



Asymptotic probabilities for conformal restriction measures for an annulus

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ARTICLE INFO

Article history:

Received 31 March 2008

Available online 29 August 2008

Submitted by S. Ruscheweyh

Keywords:

Conformal restriction

Annulus

Elliptic functions

ABSTRACT

We show that for the conformal restriction measure with exponent b in the unit disk on hulls γ connecting e^{ix} to 1 the probability of the event that γ avoids the disk of radius q centered at zero decays like $\exp(-b\pi x/(1-q))$ if either $b \in [5/8, 1] \cup [5/4, \infty)$ and $x \in (0, \pi]$, or if $b \in (1, 5/4)$, $x \in (0, \pi)$, and $bx \leq \pi$.

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1. Introduction

Denote $D \subsetneq \mathbb{C}$ a simply connected domain. If z_1, z_2 are two distinct points on the boundary of D , then a *hull* connecting z_1 and z_2 in D is a closed and connected subset γ of \bar{D} such that $\gamma \cap \partial D = \{z_1, z_2\}$ and $D \setminus \gamma$ has exactly two components. Denote P_{D, z_1, z_2}^b the conformal restriction measure of exponent b and total mass 1. P_{D, z_1, z_2}^b is a probability measure supported on hulls γ connecting z_1 and z_2 in D . For more on the definition of restriction measures, their existence and properties, see [7], and [6, Chapter 9]. It is sometimes convenient to view γ itself as a hull-valued random variable. Then P_{D, z_1, z_2}^b is the (probability-) law of γ . The measure P_{D, z_1, z_2}^b exists for all $b \geq 5/8$. For example, the chordal Schramm–Loewner evolution with parameter $\kappa = 8/3$ gives rise to the conformal restriction measure with exponent $b = 5/8$, and the *filling* of a Brownian excursion gives rise to the conformal restriction measure with exponent $b = 1$. We will assume that the boundary of D near z_1 and z_2 is analytic. Let $a < 0$, $q = e^a$, and denote A_q the annulus $\{z: q < |z| < 1\}$, and \mathbb{U} the unit disk $\{z: |z| < 1\}$. In this note we derive the asymptotics of the non-intersection probability

$$(a, b, x) \in [-\infty, 0] \times [5/8, \infty) \times [0, 2\pi] \mapsto F(a, b, x) \equiv P_{\mathbb{U}, e^{ix}, 1}^b(\gamma \subset A_q)$$

as $a \nearrow 0$. The defining properties for conformal restriction measures are *conformal invariance*, i.e. for any conformal map $f: D \rightarrow f(D)$ and any subdomain D' of D , we have

$$P_{D, z_1, z_2}^b(\gamma \subset D') = P_{f(D), f(z_1), f(z_2)}^b(\gamma \subset f(D')),$$

and *restriction* in the sense that if D' is a simply connected subdomain of D such that $D \setminus D'$ is bounded away from z_1, z_2 , then

$$P_{D, z_1, z_2}^b(\cdot \mid \gamma \subset D') = P_{D', z_1, z_2}^b(\cdot).$$

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The combination of these properties is called Γ -covariance in [7]. An immediate consequence of [7, Proposition 3.3, Remark 3.8] is

$$P_{D, z_1, z_2}^b(\gamma \subset D') = |f'(z_1)f'(z_2)|^b, \quad (1)$$

where f maps D' conformally onto D , fixing z_1 and z_2 . Note that by the Riemann mapping theorem there exists a 1-parameter family of such maps f . However, the expression $|f'(z_1)f'(z_2)|$ is independent of which member of the family is chosen.

We will also use a generalization of a result of Beffara, see [2], which says that if D' is a subdomain of D , not necessarily simply connected, such that $\partial D \subset \partial D'$, and f maps D' conformally onto $f(D')$ such that $f(D') \subset D$, $f(\partial D) = \partial D$, and $f(z_{1,2}) = z_{1,2}$, then

$$\bullet \quad P_{D, z_1, z_2}^b(\gamma \subset D') = P_{D, z_1, z_2}^b(\gamma \subset f(D')) |f'(z_1)f'(z_2)|^b.$$

We give a proof of this fact in a particular case in Lemma 2.2.

The key to our analysis is the transformation of the annulus A_q to the unit disk \mathbb{U} with a horizontal slit $[-L, L]$ along the real axis. What our estimate shows for example is that the probability that $\text{SLE}_{8/3}$ in the unit disk from i to $-i$ stays in a thin annulus A_q is, to leading order, the same as hitting at least one of the two real segments $(-1, -L), (L, 1)$ which are not part of the slit.

