



Stochastic phase field α -Navier-Stokes vesicle-fluid interaction model.Ludovic Goudenège^a, Luigi Manca^{b,*}^a*CNRS, Fédération de Mathématiques de CentraleSupélec FR 3487, Univ. Paris-Saclay, CentraleSupélec, F-91190 Gif-sur-Yvette, France*^b*LAMA, Univ. Gustave Eiffel, UPEM, Univ. Paris Est Creteil, CNRS, F-77447 Marne-la-Vallée, France*

Abstract

We consider a stochastic perturbation of the phase field α -Navier-Stokes model with vesicle-fluid interaction. It consists in a system of nonlinear evolution partial differential equations modeling the fluid-structure interaction associated to the dynamics of an elastic vesicle immersed in a moving incompressible viscous fluid. This system of equations couples a phase-field equation -for the interface between the fluid and the vesicle- to the α -Navier-Stokes equation -for the viscous fluid- with an extra nonlinear interaction term, namely the bending energy.

The stochastic perturbation is an additive space-time noise of trace class on each equation of the system. We prove the existence and uniqueness of solution in classical spaces of L^2 functions with estimates of non-linear terms and bending energy. It is based on a priori estimate about the regularity of solutions of finite dimensional systems, and tightness of the approximated solution.

Keywords: Navier-Stokes, Camassa-Holm, Lagrange Averaged α , stochastic partial differential equations, vesicle, interaction model.

2000 MSC: 60H15, 60H30, 37L55, 35Q30v 35Q35, 76D05

*Corresponding author

Email addresses: goudenège@math.cnrs.fr (Ludovic Goudenège), luigi.manca@u-pem.fr (Luigi Manca)

Introduction and main results

This paper is devoted to study a random perturbation of the equations governing the dynamic of an elastic vesicle immersed in a moving incompressible viscous fluid, whose deterministic model have been studied in [13] and [35].

According to [37], these equations are key research in the study of the dynamics of cells in fluid media. This type of models are of crucial importance in biology, where the analysis of the deformation of vesicles immersed in fluids is central topic. In particular we can refer to the articles on the biological aspects (see [1, 4, 5, 14, 15, 17, 33]). In all these articles there is a common idea about usefulness of phase field approaches. The phase field approaches, compared to sharp interface models, are natural ways to include several important physical aspects of the phenomenon being considered, without complexity of the free-boundary value problems, both in the theoretical and numerical aspects.

First consider the α -Navier-Stokes equation which reads, on the time interval $[0, T]$, on smooth, open and bounded space domain $Q \subset \mathbb{R}^N$ in dimension $N = 2$ or 3 , with ν the constant viscosity, and ρ the constant density of the incompressible fluid:

$$\begin{cases} \partial_t u + (w \cdot \nabla)u + (\nabla w)^T \cdot u + \frac{1}{\rho} \nabla p = f + \nu \Delta u, \\ u = w - \alpha^2 \Delta w - \nabla q, \\ \operatorname{div}(u) = \operatorname{div}(w) = 0. \end{cases} \quad (0.1)$$

where Δ is the Laplace operator and f is the forcing. The unknowns¹ are the random fields p and u (also q and w), which respectively represent the (modified²) pressure and the averaged velocity vector field of the point x at time t . Both unknowns u and w have homogeneous Dirichlet boundary conditions, and the pressures p and q are defined up to an additive term which could be used to stay divergence free. This model takes part of a general class of regularized models for high Reynolds number flows, firstly proposed by Leray in [31, 32] for Euler equations. Some authors stress that, from the biological point of view, the α -Navier-Stokes type equations are relevant since they are adequate for flows with high Reynolds number (like in turbulence), which may occur in some biological situations. This model is also known as viscous Camassa-Holm or the Lagrangian Averaged Navier-Stokes- α (LANS- α) model. These models have been introduced by Holm, Marsden and Ratiu in [29, 30]. It has been studied in the deterministic case by Foias, Holm and Titi (see [21, 22]) which have obtained the necessary estimations about the non-linear term in the Navier-Stokes equation in periodic domain. There are also works in alternative conditions about domain and boundary conditions in [6, 7, 25]. See also [18] for the link between Camassa-Holm and LANS- α models.

In [13] and [35] the authors have considered α -Navier-Stokes model for the fluid coupled with a phase field equation for the membrane of the vesicle. They have introduced a forcing term f which is a non-linear additive term depending of the phase field unknown. The form of the interaction is given by the variational derivative of a bending energy of the membrane of the vesicle. We obtain a system of interaction in the space-time domain $[0, T] \times Q$ between the fluid and the phase-field equations under the form :

$$\begin{cases} \partial_t(w + \alpha^2 A w) + \nu A(w + \alpha^2 A w) + \tilde{B}(w, w + \alpha^2 A w) = \mathcal{P} \left(\frac{\delta E(\phi)}{\delta \phi} \nabla \phi \right), \\ \operatorname{div}(w) = 0, \\ \phi_t + w \cdot \nabla \phi = -\gamma \frac{\delta E(\phi)}{\delta \phi}, \end{cases} \quad (0.2)$$

where ϕ is the phase field unknown/order parameter which describes the membrane of the vesicle, with the linear Stokes operator $A = -\mathcal{P}\Delta$, the Leray orthogonal projector \mathcal{P} on divergence free space H , and the non-linear operator \tilde{B} which will be described later.

¹With $u = (u_1, u_2, u_3)$ in dimension $N = 3$ or $u = (u_1, u_2)$ in dimension $N = 2$.

²Here the modified or hydrodynamic pressure satisfies $p = \pi - \frac{\rho}{2}|w|^2$ where π is the pressure.

This unknown ϕ takes the values $+1$ outside the membrane and -1 inside, with a thin transition width characterized by a small positive parameter ε . The surface of the membrane corresponds to the points where $\phi = 0$, which is actually a very complex area described by the level-set approach, but not explicitly considered in the phase field approach, or in various numerical approaches. The term $\frac{\delta E(\phi)}{\delta \phi}$ is sometimes called the *chemical potential*. It is multiplied by the constant γ which is a positive real number controlling the strength of the chemical potential. Moreover this term can be modeled using various description, depending of the physical consideration about the vesicle.

It is assumed that the energy associated with the deformation of the vesicle membrane comes mainly from the bending energy. Actually this energy is not directly well adapted to *a priori* estimate of quantities related to ϕ (like its norm in Sobolev spaces) since the vesicle tends to minimize the quantity

$$f_\varepsilon(\phi) = -\varepsilon \Delta \phi - \frac{\phi}{\varepsilon} (1 - \phi^2),$$

by minimization of the penalized bending energy given by

$$\mathcal{E}_\varepsilon(\phi) = \frac{k}{2\varepsilon} \int_Q \left(\varepsilon \Delta \phi + \frac{\phi}{\varepsilon} (1 - \phi^2) \right)^2 dx = \frac{k}{2\varepsilon} \int_Q f_\varepsilon(\phi)^2 dx,$$

with the physical parameter k of low relevance here. This bending energy is clearly not a norm or the sum of two competitive behaviors like in classical Allen-Cahn or Cahn-Hilliard equations. Although it implies only a second-order differential operator, this energy is more close to a fourth-order differential linearity as in the Cahn-Hilliard model. Thus the difficulty of the model comes from this form of energy. Moreover, knowing that the volume and the surface area of the vesicle are basically preserved, we penalize the bending energy \mathcal{E}_ε by adding extra terms to form the total energy E :

$$E(\phi) = \mathcal{E}_\varepsilon(\phi) + \frac{1}{2} M_1 (\mathcal{A}(\phi) - a)^2 + \frac{1}{2} M_2 (\mathcal{B}_\varepsilon(\phi) - b)^2,$$

where

$$\begin{aligned} \mathcal{A}(\phi) &= \int_Q \phi \, dx, \\ \mathcal{B}_\varepsilon(\phi) &= \int_Q \left(\frac{\varepsilon}{2} |\nabla \phi|^2 + \frac{1}{4\varepsilon} (\phi^2 - 1)^2 \right) dx, \end{aligned}$$

with M_1, M_2 which are (large) constants used to enforce that the volume and the surface area of the vesicle remain the same. The constants a and b are physical parameters related to the actual volume and surface area of the vesicle (see [16] for details).

Finally -and this is the novelty in the modeling- we assume that there exist two stochastic perturbations ξ_w and ξ_ϕ which are the derivative of cylindrical Wiener processes W and Z taking values in $H, L^2(Q)$ respectively, thus formally $\xi_w = dW$ and $\xi_\phi = dZ$ where

$$W = \sum_{j \in \mathbb{N}^*} \beta_j e_j \quad \text{and} \quad Z = \sum_{j \in \mathbb{N}} \beta_{-j} \eta_j$$

for any orthonormal basis $\{e_j\}_{j \in \mathbb{N}^*}$ of H , any orthonormal basis $\{\eta_j\}_{j \in \mathbb{N}}$ of $L^2(Q)$, and any sequences of independent brownian motion $\{\beta_j\}_{j \in \mathbb{Z}}$.

These perturbations are added linearly to both equations of the system of interaction via linear bounded operators Σ and Ξ acting on H and $L^2(Q)$ respectively such that $\Sigma^* \Sigma$ and $\Xi^* \Xi$ are classical covariance operators (see the next section for more details). This modeling could be extended naturally to other infinite Brownian motion as soon as some classical assumptions about regularity are satisfied. It is worth noticing that in [12] the authors have studied a simplified version of this system. They have considered an incompressible viscous fluid in the low Reynolds number regime, namely the Stokes equation instead of nonlinear Navier-Stokes equation, the interacting constant $\gamma = 0$ and $\alpha = 0$. They have also considered a stochastic forcing on the Stokes equation with operator $\Sigma = \sigma \mathcal{P} \sqrt{-\Delta}$, $\sigma \in \mathbb{R}_+$ and $\Xi = 0$, by using a mollification to obtain well defined system.

From a physical perspective, the stochastic perturbation can be seen as an unknown internal microscopic thermal agitation, or a random source. Thus it is natural to assume that these perturbations are compatible with the underlying physical system, for instance the incompressibility or the fluctuation-dissipation theorem should remain valid. In our case, the Hamiltonian is given by $\frac{1}{2}|w|_2^2 + \frac{1}{2}\alpha^2|\nabla w|_2^2 + E(\phi)$ which leads to the fluctuation-dissipation balance

$$\begin{pmatrix} \Sigma^*(I + \alpha^2 A)^{-2}\Sigma & \mathbf{0} \\ \mathbf{0} & \Xi^*\Xi \end{pmatrix} \propto \begin{pmatrix} \nu(I + \alpha^2 A)^{-1}A & \mathbf{0} \\ \mathbf{0} & \gamma I_d \end{pmatrix}.$$

We can remark that the simplified system considered in [12] respects this balance. However the balance suggests that we are not in the trace-class paradigm but we are dealing with stronger noises which require more complex mathematical work. Therefore we will restrict our study to a subcritical case (less energy is injected by noise in the system than the energy dissipated by macroscopic dynamics) assuming trace-class operators.

Hypothesis 0.1. *We assume that*

$$\text{Tr}[\Sigma^*\Sigma] < \infty, \quad \text{Tr}[\Xi^*\Delta^2\Xi] < \infty.$$

where traces are computed on H and $L^2(Q)$ respectively.

These technical assumptions of the noises permit to use the Itô-formula, which is the key to obtain *a priori* estimates depending of the trace of the operators. It seems clear that since the higher order term in the energy E is Δ , then we need a smoothing effect of Ξ . This is enlightened by our hypothesis assuming $\Delta\Xi$ is an Hilbert-Schmidt operator on $L^2(Q)$.

We obtain the abstract formulation of our studied system

$$\begin{cases} dw + \alpha^2 Aw = \left(-\nu A(w + \alpha^2 Aw) - \tilde{B}(w, w + \alpha^2 Aw) + \mathcal{P} \left(\frac{\delta E(\phi)}{\delta \phi} \nabla \phi \right) \right) dt + \Sigma dW_t \\ d\phi = \left(-w \cdot \nabla \phi - \gamma \frac{\delta E(\phi)}{\delta \phi} \right) dt + \Xi dZ_t. \end{cases} \quad (0.3)$$

Moreover this system is endowed by boundary and initial conditions

$$\begin{cases} w = 0, & Aw = 0, & \text{on } [0, T] \times \partial Q, \\ \phi = -1, & \Delta \phi = 0, & \text{on } [0, T] \times \partial Q, \\ u(0, x) = u_0(x) & & \text{on } Q, \\ \phi(0, x) = \phi_0(x) & & \text{on } Q, \end{cases}$$

with initial data u_0 and ϕ_0 .

Remark 0.2. *The apparently extra boundary condition $Aw = 0$ makes sense, since in α -Navier-Stokes model, we study a couple of unknowns w and $u = w + \alpha^2 Aw$ (the pressure q disappears with Leray's projection) which have both homogeneous Dirichlet boundary condition $u = w = 0$ on ∂Q . Thus $\alpha^2 Aw = 0$ on ∂Q .*

This system is composed of two stochastic partial differential equations which are coupled by an energy. So this is clear that the results obtained in this paper about existence and uniqueness of solution can be extended to more general forms of coupling energy, as soon as it permits a control of some norm of ϕ in Hilbert space with space regularity. Actually the studied form of energy is a mixing between fourth-order Cahn-Hilliard equation and second-order Allen-Cahn equation. These types of stochastic equations with additive noise have been studied in many works. For the Cahn-Hilliard equation there are results about existence and uniqueness in [8, 10, 19] with polynomial nonlinearity, and in [11, 23, 24] for singular nonlinearity and space-time white noises, or degenerate noises. We can also cite a result of existence for a stochastic partial differential equation with a mixing between Cahn-Hilliard and Allen-Cahn equation with multiplicative noise. It has been obtained in [2] with estimations on the Green functions in the spirit of [3, 26]. Concerning the stochastic Navier-Stokes equation, we can cite the important work present in [27, 28, 34].

Using approximated equations in finite dimensional space, we have exhibited a priori estimates and compactness of a sequence of solution of these approximated equations. It permits to prove existence (and uniqueness) of weak (martingale) solution obtained by convergence in weak topology of classical spaces $L^2(Q)$. Precisely we have proved the following:

Theorem 0.3.

Let $T > 0$ and $(w_0, \phi_0) \in D(A) \times L^2(Q)$ with $\phi_0 = -1$ on ∂Q .

Assume that the linear operators (Σ, Ξ) satisfy Hypothesis 0.1.

Then there exists a unique weak solution $((w, \phi), (\Omega, \mathcal{F}, \mathbb{P}, (\mathcal{F}_t)_{t \in [0, T]}), (W, Z))$ of problem (0.3). Moreover, for any $k \in \mathbb{N}^*$ there exists a constant $c = c(k, T, w_0, \phi_0) > 0$ such that

$$\begin{aligned} \mathbb{E} \left[\sup_{0 \leq t \leq T} (|w(t)|_2 + |\phi(t)|_2)^k \right] &\leq c, \\ \mathbb{E} \left[\int_0^T (|w|_2^2 + \alpha^2 |\nabla w|_2^2 + E(\phi))^{k-1} \left(\nu (|\nabla w|_2^2 + \alpha^2 |Aw|_2^2) + \gamma \left| \frac{\delta E(\phi)}{\delta \phi} \right|_2^2 \right) ds \right] &\leq c. \end{aligned} \quad (0.4)$$

Finally, ϕ, w are continuous in mean square, that is for any $t_0 \geq 0$ we have

$$\lim_{t \rightarrow t_0} \mathbb{E} [\|w(t) - w(t_0)\|_V^2] = 0, \quad \lim_{t \rightarrow t_0} \mathbb{E} [|\phi(t) - \phi(t_0)|_2^2] = 0.$$

In section 1, we will describe notations about spaces, classical inequalities and nonlinear estimates about the bending energy which are of crucial importance for the proof of the main result. Moreover we describe the definition of a solution of equation (0.3). In section 2, we derive a priori estimate and we prove technical lemmas which will be used in the proof of the main theorem. Finally in section 3, under the hypotheses of 0.3, we prove existence and uniqueness of solution which satisfies 0.3. This result is a corollary of a more general result obtained in Section 3 about existence and uniqueness of solution with an approximation procedure in finite dimensional spaces. In particular we prove continuity of solution with respect to time with values in Sobolev spaces, and L^p integrability of solution with respect to time with values in Sobolev spaces (H^1 for fluid unknown w and H^4 for parameter order ϕ).

1. Spaces, inequalities and nonlinear estimates

The α -Navier-Stokes equation (0.1) can be formulated in the equivalent form given in (0.2). We need to explain this equivalence, since this is the core of the variational formulation. First we introduce the following spaces:

- $\mathcal{C}_0^\infty(Q)$ is the space of infinitely differentiable functions with compact support;
- $L^2(Q)$, $L^p(Q)$, $H^k(Q)$, $H_0^k(Q)$, $W^{p,k}(Q)$ denotes the usual Sobolev spaces for integrability order $p \in \mathbb{N}$ and derivative order $k \in \mathbb{N}$; when the functions are vector-valued in dimension $N = 1, 2, 3$, we write $(L^2(Q))^N$;
- (\cdot, \cdot) denotes the inner product of the Hilbert space $(L^2(Q))^N$, with $N = 1, 2, 3$;
- $|\cdot|_p$ denotes the norm in the space $(L^p(Q))^N$, with $p \in \mathbb{N}$ and $N = 1, 2, 3$;
- $\|\cdot\|_L$ denotes the norm in a generic space L ;
- $\langle \cdot, \cdot \rangle_{L', L}$ denotes the duality between a generic space L and its dual space L' ;
- X is the space $(H_0^1(Q))^N \cap (H^2(Q))^N$;
- H is the closure in $(L^2(Q))^N$ of $\{u \in X : \operatorname{div}(u) = 0\}$;
- V is the closure in $(H^1(Q))^N$ of $\{u \in X : \operatorname{div}(u) = 0\}$;

1.1. The Stokes operator A

We denote by $\mathcal{P} : (L^2(Q))^N \rightarrow H$ the Leray orthogonal projector. The Stokes operator is then defined by

$$A := -\mathcal{P}\Delta : D(A) \rightarrow H,$$

with domain $D(A) = X \cap V \subset H$. The operator A is self adjoint and positive. Its inverse, $A^{-1} : H \rightarrow H$, is a compact self adjoint operator, thus H admits an orthonormal basis $\{e_j\}_{j \in \mathbb{N}^*}$ formed by the eigenfunctions of A , i.e. $Ae_j = \lambda_j e_j$, with $0 < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_j \rightarrow +\infty$. For $\rho \in \mathbb{R}$, the Sobolev spaces $D(A^\rho)$ are the closure of $C_0^\infty(Q)$ with respect to the norm

$$\|x\|_{D(A^\rho)} = \left(\sum_{j=1}^{\infty} (1 + \lambda_j^{2\rho}) \langle x, e_j \rangle^2 \right)^{\frac{1}{2}}.$$

As well known (see, for instance, [22]) the operator A can be continuously extended to $V = D(A^{\frac{1}{2}})$ with values in $V' = D(A^{-\frac{1}{2}})$ such that for all $u, v \in V$

$$\langle Au, v \rangle_{V', V} = (A^{1/2}u, A^{1/2}v) = \int_Q (\nabla u \cdot \nabla v) \, dx,$$

Similarly A^2 can be continuously extended to $D(A)$ with values in $D(A)'$ (the dual space of the Hilbert space $D(A)$) such that for all $u, v \in D(A)$

$$\langle A^2 u, v \rangle_{D(A)', D(A)} = (Au, Av).$$

One can show that there is a constant $c > 0$ such that for all $w \in D(A)$

$$c^{-1}|Aw|_2 \leq \|w\|_{H^2} \leq c|Aw|_2.$$

This operator A could also be used to define a stochastic convolution thanks to the *strongly* continuous semigroup $(e^{tA})_{t \geq 0}$ by the formula

$$W_A(t) = \int_0^t e^{(t-s)A} dW(s),$$

for cylindrical Wiener processes, which could be used for instance to define mild solutions. This is not the choice made here, since we have enough regularity to define solution with variational estimation.

