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A conjecture for the regularized fourth moment of Eisenstein series

Goran Djanković^{a,1}, Rizwanur Khan^{b,*}

^a *University of Belgrade, Faculty of Mathematics, Studentski Trg 16, p.p. 550, 11000 Belgrade, Serbia*

^b *Texas A&M University at Qatar, PO Box 23874, Doha, Qatar*

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ABSTRACT

We formulate a version of the Random Wave Conjecture for the fourth moment of Eisenstein series which is based on Zagier's regularized inner product. We prove an asymptotic formula expressing the regularized fourth moment as a mean value of L -functions. This is an advantage over previous work in the literature, which has approached the fourth moment problem through truncated Eisenstein series and not yielded a suitable expression in terms of L -functions.

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* Corresponding author.

E-mail addresses: djankovic@matf.bg.ac.rs (G. Djanković), rizwanur.khan@qatar.tamu.edu (R. Khan).

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1. Introduction

One of the main research themes in recent years in the theory of automorphic forms is the problem of mass distribution. Let $X = \Gamma \backslash \mathbb{H}$, where \mathbb{H} is the upper half complex plane and $\Gamma = SL_2(\mathbb{Z})$. In his PhD thesis, Spinu [Sp] obtained the following type of weak equidistribution result:

$$\int_X |E_A(z, \tfrac{1}{2} + iT)|^4 d\mu z \ll T^\epsilon, \quad (1.1)$$

where $d\mu(z) = \frac{dx dy}{y^2}$ and $E_A(z, s)$ is the truncated Eisenstein series, which on the fundamental domain equals $E(z, s)$ for $\text{Im}(z) \leq A$, and $E(z, s)$ minus its constant term for $\text{Im}(z) > A$. See the next section for a more careful definition. Spinu's result (see also [Lu] for a closely related result) is in line with a much more general conjecture, called the Random Wave Conjecture. This conjecture was made for Eisenstein series in [HR, section 7.3]. In terms of moments this implies: for any even integer $p \geq 0$ and any nice compact $\Omega \subset X$, we should have

$$\lim_{T \rightarrow \infty} \frac{1}{\text{vol}(\Omega)} \int_{\Omega} \left| \frac{E_A(z, \frac{1}{2} + iT)}{\sqrt{2 \log T}} \right|^p d\mu z = \frac{c_p}{\text{vol}(X)^{p/2}}, \quad (1.2)$$

where c_p is the p th moment of the normal distribution $\mathcal{N}(0, 1)$. The same conjecture is also made for $E(z, \frac{1}{2} + iT)$. As we will see below, $\sqrt{2 \log T}$ roughly equals $\|E_A(\cdot, \frac{1}{2} + iT)\|_2$.

One would of course like to go beyond Spinu's upper bound and prove an asymptotic for the fourth moment of Eisenstein series. In [BK], this was achieved, conditional on the Generalized Lindelöf Hypothesis, for Hecke Maass forms of large eigenvalue when $\Omega = X$, and agreement was found with the RWC. Thus in analogy one would expect (1.2) to also hold for $p = 4$ and $\Omega = X$, and one may hope that the statement in this case can be proven unconditionally. After all, such problems can be a bit easier for Eisenstein series – for example, recall that the case $p = 2$ of (1.2) was first proven for Eisenstein series [LS] before the analogue was proven for Hecke Maass forms [Li, So].

What would the proof of such an asymptotic entail? The starting point in [BK] is to relate the fourth moment of an L^2 -normalized Hecke Maass form f to L -functions. One uses the spectral decomposition and Plancherel's theorem to write

$$\langle f^2, f^2 \rangle = \sum_{j \geq 1} |\langle f^2, u_j \rangle|^2 + \dots, \quad (1.3)$$

where the inner product is the Petersson inner product, $\{u_j : j \geq 1\}$ is an orthonormal basis of Hecke Maass forms, and the ellipsis denotes the contribution of the Eisenstein spectrum and constant eigenfunction. Next one can use Watson's triple product formula to relate the squares of the inner products on the right hand side to central values of L -functions. Thus the problem is reduced to one of obtaining a mean value of L -functions.

If one tries to mimic this set up for $E(z, \frac{1}{2} + iT)$ in place of f , the first obvious difficulty encountered is that the left hand side of (1.3) does not even converge. To circumvent this, Spinu worked with the truncated $E_A(z, \frac{1}{2} + iT)$, which decays exponentially at the cusp. However a major drawback is that $E_A(z, \frac{1}{2} + iT)$ is not automorphic, so Spinu could not obtain a precise relationship with L -functions. He could only obtain an upper bound [Sp, section 4.2].

The goal of this paper is to reformulate entirely the fourth moment problem for Eisenstein series. To make sense of $\langle E^2(\cdot, \frac{1}{2} + iT), E^2(\cdot, \frac{1}{2} + iT) \rangle$, we contend that it is more natural² to use Zagier's regularized inner product [Za], which does converge. The basic idea of Zagier's method is that to kill off the growth of an automorphic form, one should not subtract off the constant term like Spinu does, but rather subtract off another Eisenstein series in such a way that the final object is square integrable and automorphic. This way we will end up with a precise relationship between a *regularized* fourth moment and L -functions. This is the first goal of our paper, and we will prove

Theorem 1.1 (*Regularized fourth moment in terms of L -functions*). *Let $\{u_j : j \geq 1\}$ denote an orthonormal basis of even and odd Hecke Maass cusp forms for Γ , ordered by Laplacian eigenvalue $\frac{1}{4} + it_j^2$, and let $\Lambda(s, u_j)$ denote the corresponding completed L -functions. Let $\xi(s)$ denote the completed Riemann ζ function. As $T \rightarrow \infty$, we have*

$$\begin{aligned} & \int_X^{reg} |E(z, 1/2 + iT)|^4 d\mu(z) \\ &= \frac{24}{\pi} \log^2 T + \sum_{j \geq 1} \frac{\cosh(\pi t_j)}{2} \frac{|\Lambda(\frac{1}{2} + 2Ti, u_j)|^2 \Lambda^2(\frac{1}{2}, u_j)}{L(1, \text{sym}^2 u_j) |\xi(1 + 2Ti)|^4} + O(\log^{5/3+\epsilon} T), \end{aligned}$$

for any $\epsilon > 0$.

This result is potentially very useful. We could try to obtain an asymptotic for the mean value of L -functions on the right hand side (and we will return to this problem in a future paper), thereby obtaining an asymptotic for the regularized fourth moment. This would be nice, but how would we know whether or not our answer is in agreement with the RWC? Thus the purpose of our second result is to translate the RWC to the setting of the regularized fourth moment. As defined in the next section, \mathcal{D}_A is the part of the fundamental domain with $\text{Im}(z) \leq A$.

Theorem 1.2 (*RWC for the regularized fourth moment of Eisenstein series*). *Suppose that (1.2) holds for $p = 4$ and $\Omega = X$, and $p = 4$ and $\Omega = \mathcal{D}_A$ for some $A = A(T)$ which*

² We thank Matthew Young for suggesting this approach to us.

tends to ∞ as $T \rightarrow \infty$. Then we have

$$\int_X^{reg} |E(z, \frac{1}{2} + iT)|^4 d\mu(z) \sim \frac{72}{\pi} \log^2 T.$$

We have already explained above why (1.2) should be expected for $p = 4$ and $\Omega = X$, even though for general p we must restrict to compact sets. The other possibility $\Omega = \mathcal{D}_A$ is already included in the RWC when A is fixed. But it is reasonable to conjecture that some effective error term will exist in (1.2), so that taking A which grows arbitrarily slowly should be permissible.

Both of our main results are based on careful calculations arising from the regularized inner product. The point is to offer a new viewpoint for the fourth moment and carefully put into place all leading constants, so that the relevant conjecture might be verified in the future using the theory of L -functions.

2. Eisenstein series

We recall the definition of Eisenstein series

$$E(z, s) = \sum_{\gamma \in \Gamma_\infty \backslash \Gamma} \text{Im}(\gamma z)^s = \frac{1}{2} y^s \sum_{\substack{c, d \in \mathbb{Z} \\ (c, d) = 1}} \frac{1}{|cz + d|^{2s}}, \quad z \in \mathbb{H},$$

where Γ_∞ is the stabilizer of the cusp ∞ in Γ . The series is absolutely convergent in the half-plane $\text{Re}(s) > 1$ where it defines an automorphic function satisfying $\Delta E(z, s) = s(1-s)E(z, s)$, for the hyperbolic Laplacian $\Delta = -y^2(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2})$.

The Eisenstein series can be meromorphically continued to the whole s -plane and $E(z, s)$ has the following Fourier expansion (for $s \neq 0, \frac{1}{2}, 1$)

$$E(z, s) = y^s + \varphi(s)y^{1-s} + \frac{2}{\xi(2s)} \sum_{n \neq 0} \tau_{s-1/2}(|n|) \sqrt{y} K_{s-1/2}(2\pi|n|y) e(nx).$$

Here for complex α , $\tau_\alpha(n) = \sum_{ab=n} (a/b)^\alpha$ is the generalized divisor sum and the scattering function $\varphi(s)$ can be explicitly expressed as

$$\varphi(s) = \frac{\xi(2s-1)}{\xi(2s)}, \quad \text{where} \quad \xi(s) = \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \zeta(s).$$

We will denote with $e(y, s) := y^s + \varphi(s)y^{1-s}$ the constant term of the Eisenstein series.

