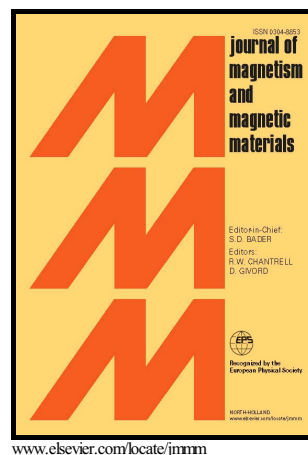


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Magnetism and charge ordering in high- and low-temperature phases of $\text{Nb}_2\text{O}_2\text{F}_3$

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Abstract

It is shown using *ab initio* band structure calculations that $\text{Nb}_2\text{O}_2\text{F}_3$ is in the orbital-selective regime in the high-temperature phase ($T > 90$ K), when two electrons occupy singlet molecular orbital, while the magnetic response comes from the remaining single electron in $\text{Nb}_2^{3.5+}$ dimer. The charge order occurs at low temperatures, resulting in the formation of $\text{Nb}^{3+}\text{-Nb}^{3+}$ and $\text{Nb}^{4+}\text{-Nb}^{4+}$ dimers, which makes this system nonmagnetic. The single electron with unpaired spin is transferred to $\text{Nb}^{3+}\text{-Nb}^{3+}$ dimer, but due to a strong splitting of the bonding xz/yz molecular orbitals the low-spin state with $S = 0$ is stabilized. We argue that the mechanism of the charge ordering in $\text{Nb}_2\text{O}_2\text{F}_3$ is the gain in kinetic energy related to the formation of molecular orbitals, which occurs due to a strong nonlinear distance dependence of the hopping parameters.

Keywords: Dimers, charge ordering.

1. Introduction

The $4d-5d$ transition metal (TM) compounds attract considerable attention in recent years. For example, a new Mott state described by the effective total angular momentum $J_{eff}=1/2$ and induced by relativistic spin-orbit coupling

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5 was found in $5d$ transition metal oxide Sr_2IrO_4 [1]. Another interesting effect thoroughly studied nowadays is a possible realization of the so-called Kitaev model [2] on hexagonal lattices, where nearest neighbor exchange Hamiltonian has a form $J \sum_{ij} \vec{S}_i^\alpha \vec{S}_j^\alpha$, where along each of metal-metal bonds only one spin projection $\alpha = x, y, z$ enters [3].

10 Going from the $3d$ to $4d$ and $5d$ transition metal compounds, we not only start to deal with more heavy elements, which implies stronger spin-orbit coupling, but also gradually decrease the on-site Coulomb repulsion (characterized by U parameter) and simultaneously increase spacial extension of d -wavefunctions [4]. This drives $4d$ and $5d$ systems to a regime where hopping parameters t_{ij}
 15 are comparable with the on-site Hubbard repulsion U . Such a growth of t_{ij} is especially pronounced in dimerized systems, where this effect may result in an orbital-selective behavior, when only a part of the d -electrons bear magnetic moments, while others form molecular-orbitals [5]. This effect, decreasing the total magnetic moment of a system and suppressing the double exchange
 20 mechanism, was found in such systems as $\text{Y}_5\text{Mo}_2\text{O}_{12}$ [6] and $\text{Ba}_5\text{AlIr}_2\text{O}_{11}$ [7, 8].

In the present paper we discuss a very similar effect of the orbital-selective behavior and partial suppression of the magnetic moment in a recently synthesized niobium oxyfluoride $\text{Nb}_2\text{O}_2\text{F}_3$ [9]. This compound presents a special interest, since here the orbital-selective behavior is observed together with a
 25 charge disproportionation. We argue that the mechanism of charge ordering in this system is of a novel type and is connected not with the gain of interaction energy, as in the conventional mechanisms of charge ordering, but is rather due to the gain in kinetic (bonding) energy due to formation of tightly bound dimers.