1.2. The bilinear form \tilde{B}

The specific form of α -Navier-Stokes equation (0.1) has been studied in [7] for bounded domains, or in [21] as the Kelvin-filtered Navier-Stokes equation. This equation is also known as the viscous version of the Camassa-Holm equation. It has been studied in [9] and for periodic domain in [22]. But the global well-posedness for the Lagrangian averaged Navier-Stokes (LANS- α) equations on bounded domains have been studied in [18] where the authors describe the equivalence between different formulations. Precisely they show that the α -Navier-Stokes equation (0.1) is equivalent to LANS- α equations under the condition $Aw = 0$ on ∂Q . We do not present all the details, but the central idea is to define a bilinear operator associated to the non-linear part of equation (0.1) in the spirit of the usual bilinear operator $B(w, u) = \mathcal{P}[(w \cdot \nabla)u]$ of Navier-Stokes equations. It is well defined for all $w, u \in (H_0^1(Q))^N$, and such that for all w, u and $v \in V \subset (H_0^1(Q))^N$

$$(B(w, u), v) = -(B(w, v), u).$$

Thus, applying \mathcal{P} to the equation (0.1) and using the identity

$$(w \cdot \nabla)u + (\nabla w)^T u = -w \times (\nabla \times u) + \nabla(u \cdot w),$$

we can see that the nonlinear term of equation (0.1) could be replaced by the bilinear operator \tilde{B} defined for all $w, u \in (H_0^1(Q))^N$ by

$$\tilde{B}(w, u) = -\mathcal{P}[w \times (\nabla \times u)]$$

since $\nabla(u \cdot w)$ is in the orthogonal of V . This operator appears clearly in the Camassa-Holm formulation. Moreover, a direct computation shows that

$$(\tilde{B}(u, v), w) = (B(u, v), w) - (B(w, v), u)$$

for $u, v, w \in V$ (see [22] for details). The next results will be crucial for many proofs:

Proposition 1.1.

(i) The operator \tilde{B} can be extended continuously to $V \times V$ with values in V' ; for all $u, v, w \in V$ it satisfies

$$\begin{aligned} \left| \left\langle \tilde{B}(u, v), w \right\rangle_{V', V} \right| &\leq c \|u\|_H^{1/2} \|u\|_V^{1/2} \|v\|_V \|w\|_V, \\ \left| \left\langle \tilde{B}(u, v), w \right\rangle_{V', V} \right| &\leq c \|u\|_V \|v\|_V \|w\|_H^{1/2} \|w\|_V^{1/2}, \\ \left\langle \tilde{B}(u, v), w \right\rangle_{V', V} &= - \left\langle \tilde{B}(w, v), u \right\rangle_{V', V}, \quad \text{and} \quad \left\langle \tilde{B}(u, v), u \right\rangle_{V', V} = 0. \end{aligned}$$

(ii) Its restriction to $D(A)$ satisfies for all $u \in V$, $v \in H$, $w \in D(A)$ it holds

$$\left| \left\langle \tilde{B}(u, v), w \right\rangle_{D(A)', D(A)} \right| \leq c \|u\|_V \|v\|_H \|w\|_{D(A)}$$

Proof. The proof of (i) is classical and can be found, for instance, on [22]. The statement (ii) follows easily by the estimate

$$\left| \left\langle \tilde{B}(u, v), w \right\rangle_{D(A)', D(A)} \right| \leq c \left(\|u\|_H^{1/2} \|u\|_V^{1/2} \|v\|_H \|Aw\|_H + \|u\|_V \|v\|_H \|w\|_V^{1/2} \|Aw\|_H^{1/2} \right)$$

which can be found, for instance, in [35]. □

1.3. Definition of solutions

We are now able to define the concept of solution of equation (0.2) or more precisely the solution of its abstract form (0.3).

Definition 1.2. Let $T > 0$ and $(w_0, \phi_0) \in D(A) \times L^2(Q)$ with $\phi_0 = -1$ on ∂Q . Assume that the linear operators (Σ, Ξ) satisfy Hypothesis 0.1. We say that $((w, \phi), (\Omega, \mathcal{F}, \mathbb{P}, (\mathcal{F}_t)_{t \in [0, T]}), (W, Z))$ is a weak solution of (0.3) if

- $(\Omega, \mathcal{F}, \mathbb{P}, (\mathcal{F}_t)_{t \in [0, T]})$ is a complete filtered probability space.
- w, ϕ are adapted to the filtration $(\mathcal{F}_t)_{t \in [0, T]}$.

Moreover, \mathbb{P} -a.s.,

- $w \in L^2([0, T]; D(A))$;
- $w + \alpha^2 Aw \in \mathcal{C}([0, T]; D(A)') \cap L^2([0, T]; H)$;
- $\tilde{B}(w, w + \alpha^2 Aw) \in L^2([0, T]; D(A)')$;
- $\phi \in L^2([0, T]; H^2(Q)) \cap \mathcal{C}([0, T]; L^2(Q))$ such that $\phi + 1 = \Delta \phi = 0$ on ∂Q ;
- $\frac{\delta E(\phi)}{\delta \phi} \in L^2([0, T]; L^2(Q))$ (this term will be defined in Section 1.4);

- $w \cdot \nabla \phi \in L^2([0, T]; L^2(Q))$;
- $\frac{\delta E(\phi)}{\delta \phi} \nabla \phi \in L^2([0, T]; D(A)')$;
- For all $\xi \in D(A)$, for all $t \in [0, T]$ we have

$$\begin{cases} \langle w(t) + \alpha^2 A w(t), \xi \rangle = \langle w_0 + \alpha^2 A w_0, \xi \rangle - \nu \int_0^t \langle w + \alpha^2 A w, A \xi \rangle ds \\ \quad - \int_0^t \langle \tilde{B}(w, w + \alpha^2 A w), \xi \rangle ds + \int_0^t \left\langle \frac{\delta E(\phi)}{\delta \phi} \nabla \phi, \xi \right\rangle ds + \langle \Sigma^* \xi, W_t \rangle \\ \phi(t) = \phi_0 - \int_0^t \left(w \cdot \nabla \phi + \gamma \frac{\delta E(\phi)}{\delta \phi} \right) ds + \Xi Z_t. \end{cases} \quad (1.1)$$

1.4. Nonlinear estimates

From now and to the end of this article we will skip the parameters ε and k (set to the value 1) because it does not bring very useful information and the reading will be clearly simplified. For this reason we will be very cautious about cancellation during subtraction of terms.

Let us choose $\{\eta_j\}_{j \in \mathbb{N}^*} \in L^2(Q)$ be the orthonormal basis in $L^2(Q)$ consisting of the eigenfunctions of the Laplacian Δ with homogeneous Dirichlet boundary conditions, and $\eta_0 = 1/\sqrt{|Q|}$.

The variational derivative of E with respect to the variable ϕ at point ϕ in the direction ψ is defined for any $\phi + 1, \psi \in C_0^\infty(Q)$. by

$$\begin{aligned} \left\langle \frac{\delta E}{\delta \phi}(\phi), \psi \right\rangle &= \lim_{h \rightarrow 0} \frac{E(\phi + h\psi) - E(\phi)}{h} = \int_Q \frac{\delta E}{\delta \phi}(\phi) \psi \, dx \\ &= \int_Q f(\phi) (f'(\phi) \psi) \, dx + M_1(\mathcal{A}(\phi) - a) \mathcal{A}(\psi) + M_2(\mathcal{B}(\phi) - b) \int_Q f(\phi) \psi \, dx \end{aligned}$$

Here we have set

$$f'(\phi) \psi = -\Delta \psi + (3\phi^2 - 1)\psi.$$

In this case the variational derivative of E can be identified with

$$\frac{\delta E}{\delta \phi}(\phi) = \Delta^2 \phi - \Delta(\phi^3 - \phi) + (3\phi^2 - 1)f(\phi) + M_1(\mathcal{A}(\phi) - a) + M_2(\mathcal{B}(\phi) - b)f(\phi). \quad (1.2)$$

Proposition 1.3. *There exists $c > 0$ such that for any $\phi + 1 \in C_0^\infty(Q)$ it holds*

$$|\Delta \phi|_2^2 + |\nabla \phi|_2^4 + |\phi \nabla \phi|_2^2 + |\phi|_4^8 + |\phi|_6^6 \leq c(1 + E(\phi)). \quad (1.3)$$

and

$$E(\phi) \leq c(1 + \|\phi\|_{H^2}^8) \quad (1.4)$$

Proof. We have

$$4\mathcal{B}(\phi) = 2|\nabla \phi|_2^2 + |\phi^2 - 1|_2^2 = 2|\nabla \phi|_2^2 + |\phi|_4^4 - 2|\phi|_2^2 + |Q|. \quad (1.5)$$

Then, since $|\phi|_2^2 \leq \frac{1}{2}|Q|^2 + \frac{1}{2}|\phi|_4^4$, there exists $c > 0$ such that

$$|\nabla \phi|_2^2 + |\phi|_4^4 + |\phi|_2^2 \leq c(1 + \mathcal{B}(\phi)). \quad (1.6)$$

Clearly, for some other constant $c > 0$ it holds

$$|\nabla \phi|_2^4 + |\phi|_4^8 + |\phi|_2^4 \leq c(1 + (\mathcal{B}(\phi))^2) \leq c(1 + E(\phi)).$$

At this point, it remains to bound the quantity $|\Delta \phi|_2^2 + |\phi \nabla \phi|_2^2 + |\phi|_6^6$. Using the expression of $f(\phi)$ we get

$$\begin{aligned} 2|f(\phi)|_2^2 &= |-\Delta \phi + \phi(\phi^2 - 1)|_2^2 \\ &= |\Delta \phi|_2^2 - 2\langle \Delta \phi, \phi(\phi^2 - 1) \rangle + |\phi(\phi^2 - 1)|_2^2 \\ &= |\Delta \phi|_2^2 + 2\langle \nabla \phi, 3\phi^2 \nabla \phi - \nabla \phi \rangle + |\phi|_6^6 - |\phi|_2^2 \\ &= |\Delta \phi|_2^2 + 6|\phi \nabla \phi|_2^2 - 2|\nabla \phi|_2^2 + |\phi|_6^6 - |\phi|_2^2 \end{aligned} \quad (1.7)$$

Here we use the fact that $\phi(\phi^2 - 1) \in C_0^\infty(Q)$ in order to perform integration by parts. Thus,

$$|\Delta\phi|_2^2 + |\phi\nabla\phi|_2^2 + |\phi|_6^6 \leq 2|f(\phi)|_2^2 + 2|\nabla\phi|_2^2 + |\phi|_2^2$$

Using the estimate (1.6) and elementary inequalities, there exist constants $c, c' > 0$ such that

$$|\Delta\phi|_2^2 + |\phi\nabla\phi|_2^2 + |\phi|_6^6 \leq 2|f(\phi)|_2^2 + c(1 + \mathcal{B}(\phi)) \leq c'(1 + E(\phi)).$$

Then, (1.3) follows easily. Let us show (1.4). By (1.7) and the embedding $H^1 \subset L^6(Q)$ we find that for some $c > 0$, independent of ϕ it holds

$$2|f(\phi)|_2^2 \leq |\Delta\phi|_2^2 + 6|\phi\nabla\phi|_2^2 + |\phi|_6^6 \leq |\Delta\phi|_2^2 + 6|\phi\nabla\phi|_2^2 + c\|\phi\|_{H^1}^6$$

Moreover, by Hölder's inequality, by Poincaré's inequality $|\phi + 1|_\infty \leq C_{\mathcal{P}}|\nabla\phi|_2$ and by Young's inequality we obtain

$$|\phi\nabla\phi|_2^2 \leq |\phi|_\infty^2 |\nabla\phi|_2^2 \leq (|\phi + 1|_\infty + 1)^2 |\nabla\phi|_2^2 \leq (C_{\mathcal{P}}|\nabla\phi|_2 + 1)^2 |\nabla\phi|_2^2 \leq 2(C_{\mathcal{P}}^2 |\nabla\phi|_2^2 + 1) |\nabla\phi|_2^2.$$

We deduce that for some $c > 0$

$$2|f(\phi)|_2^2 \leq |\Delta\phi|_2^2 + 2c|\nabla\phi|_2^4 + 2c|\nabla\phi|_2^2 + c\|\phi\|_{H^1}^6 \leq c(1 + \|\phi\|_{H^2}^6).$$

The inequality $(\mathcal{A}(\phi) - a)^2 \leq c(1 + |\phi|_2^2)$, with $c = c(a, |Q|)$ follows immediately. By (1.5) and the embedding $H^1 \subset L^4(Q)$ we get

$$4\mathcal{B}(\phi) \leq 2|\nabla\phi|^2 + |\phi|_4^4 + |Q| \leq c2|\nabla\phi|^2 + \|\phi\|_{H^1}^4 + |Q|$$

Then, for some $c > 0$ it holds

$$(\mathcal{B}(\phi) - b)^2 \leq c(1 + \|\phi\|_{H^1}^8).$$

Then, by the bounds obtained above, we deduce that (1.4) holds for some $c > 0$ independent of ϕ . \square

Proposition 1.4. *There exists a constant $c > 0$ such that for any $\phi + 1 \in C_0^\infty(Q)$ it holds*

$$|\Delta^2\phi|_2 \leq \left| \frac{\delta E}{\delta\phi}(\phi) \right|_2 + c(1 + E(\phi)^2) \quad (1.8)$$

Proof. By (1.2) we have

$$|\Delta^2\phi|_2 \leq \left| \frac{\delta E}{\delta\phi}(\phi) \right|_2 + I_1 + I_2 + I_3 + I_4,$$

where

$$\begin{aligned} I_1 &= |(-\Delta)((\phi^2 - 1)\phi)|_2, \\ I_2 &= |(3\phi^2 - 1)f(\phi)|_2, \\ I_3 &= M_1 |(\mathcal{A}(\phi) - a)|_2, \\ I_4 &= M_2 |(\mathcal{B}(\phi) - b)f(\phi)|_2. \end{aligned}$$

For I_1 we have

$$I_1 = 6\phi|\nabla\phi|^2 + 3\phi^2\Delta\phi - \Delta\phi$$

Then by basic inequality we get

$$I_1 \leq 6|\phi|_\infty |\nabla\phi|_4^2 + 3|\phi|_\infty^2 |\Delta\phi|_2 + |\Delta\phi|_2$$

The Poincaré inequality yields $|\phi|_\infty \leq |\phi + 1|_\infty + 1 \leq C_{\mathcal{P}}|\nabla\phi|_2 + 1$ where $C_{\mathcal{P}}$ is the Poincaré constant. Moreover, by the Sobolev embedding $H^1(Q) \subset L^4(Q)$ we get $|\nabla\phi|_4^2 \leq c\|\nabla\phi\|_{H^1}^2$ for some

constant $c > 0$ independent of ϕ . Then, using repeatedly the Young inequality we get that there exists $c > 0$ such that

$$\begin{aligned} I_1 &\leq 6c(C_{\mathcal{P}}|\nabla\phi|_2 + 1)\|\nabla\phi\|_{H^1}^2 + 3(C_{\mathcal{P}}|\nabla\phi|_2 + 1)^2|\Delta\phi|_2 + |\Delta\phi|_2 \\ &\leq c(1 + |\nabla\phi|_2^2 + \|\nabla\phi\|_{H^1}^4 + |\nabla\phi|_2^4 + |\Delta\phi|_2^2 + |\Delta\phi|_2). \end{aligned}$$

Notice that $\|\nabla\phi\|_{H^1} \leq \|\phi\|_{H^2} \leq c(|\phi|_2 + |\nabla\phi|_2 + |\Delta\phi|_2)$ for some $c > 0$ independent of ϕ . Then, still using Young's inequality, there exists a constant $c_1 > 0$ such that

$$I_1 \leq c_1(1 + |\phi|_4^2 + |\nabla\phi|_2^4 + |\Delta\phi|_2^4).$$

For I_2 , using the expression of $f(\phi)$ and the Poincaré inequality $|\phi + 1|_{\infty} \leq C_{\mathcal{P}}|\nabla\phi|_2$ we obtain

$$\begin{aligned} I_2 &\leq (3|\phi|_{\infty}^2 + 1)|f(\phi)|_2 \\ &\leq (3(|\phi + 1|_{\infty} + 1)^2 + 1)(|\Delta\phi|_2 + (|\phi|_6^3 + |\phi|_2)) \\ &\leq (3(C_{\mathcal{P}}|\nabla\phi|_2 + 1)^2 + 1)(|\Delta\phi|_2 + (|\phi|_6^3 + |\phi|_2)). \end{aligned}$$

By applying the inequality $(\alpha + \beta)^2 \leq 2\alpha^2 + 2\beta^2$ repeatedly, we find that there exists a constant $c_2 > 0$ such that

$$I_2 \leq c_2(|\Delta\phi|_2^2 + |\nabla\phi|_2^4 + |\phi|_6^6 + |\phi|_2^2 + 1).$$

Clearly, for I_3 there exists a constant $c_3 = c_3(a, |Q|) > 0$ such that

$$I_3 \leq c_3(|\phi|_2 + 1)$$

For I_4 we have, by the expression of $\mathcal{B}(\phi)$ and $f(\phi)$,

$$\begin{aligned} I_4 &\leq M_2(\mathcal{B}(\phi) + b)|f(\phi)|_2 \\ &\leq \left(\frac{1}{2}|\nabla\phi|_2^2 + \frac{1}{4}(|\phi|_4^4 + 2|\phi|_2^2 + |Q|) + b\right)(|\Delta\phi|_2 + (|\phi|_6^3 + |\phi|_2)) \end{aligned}$$

Using the inequality $(\alpha + \beta)^2 \leq 2\alpha^2 + 2\beta^2$ repeatedly, it is easy to show that there exists a constant $c_4 > 0$ such that

$$I_4 \leq c_4(|\phi|_4^8 + |\phi|_2^4 + |\phi|_6^6 + |\nabla\phi|_2^4 + |\Delta\phi|_2^2 + 1).$$

Taking into account the estimates on I_1, \dots, I_4 , by (1.3) we deduce that there exist $c_5, c > 0$ such that

$$I_1 + I_2 + I_3 + I_4 \leq c_5(|\phi|_2^4 + |\phi|_4^8 + |\phi|_6^6 + |\nabla\phi|_2^4 + |\Delta\phi|_2^4 + 1) \leq c(1 + (E(\phi))^2). \quad \square$$

We recall that we have made the following assumption on the operator Σ and Ξ .

$$\text{Tr}[\Sigma^*\Sigma] < \infty, \quad \text{Tr}[\Xi^*\Delta^2\Xi] < \infty.$$

Since for two separable Hilbert spaces X, Y and for a linear operator $A : X \rightarrow Y$ we have

$$\text{Tr}[A^*A] = \text{Tr}[AA^*]$$

thus $\text{Tr}[\Xi^*\Delta^2\Xi] < \infty$ if and only if $\text{Tr}[\Delta\Xi\Xi^*\Delta] < \infty$.