We denote with $\mathcal{D} = \{z \in \mathbb{H} \mid |z| \geq 1, |x| \leq \frac{1}{2}\}$ the standard fundamental domain for $\Gamma \backslash \mathbb{H}$ and recall that its volume with respect to $d\mu$ is $\text{vol}(X) = \text{vol}(\mathcal{D}) = \frac{\pi}{3}$.

For a parameter $A > 1$ we denote with $\mathcal{D}_A := \{z \in \mathcal{D} \mid \text{Im}(z) \leq A\}$ the corresponding truncated domain and with $\mathcal{C}_A = \mathcal{D} - \mathcal{D}_A$ the corresponding cuspidal region.

The truncated Eisenstein series

$$E_A(z, s) = \begin{cases} E(z, s), & z \in \mathcal{D}_A \\ E(z, s) - e(y, s), & z \in \mathcal{C}_A \end{cases}$$

is now rapidly decreasing in the cusp. Calculation of the L^2 -norm of this truncated Eisenstein series is done in [Sp], Section 2.3, both in the case of the whole fundamental domain and in the case of the cuspidal region, as follows:

$$\begin{aligned} \int_{\mathcal{D}} |E_A(z, \tfrac{1}{2} + iT)|^2 d\mu z &= -\frac{\varphi'}{\varphi}(\tfrac{1}{2} + iT) + 2 \log A + \frac{A^{2iT} \varphi(\tfrac{1}{2} - iT) - A^{-2iT} \varphi(\tfrac{1}{2} + iT)}{2iT} \\ &= 2 \log T + 2 \log A + O((\log T)^{2/3+\epsilon}) \\ \int_{\mathcal{C}_A} |E_A(z, \tfrac{1}{2} + iT)|^2 d\mu z &= \frac{6}{A\pi} \log T + O(A^{-1}(\log T)^{2/3+\epsilon}) \\ &\quad + O(A^{-1} \log A), \quad T \rightarrow \infty. \end{aligned} \tag{2.1}$$

Therefore in the compact truncated domain \mathcal{D}_A , since $E = E_A$ on \mathcal{D}_A , we have as $T \rightarrow \infty$

$$\int_{\mathcal{D}_A} |E(z, \tfrac{1}{2} + iT)|^2 d\mu z \sim \left(\frac{\pi}{3} - \frac{1}{A} \right) \frac{6 \log T}{\pi} = \text{vol}(\mathcal{D}_A) \frac{2 \log T}{\text{vol}(X)},$$

as long as $1 < A \ll \log T$ say. In other words, if we normalize the Eisenstein series as

$$\tilde{E}(z, \tfrac{1}{2} + iT) := \frac{E(z, \tfrac{1}{2} + iT)}{\sqrt{2 \log T}},$$

we have

$$\lim_{T \rightarrow \infty} \frac{1}{\text{vol}(\mathcal{D}_A)} \int_{\mathcal{D}_A} |\tilde{E}(z, \tfrac{1}{2} + iT)|^2 d\mu z = \frac{1}{\text{vol}(X)}. \tag{2.2}$$

If we denote $e^{i\theta(T)} := \frac{\xi(1+2iT)}{|\xi(1+2iT)|}$, then the function $e^{i\theta(T)} \tilde{E}(z, \tfrac{1}{2} + iT)$ is real-valued, and the Random Wave Conjecture, as extended in [HR], predicts that $e^{i\theta(T)} \tilde{E}(z, \tfrac{1}{2} + iT)$ tends to Gaussian $\mathcal{N}(0, \text{vol}(X)^{-1/2})$ in distribution, when restricted to any compact and sufficiently regular subset $\Omega \subset X$. In particular, for the fourth moment ($c_4 = 3$), the conjecture predicts

$$\lim_{T \rightarrow \infty} \frac{1}{\text{vol}(\Omega)} \int_{\Omega} |\tilde{E}(z, \tfrac{1}{2} + iT)|^4 d\mu z = \frac{3}{\text{vol}(X)^2}.$$

By heuristic considerations and numerical experiments in [HR], the same limits should hold also for the normalized truncated Eisenstein series

$$\tilde{E}_A(z, \frac{1}{2} + iT) := \frac{E_A(z, \frac{1}{2} + iT)}{\sqrt{2 \log T}}.$$

As explained, this should also include the case $\Omega = X$, in which case the conjecture is

$$\int_X |E_A(z, \frac{1}{2} + iT)|^4 d\mu z \sim \frac{36}{\pi} \log^2 T, \quad \text{as } T \rightarrow \infty. \quad (2.3)$$

3. Regularized inner product and regularized Plancherel formula

We will make use of the regularization process given by Zagier in [Za]. An adelic version with a representation theoretic interpretation and with an alternate way of defining regularization is recently given in [MV].

Let $F(z)$ be a continuous Γ -invariant function on \mathbb{H} . It is called *renormalizable* (in Zagier's terminology, or *of controlled increase* in the terminology of [MV]) if there is a function $\Phi(y)$ on $\mathbb{R}_{>0}$ of the form

$$\Phi(y) = \sum_{j=1}^l \frac{c_j}{n_j!} y^{\alpha_j} \log^{n_j} y, \quad (3.1)$$

with $c_j, \alpha_j \in \mathbb{C}$ and $n_j \in \mathbb{Z}_{\geq 0}$, such that

$$F(z) = \Phi(y) + O(y^{-N})$$

as $y \rightarrow \infty$, and for any $N > 0$.

If $F(z) = \sum_{n=-\infty}^{\infty} a_n(y) e(nx)$ is the Fourier expansion of F at the cusp ∞ , in particular if $a_0(y)$ is its 0-term, and if no α_j equals 0 or 1, then the function

$$R(F, s) := \int_0^{\infty} (a_0(y) - \Phi(y)) y^{s-2} dy,$$

where the defining integral converges for sufficiently large $\text{Re}(s)$, can be meromorphically continued to all s and has a simple pole at $s = 1$. Then one can define the regularized integral with

$$\int_{\Gamma \backslash \mathbb{H}}^{reg} F(z) d\mu(z) := \frac{\pi}{3} \text{Res}_{s=1} R(F, s). \quad (3.2)$$

Moreover, then the function $F(z)E(z, s)$ with $s \neq 0, 1$ is also renormalizable and in particular, it can be shown that

$$\int_{\Gamma \backslash \mathbb{H}}^{reg} F(z)E(z, s)d\mu(z) = R(F, s).$$

It can be shown (see [Za]) that the regularized integral can be written also as

$$\int_{\Gamma \backslash \mathbb{H}}^{reg} F(z)d\mu(z) = \int_{\mathcal{D}_A} F(z)d\mu(z) + \int_{\mathcal{C}_A} (F(z) - \Phi(y))d\mu(z) - \hat{\Phi}(A), \quad (3.3)$$

where the right-hand side is independent of the value of the parameter $A > 1$ and $\hat{\Phi}(y)$ is in the case $\alpha_j \neq 1$ for all j , given by the following explicit expression

$$\hat{\Phi}(y) = \sum_{j=1}^l c_j \frac{y^{\alpha_j-1}}{\alpha_j-1} \sum_{m=0}^{n_j} \frac{\log^m y}{m!(1-\alpha_j)^{n_j-m}}.$$

Under the assumption that no $\alpha_j = 1$, let $\mathcal{E}_\Phi(z)$ denote a linear combination of Eisenstein series $E(z, \alpha_j)$ (or suitable derivatives thereof) corresponding to all the exponents in (3.1) with $\text{Re}(\alpha_j) > 1/2$, i.e. such that $F(z) - \mathcal{E}_\Phi(z) = O(y^{1/2})$. Then the third, equivalent definition of regularization is given by

$$\int_{\Gamma \backslash \mathbb{H}}^{reg} F(z)d\mu(z) = \int_{\Gamma \backslash \mathbb{H}} (F(z) - \mathcal{E}_\Phi(z))d\mu(z). \quad (3.4)$$

For example Zagier showed in [Za] that for $s_1, s_2 \in \mathbb{C} \setminus \{0, 1\}$, $s_1 \neq s_2, 1 - s_2$, we have

$$\int_{\Gamma \backslash \mathbb{H}}^{reg} E(z, s_1)E(z, s_2)d\mu(z) = 0. \quad (3.5)$$

On the other hand, for the regularized product of the three Eisenstein series, Zagier [Za, p. 431] obtained

$$\begin{aligned} \int_{\Gamma \backslash \mathbb{H}}^{reg} E(z, \frac{1}{2} + s_1)E(z, \frac{1}{2} + s_2)E(z, \frac{1}{2} + s_3)d\mu(z) &= \\ &= \frac{\xi(\frac{1}{2} + s_1 + s_2 + s_3)\xi(\frac{1}{2} + s_1 - s_2 + s_3)\xi(\frac{1}{2} + s_1 + s_2 - s_3)\xi(\frac{1}{2} + s_1 - s_2 - s_3)}{\xi(1 + 2s_1)\xi(1 + 2s_2)\xi(1 + 2s_3)}. \end{aligned} \quad (3.6)$$

The right-hand side is of course symmetric in s_1, s_2, s_3 because of the functional equation $\xi(1-s) = \xi(s)$.