30 $\text{Nb}_2\text{O}_2\text{F}_3$ crystallizes in a monoclinic structure (space group: $\text{I}2/a$) and undergoes a structural transition at $\sim 90\text{K}$ to a triclinic structure (space group: $\text{P}-1$) in the lower temperatures [9]. The phase transition at 90K is accompanied by a strong suppression of magnetic susceptibility, which is ascribed to formation of the spin gap. Effective magnetic moment in the Curie-Weiss theory
 35 $\mu_{eff} = 2.24\mu_B$ (per dimer) in the high temperature (HT) phase is much smaller

than expected for Nb having nominal valence 3.5+. The transport measurement have been performed for $T > 90\text{K}$, where $\text{Nb}_2\text{O}_2\text{F}_3$ was found to be insulating.

The Nb ions are in octahedral surrounding, and these octahedra share their edges forming Nb-Nb dimers (see Fig. 1). There is one type of such dimers in the HT phase with crystallographically equivalent Nb ions and each Nb having 3.5+ charge state (electronic configurations $4d^{1.5}$). The transition at $\sim 90\text{K}$ results in a strong modification of the crystal structure. There appears two inequivalent types of Nb and each of them form its own type of dimers: “short” Nb1-Nb1 (2.50\AA) and “long” Nb2-Nb2 (2.66\AA) dimers. The difference in Nb-Nb bond distances was proposed to be a signature of charge disproportionated state with Nb1 being in the 3+ (electronic configurations $4d^2$) and Nb2 in the 4+ (configurations $4d^1$) charge states. The origin of the phase transition at 90K is unclear and demands thorough investigation, which we carry out in this paper using ab-initio band structure calculations.

2. Calculation details

For the band-structure calculations the pseudopotential PWscf code [11] was mainly used. We utilized the generalized gradient approximation (GGA) with Perdew-Burke-Ernzerhof version of the exchange-correlation potential [12] and ultrasoft scalar-relativistic pseudopotentials with nonlinear core correction.

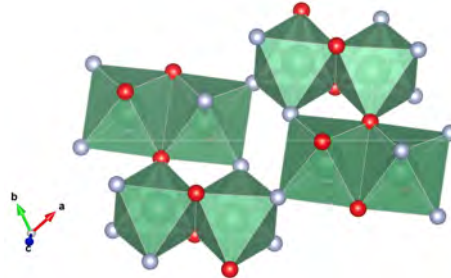


Figure 1: Crystal structure of $\text{Nb}_2\text{O}_2\text{F}_3$ (generated by the Vesta software [10]). Niobium ions are in the octahedra formed by oxygen (red balls) and fluorine (grey balls).

55 The charge density and kinetic energy cutoffs were taken as 40 Ry and 180 Ry, respectively. The mesh of 6x6x6 k -points in a full Brillouin zone was used in the calculation.

In order to estimate the crystal field splitting parameters, the Wannier function projection method [13] was used as realized in the linearized muffin-tin
60 orbitals (LMTO) method [14].

All parameters in the paper for both LT (low temperature) and HT (high temperature) phases were calculated with using of experimentally obtained structures [9].

3. Magnetic properties

65 We start with the analysis of the nonmagnetic GGA calculation results for the HT monoclinic phase. Since this approximation does not take into account strong on-site Hubbard repulsion [15], $\text{Nb}_2\text{O}_2\text{F}_3$ is metallic in this type of calculations with the Nb t_{2g} states forming both valence and conduction bands. The absence of a band gap is a standard drawback of the GGA method, which
70 can be corrected by such methods as LDA/GGA+U [15] or LDA+DMFT [16]. In spite of this fact, GGA/LDA is known to be accurate for the estimation of the on-site energies and hopping parameters, which are of most interest for us here.