Remark 1.5. The covariance operator Ξ is defined from $L^2(Q)$ into itself, but since $\text{Tr}[\Xi^*\Delta^2\Xi] < \infty$, there is a smoothing effect of Ξ . Indeed we can see that the operator $\Xi^*\Delta : H^2(Q) \cap H_0^1(Q) \rightarrow L^2(Q)$ is closable and can be extended to a bounded linear operator $\Xi^*\Delta : L^2(Q) \rightarrow L^2(Q)$.

Proof of Remark 1.5. Let $u \in H^2(Q) \cap H_0^1(Q)$. Since $\{\eta_j\}_{j \in \mathbb{N}}$ is an orthonormal basis of $L^2(Q)$ then

$$|\Xi^*\Delta u|_2^2 = \sum_{j=0}^{\infty} \langle \Xi^*\Delta u, \eta_j \rangle^2 = \sum_{j=0}^{\infty} \langle u, \Delta\Xi\eta_j \rangle^2 \leq |u|_2^2 \sum_{j=0}^{\infty} |\Delta\Xi\eta_j|_2^2 = |u|_2^2 \text{Tr}[\Xi^*\Delta^2\Xi].$$

Then by the closed graph theorem we obtain the result. \square

Proposition 1.6. *If $\text{Tr}[\Xi^* \Delta^2 \Xi] < \infty$ then $\text{Tr}[\Xi^* \Xi] < +\infty$. Moreover the sequences*

$$(|\Xi \eta_j|_\infty)_{j \in \mathbb{N}}, \quad (\|\Xi \eta_j\|_{W^{2,2}(Q)})_{j \in \mathbb{N}}, \quad (|\nabla \Xi \eta_j|_3)_{j \in \mathbb{N}}$$

are in $\ell^2(\mathbb{N})$ and there exists a constant $c > 0$ such that

$$\sum_{j=0}^{\infty} \left(|\Xi \eta_j|_\infty^2 + \|\Xi \eta_j\|_{W^{2,2}(Q)}^2 + |\nabla \Xi \eta_j|_3^2 \right) \leq c \text{Tr}[\Xi^* \Delta^2 \Xi].$$

Proof : Since η_j is a component of an orthonormal basis of $L^2(Q)$ then $\langle \Xi^* \Delta^2 \Xi \eta_j, \eta_j \rangle \leq \text{Tr}[\Xi^* \Delta^2 \Xi] < +\infty$ i.e. $\Delta \Xi \eta_j \in L^2(Q)$. This implies that

$$\begin{aligned} \text{Tr}[\Xi^* \Xi] &= |\Xi \eta_0|_2^2 + \sum_{j=1}^{\infty} |\Xi \eta_j|_2^2 \leq |\Xi \eta_0|_2^2 + \|(-\Delta)^{-1}\|_{\mathcal{L}(H)}^2 \sum_{j=1}^{\infty} |\Delta \Xi \eta_j|_2^2 \\ &\leq \max\{1, \|(-\Delta)^{-1}\|_{\mathcal{L}(H)}^2\} \text{Tr}[\Xi^* \Delta^2 \Xi], \end{aligned}$$

and in particular $\Xi \eta_j \in W^{2,2}(Q)$. Denote $M = \max\{1, \|(-\Delta)^{-1}\|_{\mathcal{L}(H)}^2\}$. Taking into account the Sobolev embedding $W^{2,2}(Q) \subset C^{0,\gamma}$ for all $\gamma < 1/2$, there exists $M' > 0$ such that

$$\sum_{j=0}^{\infty} |\Xi \eta_j|_\infty^2 \leq M' \sum_{j=0}^{\infty} \|\Xi \eta_j\|_{W^{2,2}(Q)}^2 \leq 4MM' \sum_{j=0}^{\infty} |\Delta \Xi \eta_j|_2^2 \leq 4MM' \text{Tr}[\Xi^* \Delta^2 \Xi].$$

Finally, by the Sobolev embedding $W^{1,3/2}(Q) \subset L^3(Q)$ and Hölder's inequality there exists $M'' > 0$ such that

$$\begin{aligned} \sum_{j=0}^{\infty} \|\nabla \Xi \eta_j\|_3^2 &\leq M'' \sum_{j=0}^{\infty} \|\Xi \eta_j\|_{W^{1,3/2}(Q)}^2 \leq M'' \sum_{j=0}^{\infty} \left(|\Xi \eta_j|_{3/2}^2 + |\nabla \Xi \eta_j|_{3/2}^2 \right) \\ &\leq M''(|Q|)^{1/3} \sum_{j=0}^{\infty} (|\Xi \eta_j|_2^2 + |\nabla \Xi \eta_j|_2^2) \leq M''(|Q|)^{1/3} \sum_{j=0}^{\infty} \|\Xi \eta_j\|_{W^{2,2}(Q)}^2 < +\infty. \end{aligned}$$

Proposition 1.7. *There exists a constant $c > 0$ such that for any $\phi + 1 \in C_0^\infty(Q)$*

$$\left| \Xi^* \frac{\delta E}{\delta \phi}(\phi) \right|_2 \leq c \left(1 + |\phi|_4^8 + |\phi|_6^6 + |\nabla \phi|_2^4 + |\Delta \phi|_2^2 \right) \quad (1.9)$$

Proof. By (1.2) we have

$$\Xi^* \frac{\delta E}{\delta \phi}(\phi) = I_1 + I_2 + I_3 + I_4 + I_5,$$

where (we recall that $k = \varepsilon = 1$)

$$\begin{aligned} I_1 &= \Xi^* \Delta^2 \phi, \\ I_2 &= \Xi^* (-\Delta) ((\phi^2 - 1)\phi), \\ I_3 &= \Xi^* (3\phi^2 - 1)f(\phi), \\ I_4 &= M_1 \Xi^* (\mathcal{A}(\phi) - a), \\ I_5 &= M_2 \Xi^* (\mathcal{B}(\phi) - b)f(\phi). \end{aligned}$$

By Hypothesis 0.1 and Remark 1.5, there exists $c_1 > 0$ such that

$$|I_1|_2 \leq c_1 |\Delta \phi|_2.$$

For I_2 , still by Hypothesis 0.1 and Remark 1.5 there exists $c > 0$ such that

$$|I_2|_2 \leq c |(\phi^2 - 1)\phi|_2$$

Then by basic inequality we get, for some $c_2 > 0$,

$$|I_2|_2 \leq c_2(|\phi|_6^3 + |\phi|_2).$$

Since Ξ^* is a bounded linear operator, the terms I_3, I_4, I_5 can be estimated as done for Proposition 1.4 to get

$$\begin{aligned} |I_3|_2 &\leq c_3(|\phi|_4^4 + |\Delta\phi|_2^2 + |\phi|_6^6 + |\phi|_2^2 + 1) \\ |I_4|_2 &\leq c_4(|\phi|_2 + 1) \\ |I_5|_2 &\leq c_5(|\nabla\phi|_2^4 + |\phi|_4^8 + |\phi|_2^4 + |\Delta\phi|_2^2 + |\phi|_6^6 + 1). \end{aligned}$$

for some constant c_3, c_4, c_5 independent of ϕ . Taking into account the estimates for $I_i, i = 1, \dots, 5$, the claim follows. \square

The second variational of E in ϕ is a bilinear form on $C_0^\infty(Q) \otimes C_0^\infty(Q)$ and takes the form

$$\begin{aligned} \left(\frac{\delta^2 E}{\delta\phi^2}(\phi)\right)(\psi, \rho) &= \int (f'(\phi)\psi)(f'(\phi)\rho) \, dx + \int_Q f(\phi)(f''(\phi)(\psi, \rho)) \, dx + M_1 \mathcal{A}(\psi)\mathcal{A}(\rho) \\ &\quad + M_2 \left(\int_Q f(\phi)\rho \, dx\right) \left(\int_Q f(\phi)\psi \, dx\right) + M_2(\mathcal{B}(\phi) - b) \int_Q (f'(\phi)\rho)\psi \, dx \end{aligned}$$

where

$$f''(\phi)(\psi, \rho) = 6\phi\psi\rho.$$

When $\psi = \rho$ it takes the form

$$\begin{aligned} \left(\frac{\delta^2 E}{\delta\phi^2}(\phi)\right)(\psi, \psi) &= \int (f'(\phi)\psi)^2 \, dx + \int_Q f(\phi)(f''(\phi)(\psi, \psi)) \, dx + M_1 (\mathcal{A}(\psi))^2 \\ &\quad + M_2 \left(\int_Q f(\phi)\psi \, dx\right)^2 + M_2(\mathcal{B}(\phi) - b) \int_Q (f'(\phi)\psi)\psi \, dx. \end{aligned}$$

Proposition 1.8. *There exists a constant $c > 0$ such that for any $\phi + 1, \psi \in C_0^\infty(Q)$ it holds*

$$\left(\frac{\delta^2 E}{\delta\phi^2}(\phi)\right)(\psi, \psi) \leq c(|\Delta\phi|_2^2 + |\nabla\phi|_2^4 + |\phi|_4^6 + |\phi|_6^6 + 1)(|\psi|_2^2 + |\nabla\psi|_2^2 + |\Delta\psi|_2^2)$$

Proof. Let us write

$$\frac{\delta^2 E(\phi)}{\delta\phi^2}(\psi, \psi) = I_1 + I_2 + I_3 + I_4 + I_5,$$

where

$$\begin{aligned} I_1 &= \int (f'(\phi)\psi)^2 \, dx \\ I_2 &= \int_Q f(\phi)(f''(\phi)(\psi, \psi)) \, dx \\ I_3 &= M_1 (\mathcal{A}(\psi))^2 \\ I_4 &= M_2 \left(\int_Q f(\phi)\psi \, dx\right)^2 \\ I_5 &= M_2(\mathcal{B}(\phi) - b) \int_Q (f'(\phi)\psi)\psi \, dx. \end{aligned}$$

For I_1 we have

$$\begin{aligned} I_1 &= |(f'(\phi)\psi)|_2^2 \\ &= |-\Delta\psi + (3\phi^2 - 1)\psi|_2^2 \\ &\leq (|\Delta\psi|_2 + (3|\phi|_\infty^2 + 1)|\psi|_2)^2 \end{aligned}$$

Since $|\phi|_\infty \leq |\phi + 1|_\infty + 1$ and by Poincaré's inequality there exists a constant $C_{\mathcal{P}} > 0$ such that $|\phi + 1|_\infty \leq C_{\mathcal{P}} |\nabla \phi|_2$, the right hand side is bounded by

$$(|\Delta \psi|_2 + (3(C_{\mathcal{P}} |\nabla \phi|_2 + 1) + 1) |\psi|_2)^2.$$

Then it follows that there exists a constant $d_1 > 0$ such that

$$I_1 \leq d_1 |\Delta \psi|_2^2 + d_1 (|\nabla \phi|_2^2 + 1) |\psi|_2^2.$$

For I_2 we have, using Hölder's inequality and Young's inequality,

$$\begin{aligned} I_2 &= 6 \int_Q (-\Delta \phi + (\phi^2 - 1)\phi) \phi \psi^2 \, dx \\ &\leq 6 \left(\int_Q |\Delta \phi \phi| \, dx + \int_Q |\phi^2 - 1| \phi^2 \, dx \right) |\psi|_\infty^2 \\ &\leq 6 \left(\frac{1}{2} |\Delta \phi|_2^2 + \frac{1}{2} |\phi|_2^2 + |\phi|_4^4 + |\phi|_2^2 \right) |\psi|_\infty^2. \end{aligned}$$

Then there exists d_2 such that

$$I_2 \leq d_2 (|\Delta \phi|_2^2 + |\phi|_4^4 + |\phi|_2^2) |\psi|_\infty^2$$

The term I_3 is easily bounded by

$$I_3 \leq M_1 |Q| |\psi|_2^2 = d_3 |\psi|_2^2,$$

where $d_3 = M_1 |Q|$. For I_4 we have

$$I_4 \leq M_2 |f(\phi)|_2^2 |\psi|_2^2.$$

It is easy to see that

$$|f(\phi)|_2 \leq |\Delta \phi|_2 + (|\phi|_6^3 + |\phi|_2)$$

holds. Taking into account the inequality $(a + b)^2 \leq 2a^2 + 2b^2$, there exists a constant $d_4 > 0$ such that

$$I_4 \leq d_4 (|\Delta \phi|_2^2 + |\phi|_6^6 + |\phi|_2^2) |\psi|_2^2.$$

For I_5 we can use Hölder's inequality to get

$$I_5 \leq M_2 (\mathcal{B}(\phi) + b) |f'(\phi) \psi|_2 |\psi|_2. \quad (1.10)$$

Since

$$\mathcal{B}(\phi) = \frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{4} |\phi^2 - 1|_2^2$$

using Young's inequality and Hölder's inequality we get

$$\mathcal{B}(\phi) \leq \frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{2} (|\phi|_2^2 + |Q|) = \frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{2} (|\phi|_4^4 + |Q|)$$

For the last term on the right-hand side we can argue as for I_1 to get

$$\begin{aligned} |(f'(\phi) \psi)|_2 &\leq |\Delta \psi|_2 + |(3\phi^2 - 1)\psi|_2 \\ &\leq |\Delta \psi|_2 + |3\phi^2 - 1|_2 |\psi|_\infty \\ &\leq |\Delta \psi|_2 + \left(3|\phi|_4^2 + |Q|^{\frac{1}{2}} \right) |\psi|_\infty \end{aligned} \quad (1.11)$$

Taking into account (1.10), (1.4) and (1.11), the term I_5 is bounded by

$$I_5 \leq M_2 \left(\frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{2} (|\phi|_4^4 + |Q|) + b \right) \left(|\Delta \psi|_2 |\psi|_2 + \left(3|\phi|_4^2 + |Q|^{\frac{1}{2}} \right) |\psi|_\infty |\psi|_2 \right)$$

Elementary calculus and inequality $\alpha\beta \leq \alpha^2/2 + \beta^2/2$ show that for some constant $c > 0$,

$$\left(\frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{2} (|\phi|_4^4 + |Q|) \right) \left(3|\phi|_4^2 + |Q|^{\frac{1}{2}} \right) \leq c (|\nabla \phi|_2^4 + |\phi|_4^6 + 1).$$

Using repeatedly the inequality $\alpha\beta \leq \alpha^2/2 + \beta^2/2$ there exists a constant d_5 such that

$$\begin{aligned} I_5 &\leq M_2 \left(\frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{2} (|\phi|_4^4 + |Q|) + b \right) |\Delta \psi|_2 |\psi|_2 + c (|\nabla \phi|_2^4 + |\phi|_4^6 + 1) |\psi|_\infty |\psi|_2 \\ &\leq d_5 (|\nabla \phi|_2^4 + |\phi|_4^6 + 1) (|\Delta \psi|_2^2 + |\psi|_\infty^2 + |\psi|_2^2). \end{aligned}$$

Summing up the bounds for I_1, I_2, I_3, I_4, I_5 and taking into account the Poincaré inequality $|\psi|_\infty \leq C_{\mathcal{P}} |\nabla \psi|_2$ the result follows. \square

Proposition 1.9. *Under hypothesis of Proposition 1.6, there exists a constant $c > 0$, depending on the operator Ξ , such that for any $\phi + 1 \in \mathcal{C}_0^\infty(Q)$ it holds*

$$\sum_{j \in \mathbb{N}} \left\langle \frac{\delta^2 E(\phi)}{\delta \phi^2} \Xi \eta_j, \Xi \eta_j \right\rangle := \text{Tr} \left[\Xi^* \frac{\delta^2 E(\phi)}{\delta \phi^2} \Xi \right] \leq c (|\Delta \phi|_2^2 + |\nabla \phi|_2^4 + |\phi|_4^6 + |\phi|_6^6 + 1) \text{Tr} [\Xi^* \Delta^2 \Xi]$$

Proof : Let $(\eta_j)_{j \in \mathbb{N}}$ be an orthonormal basis of $L^2(Q)$. By Proposition 1.8, we have

$$\left\langle \frac{\delta^2 E(\phi)}{\delta \phi^2} \Xi \eta_j, \Xi \eta_j \right\rangle \leq c (|\Delta \phi|_2^2 + |\nabla \phi|_2^4 + |\phi|_4^6 + |\phi|_6^6 + 1) (|\Xi \eta_j|_2^2 + |\nabla \Xi \eta_j|_2^2 + |\Delta \Xi \eta_j|_2^2)$$

By taking the sum over j and using Proposition 1.6 we get the result.

2. Existence of a solution - preliminaries

2.1. Approximated equation and a priori estimates

Before proceeding to the proof, we need an approximation of equation (0.2).

Let us choose $\{e_j\}_{j \in \mathbb{N}^*} \in \mathcal{H}$ to be the eigenfunctions of the Stokes operator with homogeneous boundary conditions, such that $\{e_j\}_{j \in \mathbb{N}^*}$ forms an orthonormal basis for \mathcal{H} . Recall that $\{\eta_j\}_{j \in \mathbb{N}} \in L^2(Q)$ is the orthonormal basis in $L^2(Q)$ consisting of the eigenfunctions of the Laplacian Δ with homogeneous Dirichlet boundary conditions.

Next, set $S_n = \text{span}\{e_1, \dots, e_n\}$, $N_n = \text{span}\{\eta_0, \dots, \eta_n\}$. Finally, we denote by $P_n : \mathcal{H} \rightarrow \mathcal{H}$ the orthogonal projection of \mathcal{H} to S_n , and by $\pi_n : L^2(Q) \rightarrow L^2(Q)$ the orthogonal projection of $L^2(Q)$ into N_n .

