



# Magnetic and submillimeter spectroscopy study of phase transitions in $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ perovskites: $T$ - $x$ phase diagram

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## Abstract

Temperature dependence of magnetization, AC magnetic susceptibility, resistance and submillimeter dielectric permittivity and dynamic conductivity were studied in colossal magnetoresistance perovskites  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ , which exhibit a lot of anomalies identified with various magnetic and structural phase transitions. As a result, an overall phase  $T$ - $x$  diagram was obtained at temperatures up to 1050 K and  $0 \leq x \leq 0.45$ . © 1999 Elsevier Science B.V. All rights reserved.

**Keywords:**  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ ; Phase transitions; Submillimeter spectroscopy

Lanthanum manganites doped by the divalent ions Sr (Ca) exhibit various interesting phenomena such as a colossal negative magnetoresistance, magnetic, structural and metal-insulator phase transitions, a charge or polaron ordering [1–3]. The parent compound  $\text{LaMnO}_3$  is a Mott insulator and has a canted antiferromagnetic layer structure. The substitution of  $\text{La}^{3+}$  ions by  $\text{Sr}^{2+}$  ions results in a transition from antiferromagnetic insulating state to a ferromagnetic metallic state at  $x \geq 0.17$ , which is stabilized by a double exchange. The crystal structure in this case is also changed from the orthorhombic to rhombohedral symmetry [1]. Complicated phase transformations are expected in the intermediate concentrations  $0.1 \leq x \leq 0.15$  [2,3]. In this work we performed complex investigations of various magnetic and structural transitions in  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$  manganites by means magnetic and electric static and submillimeter dynamic measurements.

Single crystals of  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$  ( $0 \leq x \leq 0.45$ ) were grown by the floating zone method with radiation

heating. The transmission and phase shift spectra of thin plane-parallel plates were measured using a quasi-optical submillimeter backward-wave-oscillator technique [4] in the frequency range  $\nu = 3\text{--}33\text{ cm}^{-1}$ , that allowed to obtain a dynamic dielectric permittivity  $\epsilon'$  and conductivity  $\sigma'$ . Resistance  $\rho(T)$  were measured using the four-probe method at temperatures 4.2–1050 K. Magnetization  $M(T, H)$  and AC magnetic susceptibility  $\chi_{AC}(T)$  measurements were performed at  $T = 4.2\text{--}400\text{ K}$  in magnetic field up to 12 kOe.

Temperature dependence of the resistance  $\rho(T)$  are shown in the Fig. 1. The curves exhibit several kinds of anomalies (some of them were also observed in Ref. [1] at  $T < 500\text{ K}$ ) assigned by different symbols and identified with the following transitions. (1) Ferromagnetic phase transition at  $T_C$  accompanied by a metal-semiconductor transition ( $x \geq 0.1$ ), which assigned by arrows. This transition also manifests itself distinctly as a sharp increase  $\chi_{AC}(T)$  at  $T_C$  (Fig. 2c). (2) Structural phase transitions at  $T'_s$  between low-temperature orthorhombic (Jahn–Teller)  $O'$  phase and high-temperature orthorhombic (pseudocubic)  $O^*$  phase at  $0 \leq x \leq 0.125$  (filled down triangular). The most strong anomaly in the  $\rho(T)$  is observed for the pure  $\text{LaMnO}_3$  at  $T'_s \approx 750\text{ K}$  and

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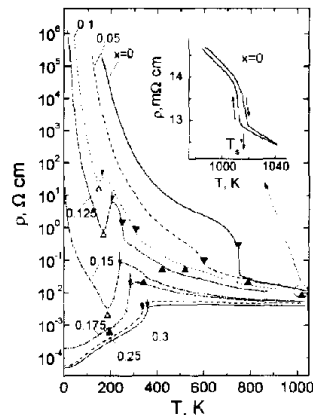


Fig. 1. Temperature dependence of the resistance in  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ . Inset:  $\rho(T)$  behavior near a new structural transition in  $\text{LaMnO}_3$ . Symbols indicate various phase transitions (see text).

