} \tag{30}$$

for some $j = 1, \dots, m$. Since there are only finitely many terms in the numerator and denominator of Δ^\pm , we see that we have the bound

$$|\Delta^\pm(n, t)| \leq K e^{2\kappa m} e^{2\Gamma} e^{-\kappa_1 \lambda(t)} \sum_{j=1}^m e^{-\kappa_1 |n-n_j(t)|}. \tag{31}$$

Summing over n then completes the proof of Lemma 2.1. \square

We now prove Corollary 2.2.

Proof. In light of the proof above, it suffices to show that $\|\partial_{\diamond_k} \Delta^\pm\|_{c_l, \pm a} \rightarrow 0$ exponentially fast where $\diamond \in \{\gamma, \kappa\}$. We work with Δ^+ , the minus case being similar. Write

$$\Delta^+ = \frac{\sum_{J \subset \{1, \dots, m\}} C^J \prod_{l \in J} \xi_l^2}{1 + \sum_{J' \subset \{1, \dots, m\}} d^{J'} \prod_{l' \in J'} \xi_{l'}^2}.$$

Here the coefficients d^J and C^J depend upon κ_j for $j \in J$. Moreover, as observed above, $d^J \neq 0$ for all J and $C^J = 0$ whenever $J = \{m-k, m-k+1, \dots, m\}$ for some k . Compute

$$\partial_{\diamond_k} \Delta^+ = \frac{(1 + \sum_{J'} d^{J'} \prod_{l' \in J'} \xi_{l'}^2) (\sum_J \partial_{\diamond_k} (C^J \prod_{l \in J} \xi_l^2)) - (\sum_J C^J \prod_{l \in J} \xi_l^2) (\sum_{J'} \partial_{\diamond_k} (d^{J'} \prod_{l' \in J'} \xi_{l'}^2))}{(\sum_{J'} d^{J'} \prod_{l' \in J'} \xi_{l'}^2)}$$

where the sums run over J and J' non-empty subsets of $\{1, \dots, m\}$. We decompose

$$\|\partial_{\diamond_l} \Delta^+\|_{c_l, \pm a} = \|\partial_{\diamond_l} \Delta^+ \chi_{n < c_1 t}\|_{c_l, \pm a} + \|\partial_{\diamond_l} \Delta^+ \chi_{n \geq c_m t}\|_{c_l, \pm a} + \sum_{j=1}^{m-1} \|\partial_{\diamond_l} \Delta^+ \chi_{c_j t \leq n < c_{j+1} t}\|_{c_l, \pm a}$$

and estimate the expression term by term.

When $n \geq c_m t$, we have $\xi_i < 1$ (modulo factors of $e^{-\gamma_i}$) for all i hence the dominant term in the denominator is 1. The dominant term in the numerator is ξ_1^2 thus

$$\begin{aligned} \|\partial_{\diamond_l} \Delta^+ \chi_{n \geq c_m t}\|_{c_l, \pm a} &\leq C \sum_{n > c_1 t} e^{\pm 2a(n-c_1 t)} e^{-2\kappa_1(n-c_1 t)} \\ &= C e^{-2\kappa_1(c_m - c_1)t} e^{\pm 2a(c_m - c_1)t} \sum_{q=0}^{\infty} e^{-2(\kappa \pm a)q} \end{aligned}$$

which decays exponentially fast.

When $n < c_1 t$ we have $\xi_i > 1$ (again modulo factors of $e^{-\gamma_i}$) for all i hence the dominant term in the denominator is $\prod_{k=1}^m (\xi_k)^4$ whereas the dominant term in the numerator is $\prod_{k=1}^m (\xi_k)^2 \prod_{k \neq 2} (\xi_k)^2$. Thus in this regime $(\partial_{\diamond_l} \Delta^+)_{n(t)} \leq C \xi_2^{-2}$. Hence

$$\begin{aligned} \|\partial_{\diamond_l} \Delta^+ \chi_{n < c_1 t}\|_{c_l, \pm a} &\leq C \sum_{n < c_1 t} e^{\pm 2a(n-c_1 t)} e^{2\kappa_2(n-c_2 t)} \\ &\leq C e^{-2(\kappa_2(c_2 - c_1) \pm a(c_1 - c_1))t} \sum_{q=0}^{\infty} e^{-2(\kappa_2 \pm a)q} \end{aligned}$$

which decays exponentially so long as $a < \frac{\kappa_2(c_2 - c_1)}{c_1 - c_1}$.