We consider the equations

$$\begin{cases} d(w_n + \alpha^2 A w_n) = P_n \Sigma dW_t + \\ \quad \left(-\nu A(w_n + \alpha^2 A w_n) - P_n \tilde{B}(w_n, w_n + \alpha^2 A w_n) + P_n \left(\pi_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \right) \nabla \phi_n \right) \right) dt, & \text{in } [0, T] \times Q, \\ d\phi_n = \left(-\pi_n (w_n \cdot \nabla \phi_n) - \gamma \pi_n \left(\frac{\delta E(\phi_n)}{\delta \phi_n} \right) \right) dt + \pi_n \Xi dZ_t, & \text{in } [0, T] \times Q, \\ w_n(0) = P_n w_0, & \text{in } Q, \\ \phi_n(0) = \pi_n(\phi_0 + 1) - 1, & \text{in } Q. \end{cases} \quad (2.1)$$

Equation (2.1) is a system of ordinary stochastic differential equations with polynomial non-linear coefficients. Therefore, there exists a unique local strong solution (w_n, ϕ_n) defined up to a blow up random time $\tau(\omega)$. In order to show global existence and uniqueness of a solution for the approximated equations, we shall show a priori estimates. First we will use the Itô formula on the

process $X := t \mapsto (u_n(t), \phi_n(t))$ for $t \in [0, \tau]$ where $u_n := (I + \alpha^2 A)w_n$ which satisfies

$$\begin{aligned} dX &:= \begin{pmatrix} du_n \\ d\phi_n \end{pmatrix} = \begin{pmatrix} dX_1 \\ dX_2 \end{pmatrix} \\ &= \begin{pmatrix} \left(-\nu AX_1 - P_n \tilde{B}((I + \alpha^2 A)^{-1} X_1, X_1) + P_n \left(\pi_n \left(\frac{\delta E(X_2)}{\delta \phi} \right) \nabla X_2 \right) \right) dt + (P_n \Sigma) dW_t \\ \left(-\pi_n ((I + \alpha^2 A)^{-1} X_1 \cdot \nabla X_2) - \gamma \pi_n \left(\frac{\delta E(X_2)}{\delta \phi} \right) \right) dt + (\pi_n \Xi) dZ_t \end{pmatrix}. \end{aligned}$$

Using the function F defined on $H \times L^2(Q)$ such that

$$\begin{aligned} F : H \times L^2(Q) &\rightarrow \mathbb{R} \\ X &\mapsto \frac{1}{2} \langle (I + \alpha^2 A)^{-1} X_1, X_1 \rangle \end{aligned}$$

with $DF(X) := \begin{pmatrix} (I + \alpha^2 A)^{-1} X_1 \\ 0 \end{pmatrix}$ and $D^2F(X) := \begin{pmatrix} (I + \alpha^2 A)^{-1} & 0 \\ 0 & 0 \end{pmatrix}$, we obtain

$$\begin{aligned} dF(X) &= \left(-\nu \langle AX_1, (I + \alpha^2 A)^{-1} X_1 \rangle - \langle P_n \tilde{B}((I + \alpha^2 A)^{-1} X_1, X_1), (I + \alpha^2 A)^{-1} X_1 \rangle \right. \\ &\quad \left. + \left\langle P_n \left(\pi_n \left(\frac{\delta E(X_2)}{\delta \phi} \right) \nabla X_2 \right), (I + \alpha^2 A)^{-1} X_1 \right\rangle + \frac{1}{2} \text{Tr}[(P_n \Sigma)^* (I + \alpha^2 A)^{-1} (P_n \Sigma)] \right) dt \\ &\quad + \langle (I + \alpha^2 A)^{-1} X_1, (P_n \Sigma) dW_t \rangle. \end{aligned}$$

This equation rewrites

$$\begin{aligned} &\frac{1}{2} d \langle (I + \alpha^2 A)^{-1} u_n, u_n \rangle = \frac{1}{2} d \langle w_n, (I + \alpha^2 A) w_n \rangle = \frac{1}{2} d (|w_n|_2^2 + \alpha^2 |\nabla w_n|_2^2) \\ &= \left(-\nu (|\nabla w_n|_2^2 + \alpha^2 |Aw_n|_2^2) + \langle P_n \tilde{B}(w_n, w_n + \alpha^2 Aw_n), w_n \rangle \right. \\ &\quad \left. + \left\langle P_n \left(\pi_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \right) \nabla \phi_n \right), w_n \right\rangle + \frac{1}{2} \text{Tr}[(P_n \Sigma)^* (I + \alpha^2 A)^{-1} (P_n \Sigma)] \right) dt \\ &\quad + \langle w_n, (P_n \Sigma) dW_t \rangle. \end{aligned}$$

Using similar computation with function F given by the energy E we obtain

$$\begin{aligned} dE(\phi_n) &= \left(-\left\langle \pi_n (\nabla \phi_n \cdot w_n), \frac{\delta E(\phi_n)}{\delta \phi} \right\rangle - \gamma \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 + \frac{1}{2} \text{Tr} \left[(\pi_n \Xi)^* \frac{\delta^2 E(\phi_n)}{\delta \phi^2} (\pi_n \Xi) \right] \right) dt \\ &\quad + \left\langle \frac{\delta E(\phi_n)}{\delta \phi}, (\pi_n \Xi) dZ_t \right\rangle. \end{aligned} \tag{2.2}$$

Notice that by Proposition 1.1 and by the fact that $w_n \in S_n$ we have

$$\langle P_n \tilde{B}(w_n, w_n + \alpha^2 Aw_n), w_n \rangle = \langle \tilde{B}(w_n, w_n + \alpha^2 Aw_n), w_n \rangle = 0$$

and

$$\left\langle P_n \left(\pi_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \right) \nabla \phi_n \right), w_n \right\rangle = \left\langle \pi_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \right), \nabla \phi_n \cdot w_n \right\rangle = \left\langle \frac{\delta E(\phi_n)}{\delta \phi}, \pi_n (\nabla \phi_n \cdot w_n) \right\rangle$$

Then, by summing up (2.2) and (2.2) we find

$$\begin{aligned} d\left(\frac{1}{2}|w_n|_2^2 + \frac{\alpha^2}{2}|\nabla w_n|_2^2 + E(\phi_n)\right) &= \left(-\nu(|\nabla w_n|_2^2 + \alpha^2|Aw_n|_2^2) - \gamma\left|\frac{\delta E(\phi_n)}{\delta\phi}\right|_2^2\right. \\ &\quad \left.+ \frac{1}{2}\text{Tr}[(P_n\Sigma)^*(I + \alpha^2 A)^{-1}(P_n\Sigma)] + \frac{1}{2}\text{Tr}\left[(\pi_n\Xi)^*\frac{\delta^2 E(\phi_n)}{\delta\phi^2}(\pi_n\Xi)\right]\right)dt \\ &\quad + \langle w_n, (P_n\Sigma)dW_t \rangle + \left\langle \frac{\delta E(\phi_n)}{\delta\phi}, (\pi_n\Xi)dZ_t \right\rangle. \end{aligned} \quad (2.3)$$

Remark 2.1. Let us explain -in an other way- how the trace term $\text{Tr}[(P_n\Sigma)^*(I + \alpha^2 A)^{-1}(P_n\Sigma)]$ appears in formula (2.2). We can write

$$|w_n|_2^2 + \alpha^2|\nabla w_n|_2^2 = \langle w_n, w_n + \alpha^2 Aw_n \rangle = \langle (I + \alpha^2 A)^{-1}(I + \alpha^2 A)w_n, (I + \alpha^2 A)w_n \rangle.$$

Notice, also, that in (2.1) the equation involving w_n can be read as $d(I + \alpha^2 A)w_n = [\dots]$. In order to compute the trace term for the Itô formula, we have to derive with respect to $(I + \alpha^2 A)w_n$. Then, the second derivative of $|w_n|_2^2 + \alpha^2|\nabla w_n|_2^2$ with respect to $(I + \alpha^2 A)w_n$ can be written as

$$D^2(|w_n|_2^2 + \alpha^2|\nabla w_n|_2^2)(h, k) = 2((I + \alpha^2 A)^{-1}h, k), \quad h, k \in D(A).$$

This leads to the trace term $\text{Tr}[(P_n\Sigma)^*(I + \alpha^2 A)^{-1}(P_n\Sigma)]$.

2.2. Existence and uniqueness for the approximated equation

Theorem 2.2. Let $(w_0, \phi_0) \in D(A) \times L^2(Q)$ and assume that Hypothesis 0.1 holds. Then, for any $n \in \mathbb{N}$, $T > 0$ there exists a solution $(w_n, \phi_n) \in L^2([0, T]; D(A)) \times L^2([0, T]; L^2(Q))$ of problem (2.1). Moreover, for any $T > 0$, $k \in \mathbb{N}^*$ there exists a constant $c = c(k, T, \phi_0, w_0) > 0$ such that for any $n \in \mathbb{N}^*$

$$\begin{aligned} \sup_{0 \leq t \leq T} \mathbb{E} \left[\left(\frac{1}{2}|w_n(t)|_2^2 + \frac{\alpha^2}{2}|\nabla w_n(t)|_2^2 + E(\phi_n(t)) \right)^k \right] &\leq c \\ \mathbb{E} \left[\int_0^T \left(\frac{1}{2}|w_n|_2^2 + \frac{\alpha^2}{2}|\nabla w_n|_2^2 + E(\phi_n) \right)^{k-1} \left(\nu(|\nabla w_n|_2^2 + \alpha^2|Aw_n|_2^2) + \gamma \left| \frac{\delta E(\phi_n)}{\delta\phi} \right|_2^2 \right) ds \right] &\leq c \end{aligned}$$

Proof. Set

$$\mathcal{F}(t) = \mathcal{F}(t, w_n, \phi_n) = \frac{1}{2}|w_n|_2^2 + \frac{\alpha^2}{2}|\nabla w_n|_2^2 + E(\phi_n) \quad (2.4)$$

For any $N > 0$, $n \in \mathbb{N}^*$ we consider the stopping time

$$\tau_N^n = \inf\{t : \mathcal{F}(t, w_n, \phi_n) > N\}.$$

As pointed out previously, (2.1) is a system of ordinary differential equations with polynomial nonlinearities. Then, there exists a local solution (w_n, ϕ_n) up to a blow up time $\tau(\omega)$. Since the functions $w_n(t \wedge \tau_N^n)$, $\phi_n(t \wedge \tau_N^n)$ are bounded by N , we can apply the Itô formula in (2.3) to obtain

$$\begin{aligned} &\mathcal{F}^k(t \wedge \tau_N^n) + 2k \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \times \left(\nu(|\nabla w_n|_2^2 + \alpha^2|Aw_n|_2^2) + \gamma \left| \frac{\delta E(\phi_n)}{\delta\phi} \right|_2^2 \right) ds \\ &= 2k \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \times \left(\frac{1}{2}\text{Tr}[(P_n\Sigma)^*(I + \alpha^2 A)^{-1}(P_n\Sigma)] + \frac{1}{2}\text{Tr}\left[(\pi_n\Xi)^*\frac{\delta^2 E(\phi_n)}{\delta\phi^2}(\pi_n\Xi)\right] \right) ds \\ &\quad + 2k \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \langle w_n, (P_n\Sigma)dW(s) \rangle + 2k \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \left\langle \frac{\delta E(\phi_n)}{\delta\phi}, (\pi_n\Xi)dW'(s) \right\rangle \\ &\quad + k(k-1) \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-2} \left(|(P_n\Sigma)^*w_n|_2^2 + \left| (\pi_n\Xi)^* \frac{\delta E(\phi_n)}{\delta\phi} \right|_2^2 \right) ds \\ &= I_1 + I_2 + M_t \end{aligned}$$

where M_t is the martingale term. Let us estimate I_1 . By Proposition 1.9 there exists $c_1 > 0$ such that

$$\frac{1}{2} \text{Tr} \left[(\pi_n \Xi)^* \frac{\delta^2 E(\phi_n)}{\delta \phi^2} (\pi_n \Xi) \right] \leq c_1 (|\Delta \phi_n|_2^2 + |\nabla \phi_n|_2^4 + |\phi_n|_4^6 + |\phi_n|_6^6 + 1) \text{Tr}[\Xi^* \Delta^2 \Xi].$$

By (1.3) and elementary inequalities there exists a positive constant c_2 such that

$$(|\Delta \phi_n|_2^2 + |\nabla \phi_n|_2^4 + |\phi_n|_4^6 + |\phi_n|_6^6 + 1) \leq c_2(1 + E(\phi_n)) \leq c_2(1 + \mathcal{F}).$$

Taking into account that $\text{Tr}[\Xi^* \Delta^2 \Xi]$ and $\text{Tr}[\Sigma^*(I + \alpha^2 A)^{-1} \Sigma]$ are bounded, there exists $c_3 > 0$ that

$$I_1 \leq c_3 \int_0^{t \wedge \tau_N^n} (\mathcal{F}^{k-1} + \mathcal{F}^k) ds$$

Let us estimate I_2 . By (1.9) there exists $c_4 > 0$ such that

$$\left| (\pi_n \Xi)^* \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 \leq c_4 (1 + |\phi_n|_4^8 + |\phi_n|_6^6 + |\nabla \phi_n|_2^4 + |\Delta \phi_n|_2^2)^2.$$

Using (1.3), the quantity on the right hand side is bounded by $c(1 + E(\phi_n)^2)$, for a suitable $c > 0$ independent of ϕ . By elementary inequalities and the fact that the operators $P_n \Sigma^*$ are uniformly bounded with respect to n , we deduce that for there exists $c_5 > 0$, independent of n, ϕ_n, w_n such that

$$I_2 \leq c_5 \int_0^{t \wedge \tau_N^n} (\mathcal{F}^{k-2} + \mathcal{F}^k) ds.$$

Finally,

$$\begin{aligned} \mathcal{F}^k(t \wedge \tau_N^n) &\leq c_3 \int_0^{t \wedge \tau_N^n} (\mathcal{F}^{k-1} + \mathcal{F}^k) ds + c_5 \int_0^{t \wedge \tau_N^n} (\mathcal{F}^{k-2} + \mathcal{F}^k) ds \\ &\quad + 2k \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \langle w_n, (P_n \Sigma) dW(s) \rangle + 2k \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \left\langle \frac{\delta E(\phi_n)}{\delta \phi}, (\pi_n \Xi) dW'(s) \right\rangle \end{aligned} \quad (2.5)$$

Before taking expectation, we need to verify that the martingale terms are integrable. Notice that since the operator Σ is bounded there exists $c_6 > 0$ such that

$$\mathcal{F}^{k-1} |(P_n \Sigma)^* w_n|_2 \leq c_6(1 + \mathcal{F}^k).$$

Then, since $\mathcal{F}^k(t \wedge \tau_n) \leq N^k$, we can take expectation to obtain

$$2k \mathbb{E} \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \langle w_n, (P_n \Sigma) dW(s) \rangle = 0.$$

Similarly, for the second term we can use estimate (1.9) and obtain, for some $c_6 > 0$

$$\mathcal{F}^{k-1} \left| (\pi_n \Xi)^* \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 \leq c_6 \mathcal{F}^{k-1} (1 + |\phi_n|_4^8 + |\phi_n|_6^6 + |\nabla \phi_n|_2^4 + |\Delta \phi_n|_2^2)$$

As we pointed out previously, by (1.3) there exists $c > 0$ such that

$$(1 + |\phi|_4^8 + |\phi_n|_6^6 + |\nabla \phi_n|_2^4 + |\Delta \phi_n|_2^2) \leq c(1 + E(\phi_n)) \leq c(1 + \mathcal{F}).$$

Then, there exists $c_7 > 0$ such that

$$\mathcal{F}^{k-1} \left| (\pi_n \Xi)^* \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 \leq c_7 \mathcal{F}^{k-1} (1 + \mathcal{F}).$$

This implies that we can take expectation to obtain

$$2k\mathbb{E} \int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \left\langle \frac{\delta E(\phi_n)}{\delta \phi}, (\pi_n \Xi) dW'(s) \right\rangle = 0.$$

Finally, by taking expectation in (2.5) we get

$$\begin{aligned} \mathbb{E}[\mathcal{F}^k(t \wedge \tau_N^n)] + 2k\mathbb{E} \left[\int_0^{t \wedge \tau_N^n} \mathcal{F}^{k-1} \times \left(\nu(|\nabla w_n|_2^2 + \alpha^2 |Aw_n|_2^2) + \gamma \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 \right) ds \right] \\ \leq c_3\mathbb{E} \left[\int_0^{t \wedge \tau_N^n} (\mathcal{F}^{k-1} + \mathcal{F}^k) ds \right] + c_5\mathbb{E} \left[\int_0^{t \wedge \tau_N^n} (\mathcal{F}^{k-2} + \mathcal{F}^k) ds \right]. \end{aligned} \quad (2.6)$$

Clearly, there exists a constant $c > 0$ such that $\mathcal{F}^{k-1} \leq c(1 + \mathcal{F}^k)$ and $\mathcal{F}^{k-2} \leq c(1 + \mathcal{F}^k)$. Then, there exists $c_7 > 0$, depending only on k, ϕ_0, w_0 , such that the right-hand side of (2.6) is bounded by

$$c_7\mathbb{E} \left[\int_0^{t \wedge \tau_N^n} (1 + \mathcal{F}^k) ds \right] \leq c_7\mathbb{E} \left[\int_0^t (1 + \mathcal{F}^k(s \wedge \tau_N^n)) ds \right].$$

Using Gronwall lemma, we find that there exists a constant $c_8 > 0$ depending on k, T, ϕ_0, w_0 , such that

$$\sup_{t \in [0, T]} \mathbb{E}[\mathcal{F}^k(t \wedge \tau_N^n)] + 2k \int_0^{T \wedge \tau_N^n} \mathcal{F}^{k-1} \times \left(\nu(|\nabla w_n|_2^2 + \alpha^2 |Aw_n|_2^2) + \gamma \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 \right) ds \leq c.$$

Letting $N \rightarrow \infty$ we conclude the proof of Theorem 2.2. \square

Theorem 2.3. *Let $(w_0, \phi_0) \in D(A) \times L^2(Q)$ and assume that Hypothesis 0.1 holds. Then for any $T > 0, k \in \mathbb{N}$ there exists $c = c(k, T, w_0, \phi_0) > 0$ such that*

$$\mathbb{E} \left[\sup_{t \in [0, T]} \left(\frac{1}{2} |w_n|_2^2 + \frac{\alpha^2}{2} |\nabla w_n|_2^2 + E(\phi_n) \right)^k \right] \leq c.$$