An estimate closely related to ours appears in [9, Lemma 18]. There, the general form of the estimate is derived from an excursion representation of $\text{SLE}_{8/3}$ and it is stated that, and briefly indicated how, the explicit values of the constants can be derived from a comparison argument. In this paper we carry through such a comparison argument in detail, and the general form of the estimate is established together with the constants at once. The upper bound is more subtle and requires the majority of the work. The upper bound, or rather our lack of finding a better one, is also the reason why there is a gap in the parameter range for which we obtain the asymptotic behavior.

A related estimate for the asymptotic behavior also appears in [3], derived using Coulomb gas techniques. In both, [9] and [3], the aim is to find—asymptotically—the weight, according to the conformally invariant measure on self-avoiding loops, of the loops which surround an annulus.

The following result and technique of proof generalize (from $b = 5/8$ to $b \in [5/8, \infty)$) those contained in our paper [1].

2. Asymptotic behavior of the non-intersection probability

For each $q \in [0, 1)$ there exists a unique $L = L(q) \in [0, 1)$ such that A_q and $\mathbb{U} \setminus [-L, L]$ are conformally equivalent. As q increases to 1, L increases to 1 as well. Denote f the conformal equivalence, normalized by $f(1) = 1$. For $x \in (0, \pi]$, let $z_1 = e^{ix/2}$, $z_2 = e^{-ix/2}$. By symmetry, if $w_{1,2} = f(z_{1,2})$, then $w_2 = \bar{w}_1$.

In what follows we will mean by $h(a) \asymp g(a)$ as $a \nearrow 0$, that

$$\lim_{a \nearrow 0} \log h(a) / \log g(a) = 1.$$

Lemma 2.1. For $x \in (0, \pi]$ and $q = e^a \in (0, 1)$, we have

$$1 - L \asymp e^{\frac{\pi^2}{4a}}, \quad \text{and} \quad |f'(z_1)| \asymp |1 - f(z_1)| \asymp e^{\frac{\pi}{4a}(\pi - x)}$$

as $a \nearrow 0$.

Proof. From [8, Chapter VI, Section 3],

$$f(z) = L \operatorname{sn} \left(\frac{2iK}{\pi} \log \frac{z}{q} + K; q^4 \right),$$

where $\operatorname{sn}(z)$ is the analytic function for which $\operatorname{sn}'(0) = 1$ and which maps the rectangle $\{z: -K < \Re z < K, 0 < \Im z < iK'\}$ onto the upper half-plane in such a way that $\operatorname{sn}(\pm K) = \pm 1$ and $\operatorname{sn}(\pm K + iK') = \pm k^{-1}$. Furthermore, $q^4 = \exp(-\pi K'/K)$, and $L = \sqrt{k}$. It is classical that $\operatorname{sn}'(z) = [(1 - \operatorname{sn}^2(z))(1 - k^2 \operatorname{sn}^2(z))]^{1/2}$. Thus

$$f'(z) = (2iK/\pi z) [(L^2 - f^2(z))(1 - L^2 f^2(z))]^{1/2}. \quad (2)$$

Define h, τ by $q^4 = h = e^{i\pi\tau}$, and set $v = \frac{i}{\pi} \log \frac{z_1}{q} + \frac{1}{2}$. Then it follows from [4, II, 3], that

$$L = \frac{\theta_2(0|\tau)}{\theta_3(0|\tau)}, \quad \text{and} \quad f(z) = \frac{\theta_1(v|\tau)}{\theta_0(v|\tau)}.$$

Here

$$\theta_1(v|\tau) = 2 \sum_{n=0}^{\infty} (-1)^n h^{(n+1/2)^2} \sin(2n+1)\pi v,$$

$$\theta_2(v|\tau) = 2 \sum_{n=0}^{\infty} h^{(n+1/2)^2} \cos(2n+1)\pi v,$$

$$\theta_3(v|\tau) = 1 + 2 \sum_{n=1}^{\infty} h^{n^2} \cos 2n\pi v,$$

$$\theta_0(v|\tau) = 1 + 2 \sum_{n=1}^{\infty} (-1)^n h^{n^2} \cos 2n\pi v.$$

Using linear transformations of theta functions we may write

$$\frac{\theta_2(0|\tau)}{\theta_3(0|\tau)} = \frac{\theta_0(0|-\frac{1}{\tau})}{\theta_3(0|-\frac{1}{\tau})}, \quad \text{and} \quad \frac{\theta_1(v|\tau)}{\theta_0(v|\tau)} = i \frac{\theta_1(\frac{v}{\tau}|-\frac{1}{\tau})}{\theta_2(\frac{v}{\tau}|-\frac{1}{\tau})}.$$

