

# Development of a Watt-level gamma-ray source based on high-repetition-rate inverse Compton scattering



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## ARTICLE INFO

### Article history:

Received 5 December 2016

Received in revised form 4 March 2017

Accepted 23 March 2017

Available online 8 April 2017

### Keywords:

Inverse Compton scattering

Gamma rays

Electron beam

Beam dynamics

## ABSTRACT

A high-brilliance ( $\sim 10^{22}$  photon  $s^{-1} mm^{-2} mrad^{-2}/0.1\%$ ) gamma-ray source experiment is currently being planned at Fermilab ( $E_\gamma \simeq 1.1$  MeV). The source implements a high-repetition-rate inverse Compton scattering by colliding electron bunches formed in a  $\sim 300$ -MeV superconducting linac with a high-intensity laser pulse. This paper describes the design rationale along with some of technical challenges associated to producing high-repetition-rate collision. The expected performances of the gamma-ray source are also presented.

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## 1. Introduction

High-flux, quasi-monochromatic,  $\gamma$ -ray sources could have widespread range of applications including in Nuclear Astrophysics, Elementary Particle Physics and Homeland security. In the later class of application, developing compact gamma-ray source capable of producing large flux could enable the rapid screening of cargos for fissile material detection. The need for monochromatic  $\gamma$  rays along with the requirement for a small-footprint source have motivated the exploration of particle-accelerator-based sources employing inverse Compton scattering (ICS) [1,2]. This development path is further supported by the increasing number of compact GeV-class electron sources based on laser-plasma wakefield accelerators (LPAs) available at various laboratories worldwide [3]. LPAs have so far been employed to generate  $\gamma$  rays with impressive brilliance but with restricted photon flux due to their low operating frequencies (typically 10 Hz) limited by the current solid-state-laser technologies [4,5]. Such a limitation of the LPAs is currently being addressed by several groups while the development of high-repetition rate interaction region could be performed at available state-of-the-art linear accelerators. Based on this observation, a collaboration between

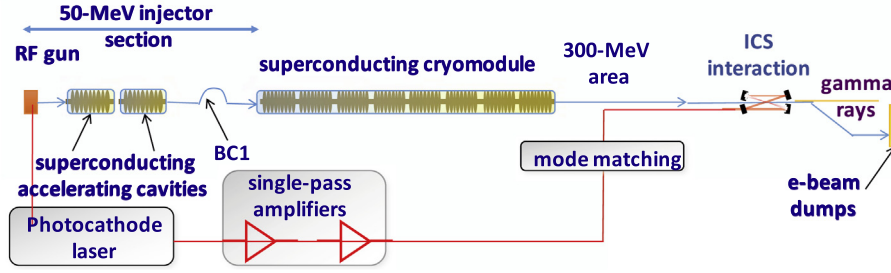
Fermilab, Northern Illinois University and RadiaBeam LLC is currently designing an experiment aimed at demonstrating high-repetition-rate ICS using the Fermilab Accelerator Science & Technology (FAST) facility based on a superconducting linear accelerator [6]. The present paper summarizes the design rationale and expected performances of the source under design.

## 2. Overview of the $\gamma$ -ray source

Fig. 1 provides an overview of the source concept to be driven by the 300-MeV electron beam available at the FAST facility. In brief the electron bunches are produced via photoemission from a semiconductor photocathode located in a  $1 + \frac{1}{2}$ -cell radiofrequency (RF) gun. The bunches are photoemitted by impinging an ultraviolet (UV) laser pulse obtained via frequency quadrupling of an amplified infrared (IR) laser pulse ( $\lambda = 1053$  nm) produced from a Nd:YLF laser system. The formed bunches are further accelerated by two TESLA-type superconducting RF (SRF) cavities. The RF gun and cavities operate at 1.3 GHz and are capable of forming 1-ms bunch trains at 5-Hz frequency and containing up to 3000 bunches (corresponding to a 3-MHz bunch frequency). The 50 MeV beam can be manipulated and diagnosed before injection in an international-linear-collider (ILC) type accelerating cryomodule. The cryomodule incorporates 8 SRF cavities with demonstrated average accelerating gradient of  $\bar{G} \simeq 31.5$  MeV/m resulting in a maximum beam energy of  $\sim 300$  MeV. The beam is further transported along a  $\sim 70$  m beamline and finally directed to the

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**Fig. 1.** Overview of the  $\gamma$ -ray source under consideration at the FAST facility. The “BC1” and “ICS” labels respectively refer to the bunch compressor and inverse-Compton scattering.