Finally, when $c_j t < n < c_{j+1} t$ we have $\xi_i > 1$ only for $i > j$. The dominant term in the denominator is $\prod_{i > j} \xi_i^4$ whereas the dominant term in the numerator is $\prod_{i > j} \xi_i^2 \prod_{i > j, i \neq j+2} \xi_i^2$. Thus in this regime we have $\partial_{\diamond} \Delta^+ \sim \xi_{j+2}^{-2}$. Hence

$$\begin{aligned} \|\partial_{\diamond_l} \Delta^+ \chi_{n < c_1 t}\|_{c_l, \pm a} &\leq C \sum_{n=c_j t}^{c_{j+1} t} e^{\pm 2a(n-c_1 t)} e^{2\kappa_{j+2}(n-c_{j+2} t)} \\ &= C \sum_{q=0}^{(c_{j+1} - c_j)t} e^{\pm 2a(q + (c_j - c_1)t)} e^{2\kappa_{j+2}(q + (c_j - c_{j+2})t)} \\ &= C e^{-2\kappa_{j+2}(c_{j+2} - c_j)t} e^{\pm 2a(c_j - c_1)t} e^{(2\kappa_{j+2} \pm 2a)(c_{j+1} - c_j)t} \\ &\leq C e^{-2(\kappa_{j+2}(c_{j+1} - c_j) \pm a(c_{j+1} - c_1))t} \end{aligned}$$

which again decays so long as a is chosen small enough. \square

We now turn to the proof of Lemma 3.6 which is an identity satisfied by the profile of the m -soliton solution.

Proof. Note that it suffices to show that the scalar equality $\partial_t r = \sum_i 2\gamma_i \sinh \kappa_i \partial_{\gamma_i} r$ for all time. This is because $p(t)$ can be recovered from $r(t)$ via $p = (S - I)^{-1} \dot{r}$ and the transformation $r \mapsto p$ commutes with time derivatives, γ_i derivatives, and linear combinations thereof. Thus if the identity holds for r , it holds also for p and hence for $U = (r, p)$.

We write $r = \Delta \log d$ where $d = \det(I + C)$ with $C_{ij} = \frac{1}{1 - e^{-(\kappa_i + \kappa_j)}} e^{-(\eta_i + \gamma_i + \eta_j + \gamma_j)}$ and $\eta_i = \kappa_i n - \sinh \kappa_i t$. Thus $\partial_\circ r = \Delta \frac{\partial_\circ d}{d}$ where ∂_\circ stands for ∂_t or ∂_{γ_i} for some i . Here we have used the fact that the discrete Laplacian Δ commutes with derivatives with respect to t and γ_i . Write $d = \sum_{\sigma \in \mathcal{S}_n} \text{sgn}(\sigma) \prod_{j=1}^n (\delta_{j\sigma(j)} + C_{j\sigma(j)})$ so that

$$\partial_\circ d = \sum_{\sigma \in \mathcal{S}_n} \text{sgn}(\sigma) \sum_{l=1}^n \partial_\circ C_{l,\sigma(l)} \prod_{j \neq l} (\delta_{j\sigma(j)} + C_{j\sigma(j)}).$$

Note that $\partial_{\gamma_i} C_{jk} = -\frac{\delta_{ij} + \delta_{ik}}{2} C_{jk}$ and that $\partial_t C_{jk} = (\sinh(\kappa_j) + \sinh(\kappa_k)) C_{jk}$. Thus

$$\begin{aligned} -\sum_{i=1}^n 2 \sinh(\kappa_i) \partial_{\gamma_i} d &= \sum_{i=1}^n \sum_{\sigma \in \mathcal{S}_n} \text{sgn}(\sigma) \sum_{l=1}^n (\delta_{i,l} + \delta_{i,\sigma(l)}) \sinh(\kappa_i) C_{l,\sigma(l)} \prod_{j \neq l} (\delta_{j\sigma(j)} + C_{j\sigma(j)}) \\ &= \sum_{l=1}^n \sum_{\sigma \in \mathcal{S}_n} \text{sgn}(\sigma) (\sinh(\kappa_l) + \sinh(\kappa_{\sigma(l)})) C_{l,\sigma(l)} \prod_{j \neq l} (\delta_{j\sigma(j)} + C_{j\sigma(j)}) \\ &= \sum_{\sigma \in \mathcal{S}_n} \text{sgn}(\sigma) \sum_{l=1}^n \partial_t C_{l,\sigma(l)} \prod_{j \neq l} (\delta_{j\sigma(j)} + C_{j\sigma(j)}) \\ &= \partial_t d. \quad \square \end{aligned}$$

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