Hence, if $h' = \exp(-i\pi/\tau)$, and using the series representation of θ_0 and θ_3 , we get

$$L = \frac{1 + 2 \sum_{n=1}^{\infty} (-1)^n (h')^{n^2}}{1 + 2 \sum_{n=1}^{\infty} (h')^{n^2}} = 1 - 4h' + O((h')^2),$$

which is the first statement of the lemma. For the second, we use the infinite product representation of θ_1 and θ_2 , giving

$$i \frac{\theta_1(\frac{v}{\tau}|-\frac{1}{\tau})}{\theta_2(\frac{v}{\tau}|-\frac{1}{\tau})} = \frac{e^{2i\pi v/\tau} - 1}{e^{2i\pi v/\tau} + 1} \prod_{n=1}^{\infty} \frac{(1 - (h')^{2n} e^{2i\pi v/\tau})(1 - (h')^{2n} e^{-2i\pi v/\tau})}{(1 + (h')^{2n} e^{2i\pi v/\tau})(1 + (h')^{2n} e^{-2i\pi v/\tau})}.$$

Since $\exp(2i\pi v/\tau) = i \exp(-(\pi/4a)(\pi - x))$, the infinite product is $1 + O(\exp(\pi^2/(4a)))$, and

$$\frac{e^{2i\pi v/\tau} - 1}{e^{2i\pi v/\tau} + 1} = 1 + 2ie^{\frac{\pi}{4a}(\pi-x)} + O(e^{\pi^2/(4a)}),$$

as $a \nearrow 0$. Using Eq. (2), the lemma now follows. \square

Recall that $z_1 = e^{ix/2}$, $w_1 = f(z_1)$, and set $u = i(1 + w_1)/(1 - w_1)$. The following result is analogous to a result in [2]. We will give a direct argument.

Lemma 2.2. *The probability $P_{\mathbb{U}, e^{ix}, 1}^b(\gamma \subset A_q)$ is equal to*

$$P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left[\frac{1-L}{1+L}, \frac{1+L}{1-L} \right] = \emptyset \right) \left| \frac{f'(z_1)(1-z_1)}{1-f(z_1)} \right|^{2b}.$$

Proof. Denote B a simple curve connecting the inner and outer boundary of A_q , so that B is bounded away from z_1 and z_2 . Denote ϕ a conformal map from $A_q \setminus B$ onto \mathbb{U} such that $\phi(z_{1,2}) = z_{1,2}$, and ψ a conformal map from $f(A_q \setminus B)$ onto \mathbb{U} such that $\psi(w_{1,2}) = w_{1,2}$. Then, by conformal restriction [7],

$$\begin{aligned} P_{\mathbb{U}, z_1, z_2}^b(\gamma \subset A_q \setminus B) &= |\phi'(z_1)\phi'(z_2)|^b, \\ P_{\mathbb{U}, w_1, w_2}^b(\gamma \subset f(A_q \setminus B)) &= |\psi'(w_2)\psi'(w_2)|^b. \end{aligned} \quad (3)$$

Since $T \equiv \phi \circ f \circ \psi^{-1}$ maps \mathbb{U} onto \mathbb{U} and sends $w_{1,2}$ to $z_{1,2}$, there is a pair $w_0, z_0 \in \partial\mathbb{U}$ such that T is the linear transformation given by

$$\frac{T(w) - w_1}{T(w) - w_2} \cdot \frac{w_0 - w_2}{w_0 - w_1} = \frac{z - z_1}{z - z_2} \cdot \frac{z_0 - z_2}{z_0 - z_1}.$$

A calculation gives

$$T'(w_1)T'(w_2) = \left(\frac{z_1 - z_2}{w_1 - w_2} \right)^2,$$

which together with $|f'(z_1)| = |f'(z_2)|$ implies

$$P_{\mathbb{U}, z_1, z_2}^b(\gamma \subset A_q \setminus B) = P_{\mathbb{U}, w_1, w_2}^b(\gamma \subset f(A_q \setminus B)) \left| \frac{f'(z_1)(z_1 - z_2)}{w_1 - w_2} \right|^{2b}. \quad (4)$$