ICS interaction point (IP). Downstream of the IP, the electron beam is then sent to a high-power beam dump.

The laser pulse used in the ICS interaction is derived from the photocathode laser system. The unspent IR laser energy downstream of the IR-to-UV conversion process is conditioned, further amplified, and directed to a passive coherent enhancement cavity for final amplification. The laser pulses collide head-on with the electron bunches with a collision frequency of 3 MHz within the 1-ms train thereby yielding  $\gamma$ -ray pulses with similar format. The energy of the backscattered  $\gamma$ -ray photon is  $E_s \approx \hbar\omega_s$  where  $\hbar$  is the reduced Planck’s constant, and the upshifted backscattered-photon frequency is  $\omega_s(\theta) = \frac{4\gamma^2\omega_l}{1+a_0^2/2+\gamma^2\theta^2}$ , with  $\hbar\omega_l \approx 1.2$  eV and  $a_0$  being respectively the laser photon energy and normalized potential, and  $\gamma$  the electron-beam Lorentz factor. The angle  $\theta$  is the direction of observation referenced to the electron-beam direction. Therefore the 300-MeV beam available at FAST will support the generation of  $\gamma$  rays with  $E_s \leq 1.5$  MeV. Considering a laser pulse energy  $\mathcal{E}_l \sim 0.5$  J and a focused transverse size  $\sigma_l \approx 30$   $\mu\text{m}$ , we obtain  $a_0^2/2 \approx 7.6 \times 10^{-3} \ll 1$  thereby confirming that nonlinear effects are insignificant for our operating parameters. Owing to its narrow bandwidth ( $\delta\lambda \sim 0.2$  nm) the laser transform-limited rms pulse duration is 3 ps close to the measured value of 3.8 ps.

### 3. Expected performances

Applications of  $\gamma$  rays generally require high photon flux, narrow spectral bandwidth and high brilliance. The brilliance can be expressed as a function of the laser and electron-beams parameters [7] as

$$B_s \propto \frac{N_l}{\sigma_l^2} \gamma^2 \frac{N_e}{\tau_e \varepsilon_\perp^2}, \quad (1)$$

where  $N_l$  (resp.  $N_e$ ) are the number of photons (resp. electron) in the laser pulse (resp. electron bunch),  $\tau_e$  the electron-bunch duration and  $\varepsilon_\perp$  its normalized transverse emittance (the electron beam and laser pulse are cylindrically symmetric). Likewise, the backscattered-photon dose  $N_s$  and relative spectral bandwidth of the scattered pulse  $\delta\omega_s/\omega_s$  can be respectively parameterized as

$$N_s \approx \frac{N_e N_l \sigma_T}{2\pi(\sigma_e^2 + \sigma_l^2)}, \quad \text{and} \quad \frac{\sigma_{\omega_s}}{\omega_s} \approx 2 \frac{\varepsilon_\perp^2}{\sigma_e^2}, \quad (2)$$

where  $\sigma_e$  is the electron-beam transverse rms size,  $\sigma_T$  the Thompson’s cross section and  $\hbar\sigma_{\omega_s}$  is the rms spectral spread of the backscattered radiation. Other sources contributing to  $\sigma_{\omega_s}/\omega_s$  includes the laser spectral bandwidth, the laser divergence (diffraction), and electron beam’s fractional momentum spread. The cumulative contribution of these effects is found to be negligible ( $\sim 0.2\%$ ) for our operating regime.

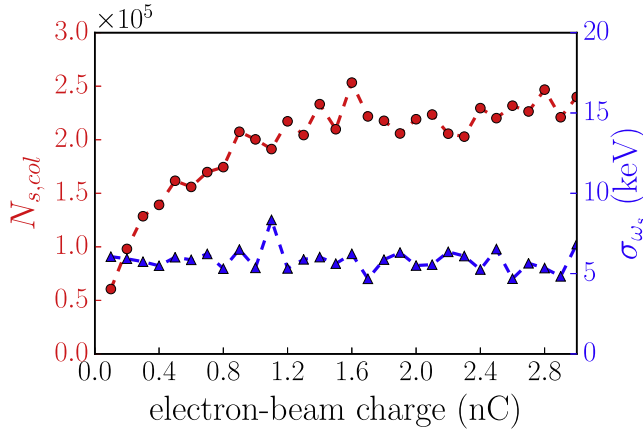
We modeled the ICS process using the program COMPTON [8] which provides a 3D treatment of the laser-electron interaction