Proof. As done for the previous Theorem, let us set \mathcal{F} as in (2.4). By Theorem 2.2 the solution (w_n, ϕ_n) is global and all moments of \mathcal{F} have finite expectation. Then by Itô formula (2.3) we get

$$\begin{aligned} \mathcal{F}^k(t) &= 2k \int_0^t \mathcal{F}^{k-1} \times \left(\left(-\nu(|\nabla w_n|_2^2 + \alpha^2 |Aw_n|_2^2) - \gamma \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 \right. \right. \\ &\quad \left. \left. + \frac{1}{2} \text{Tr}[(P_n \Sigma)^*(I + \alpha^2 A)^{-1}(P_n \Sigma)] + \frac{1}{2} \text{Tr} \left[(\pi_n \Xi)^* \frac{\delta^2 E(\phi_n)}{\delta \phi^2} (\pi_n \Xi) \right] \right) ds \right) \\ &\quad + k(k-1) \int_0^t \mathcal{F}^{k-2} \left(|(P_n \Sigma)^* w_n|_2^2 + \left| (\pi_n \Xi)^* \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 \right) ds \\ &\quad + 2k \int_0^t \mathcal{F}^{k-1} \langle w_n, (P_n \Sigma) dW(s) \rangle + 2k \int_0^t \mathcal{F}^{k-1} \left\langle \frac{\delta E(\phi_n)}{\delta \phi}, (\pi_n \Xi) dW'(s) \right\rangle \\ &= I_1 + I_2 + M_t \end{aligned}$$

Where I_1, I_2 are the integrals containing \mathcal{F}^{k-1} and \mathcal{F}^{k-2} respectively, and M_t is the martingale term. As we done for Theorem 2.2, I_1, I_2 are uniformly bounded in t by

$$I_1 + I_2 \leq c \int_0^T (1 + \mathcal{F}^k) ds,$$

where $c > 0$ is a suitable constant depending only by k, T . For the martingale part, we can use Burkholder-Davis-Gundy inequality to get for some constant $c_1, c_2 > 0$

$$\begin{aligned} \mathbb{E} \left(\sup_{t \in [0, T]} \left| \int_0^t \mathcal{F}^{k-1} \langle w_n, (P_n \Sigma) dW(s) \rangle \right| \right) &\leq c_1 \mathbb{E} \left(\int_0^T \mathcal{F}^{2(k-1)} |(P_n \Sigma)^* w_n|_2^2 ds \right)^{\frac{1}{2}} \\ &\leq c_2 \mathbb{E} \left(\int_0^T \mathcal{F}^{2k} ds \right)^{\frac{1}{2}} < \infty \end{aligned}$$

The last term is bounded thanks to Theorem 2.2. Again, by Burkholder-Davis-Gundy's inequality there exists $c_3 > 0$ such that

$$\mathbb{E} \left(\sup_{t \in [0, T]} \left| \int_0^t \mathcal{F}^{k-1} \left\langle \frac{\delta E(\phi_n)}{\delta \phi}, (\pi_n \Xi) dW'(s) \right\rangle \right| \right) \leq c_3 \mathbb{E} \left(\int_0^T \mathcal{F}^{2(k-1)} \left| (\pi_n \Xi)^* \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 ds \right)^{\frac{1}{2}}$$

By estimate (1.9) and (1.3), there exists $c_4 > 0$ such that the right-hand side is bounded by

$$c_4 \mathbb{E} \left(\int_0^T \mathcal{F}^{2(k-1)} (1 + \mathcal{F}^2) ds \right)^{\frac{1}{2}}.$$

Then, by Theorem 2.2 this integral is finite. This completes the proof. \square

2.3. Compactness argument - convergence to a solution

Let X be a Banach space with norm $\|\cdot\|_X$. For $p \geq 1$, $\theta \in]0, 1[$ we denote by $W^{\theta, p}([0, T]; X)$ the classical Sobolev space of all functions $f \in L^p([0, T]; X)$ such that

$$\int_0^T \int_0^T \frac{|f(t) - f(s)|_X^p}{|t - s|^{1+\theta p}} ds dt < \infty,$$

endowed with the norm

$$\|f\|_{W^{\theta, p}([0, T]; X)} = \left(\|f\|_{L^p([0, T]; X)}^p + \int_0^T \int_0^T \frac{\|f(t) - f(s)\|_X^p}{|t - s|^{1+\theta p}} ds dt \right)^{\frac{1}{p}}.$$

The proof of the following lemma is left to the reader.

Lemma 2.4. *Let X a Banach space. For any $\theta \in]0, 1/2[$ $p \geq 1$ there exists $c = c(\theta, p)$ such that for any $f \in L^2([0, T]; X)$ it holds*

$$\left\| \int_0^\cdot f(\tau) d\tau \right\|_{W^{\theta, p}([0, T]; X)} \leq c(\theta, p) \|f\|_{L^2([0, T]; X)}.$$

Proposition 2.5. *For any $T > 0$, $\theta \in]0, 1/2[$, $p \geq 1$ there exists $c = c(T, \theta, p) > 0$ such that for any $n \in \mathbb{N}$*

$$\mathbb{E} \left[\|w_n + \alpha^2 A w_n\|_{W^{\theta, p}([0, T]; D(A)')}^2 \right] \leq c.$$

Proof. For any $n \in \mathbb{N}$, $\xi \in D(A)$ we have

$$\begin{aligned} \langle w_n(t) + \alpha^2 A w_n(t), \xi \rangle_{D(A)', D(A)} &= -\nu \int_0^t \langle w_n(\tau) + \alpha^2 A w_n(\tau), A \xi \rangle d\tau \\ &\quad - \int_0^t \langle P_n \tilde{B}(w_n, w_n + \alpha^2 A w_n)(\tau), \xi \rangle d\tau \\ &\quad + \int_0^t \left\langle P_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \nabla \phi_n \right), \xi \right\rangle dt \\ &\quad + \langle (P_n \Sigma) W(t), \xi \rangle \\ &= J_1(t) + J_2(t) + J_3(t) + J_4(t). \end{aligned}$$

We proceed as for Proposition 2.5 by estimating each term. For J_1 we have, using Lemma 2.4 and Theorem 2.2 (with $k = 1$), that there exists $c_1 > 0$ such that

$$\mathbb{E} \left[\|J_1(\cdot)\|_{W^{\theta,p}([0,T];\mathbb{R})}^2 \right] \leq c(\theta,p) \mathbb{E} \left[\int_0^T (|w_n(\tau)|_2 + \alpha^2 |Aw_n(\tau)|_2^2) d\tau \right] |A\xi|_2^2 \leq c_1 \|\xi\|_{D(A)}^2$$

In order to estimate J_2 , observe that by (ii) of Proposition 1.1 and Young's inequality, we have

$$\begin{aligned} \left\langle P_n \tilde{B}(w_n, w_n + \alpha^2 Aw_n), \xi \right\rangle_{D(A)', D(A)} &= \left\langle \tilde{B}(w_n, w_n + \alpha^2 Aw_n), P_n \xi \right\rangle_{D(A)', D(A)} \\ &\leq c |w_n|_V (|w_n|_2 + \alpha^2 |Aw_n|_2) \|\xi\|_{D(A)} \end{aligned}$$

By Lemma 2.4 and the bound given by Theorem 2.2, we deduce that there exists $c_2 > 0$ such that

$$\mathbb{E} \left[\|J_2(\cdot)\|_{W^{\theta,p}([0,T];\mathbb{R})}^2 \right] \leq c \mathbb{E} \left[\int_0^T |w_n|_V^2 (|w_n|_2 + \alpha^2 |Aw_n|_2)^2 d\tau \right] \|\xi\|_{D(A)}^2 \leq c_2 \|\xi\|_{D(A)}^2.$$

In order to estimate J_3 , let us observe that we have, by Hölder's and Sobolev's inequalities (which works both in dimensions 2 and 3)

$$\begin{aligned} \left| \left\langle P_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \nabla \phi_n \right), \xi \right\rangle_{D(A)', D(A)} \right| &\leq \left| P_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \right) \right|_2 |\nabla \phi_n|_3 |\xi|_6 \\ &\leq \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 \|\nabla \phi_n\|_{H^1(Q)} \|\xi\|_{H^1(Q)} \\ &\leq c \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 \|\phi_n\|_{H^2} \|\xi\|_{D(A)} \\ &\leq c \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 (1 + E(\phi_n)) \|\xi\|_{D(A)}. \end{aligned}$$

In the last inequality we used (1.3). Then, by Lemma 2.4 and the estimates in Theorem 2.2 (with $k = 3$), we deduce that there exists $c_3 > 0$ such that

$$\begin{aligned} \mathbb{E} \left[\|J_3(\cdot)\|_{W^{\theta,p}([0,T];\mathbb{R})}^2 \right] &\leq \mathbb{E} \left[\int_0^T \left| \left\langle P_n \left(\frac{\delta E(\phi_n)}{\delta \phi} \nabla \phi_n \right), \xi \right\rangle_{D(A)', D(A)} \right|^2 d\tau \right] \\ &\leq c \mathbb{E} \left[\int_0^T \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 (1 + E(\phi_n))^2 d\tau \right] \|\xi\|_{D(A)}^2 \\ &\leq c_3 \|\xi\|_{D(A)}^2. \end{aligned}$$

In order to estimate J_4 , let us observe that by the gaussianity of $(P_n \Sigma)(W(t) - W(s))$ there exists a constant $c = c(p)$ such that $\mathbb{E}[(P_n \Sigma)(W(t) - W(s))_2^p] \leq c(\text{Tr}[\Sigma^* \Sigma])^{\frac{p}{2}} |t - s|^{\frac{p}{2}}$. Then, provided $\theta < 1/2$, there exists a constant $c_4 > 0$ such that

$$\mathbb{E} \left[\int_0^T \int_0^T \frac{|(P_n \Sigma)(W(t) - W(s))_2|^p}{|t - s|^{1+\theta p}} ds dt \right] \leq c(\text{Tr}[\Sigma^* \Sigma])^{\frac{p}{2}} \int_0^T \int_0^T \frac{|t - s|^p}{|t - s|^{1+\theta p}} ds dt \leq c_4. \quad (2.7)$$

Finally, the result follows by taking into account the estimates obtained for J_1, J_2, J_3, J_4 . \square

Proposition 2.6. *For any $T > 0$, $\theta \in]0, 1/2[$, $p \geq 1$ there exists $c = c(T, \theta, p) > 0$ such that for any $n \in \mathbb{N}$*

$$\mathbb{E} \left[\|\phi_n\|_{W^{\theta,p}([0,T];L^2(Q))}^2 \right] \leq c.$$

Proof. For any n we have

$$\phi_n(t) = \int_0^t \pi_n(w_n \nabla \phi_n) d\tau - \gamma \int_0^t \pi_n \left(\frac{\delta E}{\delta \phi}(\phi_n(\tau)) \right) d\tau + (\pi_n \Xi) W'(t) = K_1(t) + K_2(t) + K_3(t).$$

We proceed by estimating each term. For K_1 we have, using elementary inequalities

$$\int_0^T |\pi_n(w_n \nabla \phi_n)|_2^2 d\tau \leq T \sup_{0 \leq t \leq T} (|w_n|_\infty^2 |\nabla \phi_n|_2^2) \leq \frac{1}{2} T \sup_{0 \leq t \leq T} (|w_n|_\infty^4 + |\nabla \phi_n|_2^4).$$

Then by Lemma 2.4, Theorem 2.3 and Poincaré's inequality $|w|_\infty \leq C_P |w|_V$ we deduce that there exists $c_1 > 0$, independent of n such that

$$\begin{aligned} \mathbb{E} \left[\|K_1(\cdot)\|_{W^{\theta,p}([0,T];L^2(Q))}^2 \right] &\leq c(\theta,p) \mathbb{E} \left[\int_0^T |\pi_n(w_n \nabla \phi_n)|_2^2 d\tau \right] \\ &\leq \frac{1}{2} c(\theta,p) T \mathbb{E} \left[\sup_{0 \leq t \leq T} (|w_n|_\infty^4 + |\nabla \phi_n|_2^4) \right] \\ &\leq \frac{1}{2} c(\theta,p) T \mathbb{E} \left[\sup_{0 \leq t \leq T} (C_P^4 |w_n|_V^4 + c(1 + E(\phi_n))) \right] \leq c_1. \end{aligned}$$

In the last inequality we used (1.3).

For K_2 we have, by Lemma 2.4 and Theorem 2.2, that for some $c_2 > 0$, independent of n , it holds

$$\mathbb{E} \left[\|K_2(\cdot)\|_{W^{\theta,p}([0,T];L^2(Q))}^2 \right] \leq \mathbb{E} \left[\left\| \frac{\delta E}{\delta \phi}(\phi_n(\tau)) \right\|_{L^2([0,T];L^2(Q))}^2 \right] < c_2$$

The last term is treated as done in (2.7). Then, provided $\theta < 1/2$, there exists $c_3 > 0$ such that Then,

$$\mathbb{E} \left[\int_0^T \int_0^T \frac{|(\pi_n \Xi)(W'(t) - W'(s))|_2^p}{|t-s|^{1+\theta p}} ds dt \right] \leq c(\text{Tr}[\Xi^* \Xi])^{\frac{p}{2}} \int_0^T \int_0^T \frac{|t-s|^p}{|t-s|^{1+\theta p}} ds dt \leq c_3$$

provided $\theta < 1/2$. Taking into account the estimates on K_1 , K_2 , K_3 we obtain the result. \square

In what follows, we denote by $L_w^2([0,T];D(A)')$ the space $L^2([0,T],D(A)')$ endowed with the weak L^2 topology.

Lemma 2.7 (Tightness). *For $(w_0, \phi_0) \in D(A) \times L^2(Q)$ with $\phi_0 = -1$ on ∂Q , $T > 0$, $n \in \mathbb{N}$, let (w_n, ϕ_n) the solution of (2.1) in $[0, T]$. Then, for any $p > 2$, $\rho > 0$, the laws of $w_n, n \in \mathbb{N}$ are tight in*

$$\mathcal{C}([0, T]; D(A^{-\rho})) \cap L^p([0, T]; V) \cap L_w^2([0, T]; D(A)')$$

Moreover, for any $\sigma > 0$, the laws of $\phi_n, n \in \mathbb{N}$ are tight in

$$\mathcal{C}([0, T]; H^{-\sigma}(Q)) \cap L^p([0, T]; H^2(Q)) \cap L_w^2([0, T]; (H^4(Q))').$$

Proof. The classical interpolation inequality

$$\|w\|_{H^{1+\rho}} \leq \|w\|_{H^1}^{1-\rho} \|w\|_{H^2}^\rho, \quad \rho \in [0, 1]$$

implies

$$\|w\|_{H^{1+\frac{2}{p}}}^p \leq \|w\|_{H^1}^{p-2} \|w\|_{H^2}^2, \quad p \in [2, \infty[.$$

Then, by Theorem 2.2 and Proposition 2.5 implies that $(w_n)_n$ is bounded in

$$L^p \left(\Omega; L^p([0, T]; H^{1+\frac{2}{p}}) \right) \cap L^2 \left(\Omega; L^2([0, T]; D(A)) \right) \cap L^2 \left(\Omega; W^{\theta,p}([0, T]; H) \right)$$

for any $p \in [2, \infty[$ and $\theta < 1/2$ such that $\theta p > 1$. Taking into account Theorem [20, Theorem 2.1 and Theorem 2.2], for any $p \in [2, \infty[$ and $\theta < 1/2$ such that $\theta p > 1$ the embeddings

$$W^{\theta,p}([0, T]; H) \hookrightarrow \mathcal{C}([0, T]; D(A^{-\rho})), \quad \rho > 0$$

$$L^p([0, T]; H^{1+\frac{2}{p}}) \cap W^{\theta,p}([0, T]; H) \hookrightarrow L^p([0, T]; V)$$

are compact. Moreover, we have that $L^2([0, T]; D(A))$ is compactly embedded in the complete metrizable space $L_w^2([0, T]; D(A)')$. Then, the result follows by Prokhorov's theorem.

In order to show the tightness of the laws of ϕ_n , notice that by (1.8) there exists $c > 0$, independent of n , such that

$$\mathbb{E} \left[\int_0^T |\Delta^2 \phi_n|_2^2 dt \right] \leq c \mathbb{E} \left[\int_0^T \left(\left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2^2 + 1 + (E(\phi))^2 \right) dt \right].$$

Taking into account Theorem 2.2, this implies that the sequence $(\phi_n)_n$ is uniformly bounded in $L^2(\Omega; L^2([0, T]; H^4(Q)))$ and then the laws of $\phi_n, n \in \mathbb{N}$ are tight in the complete metrizable space $L_w^2([0, T]; (H^4(Q))')$. By the interpolation formula $\|\phi\|_{H^{2+2\rho}} \leq c_\rho \|\phi\|_{H^2}^{1-\rho} \|\phi\|_{H^4}^\rho$ we deduce that for some $c > 0$

$$\|\phi\|_{H^{2+\frac{4}{p}}}^p \leq c_\rho \|\phi\|_{H^2}^{p-2} \|\phi\|_{H^4}^2 \leq c \|\phi\|_{H^2}^{p-2} (\|\phi\|_{H^2}^2 + |\Delta^2 \phi|_2^2), \quad p \geq 2.$$

Moreover, by (1.3), (1.8), we get that for some $c > 0, p' \geq 2$ it holds

$$\|\phi_n\|_{H^{2+\frac{4}{p}}}^p \leq c(1 + E(\phi_n)^{p'}) \left(1 + \left| \frac{\delta E}{\delta \phi}(\phi_n) \right|_2^2 \right).$$

Then, thanks to Theorem 2.2, we have that for any $p \geq 2$ the sequence $(\phi_n)_n$ is uniformly bounded in $L^p(\Omega; L^p([0, T]; H^{2+\frac{4}{p}}))$, $p \geq 2$.