We calculated these parameters (hopping integrals and crystal-field splitting)
75 using so-called Wannier function projection procedure. In the edge-sharing geometry there are xy -orbitals pointing exactly to each other in the Nb-Nb dimer (we use notations of local coordinate system, where x and y axis are directed to common ligands in a dimer). This results in a strong bonding-antibonding splitting, ~ 3.3 eV, see Fig. 2. The energy difference between the bonding xy
80 and next (higher in energy) orbitals is $\Delta = 1.1$ eV. Using ionic approximation and neglecting on-site Coulomb repulsion one may expect that for a Nb-Nb dimer with 1.5 electrons per site the state with maximal spin $S_{tot} = 3/2$ will be stabilized, if $\Delta < J_H/2$, where J_H is the intra-atomic Hund's exchange, but

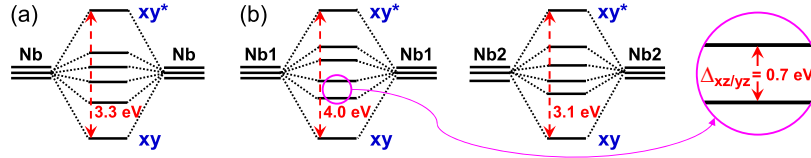


Figure 2: Formation of the molecular orbitals: (a) in the high temperature and (b) in the low temperature phases, as obtained in of the Wannier function projection method in frames of the LMTO method.

smaller net spin would exist in the opposite case [5, 7].

We estimated J_H in $\text{Nb}_2\text{O}_2\text{F}_3$ using the so-called constrained supercell method [17] in the LMTO calculations and found that $J_H = 0.7$ eV. Since the bonding-antibonding splitting is much larger than $J_H/2$, the state with minimal total spin $S_{tot} = 1/2$ per dimer is realized in the HT phase of $\text{Nb}_2\text{O}_2\text{F}_3$, see the level scheme in Fig. 2. This state is orbital-selective, since xy orbitals behave here as nonmagnetic molecular-orbital like, while xz/yz orbitals provide local magnetic moment.

The ferromagnetic GGA calculations, results of which are presented in Fig. 3a, show that total magnetic moment is $\mu = 1\mu_B$ per formula unit (f.u.). This state corresponds to $\mu_{eff} = \sqrt{3}\mu_B$ (per dimer), which is of the order of value observed experimentally [9]. The bonding and antibonding xy bands are at ~ -1 eV and ~ 3.5 eV, while remaining peaks in vicinity of the Fermi energy correspond to the xz and yz orbitals.

The electronic structure in the low temperature (LT) phase is more complex. It was proposed that there can be charge disproportionation in this phase [9]. Analysis of the hopping parameters shows that the bonding-antibonding splittings in Nb1-Nb1 dimers are even larger than in the HT phase, which is due to smaller bond distance. This is true for both xy and xz/yz orbitals. One may see from the partial density of states (DOS), shown in Fig. 3b, that both bonding xy and one of the low-lying xz/yz orbitals are occupied in the Nb1-Nb1 dimer. Therefore, one may conclude that there has to be four electrons in the Nb1-Nb1 dimer and only two electrons in the Nb2-Nb2 dimer. However,