in the time and frequency domains. The software also includes the nonlinear corrections to the Thompson cross section to correctly account for the electron recoil at the time scale of the interaction process. To gain further confidence in our modeling, we also employed the program CAIN [9] which is based on the more general Klein-Nishina cross-section for Compton-scattering process and consequently accounts for the quantum-electrodynamical (QED) effects. In our case, taking  $\gamma \approx 500$ , we find that the laser photon energy in electron-beam frame verifies  $2\hbar\gamma\omega_l \approx 0.6$  keV  $\ll m_e c^2 = 511$  keV (where  $m_e c^2$  is the electron’s rest mass) so that QED effects are expected to be insignificant. The codes COMPTON and CAIN were benchmarked and found to be in reasonable agreement over the considered range of parameters [10].

The COMPTON program was used to devise electron-beam and laser parameters such that scattered-pulse brightness, dose and bandwidth are optimized taking Eqs. 1 and 2 as guidelines. In practice only a small fraction of the scattered photons (i.e. the most energetic) are employed in front-end experiments. The selection is typically accomplished using transverse collimators exploiting the correlation between photon energy and angular spread (red shifting occurs off axis). The collimation results in the selection of a small fraction of the backscattered radiation. For example targeting  $\sigma_{\omega_s}/\omega_s \approx 0.5\%$  requires the collection angle to be  $\theta_{col} \approx \sqrt{\frac{\sigma_{\omega_s}}{2\omega_s}} \frac{1}{\gamma} \approx 0.104$  mrad [7].

For each value of the electron bunch charge the normalized emittance is evaluated with the scaling relation  $\varepsilon_\perp [\mu\text{m}] \approx 2.94 Q^{0.69} [\text{nC}]$  determined from previous beam dynamics studies [11]. This value of emittance is then used to determine the electron beam transverse size by constraining  $\frac{\varepsilon_\perp^2}{\sigma_e^2} = 2.5 \times 10^{-3}$ ; see also Eq. 2. The laser parameters are held constant and mimic the experimentally-achievable values. The number of scattered photons and associated energy bandwidth appear in Fig. 2 as functions of electron beam charge. Given the aforementioned constraint, a bandwidth  $\hbar\sigma_{\omega_s} \approx 6$  keV is maintained over the considered range of charges. The fraction of scattered photon within this bandwidth increases with charge and eventually saturate (for given our fixed set of laser parameters). Following Ref. [12] we evaluate the spectral density, defined as  $S \equiv \frac{N_{ph}}{\hbar\sigma_{\omega_s} \sigma_l}$ . Considering the case of  $Q = 0.8$  nC, we find  $S \approx \frac{1.7 \times 10^5}{6.67}$  photons/s/eV per collision. Accounting for the capability of FAST to produce 15,000 bunches per second we get  $S \sim \mathcal{O}(10^5)$  photons/s/eV comparable with the value reported in [13]. Finally, it should be pointed out that only  $\sim 0.5\%$  of the total number of backscattered photons are contained within the 0.6% relative bandwidth.

Table 1 compares the properties of the collimated and total backscattered radiation. The angular collimation leads to  $\sigma_{\omega_s}/\omega_s \approx 0.6\%$ . It should be pointed out that the Table 1 provides the scattered pulse properties for a single-bunch collision. The brightness and dose should be multiplied by 15,000 in order to



**Fig. 2.** Numbered of scattered photons collected within an angular acceptance  $\theta_{col} = 0.104$  mrad with associated rms spectral bandwidth  $\sigma_{\omega_s}$  as functions of electron pulse charge.

**Table 1**

Comparison of the single-bunch  $\gamma$ -ray parameters computed over the 0.5%-fraction of the most energetic photons and for all (“100%”) the backscattered photons. The bunch charge is 0.8 nC.

Collected radiation	0.5%	100%
$E_s$ (std. unit)	$1.8 \times 10^{18}$	$2.8 \times 10^{17}$
$N_s$ (phot.)	$1.7 \times 10^5$	$4.0 \times 10^7$
$\delta\omega_s/\omega_s$ (%)	0.60	31.3
$\theta_{max}$ (mrad)	0.104	10.0
$\sigma_{x'}/\sigma_{y'}$ (mrad)	0.032/0.032	1.18/1.56
duration $\sigma_t$ (ps)	2.92	2.92

scale these parameters to the multi-bunch collision case (i.e. when 5-Hz bunch train composed of 3000 bunches are employed). Owing to the excellent multi-bunch stability of the superconducting linac the spectral bandwidth remains unaffected in the multi-bunch operation case.