There is a countable collection of simple curves B_n connecting the inner and outer boundary of A_q as above such that $\{\gamma \subset A_q\} = \bigcup_n \{\gamma \subset A_q \setminus B_n\}$. Hence

$$P_{\mathbb{U}, z_1, z_2}^b(\gamma \subset A_q) = \lim_{N \rightarrow \infty} P_{\mathbb{U}, z_1, z_2}^b\left(\bigcup_{n=1}^N \{\gamma \subset A_q \setminus B_n\}\right).$$

Applying the inclusion/exclusion formula to the probability on the right shows that Eq. (4) also holds if $A_q \setminus B$ is replaced by A_q . Finally, by conformal invariance,

$$P_{\mathbb{U}, w_1, w_2}^b(\gamma \subset f(A_q)) = P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left[\frac{1-L}{1+L}, \frac{1+L}{1-L}\right] = \emptyset\right). \quad \square$$

Note that because $x \in (0, \pi]$ we have $\arg z_1, \arg w_1 \in (0, \pi/2]$ and so $u \leq -1$. We will use the following lower and upper bounds:

$$\begin{aligned} & P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left[\frac{1-L}{1+L}, \frac{1+L}{1-L}\right] = \emptyset\right) \\ & \geq P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) + P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left[\frac{1-L}{1+L}, \infty\right) = \emptyset\right) \\ & = P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) + P_{\mathbb{H}, \frac{1}{u}, -\frac{1}{u}}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right), \end{aligned} \quad (5)$$

and

$$\begin{aligned} & P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left[\frac{1-L}{1+L}, \frac{1+L}{1-L}\right] = \emptyset\right) \\ & \leq P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) + P_{\mathbb{H}, \frac{1}{u}, -\frac{1}{u}}^b\left(\gamma \cap i\left[\frac{1+L}{1-L}, \infty\right) = \emptyset\right) \\ & \quad + P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left(0, \frac{1-L}{1+L}\right) \neq \emptyset, \gamma \cap i\left(\frac{1+L}{1-L}, \infty\right) \neq \emptyset\right). \end{aligned} \quad (6)$$

For $c \in \mathbb{R}, d > 0$, set

$$g_{c,d}(z) = \frac{|c|}{\sqrt{c^2 + d^2}} \sqrt{z^2 + d^2}.$$

Then $g_{c,d}$ maps $\mathbb{H} \setminus i(0, d]$ conformally onto \mathbb{H} such that $g_{c,d}(\pm c) = \pm c$. Furthermore,

$$|g'_{c,d}(c)g'_{c,d}(-c)| = \frac{c^4}{(c^2 + d^2)^2},$$

and so by conformal restriction

$$P_{\mathbb{H}, c, -c}^b(\gamma \cap i(0, d] = \emptyset) = [c^2 / (c^2 + d^2)]^{2b}. \quad (7)$$

Corollary 2.3. *We have*

$$P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) + P_{\mathbb{H}, \frac{1}{u}, -\frac{1}{u}}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) \asymp e^{\frac{b\pi x}{a}}$$

as $a \nearrow 0$.

Proof. By (7),

$$P_{\mathbb{H}, u, -u}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) = \left(\frac{u(1-L)}{1+L}\right)^{4b} \left(1 + \frac{(u(1-L))^2}{(1+L)^2}\right)^{-2b},$$

and from Lemma 2.1

$$\left(\frac{u(1-L)}{1+L}\right)^{4b} \left(1 + \left(\frac{u(1-L)}{1+L}\right)^2\right)^{-2b} \asymp e^{\frac{b\pi x}{a}}.$$

Similarly,

$$P_{\mathbb{H}, \frac{1}{u}, -\frac{1}{u}}^b\left(\gamma \cap i\left(0, \frac{1+L}{1-L}\right] = \emptyset\right) \asymp e^{\frac{b\pi^2}{a} + \frac{b\pi}{a}(\pi - x)},$$

so that this term is negligible compared to the first if $0 < x < \pi$, and of the same order if $x = \pi$. \square

Lemma 2.4. *We have*

$$P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset, \gamma \cap i \left(\frac{1+L}{1-L}, \infty \right) \neq \emptyset \right) \asymp e^{\pi^2/a},$$

as $a \nearrow 0$.