Since we are interested in the regularized product of 4 Eisenstein series, one can try to apply the definition (3.2) directly. But already Zagier in [Za, p. 431] discussed that in this case there is no useful closed-form expression for the result, as is for the product of 3 Eisenstein series in (3.6). Therefore, we must proceed indirectly via a regularized Plancherel formula.

Now, let $G(z)$ be another renormalizable Γ -invariant function such that $G(z) = \Psi(y) + O(y^{-N})$ as $y \rightarrow \infty$ for any $N > 0$, where $\Psi(y) = \sum_{k=1}^{l_1} \frac{d_k}{m_k!} y^{\beta_k} \log^{m_k} y$ with $d_k, \beta_k \in \mathbb{C}$. Then the product $F(z)\overline{G(z)}$ is also a renormalizable Γ -invariant function and if $\alpha_j + \overline{\beta_k} \neq 1$, for all α_j and β_k appearing in Φ and Ψ respectively, the regularized inner product of F and G can be defined as

$$\langle F, G \rangle_{reg} := \int_{\Gamma \backslash \mathbb{H}}^{reg} F(z) \overline{G(z)} d\mu(z) = \int_{\Gamma \backslash \mathbb{H}} (F(z) \overline{G(z)} - \mathcal{E}_{\Phi \Psi}(z)) d\mu(z).$$

It is easy to see from (3.3) that this regularized product is a Hermitian form.

The regularized Plancherel formula from [MV] is much more general, but for our purposes we will state and derive it entirely in classical situation of Zagier's paper [Za], much in the spirit of Lemma 4.1 from [Yo]. Because of the cumbersome formulas, we will use the shorthand notation $E_s(z) := E(z, \frac{1}{2} + s)$, and remind the reader not to confuse this with the truncated Eisenstein series $E_A(z, s)$ which still have 2 arguments.

Proposition 3.1 ([MV]). *Let $F(z)$ and $G(z)$ be renormalizable functions on $\Gamma \backslash \mathbb{H}$ such that $F - \Phi$ and $G - \Psi$ are of rapid decay as $y \rightarrow \infty$, for some $\Phi(y) = \sum_{j=1}^l \frac{c_j}{n_j!} y^{\alpha_j} \log^{n_j} y$ and $\Psi(y) = \sum_{k=1}^{l_1} \frac{d_k}{m_k!} y^{\beta_k} \log^{m_k} y$. Moreover, let $\alpha_j \neq 1$, $\beta_k \neq 1$, $\operatorname{Re}(\alpha_j) \neq \frac{1}{2}$, $\operatorname{Re}(\beta_k) \neq \frac{1}{2}$, $\alpha_j + \overline{\beta_k} \neq 1$ and $\alpha_j \neq \overline{\beta_k}$, for all j, k . Then the following formula holds:*

$$\begin{aligned} \langle F(z), G(z) \rangle_{reg} = & \langle F, \sqrt{3/\pi} \rangle_{reg} \langle \sqrt{3/\pi}, G \rangle_{reg} + \sum_j \langle F, u_j \rangle \langle u_j, G \rangle + \frac{1}{4\pi} \int_{-\infty}^{\infty} \langle F, E_{it} \rangle_{reg} \langle E_{it}, G \rangle_{reg} dt \\ & + \langle F, \mathcal{E}_{\Psi} \rangle_{reg} + \langle \mathcal{E}_{\Phi}, G \rangle_{reg}. \end{aligned}$$

Proof. Because of the assumption $\operatorname{Re}(\alpha_j) \neq \frac{1}{2}$, $\operatorname{Re}(\beta_k) \neq \frac{1}{2}$, there exists some $\delta > 0$ such that $F_1(z) := F(z) - \mathcal{E}_{\Phi}(z) = O(y^{1/2-\delta})$ and $G_1(z) := G(z) - \mathcal{E}_{\Psi}(z) = O(y^{1/2-\delta})$. For these $F_1(z)$ and $G_1(z)$ we have $F_1(z) \overline{G_1(z)} \in L^1(\Gamma \backslash \mathbb{H})$ and hence $\langle F_1, G_1 \rangle_{reg} = \langle F_1, G_1 \rangle$ (the usual Petersson inner product), while also $F_1(z), G_1(z) \in L^2(\Gamma \backslash \mathbb{H})$ and hence one can apply the usual Plancherel formula for $\langle F_1(z), G_1(z) \rangle$, obtaining

$$\begin{aligned} \langle F(z), G(z) \rangle_{reg} &= \langle F_1(z) + \mathcal{E}_\Phi(z), G_1(z) + \mathcal{E}_\Psi(z) \rangle_{reg} = \\ &= \langle F_1, \sqrt{3/\pi} \rangle \langle \sqrt{3/\pi}, G_1 \rangle + \sum_j \langle F_1, u_j \rangle \langle u_j, G_1 \rangle + \frac{1}{4\pi} \int_{-\infty}^{\infty} \langle F_1, E_{it} \rangle \langle E_{it}, G_1 \rangle dt \\ &\quad + \langle F_1, \mathcal{E}_\Psi \rangle_{reg} + \langle \mathcal{E}_\Phi, G_1 \rangle_{reg} + \langle \mathcal{E}_\Phi, \mathcal{E}_\Psi \rangle_{reg}. \end{aligned}$$

Under our restrictions on the parameters α_j, β_k , all the inner products on the right hand side are well-defined and moreover because of (3.5), we have that $\langle \mathcal{E}_\Phi, \mathcal{E}_\Psi \rangle_{reg} = 0$, which also implies that $\langle F_1, \mathcal{E}_\Psi \rangle_{reg} = \langle F, \mathcal{E}_\Psi \rangle_{reg}$ and $\langle \mathcal{E}_\Phi, G_1 \rangle_{reg} = \langle \mathcal{E}_\Phi, G \rangle_{reg}$. Furthermore, for the products with cusp forms u_j we have $\langle F_1, u_j \rangle = \langle F, u_j \rangle - \langle \mathcal{E}_\Phi, u_j \rangle = \langle F, u_j \rangle$, the product with constant function is $\langle F_1, 1 \rangle = \langle F - \mathcal{E}_\Phi, 1 \rangle = \langle F, 1 \rangle_{reg}$ by definition of regularization and $\langle F_1, E_{it} \rangle = \langle F - \mathcal{E}_\Phi, E_{it} \rangle_{reg} = \langle F, E_{it} \rangle_{reg}$ since $\langle \mathcal{E}_\Phi, E_{it} \rangle_{reg} = 0$, by (3.5). This finishes the proof. \square

4. Proof of Theorem 1.1

We want to apply this formula for the product of four Eisenstein series. By calculating the constant term of $F(z) := E(z, \frac{1}{2} + s_1)E(z, \frac{1}{2} + s_2)$, we find that

$$\mathcal{E}_\Phi(z) = E(z, 1 + s_1 + s_2) + c_1 E(z, 1 - s_1 + s_2) + c_2 E(z, 1 + s_1 - s_2) + c_1 c_2 E(z, 1 - s_1 - s_2)$$

where

$$c_j = \varphi\left(\frac{1}{2} + s_j\right) = \frac{\xi(2s_j)}{\xi(1 + 2s_j)},$$

and we have the similar formula for $\mathcal{E}_\Psi(z)$ corresponding to $G(z) := E(z, \frac{1}{2} + s_3)E(z, \frac{1}{2} + s_4)$. Hence, under the conditions on the parameters s_j , $1 \leq j \leq 4$, described in Proposition 3.1 ($\alpha_j = 1 \pm s_1 \pm s_2$, $\beta_k = 1 \pm s_3 \pm s_4$), we get

$$\begin{aligned} \langle E_{s_1} E_{s_2}, E_{s_3} E_{s_4} \rangle_{reg} &= \tag{4.1} \\ &= \left\langle E(z, \frac{1}{2} + s_1)E(z, \frac{1}{2} + s_2), E(z, \frac{1}{2} + s_3)E(z, \frac{1}{2} + s_4) \right\rangle_{reg} = \\ &= \frac{3}{\pi} \langle E_{s_1} E_{s_2}, 1 \rangle_{reg} \langle 1, E_{s_3} E_{s_4} \rangle_{reg} + \sum_{j \geq 1} \langle E_{s_1} E_{s_2}, u_j \rangle \langle u_j, E_{s_3} E_{s_4} \rangle \\ &\quad + \frac{1}{4\pi} \int_{-\infty}^{\infty} \langle E_{s_1} E_{s_2}, E_{it} \rangle_{reg} \langle E_{it}, E_{s_3} E_{s_4} \rangle_{reg} dt \\ &\quad + \langle E_{s_1} E_{s_2}, E_{\frac{1}{2} + s_3 + s_4} + c_3 E_{\frac{1}{2} - s_3 + s_4} + c_4 E_{\frac{1}{2} + s_3 - s_4} + c_3 c_4 E_{\frac{1}{2} - s_3 - s_4} \rangle_{reg} \\ &\quad + \langle E_{\frac{1}{2} + s_1 + s_2} + c_1 E_{\frac{1}{2} - s_1 + s_2} + c_2 E_{\frac{1}{2} + s_1 - s_2} + c_1 c_2 E_{\frac{1}{2} - s_1 - s_2}, E_{s_3} E_{s_4} \rangle_{reg}. \end{aligned}$$