Table 2 provides the final set of electron- and laser-beam parameters while Fig. 3 summarizes the angular-spectral and transverse distribution and energy spectrum for all backscattered photons. The distributions displayed in Fig. 3 and parameters gathered in Table 1 are obtained from the program COMPTON and include the impact of the laser divergence and spectral bandwidth along with electron-beam emittance and energy-spread effects. The rms divergence listed in Table 1 (for the “100%” case) and associated to the distribution in Fig. 3 (b) is consistent with the full-width half maximum (FWHM) of the theoretical distribution (convolution between a cubed Lorentzian and a Lorentzian to fifth power [8])  $\theta_{FWHM} \approx 0.635/\gamma \approx 1.3$  mrad.

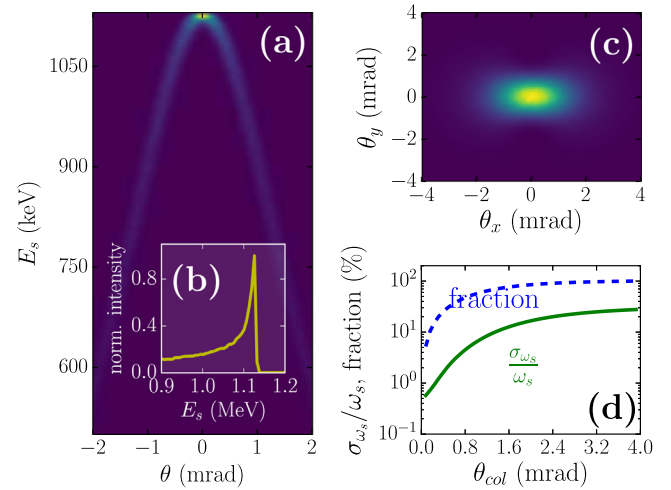
#### 4. Interaction region

At the IP, final-focus high-gradient quadrupole magnets are needed to produce tight electron-beam waist at the interaction

**Table 2**

Summary of electron-beam and laser-pulse parameters used to produce Tables 1 and Fig. 3. The duration, spot size, energy spread and bandwidth are given as rms values.

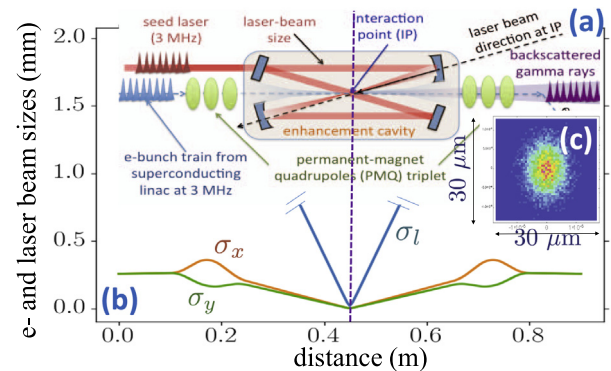
Electron beam		laser pulse	
Beam energy (MeV)	250.0	wavelength (nm)	1053
Beam charge (pC)	800.0	pulse energy (J)	0.5
Energy spread (%)	0.1	bandwidth (nm)	0.2
Duration (ps)	3.0	duration (ps)	3.0
Beam size ( $\mu\text{m}$ )	36.0	spot size ( $\mu\text{m}$ )	30.0
Emittances ( $\mu\text{m}$ )	1.80		



**Fig. 3.** Simulated angular-spectral (a) and transverse (c) distributions of the  $\gamma$ -ray generated using the parameters in Table 1 for all generated photons. Inset (b) gives the integrated spectrum associated to (a), and plot (d) presents the relative spectral bandwidth (solid line) and fraction of collected radiation (dashed line) as functions of the collection angle.

point (IP), the optical enhancement cavity and some diagnostics. The magnetic lattice combines two quadrupole-magnetic triplets (each composed of three magnets) respectively located upstream and downstream of the IP; see Fig. 4(a). The upstream triplet is used to focus the beam and match its transverse size to the laser size while the downstream triplet rematches the electron beam parameters to the downstream accelerator beamline for transport up to the beam dump. The focusing magnets are Halbach-type permanent-magnet quadrupoles (PMQs) [14] capable of producing the required field gradient of  $g \approx 300 \text{ T m}^{-1}$ . The use of such strong PMQ yields a meter-scale interaction region compatible with possible future coupling to LPAs. An example of evolution of the electron-beam transverse size appears in Fig. 4(b) and the associated transverse distribution at the IP is shown Fig. 4(c). The PMQs are transversely aligned using an in situ laser-based alignment, whereas their axial positions is adjustable to provide some flexibility. Further control on the interaction-point waist is provided by electromagnet quadrupole magnets available upstream (and downstream) of the interaction chamber.