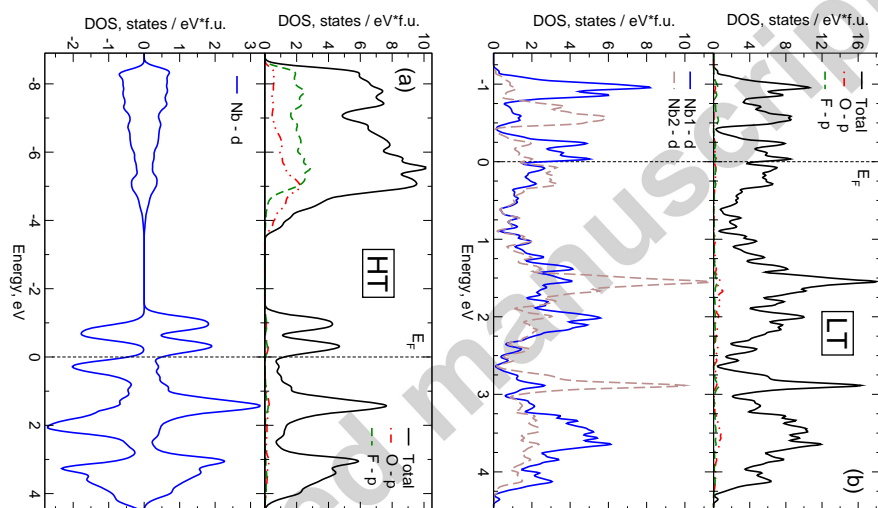


Figure 3: Total and partial densities of states (DOS) for the high-temperature (HT) monoclinic phase (a) and low-temperature (LT) triclinic phase (b) obtained in the GGA calculation. Only nonmagnetic solution survives in the LT phase, while for the HT phase the results for the ferromagnetic order are presented. Positive and negative values of DOS in lower panel of (a) correspond to different spin projections. The Fermi level is set zero.

difference in Löwdin charges (total number of d electrons) [18] for two types of Nb is minimal, $\delta n = n_{Nb1} - n_{Nb2} \sim 0.1e$. It is rather typical that the degree of charge disproportionation in a band-structure calculation turns out to be much smaller than what one would expect from formal valencies. The difference in total atomic charges or in corresponding magnetic moments in the GGA or GGA+U calculations is often an order magnitude smaller than it should formally be [19, 20, 21].

It is clear that, having formally two electrons, Nb2-Nb2 dimer stays non-magnetic since they occupy bonding xy molecular orbitals. The situation with Nb1-Nb1 dimer with four electrons per dimer is more tricky. In principle one may have three different “molecular spin-states” in this case (4 electrons per dimer) with total spin per dimer $S_{tot} = 2$, $S_{tot} = 1$, and $S_{tot} = 0$, and one could even expect magnetic transitions between these states as it occurs in MoCl_4 [22]. The final electronic configuration is determined by competition between Hund’s rule coupling, given by J_H , and splitting between molecular orbitals of the xz and yz symmetry. Simple estimates in an ionic model show that these are the states with $S_{tot} = 1$ and $S_{tot} = 0$, which can be stabilized in the LT phase. If the splitting between bonding xz/yz orbitals $\Delta_{xz/yz} > J_H/2$ (for definition see Fig.2) then $S_{tot} = 0$ is realized. The Wannier function analysis shows that $\Delta_{xz/yz} = 0.7$ eV, which is much larger than $J_H/2 = 0.35$ eV, i.e. for Nb1-Nb1 dimers a nonmagnetic $S_{tot} = 0$ state is realized. This is consistent both with the results of our GGA calculations, where magnetic solution does not survive in the LT phase, and with the experimental results showing nonmagnetic behavior for $T < 90\text{K}$ [9].

4. Charge ordering

We saw that the band structure calculations indeed show some degree of the charge disproportionation in the low-temperature phase of $\text{Nb}_2\text{O}_2\text{F}_3$. Generally speaking there can be many different origins of the charge ordering. The simplest one is the gain in long-range Coulomb (Madelung) energy, like in Wigner

crystals [23, 24]. In real transition metal compounds the electron-lattice coupling (usually to the breathing mode of MeL_6 octahedra, where L is a ligand ion) often plays even more important role. Sometimes orbital degrees or spin degrees of freedom may help to stabilize the charge order [4, 25, 19]. In all these mechanisms actually the onset of the charge ordered state is related with the
 140 tendency to gain an interaction energy. Using simple arguments we propose that the mechanism of the charge disproportionation in $\text{Nb}_2\text{O}_2\text{F}_3$ is of very different nature and is connected with the formation of the molecular orbitals in this compound, i.e. it has a two-site nature.