Consequently, by Proposition 2.6 the sequence $(\phi_n)_n$ is bounded in

$$L^p(\Omega; L^p([0, T]; H^{2+\frac{4}{p}})) \cap L^2(\Omega; W^{\theta,p}([0, T]; L^2(Q))) \cap L^2(\Omega; L^2([0, T]; H^4(Q))) \quad \theta < \frac{1}{2}, p < \infty,$$

endowed with the conditions $\phi_n = -1$ on ∂Q , $\Delta \phi_n = 0$ on ∂Q . Since by [20, Theorem 2.1 and Theorem 2.2]) we have that the embeddings

$$\begin{aligned} W^{\theta,p}([0, T]; L^2(Q)) &\hookrightarrow \mathcal{C}([0, T]; H^{-\sigma}(Q)), \quad \sigma > 0, \theta p > 1 \\ L^p([0, T]; H^{2+\frac{4}{p}}(Q)) \cap W^{\theta,p}([0, T]; L^2(Q)) &\hookrightarrow L^p([0, T]; H^2(Q)), \end{aligned}$$

are compact, the result follows by Prokhorov's Theorem. \square

Theorem 2.8. *Let $(w_0, \phi_0) \in D(A) \times L^2(Q)$ with $\phi_0 = -1$ on ∂Q . Then, there exists a probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, two cylindrical Wiener processes $\tilde{W}(t), \tilde{Z}(t)$ defined on $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, stochastic processes*

$$\begin{aligned} w &\in \mathcal{C}([0, T]; D(A^{-\rho})) \cap L^p([0, T]; V) \cap L^2([0, T]; D(A)), \quad \rho > 0, \\ \phi &\in \mathcal{C}([0, T]; H^{-\sigma}(Q)) \cap L^p([0, T]; H^2(Q)) \cap L^2([0, T]; H^4(Q)), \quad \sigma > 0, \\ \zeta &\in L^2([0, T]; L^2(Q)) \end{aligned}$$

and subsequences (for simplicity they are not relabeled) such that for any $p < \infty$ and $\tilde{\mathbb{P}}$ -a.s. the solution (w_n, ϕ_n) of problem (2.1) with $\tilde{W}(t)$ and $\tilde{Z}(t)$ instead of $W(t), Z(t)$ satisfies

- (i) $w_n \rightarrow w$ strongly in $\mathcal{C}([0, T]; D(A^{-\rho}))$, $\rho > 0$
- (ii) $w_n \rightarrow w$ strongly in $L^p([0, T]; V)$, $p \in [1, \infty[$
- (iii) $w_n \rightarrow w$ weakly in $L^2([0, T]; D(A))$
- (iv) $\phi_n \rightarrow \phi$ strongly in $\mathcal{C}([0, T]; H^{-\sigma}(Q))$, $\sigma > 0$
- (v) $\phi_n \rightarrow \phi$ strongly in $L^p([0, T]; H^2)$, $p \in [1, \infty[$
- (vi) $\Delta^2 \phi_n \rightarrow \Delta^2 \phi$ weakly in $L^2([0, T]; L^2(Q))$
- (vii) $\frac{\delta E(\phi_n)}{\delta \phi} \rightarrow \zeta$ weakly in $L^2([0, T]; L^2(Q))$
- (viii) $f(\phi_n) \rightarrow f(\phi)$ strongly in $L^2([0, T]; L^2(Q))$

Proof. Taking into account Lemma 2.7, by Skorohod's representation theorem and by a diagonal extraction argument, there exists a probability space $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, two cylindrical Wiener processes $\tilde{W}(t), \tilde{Z}(t)$ defined on $(\tilde{\Omega}, \tilde{\mathcal{F}}, \tilde{\mathbb{P}})$, two stochastic processes w, ϕ such that the convergence conditions in (i)–(vi) hold.

(vii). By Theorem 2.2, the sequence $\frac{\delta E(\phi_n)}{\delta \phi}$ are bounded in $L^2(\Omega; L^2([0, T]; L^2(Q)))$. Then, by arguing as for the previous point, the result follows by Prokhorov theorem and Skorohod theorem.

(viii) By the expression of $f(\phi_n)$ it is sufficient to show that that \mathbb{P} -almost surely $\Delta \phi_n \rightarrow \Delta \phi$ and $\phi_n^3 \rightarrow \phi^3$ strongly in $L^p([0, T]; L^2(Q))$. Indeed, the two limits follows by (v) and by standard Sobolev embedding results. \square

3. Proof of Theorem 0.3

3.1. Existence

By Theorem 2.8 we know that there exist subsequences $(w_n)_n, (\phi_n)_n$ converging $\tilde{\mathbb{P}}$ -a.s. to processes $(w, \phi) \in L^2(\Omega; L^2([0, T]; D(A))) \times L^2(\Omega; L^2([0, T]; H^2(Q)))$.

The rest of the proof will be splitted in several lemma : in Lemma 3.1, we will show that the processes (w, ϕ) satisfied (0.4). Then we will show that w, ϕ fulfill the definition 1.2 of a solution for the abstract problem.

Lemma 3.1. *Under hypothesis of Theorem 0.3, we have that (0.4) hold.*

Proof. Let us show the first bound of (0.4). Let us notice that by the definition of the norm in $D(A^\rho)$ it holds $\|w\|_{D(A^{-\rho})} \leq \|w\|_H$, for all $\rho > 0$. By Theorem 2.8,

$$\sup_{t \in [0, T]} \|w(t)\|_{D(A^{-\rho})} = \lim_{n \rightarrow \infty} \left(\sup_{t \in [0, T]} \|w_n(t)\|_{D(A^{-\rho})} \right) \leq \liminf_{n \rightarrow \infty} \left(\sup_{t \in [0, T]} \|w_n(t)\|_H \right)$$

By Fatou's lemma and Theorem 2.3 we deduce that for any $k > 0$ there exists $c > 0$ depending on k, T, w_0, ϕ_0 such that

$$\tilde{\mathbb{E}} \left[\sup_{t \in [0, T]} \|w\|_H^k \right] \leq \liminf_{n \rightarrow \infty} \tilde{\mathbb{E}} \left[\sup_{t \in [0, T]} \|w_n\|_H^k \right] \leq c.$$

With a similar argument it can be shown that for any $k > 0$ there exists $c > 0$ depending on k, T, w_0, ϕ_0 such that

$$\tilde{\mathbb{E}} \left[\sup_{t \in [0, T]} |\phi|_2^k \right] \leq \liminf_{n \rightarrow \infty} \tilde{\mathbb{E}} \left[\sup_{t \in [0, T]} |\phi_n|_2^k \right] \leq c$$

which implies that the first bound in (0.4) holds. Let us show the second bound. Notice that by Theorem 2.8 we have, $\tilde{\mathbb{P}}$ -a.s., that the limit $(\|w_n(t)\|_V + \|\phi_n(t)\|_{H^2}) \wedge M \rightarrow (\|w(t)\|_V + \|\phi(t)\|_{H^2}) \wedge M$ holds in $L^p([0, T])$, for all $p \geq 1$ and $M > 0$. Then, by Lemma 3.2 we have that the limit

$$\lim_{n \rightarrow \infty} ((\|w_n(t)\|_V + \|\phi_n(t)\|_{H^2})^p \wedge M) \frac{\delta E(\phi_n(t))}{\delta \phi} = ((\|w(t)\|_V + \|\phi(t)\|_{H^2})^p \wedge M) \frac{\delta E(\phi(t))}{\delta \phi}$$

holds weakly in $L^2([0, T] \times Q)$, for any $M > 0$. Then, for any $M > 0$,

$$\begin{aligned} & \int_0^T \left((\|w(t)\|_V + \|\phi(t)\|_{H^2})^{2p} \wedge M^2 \right) \left| \frac{\delta E(\phi(t))}{\delta \phi} \right|_2^2 dt \\ & \leq \liminf_{n \rightarrow \infty} \int_0^T \left((\|w_n(t)\|_V + \|\phi_n(t)\|_{H^2})^{2p} \wedge M^2 \right) \left| \frac{\delta E(\phi_n(t))}{\delta \phi} \right|_2^2 dt \end{aligned}$$

Letting $M \rightarrow \infty$, by monotone convergence we obtain

$$\int_0^T (\|w(t)\|_V + \|\phi(t)\|_{H^2})^{2p} \left| \frac{\delta E(\phi(t))}{\delta \phi} \right|_2^2 dt \leq \liminf_{n \rightarrow \infty} \int_0^T (\|w_n(t)\|_V + \|\phi_n(t)\|_{H^2})^{2p} \left| \frac{\delta E(\phi_n(t))}{\delta \phi} \right|_2^2 dt$$

Finally, by Fatou's Lemma we get

$$\begin{aligned} \tilde{\mathbb{E}} \left[\int_0^T (|w(t)|_V + |\phi(t)|_{H^2})^{2p} \left| \frac{\delta E(\phi(t))}{\delta \phi} \right|_2^2 dt \right] \\ \leq \liminf_{n \rightarrow \infty} \tilde{\mathbb{E}} \left[\int_0^T (\|w_n(t)\|_V + \|\phi_n(t)\|_{H^2})^{2p} \left| \frac{\delta E(\phi_n(t))}{\delta \phi} \right|_2^2 dt \right] \leq c \end{aligned}$$

where $c > 0$ is given by Theorem 2.2. By similar arguments we can show that there exists $c > 0$ such that

$$\tilde{\mathbb{E}} \left[\int_0^T (\|w(t)\|_V + \|\phi(t)\|_{H^2})^{2p} (|\nabla w|_2^2 + \alpha^2 |Aw|_2^2) dt \right] \leq c.$$

To conclude the proof, it is sufficient to notice that thanks to (1.4) there exists $c > 0$ such that $E(\phi) \leq c(1 + |\phi(t)|_{H^2}^8)$. \square

Next, we prove the convergence of nonlinear term.

Lemma 3.2. *We have, $\tilde{\mathbb{P}}$ -a.s.*

$$\lim_{n \rightarrow \infty} \frac{\delta E(\phi_n(t))}{\delta \phi} = \frac{\delta E(\phi(t))}{\delta \phi} \quad \text{weakly in } L^2([0, T]; L^2(Q))$$

and

$$\lim_{n \rightarrow \infty} P_n \left(\frac{\delta E(\phi_n(t))}{\delta \phi} \right) = \frac{\delta E(\phi(t))}{\delta \phi} \quad \text{weakly in } L^2([0, T]; L^2(Q)).$$

Proof. Let us prove the first limit. By (vii) of Theorem 2.8 we have to show that $\zeta = \frac{\delta E(\phi(t))}{\delta \phi}$. Let $g \in C_0^\infty([0, T] \times Q; \mathbb{R})$. We shall show that

$$\lim_{n \rightarrow \infty} \int_0^T \left\langle \frac{\delta E(\phi_n(t))}{\delta \phi}, g(t) \right\rangle dt = \int_0^T \left\langle \frac{\delta E(\phi(t))}{\delta \phi}, g(t) \right\rangle dt.$$

By the expression (1.2) of $\frac{\delta E(\phi_n)}{\delta \phi}$ we have to identify each limit. Indeed, we have

$$\int_0^T \langle \Delta^2 \phi_n(t), g(t) \rangle dt \rightarrow \int_0^T \langle \Delta^2 \phi(t), g(t) \rangle dt$$

by (vi) of Theorem 2.8. Similarly,

$$\begin{aligned} \lim_{n \rightarrow \infty} \int_0^T \langle \Delta(\phi_n^3(t) - \phi_n(t)), g(t) \rangle dt &= \int_0^T \langle \phi_n^3(t) - \phi_n(t), \Delta g(t) \rangle dt \\ &= \int_0^T \langle \phi^3(t) - \phi(t), \Delta g(t) \rangle dt = \int_0^T \langle \Delta(\phi^3(t) - \phi(t)), g(t) \rangle dt \end{aligned}$$

thanks to (v) of Theorem 2.8. Moreover, by Theorem 2.8, (v), (viii) and the Sobolev embedding $L^6(Q) \subset H^2$ the convergences $\phi_n^3 \rightarrow \phi^3$ and $f(\phi_n)g \rightarrow f(\phi)g$ holds in $L^2([0, T]; L^2(Q))$. We deduce that

$$\lim_{n \rightarrow \infty} \int_0^T \langle (3\phi_n^3 - 1)f(\phi_n), g(t) \rangle dt = \int_0^T \langle (3\phi^3 - 1)f(\phi), g(t) \rangle dt$$

holds. For the last term, we have to show that

$$\lim_{n \rightarrow \infty} \int_0^T \mathcal{B}(\phi_n(t)) \langle f(\phi_n(t)), g(t) \rangle dt = \int_0^T \mathcal{B}(\phi(t)) \langle f(\phi(t)), g(t) \rangle dt \quad (3.1)$$

Since $\mathcal{B}(\phi_n) = \frac{1}{2}|\nabla \phi_n|_2^2 + \frac{1}{4}|\phi_n^2 - 1|_2^2$, by (v) of Theorem 2.8 we deduce that $\mathcal{B}(\phi_n) \rightarrow \mathcal{B}(\phi)$ in $L^p([0, T]; \mathbb{R})$, for any $p \in [1, \infty[$. On the other side, by (v) of Theorem 2.8 we have $f(\phi_n)g \rightarrow f(\phi)g$

as $n \rightarrow \infty$ in $L^p([0, T]; \mathbb{R})$. Then, we deduce that (3.1) holds. Since we have identified all the terms, the first limit of the Lemma is proved.

The second limit is obvious since for any $g \in C_0^\infty([0, T] \times Q; \mathbb{R})$, we have $P_n g \rightarrow g$ strongly in $L^2([0, T]; L^2(Q))$ and then

$$\begin{aligned} \lim_{n \rightarrow \infty} \int_0^T \left\langle P_n \left(\frac{\delta E(\phi_n(t))}{\delta \phi} \right), g(t) \right\rangle dt &= \lim_{n \rightarrow \infty} \int_0^T \left\langle \frac{\delta E(\phi_n(t))}{\delta \phi}, P_n g(t) \right\rangle dt \\ &= \int_0^T \left\langle \frac{\delta E(\phi(t))}{\delta \phi}, g(t) \right\rangle dt. \quad \square \end{aligned}$$

Finally we prove that the limit processes are solutions by identifying all the terms in the equation.

Lemma 3.3. *Under hypothesis of Theorem 0.3, the limit processes (w, ϕ) solve (0.3) in the sense of Definition 1.2*

Proof. Let us first show that (w, ϕ) solve (0.3). Since w_n, ϕ_n solves (2.1), it is sufficient to show that the right-hand side of (2.1) converges to the right-hand side of (0.3).

Let $\xi \in L^2([0, T]; D(A))$. By Theorem 2.8, (iii) we have

$$\lim_{n \rightarrow \infty} \int_0^T \langle w_n + \alpha^2 A w_n, \xi(t) \rangle dt = \int_0^T \langle w + \alpha^2 A w, \xi(t) \rangle dt$$

and

$$\begin{aligned} \lim_{n \rightarrow \infty} \nu \int_0^T \left\langle \int_0^t (w_n(\tau) + \alpha^2 A w_n(\tau)) d\tau, A \xi(t) \right\rangle dt &= \lim_{n \rightarrow \infty} \nu \int_0^T \int_0^t \langle w(\tau) + \alpha^2 A w(\tau), A \xi(t) \rangle d\tau dt \\ &= \nu \int_0^T \left\langle \int_0^t (w(\tau) + \alpha^2 A w(\tau)) d\tau, A \xi(t) \right\rangle dt. \end{aligned}$$

Observe that by Hölder's inequality and Proposition 1.1 (ii) it holds

$$\left| \int_0^T \int_0^t \langle \tilde{B}(w(\tau), u(\tau)), \xi(t) \rangle d\tau dt \right| \leq c \left(\int_0^T \|w(\tau)\|_V^2 d\tau \right)^{\frac{1}{2}} \left(\int_0^T |u(\tau)|_2^2 d\tau \right)^{\frac{1}{2}} \left(\int_0^T \|\xi(t)\|_{D(A)}^2 dt \right)^{\frac{1}{2}}$$

where $c > 0$ does not depend on w, u, ξ . This implies that the trilinear form

$$\begin{aligned} &L^2([0, T]; V) \times L^2([0, T]; L^2(Q)) \times L^2([0, T]; D(A)) \rightarrow \mathbb{R} \\ &(w, u, \xi) \mapsto \int_0^T \int_0^t \langle \tilde{B}(w(\tau), u(\tau)), \xi(t) \rangle d\tau dt \end{aligned}$$

is continuous. Since by Theorem 2.8 we have that \mathbb{P} -a.s. $w_n \rightarrow w$ strongly in $L^2([0, T]; V)$, that $w_n + \alpha^2 A w_n \rightarrow w + \alpha^2 A w$ weakly in $L^2([0, T]; L^2(Q))$ and clearly $P_n \xi \rightarrow \xi$ strongly in $L^2([0, T]; D(A))$, we deduce that

$$\begin{aligned} &\lim_{n \rightarrow \infty} \int_0^T \left\langle \int_0^t P_n \tilde{B}(w_n, w_n + \alpha^2 A w_n)(\tau) d\tau, \xi(t) \right\rangle_{D(A)', D(A)} dt \\ &= \lim_{n \rightarrow \infty} \int_0^T \int_0^t \langle \tilde{B}(w_n, w_n + \alpha^2 A w_n)(\tau), P_n \xi(t) \rangle_{D(A)', D(A)} d\tau dt \\ &= \int_0^T \int_0^t \langle \tilde{B}(w, w + \alpha^2 A w)(\tau), \xi(t) \rangle_{D(A)', D(A)} d\tau dt. \end{aligned}$$

Finally, it is easy to see that $\tilde{\mathbb{P}}$ -a.s. it holds

$$\lim_{n \rightarrow \infty} \int_0^T \left\langle \int_0^t (P_n \Sigma) d\tilde{W}(\tau), \xi(t) \right\rangle dt = \int_0^T \left\langle \int_0^t \Sigma d\tilde{W}(\tau), \xi(t) \right\rangle dt.$$

By the previous Lemma 3.2 and by (vii) of Theorem 2.8 we get that

$$\int_0^t P_n \left(\frac{\delta E(\phi_n(\tau))}{\delta \phi} \right) \nabla \phi_n(\tau) d\tau \rightarrow \int_0^t \frac{\delta E(\phi(\tau))}{\delta \phi} \nabla \phi(\tau) d\tau$$

strongly in $L^2([0, T]; L^2(Q))$. Then, in particular, the convergence holds weakly in $L^2([0, T]; D(A))$. So, we have shown that w solves the first equation of (0.3). Let us show that ϕ solves the second one. Let us observe that by Theorem 2.8, $\phi_n \rightarrow \phi$ strongly in $L^p([0, T]; H^2(Q))$. Moreover, since $w_n \rightarrow w$ strongly in $L^p([0, T]; V)$, it is easy to show that the limit

$$\lim_{n \rightarrow \infty} \int_0^t \pi_n(w_n \nabla \phi_n)(\tau) d\tau = \int_0^t (w \nabla \phi)(\tau) d\tau$$

holds in $L^2([0, T]; L^2(Q))$. Finally, it is clear that

$$\lim_{n \rightarrow \infty} \int_0^T \left\langle \int_0^t (\pi_n \Xi) d\tilde{Z}(\tau), \xi(t) \right\rangle dt = \int_0^T \left\langle \int_0^t \Xi d\tilde{Z}(\tau), \xi(t) \right\rangle dt$$

holds $\tilde{\mathbb{P}}$ -a.s. Then, (w, ϕ) is a solution of (0.3).

It remains to verify that (w, ϕ) satisfy all the other conditions of Definition 1.2.

Continuity of $w + \alpha^2 Aw$, ϕ .

Notice that since the processes $w + \alpha^2 Aw$, ϕ solve the stochastic differential equation (1.1), then $w + \alpha^2 Aw \in L^2(\Omega; \mathcal{C}([0, T]; D(A)'))$ and $\phi \in L^2(\Omega; \mathcal{C}([0, T]; L^2(Q)))$. The fact that ϕ, w are adapted to the filtration \mathcal{F}_t is obvious, since ϕ, w are a.s. limits of adapted processes.