As previously pointed out, the laser beam used for the ICS process originates from the IR photocathode laser system. Our current



**Fig. 4.** Overview of the final-focus area (a) with associated rms beam sizes evolution (b). In (b)  $\sigma_{x,y}$  are the electron-beam transverse beam sizes while  $\sigma_l$  denotes the laser spot size evolution. The origin of the distance ordinate is  $z = 120$  m from the photocathode surface. Inset (c) shows a density plot of the electron beam transverse density at the IP [located at distance of 0.43 m in (b)].

approach is to condition and to amplify the IR pulses using single-pass amplifiers to bring the laser pulse energy from  $\sim 50 \mu\text{J}$  to  $\sim 5 \text{ mJ}$  eventually and then to further increase the pulse energy to Joule level using a passive enhancement cavity. This type of cavities were intensively studied in the recent years and intensity amplification factors of several thousands were reported [15,16]. An amplified IR laser pulse enter in such a cavity, illustrated in Fig. 4(a), from the left side and undergo an integer number of round trips, a new laser pulse arrives at the cavity entrance. Assuming a optimal mode matching the amplitude of the two pulses add up coherently. Ideally, the intensity build-up is only limited by the losses in the mirrors. The process is repeated with the subsequent IR pulses. Once the steady state is reached, the intra-cavity pulse intensity is  $I = \frac{1-R_1}{1-\sqrt{(R_1 R_2)^n}} I_L$  where  $R_i$  ( $i = 1, 2$ )

are the reflectivity coefficients of the two cavity mirrors,  $n$  is the number of roundtrips performed by the laser pulse inside the cavity during the time interval between consecutive bunches coming from the laser system and  $I_L$  is the intensity of the laser pulses just before the entrance in the cavity. Consider the case of a conventional Perot-Fabry resonator with commercially-available mirrors with  $R_1 = 0.999, R_2 = 0.99995$  and  $n = 1$  the total gain is  $\Gamma \equiv \frac{I}{I_L} \approx 3626$ . In such a two-mirror configuration, the injection mirror, must have a much lower reflectivity in order to allow the outside laser pulses to couple to the cavity. This later feature results in significant losses when a large  $n$  is required. In our case where the pulse repetition rate is 3 MHz (corresponding to a pulse separation of  $\sim 100 \text{ m}$ ), requiring the cavity to fit within a 2-m footprint would imply  $n = 25$  yielding a poor gain  $\Gamma \approx 6$ . Consequently an alternative approach is required. One candidate being explored is a Herriott cell [17] combined with a four-mirror bow-tie cavity. Calculations performed for such a configuration indicate a gain  $\Gamma \sim 100$  is attainable but further optimization is needed. The cavity will operate in burst mode with 1.2-ms pulses yielding  $\sim 10 \text{ kW}$  of stored power.

## 5. Summary

The coupling of an electron beam originating from a 250-MeV SRF pulsed linac with a high-repetition rate infrared laser is expected to provide  $\gamma$  rays with maximum energy up to  $E_s \sim 1.1 \text{ MeV}$ , average brilliance of  $\langle B_s \rangle \simeq 2.7 \times 10^{22} \text{ photon s}^{-1} \text{ mm}^{-2} \text{ mrad}^{-2}/0.1\%$ , photon flux of  $\langle N_s \rangle \simeq 2.5 \times 10^9 \text{ photon/s}$  and relative spectral bandwidth of  $\delta\omega_s/\omega_s \simeq 6.0 \times 10^{-3}$ . These predicted parameters corresponds to an average power of  $\sim 0.44 \text{ mW}$  and are expected to foster a wide range of applications

such as fissile-material detection, exploration and development of gamma-ray optics, or possibly enable the quantum control of nuclear states.

## Acknowledgement

We are grateful to Dr. W. Brown for sharing his inverse Compton-scattering simulation software COMPTON. This work was sponsored by the DNDO award 2015-DN-077-ARI094 to Northern Illinois University. Fermilab is operated by Fermi Research Alliance, LLC. for the U.S. Department of Energy under contract DE-AC02-07CH11359.

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