145 The bonding-antibonding splitting for the molecular orbitals in $\text{Nb}_2\text{O}_2\text{F}_3$ is defined by corresponding hopping integrals, or overlap between atomic d -orbitals. These hoppings are strongly nonlinear $t \sim 1/r^5$ [26], where r is the distance in a dimer. Therefore making short and long dimers instead of two Nb-Nb dimers having the same bond length we win a lot of kinetic energy. Similar idea was
 150 used before to explain isotope effect in manganites [27] and cobaltates [28]. Moreover, it turns out that one gains a kinetic energy not only due to the xy -orbitals having direct overlap in the common edge geometry. Other Nb $4d$ -orbitals also contribute to lowering of the charge ordered state. One can construct molecular orbitals of the xz/yz symmetry. By making a linear combination $xz + yz$ we
 155 maximize direct hopping keeping overlap via ligand p_z orbital nearly the same [29].

One may estimate the total energies of the charge ordered and homogeneous states in a simple ionic approximation using Kanamori parametrization for the Coulomb interaction [30] and calculated with the LMTO hopping parameters.
 160 Since the Nb-Nb distance in the HT phase is nearly exactly the arithmetic mean of short and long Nb-Nb bonds in the LT phase, one may use these data to calculate the difference in hopping parameters for homogeneous (H) and charge ordered (CO) states δt_{xy} and δt_{xz+yz} for the xy and $xz + yz$ orbitals, correspondingly. Then the total energy difference between these two states in

165 the ionic approximation can be estimated as

$$\delta E = E_H - E_{CO} = \delta t_{xy} + \delta t_{xz+yz} - U/2. \quad (1)$$

The direct calculation based on the results of the Wannier function analysis shows that $\delta t_{xy} = 0.8$ eV and $\delta t_{xz+yz} = 1$ eV. One may expect that U for Nb is smaller than 3 eV (typical value for, e.g., Ru^{4+} having much larger number of electrons [31]). Thus, we see that at zero temperature indeed a gain in the kinetic
170 energy due to formation of molecular orbitals exceeds a loss in the Coulomb energy because of the charge ordering, which in its turn is even overestimated in the simple Hubbard-like treatment used above: when we take into account $p-d$ covalency, the actual charge transfer in transition metal system with charge ordering is never complete, as was assumed in the crude estimate above, but is
175 much smaller, typically near $\pm 0.2e$ [4]. At finite temperature T instead of the total energies, E , one has to compare free energies $F = E - TS$ and increasing temperature the entropy term will dominate, which stabilizes homogeneous state and finally leads to the transition to this state.

However, one has to note that there are two important ingredients, which
180 we did not take into account in aforementioned model. First of all, Coulomb interaction not simply adds a penalty due to inhomogeneous charge distribution, but also modifies the wavefunctions of the corresponding states increasing contributions from the ionic terms (i.e. changing them from molecular-orbital like to Heitler-London like). This will decrease energy gain due to the formation
185 of the molecular orbitals [7].

Second, there is a very important contribution due to the elastic energy, which normally stabilizes homogeneously states. This correction is not easy to estimate on a model level for such a complicated structure as in $\text{Nb}_2\text{O}_2\text{F}_3$. However, there is an extra feature of the $\text{Nb}_2\text{O}_2\text{F}_3$ crystal structure, which has
190 to be mentioned and which helps the observed charge ordering. Dimerization of one class of Nb drives another one to change its charge state. Indeed, if, e.g., we decrease Nb1-Nb1 bond distance and want to keep Nb1 valence as it is (i.e. keep the same average Nb1-O and Nb1-F bond lengths), then oxygen ions linking

