



# All-fiberized, mode-locked laser at 1.95 $\mu\text{m}$ using copper chalcogenide $\text{Cu}_2\text{Te}$ -based evanescent field interaction

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## ABSTRACT

For the first time, a copper telluride ( $\text{Cu}_2\text{Te}$ ) based saturable absorber (SA) was demonstrated for mode-locking in a thulium-doped fiber laser. Layers of  $\text{Cu}_2\text{Te}$  were deposited onto a side polished fiber (SPF) using the optical deposition technique, forming an all-fiber SA device. The nonlinear optical properties of the SA device were investigated based on the balanced twin-detector measurement scheme. The  $\text{Cu}_2\text{Te}$ -deposited SPF SA was capable of generating stable soliton mode-locked pulses with a pulse duration of 1.58 ps at a center wavelength of 1951 nm with a signal to noise ratio (SNR) of more than 60 dB, as well as a peak power of 250 mW and a pulse energy of 0.39 nJ.

## 1. Introduction

Soliton lasers operating at the 2  $\mu\text{m}$  region are promising light sources and have received wide research interest due to their versatile applications in long-range light detection and ranging (LIDAR), laser surgery, material processing and laser spectroscopy [1–4]. Mode-locked fiber lasers are simple and exhibit significant reliability and environmental stability as compared to free-space optics-based lasers. Optical fiber-based 2  $\mu\text{m}$  lasers use thulium (Tm) [5], holmium (Ho) [6] or Tm–Ho co-doped fibers [7] as the primary gain medium. Various active/passive schemes have been employed in fiber lasers to generate mode-locked pulses in the picoseconds and femtosecond range [8]. In contrast to active scheme which requires external and complicated device [9–11], passive scheme offers advantages of a compact, flexible and cost-effective mode-locked system through the use of nonlinear optical element, called a saturable absorber (SA).

Two-dimensional (2D) materials such as graphene [12–15], black phosphorus (BP) [16–18] and topological insulators (TIs) [19–21] have been widely employed as SAs for demonstrating passively mode-locked fiber laser owing to their high optical nonlinearity property, broadband saturable absorption and ultrafast carrier dynamics. Nevertheless, the aforementioned materials have some drawbacks. For example, graphene possesses weak light–matter interaction due to its low modulation depth (typically  $\sim 1.2\%$  per layer [22]), whereas BP can be easily oxidized when exposed to a humid environment [23]. Therefore, the exploration of new SAs has always attracted a lot of attention. Recently,

transition metal dichalcogenides (TMDs), including  $\text{MoS}_2$ ,  $\text{WS}_2$ ,  $\text{MoSe}_2$ ,  $\text{WSe}_2$  and  $\text{ReS}_2$  [24–29] have been widely employed as SAs due to their advantages of wide absorption range and its exotic optical and electrical properties which are different from their bulk counterpart. More recently, binary chalcogenide semiconductors which arises as a result of quantum confinement effect have received substantial attention due to their size-dependent optical and physical properties. Copper chalcogenides have been widely used as solar cells, optical recording materials and optical filters. This group of materials possesses an interesting localized surface plasmon resonance (LSPR) effect in the near infrared (NIR), which is very useful for surface-enhanced Raman scattering (SERS) and photoacoustic imaging applications [30]. Of these copper chalcogenides, copper telluride ( $\text{Cu}_x\text{Te}$ ) in particular is attractive for a wide range of applications owing to its direct bandgap ranging from 1.1 eV to 1.5 eV [29] as well as large thermoelectric power.  $\text{Cu}_x\text{Te}$  has a wide range of compositions with each composition having a different crystal structure, depending on the value of  $x$  ( $1 < x < 2$ ). For instance, the  $\text{Cu}_2\text{Te}$  has an optical band gap of 1.1 eV [31]. The SA device based on the evanescent wave interaction has drawn extensive attention in demonstrating mode-locked fiber laser due to its robustness, large nonlinear interaction length, high degree of compatibility to the SMF and low optically induced thermal damage [32]. Various 2D materials have been utilized to enable mode-locking operation through the nonlinear interaction with evanescent field [25,33–39]. However, the saturable absorption properties of copper chalcogenides have been