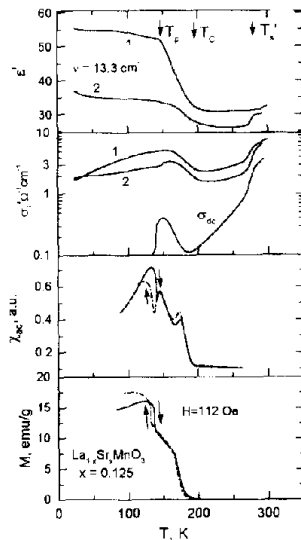


Fig. 2. Temperature dependence of the dielectric permittivity  $\epsilon'(13.3 \text{ cm}^{-1})$  (a), dynamic ( $\sigma'(13.3 \text{ cm}^{-1})$ ) and static ( $\sigma_{\text{DC}}$ ) conductivity (b), AC magnetic susceptibility (c) and magnetization (d) in  $\text{La}_{0.875}\text{Sr}_{0.125}\text{MnO}_3$ . Indices 1,2 for the  $\epsilon'$  and  $\sigma'$  correspond to different radiation polarizations.

corresponds to the known Jahn–Teller transition accompanied by an orbital ordering and significant lattice distortions (see also Ref. [5]). In order to identify these transitions we used results of a neutron diffraction study [3] for  $x = 0.125$ . (3) Structural orthorhombic ( $\text{O}^*$ ) to rhombohedral (R) phase transitions at  $T_s$  and  $0 \leq x \leq 0.22$  (filled up triangular). The  $\rho(T)$  anomaly at  $T_s$  is too weak in this case and can be seen only for a suitable scale (inset in Fig. 1). (4) Transitions to a polaron ordered state (P) at  $T_p$  and  $0.1 \leq x \leq 0.15$ , determined

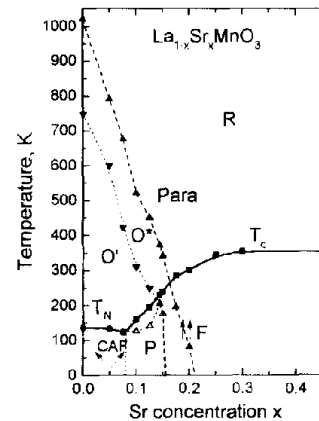


Fig. 3. Phase  $T$ - $x$  diagram of  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ : (F) ferromagnetic phase, (CAF) canted antiferromagnetic phase, ( $\text{O}'$ ) Jahn–Teller orthorhombic phase, ( $\text{O}^*$ ) orthorhombic (pseudocubic) phase, (P) polaron ordered state.

by freezing of holes on the lattice sites (open up triangular). Such transitions were observed by neutron scattering in  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$  for  $x = 0.1$  and  $0.15$  [2].

The results of the submillimeter measurements of the dielectric permittivity  $\epsilon'(T)$  and conductivity  $\sigma'(T)$  at the frequency  $13.3 \text{ cm}^{-1}$  for  $x = 0.125$ , combined with static conductivity  $\sigma_{\text{DC}}(T)$ , AC magnetic susceptibility  $\chi_{\text{AC}}(T)$  and magnetization  $M(T)$  measurements, are displayed in Fig. 2. Curves 1,2 for the  $\epsilon'$  and  $\sigma'$  correspond to the radiation polarization for a minimum and maximum transmission, respectively, that occurs due to a noticeable anisotropy of a crystal lattice. Observed features in the  $\epsilon'(T)$  and  $\sigma'(T)$  at  $T_s$ ,  $T_c$  and  $T_p$ , indicated by arrows, are in a reasonable agreement with corresponding static data. A significant increase of the  $\epsilon'(T)$  in the polaron ordered state (Fig. 2a) indicates on a noticeable transformation of a crystal lattice and, probably, a change of electron structure, which was observed recently by optical measurements [6]. We note also that  $\sigma'$  remains large enough at low temperatures in spite of localization of charge carriers which results in a significant lowering of the  $\sigma_{\text{DC}}$  (Fig. 2b). A similar behavior of the  $\epsilon'(T)$  and  $\sigma'(T)$  was also observed for  $x = 0.1$  and  $0.15$  (see also Ref. [7]).

The observed anomalies in  $\rho$ ,  $\chi_{\text{AC}}$ ,  $M$ ,  $\epsilon'$  and  $\sigma'$  at various phase transitions are displayed in the form of the  $T$ - $x$  phase diagram in Fig. 3, where solid lines correspond to magnetic transitions and dotted and dashed lines to the structural ones. The polaron ordered phase P is located approximately between  $x = 0.08$  and  $0.15$ . The phase diagram presents a general picture of the phase transitions in  $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$  and in the whole, it agrees with corresponding data of Ref. [2,3].

This work was supported in part by the RFBR (96-02-18091, 97-02-17325, 96-15-96577).

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