For the cusp forms u_j , the triple products $\langle E_{s_1} E_{s_2}, u_j \rangle$ can be evaluated by the standard unfolding argument (see Section 2 of [LS]):

Lemma 4.1. *Let $u_j(z)$ be a Hecke–Maass cusp form for the group Γ , that is an eigenvalue of the Laplace operator $\Delta u_j = (\frac{1}{4} + t_j^2)u_j$ and of all Hecke operators $T_n u_j = \lambda_j(n)u_j$, for all $n \geq 1$, which satisfies also $T_{-1} u_j(z) = \overline{u_j(-\bar{z})} = \epsilon_j u_j(z)$, $\epsilon_j = \pm 1$. Then it has the Fourier expansion $u_j(z) = \rho_j(1) \sum_{n \neq 0} \lambda_j(n) \sqrt{y} K_{it_j}(2\pi|n|y) e(nx)$ with $\lambda_j(-n) = \epsilon_j \lambda_j(n)$. Let $L(s, u_j)$ be the L -function associated to u_j , defined by analytic continuation from the Dirichlet series $\sum_{n \geq 1} \lambda_j(n) n^{-s}$. Then for $s_1, s_2 \neq \pm 1/2$, if u_j is even ($\epsilon_j = 1$) we have*

$$\langle E(\cdot, 1/2 + s_1) E(\cdot, 1/2 + s_2), u_j \rangle = \frac{\overline{\rho_j(1)}}{2} \frac{\Lambda(\frac{1}{2} + s_1 + s_2, u_j) \Lambda(\frac{1}{2} + s_1 - s_2, u_j)}{\xi(1 + 2s_1) \xi(1 + 2s_2)}, \quad (4.2)$$

where $\Lambda(s, u_j) := \pi^{-s} \Gamma(\frac{s+it_j}{2}) \Gamma(\frac{s-it_j}{2}) L(s, u_j)$ is the completed L -function corresponding to u_j . In the case of odd u_j (i.e. $\epsilon_j = -1$), the triple product is 0.

Remark. The right hand side in (4.2) is symmetric in s_1, s_2 , since for even u_j , we have the functional equation $\Lambda(s, u_j) = \Lambda(1 - s, u_j)$. Moreover, we have the following formula relating the normalizing factor $\rho_j(1)$ with the symmetric square L -function:

$$|\rho_j(1)|^2 = \frac{2 \cosh(\pi t_j)}{L(1, \text{sym}^2 u_j)}.$$

Further, for $s_1 \neq \pm s_2$ and $s_3 \neq \pm s_4$, by (3.5) the first term on the right-hand side in (4.1) vanishes. Finally, using (3.6) for all regularized triple products of Eisenstein series, we arrive at

$$\begin{aligned} & \langle E_{s_1} E_{s_2}, E_{s_3} E_{s_4} \rangle_{reg} \quad (4.3) \\ &= \sum_{\substack{j \geq 1 \\ \epsilon_j = 1}} \frac{\cosh(\pi t_j)}{2} \frac{\Lambda(\frac{1}{2} + s_1 + s_2, u_j) \Lambda(\frac{1}{2} + s_1 - s_2, u_j) \Lambda(\frac{1}{2} + \bar{s}_3 + \bar{s}_4, u_j) \Lambda(\frac{1}{2} + \bar{s}_3 - \bar{s}_4, u_j)}{L(1, \text{sym}^2 u_j) \xi(1 + 2s_1) \xi(1 + 2s_2) \xi(1 + 2\bar{s}_3) \xi(1 + 2\bar{s}_4)} \\ &+ \frac{1}{4\pi} \int_{-\infty}^{\infty} \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(\frac{1}{2} + ti + \delta_1 s_1 + \delta_2 s_2) \xi(\frac{1}{2} + ti + \delta_1 \bar{s}_3 + \delta_2 \bar{s}_4)}{|\xi(1 + 2ti)|^2 \xi(1 + 2s_1) \xi(1 + 2s_2) \xi(1 + 2\bar{s}_3) \xi(1 + 2\bar{s}_4)} dt \\ &+ \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 + \delta_1 s_1 + \delta_2 s_2 + \bar{s}_3 + \bar{s}_4)}{\xi(1 + 2s_1) \xi(1 + 2s_2) \xi(2 + 2\bar{s}_3 + 2\bar{s}_4)} \\ &+ \bar{c}_3 \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 + \delta_1 s_1 + \delta_2 s_2 - \bar{s}_3 + \bar{s}_4)}{\xi(1 + 2s_1) \xi(1 + 2s_2) \xi(2 - 2\bar{s}_3 + 2\bar{s}_4)} \\ &+ \bar{c}_4 \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 + \delta_1 s_1 + \delta_2 s_2 + \bar{s}_3 - \bar{s}_4)}{\xi(1 + 2s_1) \xi(1 + 2s_2) \xi(2 + 2\bar{s}_3 - 2\bar{s}_4)} \end{aligned}$$

$$\begin{aligned}
& + \bar{c}_3 \bar{c}_4 \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 + \delta_1 s_1 + \delta_2 s_2 - \bar{s}_3 - \bar{s}_4)}{\xi(1 + 2s_1) \xi(1 + 2s_2) \xi(2 - 2\bar{s}_3 - 2\bar{s}_4)} \\
& + \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 + s_1 + s_2 + \delta_1 \bar{s}_3 + \delta_2 \bar{s}_4)}{\xi(2 + 2s_1 + 2s_2) \xi(1 + 2\bar{s}_3) \xi(1 + 2\bar{s}_4)} \\
& + c_1 \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 - s_1 + s_2 + \delta_1 \bar{s}_3 + \delta_2 \bar{s}_4)}{\xi(2 - 2s_1 + 2s_2) \xi(1 + 2\bar{s}_3) \xi(1 + 2\bar{s}_4)} \\
& + c_2 \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 + s_1 - s_2 + \delta_1 \bar{s}_3 + \delta_2 \bar{s}_4)}{\xi(2 + 2s_1 - 2s_2) \xi(1 + 2\bar{s}_3) \xi(1 + 2\bar{s}_4)} \\
& + c_1 c_2 \frac{\prod_{\delta_1, \delta_2 \in \{\pm 1\}} \xi(1 - s_1 - s_2 + \delta_1 \bar{s}_3 + \delta_2 \bar{s}_4)}{\xi(2 - 2s_1 - 2s_2) \xi(1 + 2\bar{s}_3) \xi(1 + 2\bar{s}_4)}.
\end{aligned}$$

Let us denote the last eight terms (quotients of products of ξ -functions, coming from the regularization process) on the right hand side of (4.3) with Ξ_j , $1 \leq j \leq 8$, respectively with the order of appearance in (4.3).

Now, let us choose for s_j the following values: $s_1 = iT$, $s_2 = iT + \nu$, $s_3 = iT$ and $s_4 = iT + \bar{\eta}$, with complex parameters ν and η satisfying $0 < \operatorname{Re}(\nu) < \operatorname{Re}(\eta) < \frac{1}{4}$. For these values all the conditions from Proposition 3.1 are satisfied. From (3.4), we see that $\langle E_{iT} E_{iT+\nu}, E_{iT} E_{iT+\bar{\eta}} \rangle_{reg}$ is continuous in ν, η , and therefore, if we first let $\nu \rightarrow 0$ in (4.3), keeping η fixed, we get

$$\begin{aligned}
& \langle E_{iT}^2, E_{iT} E_{iT+\bar{\eta}} \rangle_{reg} \\
& = \sum_{\substack{j \geq 1 \\ \epsilon_j = 1}} \frac{\cosh(\pi t_j)}{2} \frac{\Lambda(\frac{1}{2} + 2Ti, u_j) \Lambda(\frac{1}{2} - 2Ti + \eta, u_j) \Lambda(\frac{1}{2}, u_j) \Lambda(\frac{1}{2} - \eta, u_j)}{L(1, \operatorname{sym}^2 u_j) \xi^2(1 + 2Ti) \xi(1 - 2Ti) \xi(1 - 2Ti + 2\eta)} \\
& \quad + \frac{1}{4\pi} \int_{-\infty}^{\infty} \frac{\xi^2(\frac{1}{2} + ti) \prod_{\pm} \xi(\frac{1}{2} + ti \pm 2Ti) \xi(\frac{1}{2} + ti \pm 2Ti \mp \eta) \xi(\frac{1}{2} + ti \pm \eta)}{|\xi(1 + 2ti)|^2 \xi^2(1 + 2Ti) \xi(1 - 2Ti) \xi(1 - 2Ti + 2\eta)} dt \\
& \quad + \sum_{j=1}^8 \Xi_j(T, \eta),
\end{aligned} \tag{4.4}$$