It remains to show that ϕ, w are continuous in mean square. Indeed, by Itô's formula (2.2) we deduce

$$\begin{aligned} \frac{1}{2} \mathbb{E} [|w_n(t) - w_n(t_0)|_2^2 + \alpha^2 |\nabla(w_n(t) - w_n(t_0))|_2^2] &\leq -\mathbb{E} \int_{t_0}^t \nu(|\nabla w_n|_2^2 + \alpha^2 |Aw_n|_2^2) ds \\ &\quad + \mathbb{E} \int_{t_0}^t \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 |\nabla \phi_n|_2 |w_n|_\infty ds + \text{Tr}[(\pi_n \Sigma)^*(I + \alpha^2 A)^{-1}(\pi_n \Sigma)](t - t_0) \\ &\leq -\mathbb{E} \int_{t_0}^t \nu(|\nabla w_n|_2^2 + \alpha^2 |Aw_n|_2^2) ds \\ &\quad + C_{\mathcal{P}} \mathbb{E} \int_{t_0}^t \left| \frac{\delta E(\phi_n)}{\delta \phi} \right|_2 |\nabla \phi_n|_2 \|w_n\|_V ds + \text{Tr}[\Sigma^*(I + \alpha^2 A)^{-1}\Sigma](t - t_0). \end{aligned}$$

Notice that we have used the property $\langle \tilde{B}(w, w + \alpha^2 Aw), w \rangle = 0$ and the Poincaré inequality $|w_n|_\infty \leq C_{\mathcal{P}} \|w_n\|_V$. Moreover, by Theorem 2.8 and the bounds in (0.4) we can apply Fatou's Lemma to get, as $n \rightarrow \infty$

$$\begin{aligned} \frac{1}{2} \mathbb{E} [|w(t) - w(t_0)|_2^2 + \alpha^2 |\nabla(w(t) - w(t_0))|_2^2] &\leq -\mathbb{E} \int_{t_0}^t \nu(|\nabla w|_2^2 + \alpha^2 |Aw|_2^2) ds \\ &\quad + C_{\mathcal{P}} \mathbb{E} \int_{t_0}^t \left| \frac{\delta E(\phi)}{\delta \phi} \right|_2 |\nabla \phi|_2 \|w\|_V ds + \text{Tr}[\Sigma^*(I + \alpha^2 A)^{-1}\Sigma](t - t_0). \end{aligned}$$

Then, the continuity in mean square for w follows. In a similar way (we omit the calculus, which are standard) we get the continuity in mean square for the process ϕ . \square

Corollary 3.4. *Under hypothesis of Theorem 0.3, we have*

$$\lim_{n \rightarrow \infty} \int_0^t P_n \left(\frac{\delta E(\phi_n(\tau))}{\delta \phi} \right) d\tau = \int_0^t \frac{\delta E(\phi(\tau))}{\delta \phi} d\tau \quad \text{in } L^p([0, T]; L^2(Q)), p \in [1, \infty[.$$

3.2. Uniqueness

By the Yamada-Watanabe theorem for spdes (see, for instance, [36, 38]), it is sufficient to show pathwise uniqueness of the solution.

Theorem 3.5. *Under Hypothesis 0.1 for any initial condition $(w_0, \phi_0) \in D(A) \times L^2(Q)$ there exists a unique solution (w, ϕ) to equation (0.3) such that for any $T > 0$ and \mathbb{P} -a.s.*

$$\int_0^T \left(\|w(t)\|_V^2 + |Aw|_2^2 + \left| \frac{\delta E}{\delta \phi}(\phi(t)) \right|_2^2 + |\phi|_{H^2}^8 + |\Delta^2 \phi|_2^2 \right) dt < \infty \quad (3.2)$$

Since the proof of this result is quite the same as in [35], for the reader's convenience we only give here the main ideas.

Proof. By Theorem 2.8 and Theorem 3.1, there exists at least a solution (w, ϕ) satisfying (3.2). As usual, consider two solutions of the system (w_1, ϕ_1) and (w_2, ϕ_2) with the expected regularity stated before, and consider the difference $(w, \phi) = (w_1, \phi_1) - (w_2, \phi_2)$ between these two solutions. We shall show that $(w_1, \phi_1) = (w_2, \phi_2)$ on the full measure set

$$\left\{ \sum_{i=1}^2 \int_0^T \left(\|w_i(t)\|_V^2 + |Aw_i|_2^2 + \left| \frac{\delta E}{\delta \phi}(\phi_i(t)) \right|_2^2 + |\phi_i|_{H^2}^8 + |\Delta^2 \phi_i|_2^2 \right) dt < \infty \right\}. \quad (3.3)$$

As in [35], we write

$$\frac{\delta E(\phi)}{\delta \phi}(\phi) = M(\phi) + N(\phi),$$

where

$$M(\phi) = \Delta^2 \phi - \Delta \phi + \phi$$

$$N(\phi) = \frac{\delta E}{\delta \phi}(\phi) - M(\phi).$$

Let us set $G(\phi) = |\Delta \phi|_2^2 + |\nabla \phi|_2^2 + |\phi|_2^2$. The proof of the following lemma is easy and it is left to the reader.

Lemma 3.6. *The function $G(\phi)$ defines a norm equivalent to the $H^2(Q)$ norm. That is, there exists $C > 0$ such that it holds*

$$\frac{1}{C} \|\phi\|_{H^2}^2 \leq G(\phi) \leq C \|\phi\|_{H^2}^2, \quad \forall \phi \in H^2(Q).$$

Moreover, it holds

$$\int_0^T G(\phi) dt = \int_0^T \langle M(\phi), \phi \rangle dt, \quad \forall \phi \in L^2([0, T]; H^4(Q)) \cap \{\phi : \phi = \Delta \phi = 0 \text{ on } \partial Q\}.$$

For any $\tilde{v} \in L^2([0, T]; D(A))$, the couple (w, ϕ) satisfies

$$\begin{cases} d \langle w + \alpha^2 Aw, \tilde{v} \rangle = \left(\langle -\nu A(w + \alpha^2 Aw), \tilde{v} \rangle \right. \\ \quad \left. + \langle \tilde{B}(w_2, w_2 + \alpha^2 Aw_2) - \tilde{B}(w_1, w_1 + \alpha^2 Aw_1), \tilde{v} \rangle \right. \\ \quad \left. + \langle M(\phi_1), \nabla \phi_1 \cdot \tilde{v} \rangle - \langle M(\phi_2), \nabla \phi_2 \cdot \tilde{v} \rangle + \langle N(\phi_1), \nabla \phi_1 \cdot \tilde{v} \rangle - \langle N(\phi_2), \nabla \phi_2 \cdot \tilde{v} \rangle \right) dt & \text{in } [0, T] \times Q, \\ d\phi = (-w_1 \cdot \nabla \phi_1 + w_2 \cdot \nabla \phi_2 - \gamma M(\phi) - \gamma N(\phi_1) + \gamma N(\phi_2)) dt & \text{in } [0, T] \times Q, \\ w(0) = 0 & \text{in } Q, \\ \phi(0) = 0 & \text{in } Q. \end{cases}$$

Let us look at the second equation. By multiplying with $M(\phi)$ and integrating over $[0, t] \times Q$ we find

$$\frac{1}{2} G(\phi(t)) = - \int_0^t \langle w_1 \cdot \nabla \phi_1 - w_2 \cdot \nabla \phi_2, M(\phi) \rangle ds - \gamma \int_0^t \left(\left\langle \frac{\delta E(\phi_1)}{\delta \phi} - \frac{\delta E(\phi_2)}{\delta \phi}, M(\phi) \right\rangle \right) ds \quad (3.4)$$

Here we used the fact that $E(\phi_i) < \infty$ implies $\phi_i \in H^2(Q)$ (see (1.3)). Moreover, notice that $\phi = \Delta\phi = 0$ on ∂Q and (3.2) holds, then we can apply the integration by parts in Lemma 3.6.

Since $w_i \in L^2([0, T]; D(A))$, we can set $\tilde{v} = w$ in the first equation. By integrating over $[0, t] \times Q$ we find

$$\begin{aligned} & \frac{1}{2}(|w(t)|_2^2 + \alpha^2|\nabla w(t)|_2^2) + \nu \int_0^t (|\nabla w|_2^2 + \alpha^2|\Delta w|_2^2) \, ds \\ &= \int_0^t \left(\left\langle -\tilde{B}(w_1, w_1 + \alpha^2 Aw_1) + \tilde{B}(w_2, w_2 + \alpha^2 Aw_2), w \right\rangle \right) \, ds \\ &+ \int_0^t \left(\left\langle \frac{\delta E(\phi_1)}{\delta \phi} \nabla \phi_1 - \frac{\delta E(\phi_2)}{\delta \phi} \nabla \phi_2, w \right\rangle \right) \, ds \\ &= \int_0^t \left\langle \tilde{B}(w, w + \alpha^2 Aw), w_2 \right\rangle \, ds + \int_0^t (\langle M(\phi_1) \nabla \phi_1 - M(\phi_2) \nabla \phi_2, w \rangle) \, ds \\ &+ \int_0^t (\langle N(\phi_1) \nabla \phi_1 - N(\phi_2) \nabla \phi_2, w \rangle) \, ds \end{aligned} \quad (3.5)$$

Here we have used the properties of \tilde{B} (see Proposition 1.1) which yield

$$\left\langle \tilde{B}(w_1, w_1 + \alpha^2 Aw_1), w \right\rangle - \left\langle \tilde{B}(w_2, w_2 + \alpha^2 Aw_2), w \right\rangle = - \left\langle \tilde{B}(w, w + \alpha^2 Aw), w_2 \right\rangle.$$

By adding (3.5) and (3.4) we get

$$\frac{1}{2}(|w(t)|_2^2 + \alpha^2|\nabla w(t)|_2^2 + |G(\phi(t))|_2^2) + \nu \int_0^t (|\nabla w|_2^2 + \alpha^2|\Delta w|_2^2) \, ds + \gamma \int_0^t |M(\phi)|_2^2 \, ds = \int_0^t \mathcal{F}(w, \phi) \, ds$$

where

$$\begin{aligned} \mathcal{F}(w, \phi) &= \left\langle \tilde{B}(w, w + \alpha^2 Aw), w_2 \right\rangle + \langle M(\phi_1) \nabla \phi_1 - M(\phi_2) \nabla \phi_2, w \rangle \\ &- \langle w_1 \cdot \nabla \phi_1 - w_2 \cdot \nabla \phi_2, M(\phi) \rangle + \langle N(\phi_1) \nabla \phi_1 - N(\phi_2) \nabla \phi_2, w \rangle \\ &- \gamma \langle N(\phi_1) - N(\phi_2), M(\phi) \rangle \end{aligned}$$

As in [35], we have to estimate each term of $\mathcal{F}(w, \phi)$. A key role is played by the following result, which is similar to Lemma 5.2 of [35]. The main difference is that in [35] the solution ϕ belongs to $\mathcal{C}^0([0, T]; H^2)$. In our case, we are able to prove only $\phi \in L^p([0, T]; H^2)$.

Lemma 3.7. *Let $\phi_1, \phi_2 \in H^2(Q)$ such that $\phi_i + 1 = \Delta\phi_i = 0$ on ∂Q , $i = 1, 2$. Then there exists $c > 0$, independent of ϕ_1, ϕ_2 such that*

$$|N(\phi_1) - N(\phi_2)|_2 \leq c(1 + \|\phi_1\|_{H^2}^6 + \|\phi_2\|_{H^2}^6) \|\phi_1 - \phi_2\|_{H^2}. \quad (3.6)$$

Proof. By (1.2) we can write

$$N(\phi) = -\Delta\phi^3 + 2\Delta\phi + 3\phi^2 f(\phi) - f(\phi) - \phi + M_1(\mathcal{A}(\phi) - a) + M_2(\mathcal{B}(\phi) - b)f(\phi).$$

Then,

$$\begin{aligned} |N(\phi_1) - N(\phi_2)|_2 &\leq |\Delta(\phi_1^3 - \phi_2^3)|_2 + 2|\Delta(\phi_1 - \phi_2)|_2 + 3|\phi_1^2 f(\phi_1) - \phi_2^2 f(\phi_2)|_2 + \\ &+ (1 + M_2 b) |f(\phi_1) - f(\phi_2)|_2 + |\phi_1 - \phi_2|_2 + M_1 |\mathcal{A}(\phi_1) - \mathcal{A}(\phi_2)|_2 + \\ &+ M_2 |\mathcal{B}(\phi_1) f(\phi_1) - \mathcal{B}(\phi_2) f(\phi_2)|_2 = I_1 + \dots + I_7. \end{aligned}$$

Let us proceed by estimating each term. For I_1 , we set $\tilde{\phi} = \phi_1^2 + \phi_1 \phi_2 + \phi_2^2$. Using Poincaré inequality, it holds $|\phi_i|_\infty \leq (1 + C_P \|\phi_i\|_{H^1})$, $i = 1, 2$ where $C_P > 0$ is the Poincaré constant. We deduce that there exists $c > 0$ such that

$$|\tilde{\phi}|_\infty \leq \sum_{i,j=1}^2 |\phi_i|_\infty |\phi_j|_\infty \leq \sum_{i,j=1}^2 (1 + C_P \|\phi_i\|_{H^1})(1 + C_P \|\phi_j\|_{H^1}) \leq c(1 + \|\phi_1\|_{H^1}^2 + \|\phi_2\|_{H^1}^2). \quad (3.7)$$

Similarly, since $|\nabla \tilde{\phi}| \leq \sum_{i,j=1}^2 |\phi_i| |\nabla \phi_j|$ by Poincaré inequality and the Sobolev embedding $H^1(Q) \subset L^4(Q)$ there exists $c > 0$ such that

$$\begin{aligned} |\nabla \tilde{\phi}|_4 &\leq \sum_{i,j=1}^2 |\phi_i|_\infty |\nabla \phi_j|_4 \leq \sum_{i,j=1}^2 (1 + C_{\mathcal{P}} \|\nabla \phi_i\|_{H^1}) \|\nabla \phi_j\|_{H^1} \\ &\leq \sum_{i,j=1}^2 (1 + C_{\mathcal{P}} \|\phi_i\|_{H^2}) \|\phi_j\|_{H^2} \leq (1 + \|\phi_1\|_{H^2}^2 + \|\phi_2\|_{H^2}^2). \end{aligned} \quad (3.8)$$

Moreover, since $\Delta \tilde{\phi} = \sum_{i,j=1}^2 \Delta(\phi_i \phi_j) = \sum_{i,j=1}^2 ((\Delta \phi_i) \phi_j + \nabla \phi_i \nabla \phi_j + \phi_i (\Delta \phi_j))$, still using the Poincaré inequality and the Sobolev embedding $H^1(Q) \subset L^4(Q)$ we get

$$\begin{aligned} |\Delta \tilde{\phi}|_2 &\leq \sum_{i,j=1}^2 (|\Delta \phi_i|_2 |\phi_j|_\infty + |\nabla \phi_i|_4 |\nabla \phi_j|_4 + |\phi_i|_\infty |\Delta \phi_j|_2) \\ &\leq \sum_{i,j=1}^2 (|\Delta \phi_i|_2 (C_{\mathcal{P}} \|\phi_j\|_{H^1} + 1) + \|\phi_i\|_{H^2} \|\phi_j\|_{H^2} + (C_{\mathcal{P}} \|\phi_i\|_{H^1} + 1) |\Delta \phi_j|_2) \\ &\leq c(1 + \|\phi_1\|_{H^2}^2 + \|\phi_2\|_{H^2}^2) \end{aligned} \quad (3.9)$$

where $c > 0$ is independent of ϕ_1, ϕ_2 . By taking in mind (3.7), (3.8), (3.9) there exists $c > 0$ such that

$$\begin{aligned} I_1 = |\Delta(\phi_1^3 - \phi_2^3)|_2 &= |\Delta(\phi \tilde{\phi})|_2 \\ &\leq |(\Delta \phi) \tilde{\phi}|_2 + 2|\nabla \phi \cdot \nabla \tilde{\phi}|_2 + |\phi(\Delta \tilde{\phi})|_2 \\ &\leq |\Delta \phi|_2 |\tilde{\phi}|_\infty + 2|\nabla \phi|_4 |\nabla \tilde{\phi}|_4 + |\phi|_\infty |\Delta \tilde{\phi}|_2 \\ &\leq c(1 + \|\phi_1\|_{H^2}^2 + \|\phi_2\|_{H^2}^2) (|\Delta \phi|_2 + |\nabla \phi|_4 + |\phi|_\infty) \\ &\leq c(1 + \|\phi_1\|_{H^2}^2 + \|\phi_2\|_{H^2}^2) \|\phi_1 - \phi_2\|_{H^2}. \end{aligned}$$

In the last inequality we have used the Sobolev embedding $H^1(Q) \subset L^4(Q)$ and the Poincaré inequality $|\phi_1 - \phi_2|_\infty \leq C_{\mathcal{P}} |\phi_1 - \phi_2|_{H^1}$. For I_2 , we have clearly $I_2 \leq c |\phi_1 - \phi_2|_{H^2}$. For I_3 we can write

$$I_3 = 3|\phi_1^2 f(\phi_1) - \phi_2^2 f(\phi_2)|_2 \leq 3|\phi_1^2 (f(\phi_1) - f(\phi_2))|_2 + 3|(\phi_1^2 - \phi_2^2) f(\phi_2)|_2 = J_1 + J_2.$$

For J_1 , we have

$$J_1 \leq 3|\phi_1|_\infty^2 |f(\phi_1) - f(\phi_2)|_2 \leq 3|\phi_1|_\infty^2 (|\Delta(\phi_1 - \phi_2)|_2 + |\phi_1^3 - \phi_2^3|_2 + |\phi_1 - \phi_2|_2).$$

With a similar calculus done for I_1 , we have $|\phi_1^3 - \phi_2^3|_2 \leq |\tilde{\phi}|_\infty |\phi_1 - \phi_2|_2$. Then, using Poincaré's inequality there exists $c > 0$ such that $|\phi_1^3 - \phi_2^3|_2 \leq c(1 + \|\phi_1\|_{H^1}^2 + \|\phi_2\|_{H^1}^2) |\phi_1 - \phi_2|_2$. Then, for some $c > 0$ independent of ϕ_1, ϕ_2 we obtain

$$J_1 \leq c(1 + \|\phi_1\|_{H^1}^2) (1 + \|\phi_1\|_{H^1}^2 + \|\phi_2\|_{H^1}^2) \|\phi_1 - \phi_2\|_{H^2}.$$

For J_2 , since $H^1(Q) \subset L^6(Q)$ we have

$$\begin{aligned} J_2 &\leq 3|\phi_1 - \phi_2|_\infty |\phi_1 + \phi_2|_\infty (|\Delta \phi_2|_2 + |\phi_2|_6^3 + |\phi_2|_2) \\ &\leq c \|\phi_1 - \phi_2\|_{H^1} (1 + \|\phi_1\|_{H^1} + \|\phi_2\|_{H^1}) (\|\phi_2\|_{H^2} + \|\phi_2\|_{H^1}^3 + |\phi_2|_2) \end{aligned}$$

Finally, using Young's inequality repeatedly, we find that for some $c > 0$ it holds

$$I_3 \leq c(1 + \|\phi_1\|_{H^2}^6 + \|\phi_2\|_{H^2}^6) \|\phi_1 - \phi_2\|_{H^2}.$$