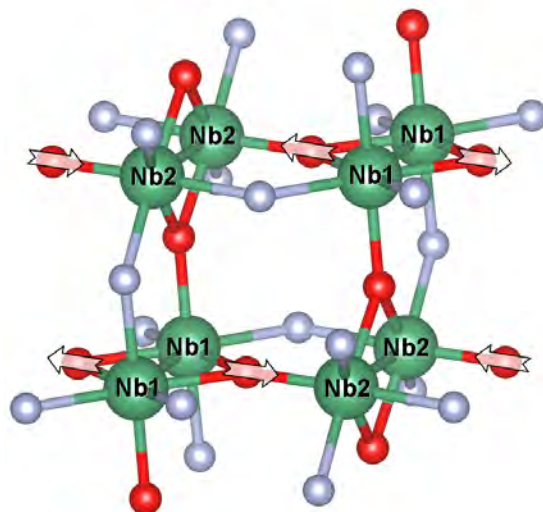


Figure 4: Illustration of cooperative distortions, which support formation of the charge ordered state.

Nb1-Nb1 and Nb2-Nb2 dimers should move towards Nb2 as it is shown in Fig. 4.

195 This should lead to decrease of corresponding Nb2-O bond distance and hence to the increase of the valence of Nb2 as a result. It is interesting to note that the low temperature crystal structure of $\text{Nb}_2\text{O}_2\text{F}_3$ indeed shows this effect with shortest Nb-O bond (1.89 Å, while other Nb2-O bond lengths are 2.03 Å and Nb2-F are $\sim 2.07\text{-}3.0$ Å) linking two dimers together. Thus, we see that in this
200 sense dimerization goes hand in hand with the charge disproportionation in $\text{Nb}_2\text{O}_2\text{F}_3$. This situation actually strongly resembles that in M_1 and M_2 phases of VO_2 [32, 33, 4].

5. Summary

To sum up, we show in the present paper that magnetic response in $\text{Nb}_2\text{O}_2\text{F}_3$
205 at the high-temperatures ($T > 90$ K) is related to the orbital selective regime, when part of the electrons form molecular orbitals while other electrons have local magnetic moments. The charge disproportionation, which occurs at $T \sim 90$

K is seen in the GGA calculations, but its degree ($\delta n \sim 0.1$ electron) is far from what one would expect from naive expectations based on the formal ionic valences. The mechanism of the charge ordering is argued to be related with a sizable kinetic energy gain due to formation of two molecular orbitals in short Nb^{3+} - Nb^{3+} dimers caused by a strong nonlinearity of the distance dependence on electron hopping. We think that this mechanism of charge ordering, stabilized not by decrease of interaction energy, but rather by the gain in kinetic energy, may be operative in many other systems, especially consisting of structural dimers.

Acknowledgments

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References

- [1] B. J. Kim, H. Jin, S. J. Moon, J.-Y. Kim, B.-G. Park, C. S. Leem, J. Yu, T. W. Noh, C. Kim, S.-J. Oh, J.-H. Park, V. Durairaj, G. Cao, E. Rotenberg, Novel $j_{eff}=1/2$ mott state induced by relativistic spin-orbit coupling in Sr_2IrO_4 , Phys. Rev. Lett. 101 (2008) 076402. doi:10.1103/PhysRevLett.101.076402.
- [2] A. Kitaev, Anyons in an exactly solved model and beyond, Ann. Phys. (N.Y.) 321 (2006) 2–111. doi:10.1016/j.aop.2005.10.005.
- [3] G. Jackeli, G. Khaliullin, Mott insulators in the strong spin-orbit coupling limit: From heisenberg to a quantum compass and kitaev models, Phys. Rev. Lett. 102 (2009) 017205. doi:10.1103/PhysRevLett.102.017205.
- [4] D. I. Khomskii, Transition metal compounds, Cambridge University Press.