where

$$\begin{aligned}
\Xi_1(T, \eta) &= \frac{\xi(1 + \eta) \xi^2(1 - 2Ti + \eta) \xi(1 - 4Ti + \eta)}{\xi^2(1 + 2Ti) \xi(2 - 4Ti + 2\eta)} =: \xi(1 + \eta) F_1(\eta), \\
\Xi_2(T, \eta) &= \frac{\xi^2(1 + \eta) \xi(1 + 2Ti + \eta) \xi(1 - 2Ti + \eta)}{\xi(1 - 2Ti) \xi(1 + 2Ti) \xi(2 + 2\eta)} =: \xi^2(1 + \eta) F_2(\eta), \\
\Xi_3(T, \eta) &= \frac{\xi^2(1 - \eta) \xi(1 + 2Ti - \eta) \xi(1 + 2Ti - 2\eta) \xi(1 - 2Ti - \eta)}{\xi^2(1 + 2Ti) \xi(1 - 2Ti + 2\eta) \xi(2 - 2\eta)} =: \xi^2(1 - \eta) F_3(\eta), \\
\Xi_4(T, \eta) &= \frac{\xi(1 - \eta) \xi^2(1 + 2Ti - \eta) \xi(1 + 2Ti - 2\eta) \xi(1 + 4Ti - \eta)}{\xi(1 + 2Ti) \xi(1 - 2Ti) \xi(1 - 2Ti + 2\eta) \xi(2 + 4Ti - 2\eta)} =: \xi(1 - \eta) F_4(\eta),
\end{aligned}$$

$$\begin{aligned}
\Xi_5(T, \eta) &= \frac{\xi(1+\eta)\xi(1+2Ti-\eta)\xi(1+2Ti+\eta)\xi(1+4Ti-\eta)}{\xi(2+4Ti)\xi(1-2Ti)\xi(1-2Ti+2\eta)} \\
&=: \xi(1+\eta)F_5(\eta), \\
\Xi_6(T, \eta) &= \Xi_7(T, \eta) = \frac{\xi(1-\eta)\xi(1+\eta)\xi(1+2Ti-\eta)\xi(1-2Ti+\eta)}{\xi(2)\xi(1+2Ti)\xi(1-2Ti+2\eta)} \\
&=: \xi(1-\eta)\xi(1+\eta)F_6(\eta), \\
\Xi_8(T, \eta) &= \frac{\xi(1-\eta)\xi(1-2Ti-\eta)\xi(1-2Ti)\xi(1-2Ti+\eta)\xi(1-4Ti+\eta)}{\xi^2(1+2Ti)\xi(1-2Ti+2\eta)\xi(2-4Ti)} \\
&=: \xi(1-\eta)F_8(\eta).
\end{aligned}$$

Each of Ξ_j has a pole at $\eta = 0$, but the whole sum $\sum_{j=1}^8 \Xi_j$ has a removable singularity at $\eta = 0$. This can be seen by grouping together Ξ_1 with Ξ_8 , Ξ_4 with Ξ_5 , and $\Xi_2 + \Xi_3$ with $2\Xi_6 = \Xi_6 + \Xi_7$. More explicitly, if we denote with

$$\xi(s) = \frac{1}{s-1} + a + b(s-1) + O((s-1)^2) \quad (4.5)$$

the Laurent expansion of $\xi(s)$ around $s = 1$, we get the following expansions of $\Xi_j(\eta)$:

$$\begin{aligned}
\Xi_1(\eta) &= \frac{F_1(0)}{\eta} + F'_1(0) + aF_1(0) + O(\eta), \\
\Xi_2(\eta) &= \frac{F_2(0)}{\eta^2} + \frac{F'_2(0) + 2aF_2(0)}{\eta} + (a^2 + 2b)F_2(0) + 2aF'_2(0) + \frac{1}{2}F''_2(0) + O(\eta), \\
\Xi_3(\eta) &= \frac{F_3(0)}{\eta^2} + \frac{F'_3(0) - 2aF_3(0)}{\eta} + (a^2 + 2b)F_3(0) - 2aF'_3(0) + \frac{1}{2}F''_3(0) + O(\eta), \\
\Xi_4(\eta) &= -\frac{F_4(0)}{\eta} - F'_4(0) + aF_4(0) + O(\eta), \\
\Xi_5(\eta) &= \frac{F_5(0)}{\eta} + F'_5(0) + aF_5(0) + O(\eta), \\
\Xi_6(\eta) = \Xi_7(\eta) &= -\frac{F_6(0)}{\eta^2} - \frac{F'_6(0)}{\eta} + (a^2 - 2b)F_6(0) - \frac{1}{2}F''_6(0) + O(\eta), \\
\Xi_8(\eta) &= -\frac{F_8(0)}{\eta} - F'_8(0) + aF_8(0) + O(\eta).
\end{aligned}$$

But, $F_2(0) = F_3(0) = F_6(0) = \frac{1}{\xi(2)}$, so the polar terms with $\frac{1}{\eta^2}$ cancel out in the sum. Further, $F_1(0) = F_8(0)$, $F_4(0) = F_5(0)$ and one calculates

$$F'_2(0) = \frac{2}{\xi(2)} \left[\operatorname{Re} \frac{\xi'}{\xi}(1+2Ti) - \frac{\xi'}{\xi}(2) \right], \quad F'_3(0) = \frac{2}{\xi(2)} \left[\frac{\xi'}{\xi}(2) - 3\operatorname{Re} \frac{\xi'}{\xi}(1+2Ti) \right]$$

and

$$F'_6(0) = -\frac{2}{\xi(2)} \operatorname{Re} \frac{\xi'}{\xi}(1+2Ti),$$

from which it follows that $F'_2(0) + F'_3(0) - 2F'_6(0) = 0$, a.e. the coefficient in front of $\frac{1}{\eta}$ also vanishes. Therefore we can take $\eta \rightarrow 0$ in (4.4) and after calculation of all the other required derivatives appearing in

$$\lim_{\eta \rightarrow 0} \sum_{j=1}^8 \Xi_j(T, \eta) = a(F_1(0) + F_4(0) + F_5(0) + F_8(0)) + a^2(F_2(0) + F_3(0) + 2F_6(0)) \\ + F'_1(0) - F'_4(0) + F'_5(0) - F'_8(0) + 2a(F'_2(0) - F'_3(0)) + \frac{1}{2}F''_2(0) + \frac{1}{2}F''_3(0) - F''_6(0),$$

we obtain the following exact evaluation of the regularized fourth power of Eisenstein series:

Proposition 4.2. *For any nonzero real T , we have:*

$$\int_{\Gamma \backslash \mathbb{H}}^{reg} |E(z, 1/2 + iT)|^4 d\mu(z) = \langle E_{iT}^2, E_{iT}^2 \rangle_{reg} \quad (4.6) \\ = \sum_{\substack{j \geq 1 \\ \epsilon_j = 1}} \frac{\cosh(\pi t_j)}{2} \frac{\Lambda(\frac{1}{2} + 2Ti, u_j) \Lambda(\frac{1}{2} - 2Ti, u_j) \Lambda^2(\frac{1}{2}, u_j)}{L(1, \text{sym}^2 u_j) |\xi(1 + 2Ti)|^4} \\ + \frac{1}{4\pi} \int_{-\infty}^{\infty} \frac{\xi^4(\frac{1}{2} + ti) \xi^2(\frac{1}{2} + ti + 2Ti) \xi^2(\frac{1}{2} + ti - 2Ti)}{|\xi(1 + 2ti)|^2 |\xi(1 + 2Ti)|^4} dt \\ + \frac{4}{\xi(2)} \left[\text{Re} \frac{\xi''}{\xi} (1 + 2Ti) + 2 \left| \frac{\xi'}{\xi} (1 + 2Ti) \right|^2 + \text{Re} \frac{(\xi')^2}{\xi^2} (1 + 2Ti) \right. \\ \left. + 4(a - \frac{\xi'}{\xi}(2)) \text{Re} \frac{\xi'}{\xi} (1 + 2Ti) + 2 \frac{(\xi')^2}{\xi^2} (2) - \frac{\xi''}{\xi} (2) - 2a \frac{\xi'}{\xi} (2) + a^2 \right] \\ + \frac{\xi^2(1 + 2Ti) \xi(1 + 4Ti)}{\xi^2(1 - 2Ti) \xi(2 + 4Ti)} \left[2a + 4 \frac{\xi'}{\xi} (1 + 2Ti) - 2 \frac{\xi'}{\xi} (2 + 4Ti) \right] \\ + \frac{\xi^2(1 - 2Ti) \xi(1 - 4Ti)}{\xi^2(1 + 2Ti) \xi(2 - 4Ti)} \left[2a + 4 \frac{\xi'}{\xi} (1 - 2Ti) - 2 \frac{\xi'}{\xi} (2 - 4Ti) \right],$$

where $a = \lim_{s \rightarrow 1} (\xi(s) - (s - 1)^{-1})$.