Proof. First,

$$\begin{aligned} & P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset, \gamma \cap i \left(\frac{1+L}{1-L}, \infty \right) \neq \emptyset \right) \\ &= P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset \right) + P_{\mathbb{H}, \frac{1}{u}, -\frac{1}{u}}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset \right) \\ &\quad - P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left(\left(0, \frac{1-L}{1+L} \right) \cup \left(\frac{1+L}{1-L}, \infty \right) \right) \neq \emptyset \right). \end{aligned} \quad (8)$$

The last probability on the right equals

$$P_{\mathbb{U}, w_1, w_2}^b (\gamma \cap ((-1, -L] \cup [L, 1)) \neq \emptyset).$$

To calculate this probability, note that

$$g_L(w) \equiv \frac{1 + w^2 - \sqrt{(1 + w^2)^2 - 4p^2 w^2}}{2pw}$$

maps $\mathbb{U} \setminus ((-1, -L] \cup [L, 1))$ onto \mathbb{U} if $2p = (L + 1/L)$, see [5, Chapter 3]. Here, the square root is chosen so that $g_L(i) = i$. Setting $w = e^{i\varphi}$, this can be written

$$g_L(w) = \begin{cases} \frac{1}{p} \cos \varphi + i \sqrt{1 - \frac{1}{p^2} \cos^2 \varphi}, & \text{if } \varphi \in (0, \pi/2]; \\ \frac{1}{p} \cos \varphi - i \sqrt{1 - \frac{1}{p^2} \cos^2 \varphi}, & \text{if } \varphi \in [-\pi/2, 0). \end{cases} \quad (9)$$

Then

$$g'_L(w) g'_L(\bar{w}) = -\frac{\sin^2 \varphi}{p^2 - 1 + \sin^2 \varphi}.$$

Denote T a (fractional) linear transformation from \mathbb{U} onto \mathbb{U} sending $g_L(w_{1,2})$ onto $w_{1,2}$. Then, as in the proof of Lemma 2.2,

$$T'(g_L(w_1)) T'(g_L(w_2)) = \frac{\sin^2 \varphi}{1 - \frac{1}{p^2} \cos^2 \varphi},$$

where now $\varphi = \arg w_1$. Thus, by conformal restriction,

$$P_{\mathbb{U}, w_1, w_2}^b (\gamma \cap ((-1, -L] \cup [L, 1)) \neq \emptyset) = 1 - \left[\frac{p \sin^2 \varphi}{p^2 - 1 + \sin^2 \varphi} \right]^{2b}. \quad (10)$$

Finally, from the definition of u and φ in terms of w_1 , it follows that $u = -\cot(\varphi/2)$ and so $4/\sin^2 \varphi = (u + 1/u)^2$. A calculation now gives

$$\frac{p^2 - 1 + \sin^2 \varphi}{p \sin^2 \varphi} = 1 + \left(\frac{1-L}{1+L} \right)^2 \left(u^2 + \frac{1}{u^2} \right) + \frac{(1-L)^4}{8(L+L^3)} \left[2 + \left(\frac{1-L}{1+L} \right)^2 \left(u^2 + \frac{1}{u^2} \right) \right]. \quad (11)$$

On the other hand, (7) implies

$$P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset \right) = 1 - \left(1 + \left(\frac{1-L}{1+L} \right)^2 \frac{1}{u^2} \right)^{-2b} \quad (12)$$

and

$$P_{\mathbb{H}, \frac{1}{u}, -\frac{1}{u}}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset \right) = 1 - \left(1 + \left(\frac{1-L}{1+L} \right)^2 u^2 \right)^{-2b}. \quad (13)$$

Combining (12), (13), (10), and (8), we get

$$\begin{aligned}
P_{\mathbb{H}, u, -u}^b \left(\gamma \cap i \left(0, \frac{1-L}{1+L} \right) \neq \emptyset, \gamma \cap i \left(\frac{1+L}{1-L}, \infty \right) \neq \emptyset \right) \\
= 1 - \left(1 + \left(\frac{1-L}{1+L} \right)^2 \frac{1}{u^2} \right)^{-2b} + 1 - \left(1 + \left(\frac{1-L}{1+L} \right)^2 u^2 \right)^{-2b} - 1 + \left(\frac{p^2 - 1 + \sin^2 \varphi}{p \sin^2 \varphi} \right)^{-2b}.
\end{aligned} \quad (14)$$