- [5] S. V. Streltsov, D. I. Khomskii, Orbital-dependent singlet dimers and orbital-selective peierls transitions in transition-metal compounds, Phys. Rev. B 89 (2014) 161112(R). doi:10.1103/PhysRevB.89.161112.
- [6] S. V. Streltsov, Orbital-selective behavior in $\text{Y}_5\text{Mo}_2\text{O}_{12}$ and $(\text{Cd,Zn})\text{V}_2\text{O}_4$, J. Magn. Magn. Mater. 383 (2015) 27–29. doi:10.1016/j.jmmm.2014.10.150.
- [7] S. V. Streltsov, D. I. Khomskii, Covalent bonds against magnetism in transition metal compounds, arXiv:1602.06425.
- [8] J. Terzic, J. C. Wang, F. Ye, W. H. Song, S. J. Yuan, S. Aswartham, L. E. DeLong, S. V. Streltsov, D. I. Khomskii, G. Cao, Coexisting charge and magnetic orders in the dimer-chain iridate $\text{Ba}_5\text{AlIr}_2\text{O}_{11}$, Phys. Rev. B 9 (2015) 235147. doi:10.1103/PhysRevB.91.235147.
- [9] T. Tran, M. Gooch, B. Lorenz, A. Litvinchuk, M. n. Sorolla, B. J., P. Chu, A. Guloy, $\text{Nb}_2\text{O}_2\text{F}_3$: A reduced niobium (iii/iv) oxyfluoride with a complex structural, magnetic, and electronic phase transition, J. Am. Chem. Soc. 137 (2015) 636. doi:10.1021/ja511745q.
- [10] K. Momma, F. Izumi, Vesta 3 for three-dimensional visualization of crystal, volumetric and morphology data, J. Appl. Crystallogr. 44 (2011) 1272. doi:10.1107/S0021889811038970.
- [11] P. Giannozzi, et al, Quantum espresso: a modular and open-source software project for quantum simulations of materials, J. Phys.: Condens. Matter 21 (2009) 395502. doi:10.1088/0953-8984/21/39/395502.
- [12] J. P. Perdew, K. Burke, M. Ernzerhof, Generalized gradient approximation made simple, Phys. Rev. Lett. 77 (1996) 3865. doi:10.1103/PhysRevLett.77.3865.
- [13] S. V. Streltsov, A. S. Mylnikova, A. O. Shorikov, Z. V. Pchelkina, D. I. Khomskii, A. V. I., Crystal-field splitting for low symmetry systems in

- 260 ab initio calculations, Phys. Rev. B 71 (2005) 245114. doi:10.1103/PhysRevB.71.245114.
- [14] O. K. Andersen, O. Jepsen, Explicit, first-principles tight-binding theory, Phys. Rev. Lett. 53 (1984) 2571. doi:10.1103/PhysRevLett.53.2571.
- [15] V. Anisimov, F. Aryasetiawan, A. Lichtenstein, First-principles calculations
265 of the electronic structure and spectra of strongly correlated systems: the lda+ u method, Journal of Physics: Condensed Matter 9 (1997) 767. doi:10.1088/0953-8984/9/4/002.
- [16] V. I. Anisimov, A. I. Poteryaev, M. A. Korotin, A. O. Anokhin, G. Kotliar, First-principles calculations of the electronic structure and spectra of
270 strongly correlated systems: dynamical mean-field theory, Journal of Physics: Condensed Matter 9 (1997) 7359. doi:10.1088/0953-8984/9/35/010.
- [17] V. Anisimov, O. Gunnarsson, Density-functional calculation of effective coulomb interactions in metals, Phys. Rev. B 43 (1991) 7570. doi:10.1103/PhysRevB.43.7570.
275
- [18] P.-O. Lowdin, On the non-orthogonality problem connected with the use of atomic wave functions in the theory of molecules and crystals, J. Chem. Phys. 18 (1950) 365. doi:10.1063/1.1747632.
- [19] S. V. Streltsov, D. I. Khomskii, Jahn-teller distortion and charge, orbital, and magnetic order in LaMnO_3 , Phys. Rev. B 89 (2014) 201115. doi:10.1103/PhysRevB.89.201115.
280
- [20] I. Leonov, A. N. Yaresko, V. N. Antonov, M. A. Korotin, V. I. Anisimov, Charge and orbital order in Fe_3O_4 , Phys. Rev. Lett. 93 (2004) 146404. doi:10.1103/PhysRevLett.93.146404.
- [21] W. E. Pickett, Y. Quan, V. Pardo, Charge states of ions, and mechanisms of charge ordering transitions, J. Phys. Cond. Matt. 26 (2014) 274203. doi:10.1088/0953-8984/26/27/274203.
285