Remark. The exact value of the constant a is $\frac{C_0}{2} - \frac{\ln \pi}{2} - \ln 2 = -0.9769 \dots$, where $C_0 = 0.57721 \dots$ is Euler's constant. This is a consequence of the following two formulas: $\zeta(s) = \frac{1}{s-1} + C_0 + O((s-1))$ and $\frac{\Gamma'}{\Gamma}(1/2) = -C_0 - 2 \ln 2$. Also we recall that $\xi(2) = \frac{\pi}{6}$.

Using Stirling's approximations $|\Gamma(\sigma + it)| = e^{-\pi|t|/2} |t|^{\sigma - \frac{1}{2}} \sqrt{2\pi} \{1 + O(|t|^{-1})\}$ and $\frac{\Gamma'}{\Gamma}(s) = \log s + O(|s|^{-1})$ valid in a fixed vertical strip when $|t| \rightarrow \infty$, and classical estimates for the Riemann zeta-function on the edge of the critical strip [MoV, section 6.3],

$$(\log t)^{-2/3}(\log \log t)^{-1/3} \ll \zeta(s) \ll (\log t)^{2/3} \quad \text{and} \\ \frac{\zeta'}{\zeta}(s) \ll (\log t)^{2/3}(\log \log t)^{1/3} \quad (4.7)$$

for

$$s = \sigma + it, \quad 1 - \sigma \ll (\log t)^{-2/3}(\log \log t)^{-1/3},$$

one obtains first that $\frac{\xi'}{\xi}(1 \pm 2Ti) \ll \log T$ and then that the contribution of the terms in the last two lines in the formula (4.6) is $O(\frac{\log^2 T}{T^{1/2}})$. Therefore the contribution on the right-hand side of (4.6) coming from the regularization process is

$$\frac{24}{\pi} \left[\operatorname{Re} \frac{\xi''}{\xi}(1 + 2Ti) + 2 \left| \frac{\xi'}{\xi}(1 + 2Ti) \right|^2 + \operatorname{Re} \frac{(\xi')^2}{\xi^2}(1 + 2Ti) \right] + O(\log T).$$

Since $\frac{\xi'}{\xi}(s) = -\frac{\log \pi}{2} + \frac{1}{2} \frac{\Gamma'}{\Gamma}(\frac{s}{2}) + \frac{\zeta'}{\zeta}(s)$, from Stirling's approximation and (4.7) we obtain further that when $T \rightarrow \infty$

$$2 \left| \frac{\xi'}{\xi}(1 + 2Ti) \right|^2 + \operatorname{Re} \frac{(\xi')^2}{\xi^2}(1 + 2Ti) = \frac{3}{4} \log^2 T + O(\log^{5/3+\epsilon} T),$$

for any $\epsilon > 0$.

Lemma 4.3. *As $t \rightarrow \infty$, we have*

$$\frac{\zeta''}{\zeta}(1 + ti) \ll \log^{4/3+\epsilon} t.$$

Proof. We can use (4.7) and the Borel–Carathéodory lemma [MoV, Lemma 6.2] to get the bound

$$\left(\frac{\zeta'}{\zeta} \right)'(1 + ti) \ll \log^{4/3+\epsilon} t.$$

This and (4.7) again imply the stated bound for $\frac{\zeta''}{\zeta}(1 + ti) = (\frac{\zeta'}{\zeta})'(1 + ti) + (\frac{\zeta'}{\zeta}(1 + ti))^2$. \square

From

$$\frac{\xi''}{\xi}(s) = \frac{\log^2 \pi}{4} + \frac{1}{4} \frac{\Gamma''}{\Gamma}(\frac{s}{2}) + \frac{\zeta''}{\zeta}(s) - \frac{\log \pi}{2} \frac{\Gamma'}{\Gamma}(\frac{s}{2}) - (\log \pi) \frac{\zeta'}{\zeta}(s) + \frac{\Gamma'}{\Gamma}(\frac{s}{2}) \frac{\zeta'}{\zeta}(s),$$

using (4.7), Lemma 4.3 and another well-known approximation $\frac{\Gamma''}{\Gamma}(\frac{1}{2} + Ti) = \log^2 T + O(\log T)$, we obtain also the asymptotic

$$\operatorname{Re} \frac{\xi''}{\xi}(1 + 2Ti) = \frac{1}{4} \log^2 T + O(\log^{5/3+\epsilon} T).$$

On the other hand, the contribution of the continuous spectrum in (4.6) i.e. the integral on the right-hand side is of a smaller size, being bounded by the integral in the following Lemma:

Lemma 4.4. *For $T \geq 1$ we have:*

$$\int_{-\infty}^{\infty} \frac{|\xi(\frac{1}{2} + ti)|^4 |\xi(\frac{1}{2} + ti + 2Ti)|^2 |\xi(\frac{1}{2} + ti - 2Ti)|^2}{|\xi(1 + 2ti)|^2 |\xi(1 + 2Ti)|^4} dt \ll T^{-1/6},$$

for some absolute implicit constant.

Proof. This is exactly Proposition 3.4 in [Sp]. For completeness we briefly repeat here the argument. After employing Stirling's asymptotic formula for Gamma functions and after splitting the integral $\int_{-\infty}^{+\infty} = 2 \int_0^{3T} + 2 \int_{3T}^{+\infty}$, one can see easily that the contribution in the range $\int_{3T}^{+\infty}$ decays exponentially with T . Therefore, one needs to bound the integral

$$\int_0^{3T} \frac{|\zeta(\frac{1}{2} + ti)|^4}{(1 + |t|)} \cdot \frac{|\zeta(\frac{1}{2} + (t + 2T)i)|^2}{(1 + |t + 2T|)^{1/2}} \cdot \frac{|\zeta(\frac{1}{2} + (t - 2T)i)|^2}{(1 + |t - 2T|)^{1/2}} \cdot \frac{e^{\frac{\pi}{2}(4T - |t - 2T| - |t + 2T|)}}{|\zeta(1 + 2ti)|^2 |\zeta(1 + 2Ti)|^4} dt.$$

By (4.7) the fourth ratio can be bounded by T^ε , the third ratio can be bounded by convexity bound, while the second ratio can be bounded using the subconvexity bound $\zeta(\frac{1}{2} + ti) \ll (1 + |t|)^{\theta + \varepsilon}$ for some $\theta < \frac{1}{6}$, which is available and sufficient. The bound follows by the fourth moment estimate $\int_0^{3T} \frac{|\zeta(\frac{1}{2} + ti)|^4}{1 + |t|} dt \ll T^\varepsilon$, for any $\varepsilon > 0$. \square

Therefore after putting together everything in this section, we obtain the asymptotic formula in Theorem 1.1. Note that the theorem has dropped the condition $\epsilon_j = 1$. This is fine because when $\epsilon_j = -1$, we have $\Lambda(\frac{1}{2}, u_j) = 0$ and the summand vanishes.

5. Proof of Theorem 1.2

The regularized fourth moment of Eisenstein series $E(z, 1/2 + iT)$

$$\langle E_{iT}^2, E_{iT}^2 \rangle_{reg} = \int_{\Gamma \backslash \mathbb{H}}^{reg} |E(z, \frac{1}{2} + iT)|^4 d\mu(z)$$

can also be expressed directly using (3.3). The corresponding function $\Phi(y)$ is given by

$$\Phi(y) = |e(y, 1/2 + iT)|^4 = c^2 y^{2-4Ti} + 4cy^{2-2Ti} + 6y^2 + 4\bar{c}y^{2+2Ti} + \bar{c}^2 y^{2+4Ti}$$

where $c = \varphi(1/2 + iT) = \frac{\xi(1-2Ti)}{\xi(1+2Ti)}$, so in particular $|c| = 1$. Hence we get

$$\begin{aligned}
& \int_{\Gamma \backslash \mathbb{H}}^{reg} |E(z, \frac{1}{2} + iT)|^4 d\mu(z) = \\
& = \int_{\mathcal{D}_A} |E(z, 1/2 + iT)|^4 d\mu(z) + \int_{\mathcal{C}_A} (|E(z, 1/2 + iT)|^4 - |e(y, 1/2 + iT)|^4) d\mu(z) - \hat{\Phi}(A),
\end{aligned}$$

with $\hat{\Phi}(A)$ given explicitly by

$$\hat{\Phi}(A) = c^2 \frac{A^{1-4Ti}}{1-4Ti} + 4c \frac{A^{1-2Ti}}{1-2Ti} + 6A + 4\bar{c} \frac{A^{1+2Ti}}{1+2Ti} + \bar{c}^2 \frac{A^{1+4Ti}}{1+4Ti}.$$

In particular, $\hat{\Phi}(A) \ll A$. This will be an admissible error for all the values of the truncation parameter in the range $1 < A \ll \log T$.