Using (11), straightforward expansion of the right-hand side of (14) shows it to be equal to

$$\frac{b(2b-1)}{8}(1-L)^4 + \frac{b(2b-1)}{4}(1-L)^5 + (1-L)^4 O(u^2(1-L)^2). \quad \square$$

Theorem 2.5. Let $q = e^a \in (0, 1)$. Then

$$F(a, b, x) = P_{\mathbb{U}, e^{ix}, 1}^b (\gamma \subset A_q) \asymp \exp \left(\frac{b\pi x}{a} \right) \quad (15)$$

as $a \nearrow 0$, if either

$$(b, x) \in ([5/8, 1] \cup [5/4, \infty)) \times (0, \pi], \quad (16)$$

or

$$(b, x) \in (1, 5/4) \times (0, \pi) \quad \text{and} \quad bx \leq \pi. \quad (17)$$

Proof. Eq. (15) holds as long as the difference between the upper and lower bound, estimated in Lemma 2.4, is not bigger than the lower bound from Corollary 2.3, i.e. for $(b, x) \in [5/8, 1] \times (0, \pi]$ or for $(b, x) \in [5/4, \infty) \times [0, \pi]$ such that $bx \leq \pi$.

For the remaining cases we use the following property of restriction measures: If γ and γ' are independent and with respective laws P_{D, z_1, z_2}^b and $P_{D, z_1, z_2}^{b'}$, then the filling of $\gamma \cup \gamma'$ has law $P_{D, z_1, z_2}^{b+b'}$, see [7, Remark 3.6]. Here the filling of $\gamma \cup \gamma'$ is the smallest simply connected subdomain of \bar{D} containing $\gamma \cup \gamma'$. This property and the definition extend to any finite number of hulls $\gamma_1, \dots, \gamma_n$ by induction.

Let now $b \in [5/4, \infty)$. Then $b = b_1 + \dots + b_n$ for some $n \in \mathbb{Z}^+$ and $b_1, \dots, b_n \in [5/8, 1]$. Corollary 2.3 provides a lower bound for $F(a, b, x)$ which is $\asymp \exp(b\pi x/a)$ as $a \nearrow 0$ for all $x \in (0, \pi]$. For the upper bound, let $\gamma_1, \dots, \gamma_n$ be independent with respective laws $P_{\mathbb{U}, e^{ix}, 1}^{b_k}$, $1 \leq k \leq n$. Then the event that the filling of $\gamma_1 \cup \dots \cup \gamma_n$ is contained in A_q is a subset of the event $\{\gamma_1 \cup \dots \cup \gamma_n \subset A_q\}$. Hence

$$P_{\mathbb{U}, e^{ix}, 1}^b (\gamma \subset A_q) \leq \prod_{k=1}^n P_{\mathbb{U}, e^{ix}, 1}^{b_k} (\gamma \subset A_q) \asymp e^{b\pi x/a}. \quad \square$$

Remark 2.6. We believe the estimate (15) holds for all $b \geq 5/8$ and $x \in (0, \pi]$ but a proof of this statement likely requires an upper bound closer to the lower bound than the upper bound we use.

References

- [1] Robert O. Bauer, Restricting SLE(8/3) to an annulus, *Stochastic Process. Appl.* 117 (9) (2007) 1165–1188.
- [2] Vincent Beffara, Thesis, University of Paris, Orsay, 2003.
- [3] John Cardy, The $O(n)$ model on the annulus, preprint, math-ph/0604043.
- [4] Adolf Hurwitz, Vorlesungen über allgemeine Funktionentheorie und elliptische Funktionen, fourth ed., Springer-Verlag, Berlin, 1964.
- [5] V.I. Ivanov, M.K. Trubetskoy, Handbook of Conformal Mapping with Computer-aided Visualization, CRC Press, Boca Raton, Florida, 1995.
- [6] Gregory F. Lawler, Conformally Invariant Processes in the Plane, *Math. Surveys Monogr.*, vol. 114, American Mathematical Society, Providence, RI, 2005.
- [7] Gregory Lawler, Oded Schramm, Wendelin Werner, Conformal restriction: The chordal case, *J. Amer. Math. Soc.* 16 (4) (2003) 917–955 (electronic).
- [8] Zeev Nehari, Conformal Mapping, McGraw-Hill, New York, 1952.
- [9] Wendelin Werner, The conformally invariant measure on self-avoiding loops, preprint, math.PR/0511605.