For I_4 , we can perform a calculus as done for J_1 to obtain, for some $c > 0$

$$I_4 = (1 + M_2 b) |f(\phi_1) - f(\phi_2)|_2 \leq c(1 + \|\phi_1\|_{H^1}^2 + \|\phi_2\|_{H^1}^2) \|\phi_1 - \phi_2\|_{H^2}.$$

Clearly, for I_5 and I_6 we have

$$I_5 + I_6 = |\phi_1 - \phi_2|_2 + M_1 |A(\phi_1) - A(\phi_2)|_2 \leq |\phi_1 - \phi_2|_2 + M_1 |Q|^{1/2} |A(\phi_1 - \phi_2)| \leq (1 + M_1 |Q|) |\phi_1 - \phi_2|_2.$$

For I_7 ,

$$\begin{aligned} I_7 &= M_2 |\mathcal{B}(\phi_1) f(\phi_1) - \mathcal{B}(\phi_2) f(\phi_2)|_2 \\ &\leq \mathcal{B}(\phi_1) |f(\phi_1) - f(\phi_2)|_2 + |\mathcal{B}(\phi_1) - \mathcal{B}(\phi_2)| |f(\phi_2)|_2 \\ &= K_1 + K_2 \end{aligned}$$

For K_1 , since $\mathcal{B}(\phi) = \frac{1}{2} |\nabla \phi|_2^2 + \frac{1}{4} |\phi^2 - 1|_2^2$ we have

$$\mathcal{B}(\phi_1) \leq \frac{1}{2} |\nabla \phi_1|_2^2 + \frac{1}{4} (|\phi_1|_4^4 + |Q|) \leq \frac{1}{2} |\nabla \phi_1|_2^2 + \frac{1}{4} (\|\phi_1\|_{H^1}^4 + |Q|).$$

Then, there exists $c > 0$ such that $\mathcal{B}(\phi_1) \leq c(1 + \|\phi_1\|_{H^1}^4)$. In order to estimate $|f(\phi_1) - f(\phi_2)|_2$ we can argue as done before for the term J_1 to obtain

$$|f(\phi_1) - f(\phi_2)|_2 \leq c(1 + \|\phi_1\|_{H^1}^2 + \|\phi_2\|_{H^1}^2) \|\phi_1 - \phi_2\|_{H^2}$$

Then, by using Young's inequality repeatedly, there exists $c > 0$ such that

$$K_1 \leq c(1 + \|\phi_1\|_{H^1}^6 + \|\phi_2\|_{H^1}^6) \|\phi_1 - \phi_2\|_{H^2}.$$

Before consider K_2 , let us observe that by the expression of $\mathcal{B}(\phi_1)$ we have

$$\begin{aligned} |\mathcal{B}(\phi_1) - \mathcal{B}(\phi_2)| &\leq \frac{1}{2} |\langle \nabla(\phi_1 - \phi_2), \nabla(\phi_1 + \phi_2) \rangle| + \frac{1}{4} |\langle \phi_1 - \phi_2, (\phi_1 + \phi_2)(\phi_1^2 + \phi_2^2 - 2) \rangle| \\ &\leq \frac{1}{2} \|\phi_1 - \phi_2\|_{H^1} (\|\phi_1\|_{H^1} + \|\phi_2\|_{H^1}) + \frac{1}{4} |\phi_1 - \phi_2|_2 |(\phi_1 + \phi_2)(\phi_1^2 + \phi_2^2 - 2)|_2. \end{aligned}$$

By Young inequality $ab^2 \leq a^2/3 + 2b^3/3$ we get $(\phi_1 + \phi_2)(\phi_1^2 + \phi_2^2 - 2) \leq 2\phi_1^3 + 2\phi_2^3 + \phi_1 + \phi_2$. Therefore, the last expression is bounded by

$$|\mathcal{B}(\phi_1) - \mathcal{B}(\phi_2)| \leq \|\phi_1 - \phi_2\|_{H^1} \left(\frac{\|\phi_1\|_{H^1} + \|\phi_2\|_{H^1}}{2} + \frac{|Q|^{1/2}}{4} (2|\phi_1|_\infty^3 + 2|\phi_2|_\infty^3 + |\phi_1|_\infty + |\phi_2|_\infty) \right).$$

Since by Poincaré's inequality we have $|\phi_i|_\infty \leq |\phi_i + 1|_\infty + 1 \leq C_P \|\phi_i\|_{H^1} + 1$, we deduce that there exists $c > 0$ such that

$$|\mathcal{B}(\phi_1) - \mathcal{B}(\phi_2)| \leq c \|\phi_1 - \phi_2\|_{H^1} (1 + \|\phi_1\|_{H^1}^3 + 2\|\phi_2\|_{H^1}^3)$$

Moreover, since $f(\phi) = -\Delta \phi + \phi(\phi^2 - 1)$ and the continuous embedding $H^1(Q) \subset L^6(Q)$ holds, there exists $c > 0$ such that

$$|f(\phi_2)|_2 \leq |\Delta \phi_2|_2 + |\phi_2|_6^3 + |\phi_2|_2 \leq \|\phi_2\|_{H^2} + |\phi_2|_{H^1}^3 + |\phi_2|_2 \leq c(1 + \|\phi_2\|_{H^2}^3)$$

By the previous results, we deduce that for K_2 we have

$$K_2 \leq c \|\phi_1 - \phi_2\|_{H^1} (1 + \|\phi_1\|_{H^1}^3 + 2\|\phi_2\|_{H^1}^3) (1 + \|\phi_2\|_{H^2}^3)$$

Then, for some $c > 0$ independent of ϕ_1, ϕ_2 we obtain the bound

$$K_2 \leq c \|\phi_1 - \phi_2\|_{H^1} (1 + \|\phi_1\|_{H^2}^6 + \|\phi_2\|_{H^2}^6).$$

Taking into account the estimates on K_1 and K_2 we get that for some $c > 0$ we have

$$I_7 \leq c(1 + \|\phi_1\|_{H^2}^6 + \|\phi_2\|_{H^2}^6) \|\phi_1 - \phi_2\|_{H^2}.$$

Finally, taking into account the estimates on I_1, \dots, I_7 , we get that there exists $c > 0$ such that (3.6) holds. \square

By arguing as in [35] (see equations (60)–(67)), the term $\mathcal{F}(w, \phi)$ is bounded by

$$\begin{aligned} \mathcal{F}(w, \phi) \leq & C_{\tilde{\varepsilon}} |\nabla w|_2^2 |Aw_2|_2^2 + \tilde{\varepsilon} |w + \alpha^2 Aw|_2^2 \\ & + \tilde{\varepsilon} \|w\|_V^2 + C_{\tilde{\varepsilon}} |M(\phi_1)|_2^2 \|\phi\|_{H^2}^2 \\ & + \tilde{\varepsilon} |M(\phi)|_2^2 + C_{\tilde{\varepsilon}} \|w_1\|_V^2 \|\phi\|_{H^2}^2 \\ & + \tilde{\varepsilon} \|w\|_V^2 + C_{\tilde{\varepsilon}} |N(\phi_1)|_2^2 \|\phi\|_{H^2}^2 \\ & + \tilde{\varepsilon} \|w\|_V^2 + C_{\tilde{\varepsilon}} |N(\phi_1) - N(\phi_2)|_2^2 \|\phi_2\|_{H^2}^2 \\ & + \tilde{\varepsilon} |M(\phi)|_2^2 + C_{\tilde{\varepsilon}} |N(\phi_1) - N(\phi_2)|_2^2, \end{aligned}$$

where $\tilde{\varepsilon} > 0$ can be chosen arbitrarily and $C_{\tilde{\varepsilon}} > 0$ depends only on $\tilde{\varepsilon} > 0$. By (3.6) there exists $c > 0$ such that

$$\begin{aligned} \mathcal{F}(w, \phi) \leq & 3\tilde{\varepsilon} \|w\|_V^2 + 2\tilde{\varepsilon} |M(\phi)|_2^2 + \tilde{\varepsilon} |w + \alpha^2 Aw|_2^2 + C_{\tilde{\varepsilon}} |\nabla w|_2^2 |Aw_2|_2^2 \\ & + (C_{\tilde{\varepsilon}} (|M(\phi_1)|_2^2 + \|w_1\|_V^2 + |N(\phi_1)|_2^2 + c(1 + \|\phi_2\|_{H^2}^2)(1 + \|\phi_1\|_{H^2}^6 + \|\phi_2\|_{H^2}^6))) \|\phi\|_{H^2}^2. \end{aligned}$$

Since $|N(\phi)|_2 \leq \left| \frac{\delta E}{\delta \phi}(\phi) \right|_2 + |M(\phi)|_2$ and $|M(\phi)|_2 \leq c(|\Delta^2 \phi|_2 + \|\phi\|_{H^2})$, there exists $c_1, c_2 > 0$, depending only on $\tilde{\varepsilon}$ such that

$$\begin{aligned} \mathcal{F}(w, \phi) \leq & 3\tilde{\varepsilon} \|w\|_V^2 + 2\tilde{\varepsilon} |M(\phi)|_2^2 + \tilde{\varepsilon} |w + \alpha^2 Aw|_2^2 + C_{\tilde{\varepsilon}} |\nabla w|_2^2 |Aw_2|_2^2 \\ & + c_1 \left(|\Delta^2 \phi_1|_2^2 + \|w_1\|_V^2 + \left| \frac{\delta E}{\delta \phi}(\phi_1) \right|_2^2 + 1 + \|\phi_1\|_{H^2}^8 + \|\phi_2\|_{H^2}^8 \right) \|\phi\|_{H^2}^2 \\ \leq & 3\tilde{\varepsilon} \|w\|_V^2 + 2\tilde{\varepsilon} |M(\phi)|_2^2 + \tilde{\varepsilon} |w + \alpha^2 Aw|_2^2 \\ & + c_2 \left(|\Delta^2 \phi_1|_2^2 + \|w_1\|_V^2 + |Aw_2|_2^2 + \left| \frac{\delta E}{\delta \phi}(\phi_1) \right|_2^2 + \|\phi_1\|_{H^2}^8 + \|\phi_2\|_{H^2}^8 + 1 \right) \times \\ & \times (G(\phi) + |w|_2^2 + \alpha^2 |\nabla w|_2^2) \end{aligned}$$

Consequently, for $\tilde{\varepsilon}$ small enough, it holds

$$\begin{aligned} \frac{1}{2} (|w(t)|_2^2 + \alpha^2 |\nabla w(t)|_2^2 + G(\phi)(t)) + \frac{\gamma}{2} \int_0^T |M(\phi)|_2^2 dt + \frac{\nu}{2} \int_0^T (|\nabla w(t)|_2^2 + \alpha^2 |Aw(t)|_2^2) dt \\ \leq \frac{1}{2} \int_0^T H(t) (|w(t)|_2^2 + \alpha^2 |\nabla w(t)|_2^2 + G(\phi)(t)) dt. \end{aligned}$$

Here, H (up to a multiplicative constant) is explicitly given by

$$|\Delta^2 \phi_1|_2^2 + \|w_1\|_V^2 + |Aw_2|_2^2 + \left| \frac{\delta E}{\delta \phi}(\phi_1) \right|_2^2 + \|\phi_1\|_{H^2}^8 + \|\phi_2\|_{H^2}^8 + 1.$$

By the conditions (3.2), the quantity $\int_0^T H(t) dt$ is bounded. Then we can apply Gronwall's lemma to deduce

$$|w(t)|_2^2 + \alpha^2 |\nabla w(t)|_2^2 + G(\phi)(t) \leq 0$$

which implies $(w_1, \phi_1) = (w_2, \phi_2)$ on the full measure set defined in (3.3). \square

Acknowledgement

The authors would like to thank the anonymous reviewers for their interesting remarks, corrections and helpful comments.

References

- [1] M. Abkarian, C. Lartigue, and A. Viallat. Tank treading and unbinding of deformable vesicles in shear flow: Determination of the lift force. *Phys. Rev. Lett.*, 88(6), 2002.
- [2] D. C. Antonopoulou, G. Karali, and A. Millet. Existence and regularity of solution for a stochastic Cahn-Hilliard/Allen-Cahn equation with unbounded noise diffusion. *J. Differential Equations*, 260(3):2383–2417, 2016.
- [3] J. B. Walsh. An introduction to stochastic partial differential equations. *École d’Été de Probabilités de Saint Flour XIV-1984*, 1180, January 1986.
- [4] J. Beaucourt, F. Rioual, T. Séon, T. Biben, and C. Misbah. Steady to unsteady dynamics of a vesicle in a flow. *Physical review. E, Statistical, nonlinear, and soft matter physics*, 69:011906, February 2004.
- [5] T. Biben, K. Kassner, and C. Misbah. Phase-field approach to three-dimensional vesicle dynamics. *Physical review. E, Statistical, nonlinear, and soft matter physics*, 72:041921, November 2005.
- [6] C. Bjorland and M. Schonbek. On questions of decay and existence for the viscous camassa-holm equations. *Annales de l’Institut Henri Poincaré (C) Non Linear Analysis*, 25:907–936, September 2006.
- [7] A. Çağlar. Convergence analysis of the navier-stokes alpha model. *Numerical Methods for Partial Differential Equations*, 26:1154 – 1167, July 2009.
- [8] C. Cardon-Weber. Cahn-hilliard stochastic equation: Existence of the solution and of its density. *Bernoulli*, 7, October 2001.
- [9] S. Chen, C. Foias, D. Holm, E. Olson, E. Titi, and S. Wynne. A connection between the camassa-holm equations and turbulent flows in channels and pipes. *Physics of Fluids*, 11, March 1999.
- [10] G. Da Prato and A. Debussche. Stochastic Cahn-Hilliard equation. *Nonlinear Anal.*, 26(2):241–263, 1996.
- [11] A. Debussche and L. Goudenège. Stochastic Cahn-Hilliard equation with double singular nonlinearities and two reflections. *SIAM J. Math. Anal.*, 43(3):1473–1494, 2011.
- [12] Q. Du and M. Li. Analysis of a stochastic implicit interface model for an immersed elastic surface in a fluctuating fluid. *Archive for rational mechanics and analysis*, 199(1):329–352, 2011.
- [13] Q. Du, M. Li, and C. Liu. Analysis of a phase field navier-stokes vesicle-fluid interaction model. *Discrete and Continuous Dynamical Systems - Series B*, 8, October 2007.
- [14] Q. Du, C. Liu, R. Ryham, and X. Wang. Modeling the spontaneous curvature effects in static cell membrane deformations by a phase field formulation. *Communications on Pure and Applied Analysis*, 4:537–548, September 2005.
- [15] Q. Du, C. Liu, R. Ryham, and X. Wang. A phase field formulation of the willmore problem. *Nonlinearity*, 18:1249–1267, May 2005.
- [16] Q. Du, C. Liu, and X. Wang. A phase field approach in the numerical study of the elastic bending energy for vesicle membranes. *Journal of Computational Physics*, 198:450–468, August 2004.
- [17] Q. Du, C. Liu, and X. Wang. Retrieving topological information for phase field models. *SIAM Journal of Applied Mathematics*, 65:1913–1932, January 2005.

- [18] J. E. Marsden and S. Shkoller. Global well-posedness for the lagrangian averaged navier-stokes (lans-) equations on bounded domains. *Philosophical Transactions of The Royal Society B: Biological Sciences*, 359, July 2001.
- [19] N. Elezović and A. Mikelić. On the stochastic Cahn-Hilliard equation. *Nonlinear Anal.*, 16(12):1169–1200, 1991.
- [20] F. Flandoli and D. Gatarek. Martingale and stationary solutions for stochastic navier-stokes equations. *Probability Theory and Related Fields*, 102:367–391, September 1995.
- [21] C. Foias, D. Holm, and E. Titi. The navier-stokes-alpha model of fluid turbulence. *Physica D: Nonlinear Phenomena*, 152:505–519, May 2001.
- [22] C. Foias, D. Holm, and E. Titi. The three dimensional viscous Camassa-Holm equations, and their relation to the navier-stokes equations and turbulence theory. *Journal of Dynamics and Differential Equations*, 14:1–35, January 2002.
- [23] L. Goudenège. Stochastic Cahn-Hilliard equation with singular nonlinearity and reflection. *Stochastic Process. Appl.*, 119(10):3516–3548, 2009.
- [24] L. Goudenège and L. Manca. Asymptotic properties of stochastic Cahn-Hilliard equation with singular nonlinearity and degenerate noise. *Stochastic Processes and their Applications*, 125(10):3785 – 3800, 2015.
- [25] J. Guermond, J. Oden, and S. Prudhomme. An interpretation of the ns alpha model as a frame indifferent leray regularization. *Physica D: Adv. Math. Phys. Fluids*, 177:23–30, March 2003.
- [26] I. Gyongy. Existence and uniqueness results for semilinear stochastic partial differential equations. *Stochastic Processes and their Applications*, 73:271–299, March 1998.
- [27] M. Hairer and J. C. Mattingly. Ergodicity of the 2D Navier-Stokes equations with degenerate stochastic forcing. *Ann. of Math.*, 164(3):993–1032, 2006.
- [28] M. Hairer and J. C. Mattingly. A theory of hypoellipticity and unique ergodicity for semilinear stochastic PDEs. *Electron. J. Probab.*, 16:no. 23, 658–738, 2011.
- [29] D. Holm, J. E. Marsden, and T. Ratiu. The euler-poincaré equations and semidirect products with applications to continuum theories. *Advances in Mathematics*, 137:1–81, February 1998.
- [30] D. Holm, J. E. Marsden, and T. Ratiu. Euler-poincaré models of ideal fluids with nonlinear dispersion. *Physical Review Letters*, 80:4173–4176, May 1998.
- [31] J. Leray. Essai sur le mouvement plan d’un liquide visqueux que limitent des parois. *Journal de Mathématiques Pures et Appliquées. Neuvième Série*, 13, January 1934.
- [32] J. Leray. Sur le mouvement d’un liquide visqueux emplissant l’espace. *Acta Mathematica*, 63(1):193–248, December 1934.
- [33] Y. Liu, T. Takahashi, and T. Marius. Strong solutions for a phase field navier-stokes vesicle-fluid interaction model. *J. Math. Fluid Mech. c*, 14, March 2011.
- [34] C. Odasso. Exponential mixing for the 3D stochastic Navier-Stokes equations. *Comm. Math. Phys.*, 270(1):109–139, 2007.
- [35] A. Piovezan Entringer and J. Luiz Boldrini. A phase field alpha-navier-stokes vesicle-fluid interaction model: Existence and uniqueness of solutions. *Discrete and Continuous Dynamical Systems - Series B*, 20:397–422, March 2015.
- [36] M. Rockner, B. Schmulland, and X. Zhang. Yamada-watanabe theorem for stochastic evolution equations in infinite dimensions. *Condensed Matter Physics*, 54(11):247, 2008.

- [37] U. Seifert. Configurations of fluid membranes and vesicles. *Adv. Phys.*, 46:13–137, February 1997.
- [38] S. Tappe. The Yamada-Watanabe theorem for mild solutions to stochastic partial differential equations. *Electron. Commun. Probab.*, 18:no. 24, 13, 2013.