Hence the difference between the regularized integral of $|E|^4$ and the integral $\int_X |E_A|^4 d\mu$ with the truncated Eisenstein series considered in [Sp] is

$$\begin{aligned}
& \int_{\Gamma \backslash \mathbb{H}}^{reg} |E(z, \frac{1}{2} + iT)|^4 d\mu(z) - \|E_A(\cdot, \frac{1}{2} + iT)\|_4^4 \\
& = \int_{\mathcal{D}_A} |E|^4 d\mu + \int_{\mathcal{C}_A} (|E_A + e|^4 - |e|^4) d\mu - \hat{\Phi}(A) - \int_{\mathcal{D}_A} |E|^4 d\mu - \int_{\mathcal{C}_A} |E_A|^4 d\mu \\
& = \int_{\mathcal{C}_A} (e^2 \bar{E}_A^2 + \bar{e}^2 E_A^2 + 4|e|^2 |E_A|^2) d\mu + 2 \int_{\mathcal{C}_A} (|E_A|^2 \bar{E}_A e + |E_A|^2 E_A \bar{e}) d\mu - \hat{\Phi}(A), \quad (5.1)
\end{aligned}$$

since $\int_{\mathcal{C}_A} E_A \bar{e} |e|^2 d\mu = \int_{\mathcal{C}_A} \bar{E}_A e |e|^2 d\mu = 0$.

Here, the first integral in the cuspidal region can be explicitly computed. From the integral representation

$$K_\nu(y) = \int_0^{+\infty} e^{-y \cosh t} \cosh(\nu t) dt, \quad \operatorname{Re}(\nu) > -\frac{1}{2},$$

we see that $K_{iT}(y)$ is real for $y > 0$, $T \in \mathbb{R}$ and hence for $z \in \mathcal{C}_A$

$$\xi(1 + 2Ti) E_A(z, 1/2 + iT) = 4 \sum_{n=1}^{\infty} \tau_{iT}(n) y^{\frac{1}{2}} K_{iT}(2\pi n y) \cos(2\pi n x) \quad (5.2)$$

is also real-valued. Using this and the functional equation for $\xi(s)$, after a short calculation one gets that the first integral in (5.1) is equal to

$$\int_{\mathcal{C}_A} \left(\frac{12y}{|\xi(1 + 2Ti)|^2} + \frac{6y^{1+2Ti}}{\xi^2(1 - 2Ti)} + \frac{6y^{1-2Ti}}{\xi^2(1 + 2Ti)} \right) \xi^2(1 + 2Ti) E_A^2(z, 1/2 + iT) d\mu(z).$$

Therefore we need to calculate the twisted integrals of the second moment of the truncated Eisenstein series in the cuspidal region

$$I_\eta := \int_{\mathcal{C}_A} y^{1+\eta} \xi^2 (1 + 2Ti) E_A^2(z, 1/2 + iT) d\mu(z),$$

for the values of parameter $\eta \in \{0, \pm 2Ti\}$. Substituting here the Fourier expansion (5.2) we obtain

$$\begin{aligned} I_\eta &= 16 \int_A^{\infty} \int_0^1 y^{1+\eta} \left(\sum_{n=1}^{\infty} \tau_{iT}(n) y^{\frac{1}{2}} K_{iT}(2\pi ny) \cos(2\pi nx) \right)^2 \frac{dx dy}{y^2} \\ &= 8 \sum_{n=1}^{\infty} \tau_{iT}^2(n) \int_A^{\infty} K_{iT}^2(2\pi ny) y^\eta dy = 8 \sum_{n=1}^{\infty} \tau_{iT}^2(n) (2\pi n)^{-1-\eta} g(2\pi An), \end{aligned}$$

where

$$g(x) := \int_x^{\infty} K_{iT}^2(y) y^\eta dy.$$

The Mellin transform of this function is equal to

$$\begin{aligned} G(s) &:= \int_0^{\infty} g(x) x^s \frac{dx}{x} = \frac{1}{s} \int_0^{\infty} K_{iT}^2(x) x^{\eta+s} dx \\ &= \frac{2^{\eta-2+s}}{s\Gamma(1+\eta+s)} \Gamma^2\left(\frac{1+\eta+s}{2}\right) \Gamma\left(\frac{1+\eta+s}{2} + iT\right) \Gamma\left(\frac{1+\eta+s}{2} - iT\right) \end{aligned}$$

by integration by parts and the Mellin–Barnes formula [GR], 6.576.4

$$\int_0^{\infty} K_\mu(y) K_\nu(y) y^s \frac{dy}{y} = \frac{2^{s-3}}{\Gamma(s)} \prod_{\pm, \pm} \Gamma\left(\frac{s \pm \mu \pm \nu}{2}\right).$$

By the inverse Mellin transform we have $g(x) = \frac{1}{2\pi i} \int_{(3)} G(s) x^{-s} ds$ (where the integration is over the line $\operatorname{Re}(s) = 3$) and so we get

$$I_\eta = \frac{8}{(2\pi)^{1+\eta}} \frac{1}{2\pi i} \int_{(3)} G(s) (2\pi A)^{-s} \sum_{n=1}^{\infty} \frac{\tau_{iT}^2(n)}{n^{s+1+\eta}} ds.$$

Here, since $\tau_{iT}(n) = \sigma_{2iT}(n) n^{-iT}$, we have by Ramanujan's identity

$$\sum_{n=1}^{\infty} \frac{\tau_{iT}^2(n)}{n^{s+1+\eta}} = \frac{\zeta^2(s+1+\eta) \zeta(s+1+\eta+2Ti) \zeta(s+1+\eta-2Ti)}{\zeta(2s+2+2\eta)},$$

which then gives

$$I_\eta = \frac{1}{2\pi i} \int_{(4)} \frac{A^{1-s}}{s-1} \frac{\xi^2(s+\eta)\xi(s+\eta+2Ti)\xi(s+\eta-2Ti)}{\xi(2s+2\eta)} ds.$$

The integrand is rapidly decreasing in vertical strips and it is regular on the line $\operatorname{Re}(s) = \frac{1}{2}$ (for all three values of the parameter η), so we can shift the line of integration from $\operatorname{Re}(s) = 4$ to $\operatorname{Re}(s) = \frac{1}{2}$:

$$I_\eta = \mathcal{R}_\eta + \frac{1}{2\pi i} \int_{(1/2)} \frac{A^{1-s}}{s-1} \frac{\xi^2(s+\eta)\xi(s+\eta+2Ti)\xi(s+\eta-2Ti)}{\xi(2s+2\eta)} ds, \quad (5.3)$$

where $\mathcal{R}_\eta = \sum_P R_{\eta,P}$ is the sum of residues $R_{\eta,P}$ of the poles P that we encounter.

In the case $\eta = 0$, the integrand has two simple poles at $s = 1 \pm 2Ti$ with the residues

$$R_{0,1-2Ti} = -\frac{A^{2Ti}\xi^2(1-2Ti)\xi(1-4Ti)}{2Ti\xi(2-4Ti)} \quad \text{and} \quad R_{0,1+2Ti} = \frac{A^{-2Ti}\xi^2(1+2Ti)\xi(1+4Ti)}{2Ti\xi(2+4Ti)}$$

and the triple pole at $s = 1$ with residue

$$\begin{aligned} R_{0,1} = & \frac{|\xi(1+2Ti)|^2}{\xi(2)} \cdot \left[\left| \frac{\xi'}{\xi}(1+2Ti) \right|^2 + \operatorname{Re} \frac{\xi''}{\xi}(1+2Ti) + \frac{1}{2} \log^2 A \right. \\ & - 2(\log A) \operatorname{Re} \frac{\xi'}{\xi}(1+2Ti) + 2 \left(\frac{\xi'}{\xi}(2) - a \right) \left(\log A - 2 \operatorname{Re} \frac{\xi'}{\xi}(1+2Ti) \right) \\ & \left. + 4 \left(\frac{\xi'}{\xi}(2) \right)^2 - 4a \frac{\xi'}{\xi}(2) + a^2 + 2b - 2 \frac{\xi''}{\xi}(2) \right], \end{aligned}$$

where the constants a and b are as in (4.5).