- [22] D. Korotin, V. Anisimov, S. Streltsov, Pressure-induced magnetic transitions with change of the orbital configuration in dimerised systems, Sci. Rep. 6 (2016) 25831. doi:10.1038/srep25831.
- [23] E. Wigner, Effects of the electron interaction on the energy levels of electrons in metals, Trans. Faraday Soc. 34 (1938) 678–685. doi:10.1039/TF9383400678.
- [24] J. Hubbard, Generalized wigner lattices in one dimension and some applications to tetracyanoquinodimethane (tcnq) salts, Phys. Rev. B 17 (1978) 494. doi:10.1103/PhysRevB.17.494.
- [25] I. V. Solovyev, K. Terakura, Magnetic spin origin of the charge-ordered phase in manganites, Phys. Rev. Lett. 83 (1999) 2825. doi:10.1103/PhysRevLett.83.2825.
- [26] W. A. Harrison, Elementary electronic structure, World Scientific, Singapore.
- [27] N. A. Babushkina, L. M. Belova, V. I. Ozhogin, O. Y. Gorbenko, A. R. Kaul, A. A. Bosak, D. I. Khomskii, K. I. Kugel, N. V. Ozhogin, Metal-insulator transition induced by ^{16}O - ^{18}O oxygen isotope exchange in colossal negative magnetoresistance manganites, J. Appl. Phys 83 (1998) 7369–7371. doi:10.1063/1.367636.
- [28] N. A. Babushkina, A. N. Taldenkov, S. V. Streltsov, T. G. Kuzmova, A. A. Kamenev, A. R. Kaul, D. I. Khomskii, K. Kugel, Effect of eu doping and partial oxygen isotope substitution on magnetic phase transitions in $(\text{Pr}_{1-y}\text{Eu}_y)_{0.7}\text{Ca}_{0.3}\text{CoO}_3$ cobaltites, JETP 118 (2014) 266. doi:10.1134/S1063776114010026.
- [29] S. A. J. Kimber, I. I. Mazin, J. Shen, H. O. Jeschke, S. V. Streltsov, D. N. Argyriou, R. Valenti, D. I. Khomskii, Valence bond liquid phase in the honeycomb lattice material Li_2RuO_3 , Phys. Rev. B 89 (2014) 081408(R). doi:10.1103/PhysRevB.89.081408.

- [30] J. Kanamori, Electron correlation and ferromagnetism of transition metals, Prog. Theor. Phys. 30 (1963) 275–289. doi:10.1143/PTP.30.275.
- [31] S. Lee, J.-G. Park, D. T. Adroja, D. I. Khomskii, S. V. Streltsov, K. A. McEwen, H. Sakai, K. Yoshimura, V. I. Anisimov, D. Mori, R. Kanno, R. Ibberson, Spin gap in $\text{Ti}_2\text{Ru}_2\text{O}_7$ and the possible formation of haldane chains in three-dimensional crystals, Nat. Mater. 5 (2006) 471. doi:10.1038/nmat1605.
- [32] M. Ghedira, H. Vincent, M. Marezio, J. Launay, Structural aspects of the metal-insulator transitions in $\text{V}_{0.985}\text{Al}_{0.015}\text{O}_2$, J. Solid State Chem. 22 (1977) 423. doi:10.1016/0022-4596(77)90020-2.
- [33] L. Whittaker, C. J. Partridge, S. Banerjee, Microscopic and nanoscale perspective of the metalinsulator phase transitions of VO_2 : Some new twists to an old tale, J. Phys .Chem. Lett. 2 (2011) 745. doi:10.1021/jz101640n.