In the case $\eta = 2Ti$, the integrand has the simple pole at $s = 1 - 4Ti$ with the residue

$$R_{2Ti,1-4Ti} = -\frac{A^{4Ti}\xi^2(1-2Ti)\xi(1-4Ti)}{4Ti\xi(2-4Ti)}$$

and two double poles at $s = 1$ and $s = 1 - 2Ti$ with the corresponding residues

$$\begin{aligned} R_{2Ti,1} = & \frac{\xi^2(1+2Ti)\xi(1+4Ti)}{\xi(2+4Ti)} \\ & \times \left[a - \log A + \frac{\xi'}{\xi}(1+4Ti) + 2 \frac{\xi'}{\xi}(1+2Ti) - 2 \frac{\xi'}{\xi}(2+4Ti) \right] \end{aligned}$$

and

$$R_{2Ti, 1-2Ti} = -\frac{A^{2Ti}\xi(1+2Ti)\xi(1-2Ti)}{2Ti\xi(2)} \\ \times \left[2\operatorname{Re}\frac{\xi'}{\xi}(1+2Ti) - \log A + \frac{1}{2Ti} + 2a - 2\frac{\xi'}{\xi}(2) \right].$$

In the case $\eta = -2Ti$, the integrand has the simple pole at $s = 1 + 4Ti$ with the residue

$$R_{-2Ti, 1+4Ti} = \frac{A^{-4Ti}\xi^2(1+2Ti)\xi(1+4Ti)}{4Ti\xi(2+4Ti)}$$

and two double poles at $s = 1$ and $s = 1 + 2Ti$ with the corresponding residues

$$R_{-2Ti, 1} = \frac{\xi^2(1-2Ti)\xi(1-4Ti)}{\xi(2-4Ti)} \\ \times \left[a - \log A + \frac{\xi'}{\xi}(1-4Ti) + 2\frac{\xi'}{\xi}(1-2Ti) - 2\frac{\xi'}{\xi}(2-4Ti) \right]$$

and

$$R_{-2Ti, 1+2Ti} = \frac{A^{-2Ti}\xi(1+2Ti)\xi(1-2Ti)}{2Ti\xi(2)} \\ \times \left[2\operatorname{Re}\frac{\xi'}{\xi}(1+2Ti) - \log A - \frac{1}{2Ti} + 2a - 2\frac{\xi'}{\xi}(2) \right].$$

In particular, we have

$$\frac{1}{\xi^2(1 \pm 2Ti)} \mathcal{R}_{\mp 2Ti} \ll (\log T)^2 T^{-1/2},$$

when $T \rightarrow \infty$.

The contribution of the integrals on the shifted line in (5.3) is bounded in the following Lemma:

Lemma 5.1. *For any $\eta \in \{0, \pm 2Ti\}$ with $T > 1$, we have*

$$\int_{-\infty}^{\infty} \left| \frac{A^{\frac{1}{2}-ti}}{\frac{1}{2}-ti} \cdot \frac{\xi^2(\frac{1}{2}+ti+\eta)\xi(\frac{1}{2}+ti+\eta+2Ti)\xi(\frac{1}{2}+ti+\eta-2Ti)}{\xi^2(1+2Ti)\xi(1+2ti+2\eta)} \right| dt \ll A^{1/2} T^{-1/6},$$

with an absolute implicit constant.

Proof. The analysis is similar to that in Lemma 4.4. The case $\eta = 0$ was treated in [Sp], section 4.3.2, where the bound $O(A^{1/2}T^{-1/6})$ is obtained. Here, we treat the case $\eta = 2Ti$ (for $\eta = -2Ti$, the value of the integral is the same). Using Stirling's formula, we see that the integrand is bounded by

$$A^{1/2} \frac{e^{\frac{\pi}{4}(4T-|t|-|t+4T|)}}{(1+|t|)^{5/4}(1+|t+2T|)^{1/2}(1+|t+4T|)^{1/4}} \times \left| \frac{\zeta(\frac{1}{2}+ti) \zeta^2(\frac{1}{2}+(t+2T)i) \zeta(\frac{1}{2}+(t+4T)i)}{\zeta^2(1+2Ti) \zeta(1+(2t+4T)i)} \right|.$$

Using subconvexity estimate $\zeta(1/2+ti) \ll (1+|t|)^{\theta+\epsilon}$, for all $\epsilon > 0$ and some $\theta < \frac{1}{6}$ for the zeta-functions in the numerator and (4.7) for the zeta-functions in the denominator, this is further bounded by

$$A^{1/2}T^\epsilon \frac{e^{\frac{\pi}{4}(4T-|t|-|t+4T|)}}{(1+|t|)^{\frac{5}{4}-\theta-\epsilon}(1+|t+2T|)^{\frac{1}{2}-2\theta-\epsilon}(1+|t+4T|)^{\frac{1}{4}-\theta-\epsilon}}.$$

We split the integration into 3 ranges: $\int_{-\infty}^{\infty} = \int_{-\infty}^{-4T} + \int_{-4T}^0 + \int_0^{\infty}$. In the first and the third range we have an exponential decay of the integrand and so we have that in these ranges the integrals are bounded respectively by

$$\int_0^{\infty} \ll A^{1/2}T^\epsilon \int_0^{\infty} \frac{e^{-\frac{\pi}{2}t}}{(1+|t|)^{\frac{5}{4}-\theta-\epsilon}T^{\frac{1}{2}-2\theta+\frac{1}{4}-\theta}} dt \ll A^{1/2}T^{-1/4}$$

and

$$\int_{-\infty}^{-4T} \ll A^{1/2}T^\epsilon \int_{-\infty}^{-4T} \frac{e^{\frac{\pi}{2}(t+4T)}}{T^{\frac{5}{4}-\theta+\frac{1}{2}-2\theta}(1+|t+4T|)^{\frac{1}{4}-\theta-\epsilon}} dt \ll A^{1/2}T^{-5/4}.$$

In the middle range the integral is bounded by

$$\begin{aligned} &\ll A^{1/2}T^\epsilon \int_{-4T}^0 \frac{dt}{(1+|t|)^{\frac{5}{4}-\theta}(1+|t+2T|)^{\frac{1}{2}-2\theta}(1+|t+4T|)^{\frac{1}{4}-\theta}} \\ &\ll A^{1/2}T^{\theta-\frac{1}{4}+\epsilon} \int_{-2T}^0 \frac{dt}{(1+|t|)^{\frac{5}{4}-\theta}(1+|t+2T|)^{\frac{1}{2}-2\theta}} \ll A^{1/2}T^{3\theta-\frac{3}{4}+\epsilon} \ll A^{1/2}T^{-1/4}. \end{aligned}$$

Therefore, in the cases $\eta = \pm 2Ti$, we get an even better bound $O(A^{1/2}T^{-1/4})$. \square

After we collect everything together, and use asymptotic formulas for $\frac{\xi'}{\xi}(1+2iT)$ and $\frac{\xi''}{\xi}(1+2iT)$ already seen in the previous section, we get that the contribution of the first integral in (5.1) is

$$\begin{aligned}
& \frac{12}{|\xi(1+2Ti)|^2} I_0 + \frac{6}{\xi^2(1-2Ti)} I_{2Ti} + \frac{6}{\xi^2(1+2Ti)} I_{-2Ti} \\
&= \frac{12}{|\xi(1+2Ti)|^2} \mathcal{R}_0 + \frac{6}{\xi^2(1-2Ti)} \mathcal{R}_{2Ti} + \frac{6}{\xi^2(1+2Ti)} \mathcal{R}_{-2Ti} + O(A^{1/2} T^{-1/6}) \\
&= \frac{36}{\pi} \log^2 T + O(\log^{\frac{5}{3}+\epsilon} T),
\end{aligned} \tag{5.4}$$

for the range $1 < A \ll \log T$. The main contribution is coming from $R_{0,1}$.

Putting together all our calculations in this section, we end up with the following Proposition:

Proposition 5.2. *When $T \rightarrow \infty$, for any value of the truncation parameter $1 < A \ll \log T$ we have*

$$\begin{aligned}
& \int_X^{reg} |E(z, \frac{1}{2} + iT)|^4 d\mu(z) = \\
& \int_X |E_A(z, \frac{1}{2} + iT)|^4 d\mu(z) + \frac{36}{\pi} \log^2 T + 2 \int_{\mathcal{C}_A} (|E_A|^2 \overline{E}_A e + |E_A|^2 E_A \overline{e}) d\mu + O(\log^{\frac{5}{3}+\epsilon} T).
\end{aligned} \tag{5.5}$$

The first integral on the right hand side of (5.5) is asymptotic to $\frac{36}{\pi} \log^2 T$ under assumption (2.3). The integral over cuspidal region \mathcal{C}_A in (5.5) is bounded by

$$4 \int_{\mathcal{C}_A} \left| E_A^3(z, \frac{1}{2} + Ti) e(y, \frac{1}{2} + Ti) \right| d\mu(z) \leq 4 \left(\int_{\mathcal{C}_A} |E_A|^4 d\mu \right)^{1/2} \left(\int_{\mathcal{C}_A} |e E_A|^2 d\mu \right)^{1/2}. \tag{5.6}$$

The second integral on the right hand side of (5.6) is $\sim \frac{6}{\pi} \log^2 T$; this is implicit in the calculation of the first integral in (5.1). Under the RWC, the first integral on the right hand side of (5.6) can be bounded by

$$\int_X |E_A|^4 d\mu - \int_{\mathcal{D}_A} |E_A|^4 d\mu \ll (\text{vol}(X) - \text{vol}(\mathcal{D}_A)) \log^2 T \ll A^{-1} \log^2 T,$$

which is $o(\log^2 T)$ if A grows arbitrary slowly to infinity as $T \rightarrow \infty$. This way the right hand side of (5.5) is asymptotic to $\frac{72}{\pi} \log^2 T$.

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