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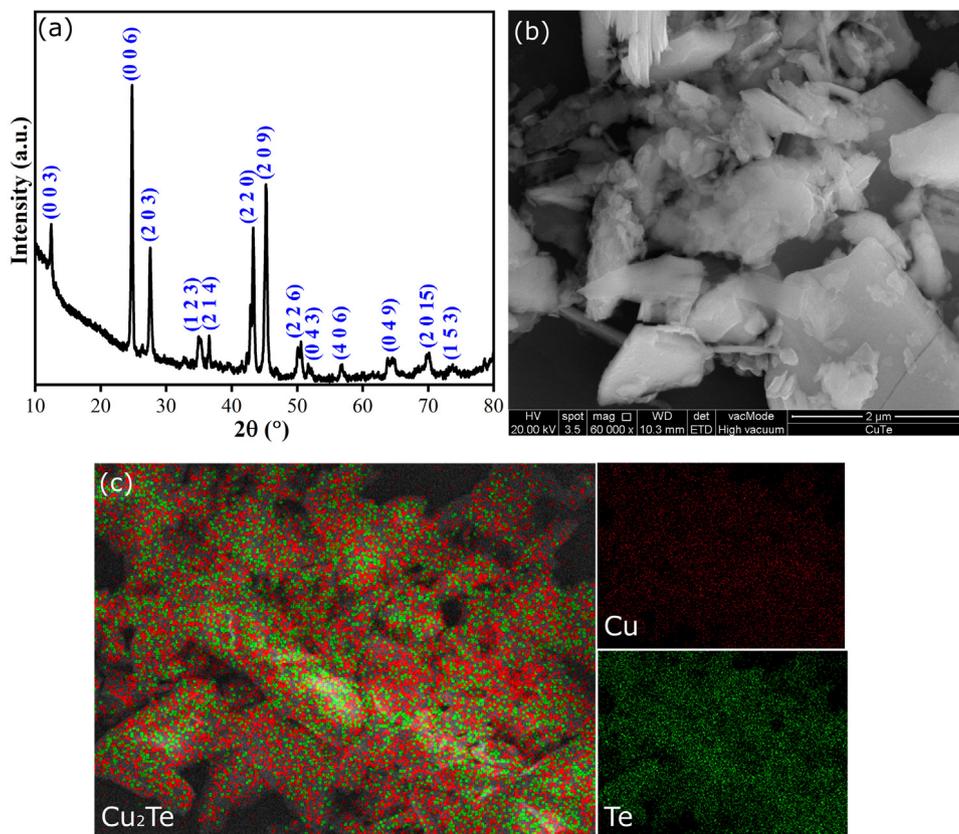


Fig. 1. Characterization of  $\text{Cu}_2\text{Te}$ . (a) XRD spectrum; (b) FESEM image; and (c) selected area EDX elemental mapping.

investigated. Thus, in this work, a  $\text{Cu}_2\text{Te}$ -based evanescent field interaction is demonstrated to produce mode-locked pulses. The  $\text{Cu}_2\text{Te}$  is synthesized using simple solution method and deposited on the side polished fiber (SPF) using optical deposition method. By utilizing the  $\text{Cu}_2\text{Te}$ -deposited SPF, stable picoseconds soliton pulses are delivered from the Tm-doped fiber laser (TDFL) with >60 dB of signal-to-noise ratio (SNR). At maximum pump power of 462 mW, the laser average output power is 3.2 mW, with a corresponding pulse energy and peak power of about 0.39 nJ and 250 W. This is the first time, to the best of the author's knowledge, of the use of a  $\text{Cu}_2\text{Te}$  on a SPF as an SA. This work proposed a practical photonic device which may find wide applications in ultrafast photonics.

## 2. $\text{Cu}_2\text{Te}$ -deposited side polished fiber fabrication and characterization

A  $\text{Cu}_2\text{Te}$  powder was purchased from 2D Semiconductors, while polyvinyl alcohol (PVA,  $M_w \sim 31,000$ ) was bought from Sigma Aldrich. All chemicals are used as received without further purification. The  $\text{Cu}_2\text{Te}$  was synthesized using a simple solution method. Initially, 100 mg of  $\text{Cu}_2\text{Te}$  powder was dispersed in 10 ml of deionized (DI) water. The dispersion was sonicated in a bath sonicator for 10 h at room temperature to form a homogeneous solution. In order to remove large agglomerations, the  $\text{Cu}_2\text{Te}$  dispersion was then centrifuged at 4000 rpm for 10 min and the supernatant was collected.

The X-ray diffraction (XRD) investigation was carried out using a Malvern Panalytical Empyrean, recorded in the range of  $2\theta = 5 - 70^\circ$ . The XRD pattern obtained from the  $\text{Cu}_2\text{Te}$  powder is depicted in Fig. 1(a). A few sharp peaks can be observed in the XRD spectrum of  $\text{Cu}_2\text{Te}$  located at  $2\theta$  of 12.45, 24.80, 27.55, 43.48, and  $45.23^\circ$  which corresponds to the (0 0 3), (0 0 6), (2 0 3), (2 2 0), and (2 0 9) planes that matched with the standard diffraction pattern of hexagonal  $\text{Cu}_2\text{Te}$  (JCPDS 96-901-1862). The existence of the sharp peaks and the

absence of impurity peaks indicate the high crystallinity and purity of  $\text{Cu}_2\text{Te}$ . The surface morphology and the energy dispersive X-ray (EDX) elemental analysis of  $\text{Cu}_2\text{Te}$  were investigated using a JEOL JSM 7600F field emission scanning electron microscope (FESEM) together with an EDX. It is evident from the FESEM image presented in Fig. 1(b) that the  $\text{Cu}_2\text{Te}$  is composed of few flake-like sheets that overlap with one another and the flakes size is mainly  $\sim 1 \mu\text{m}$ . The EDX elemental mapping images of individual copper (Cu) and tellurium (Te) as well as the overlay image of these two elements are illustrated in Fig. 1(c). It is noteworthy that the  $\text{Cu}_2\text{Te}$  consists of Cu and Te elements, hence further confirming the formation of  $\text{Cu}_2\text{Te}$ .

After the successful exfoliation of the  $\text{Cu}_2\text{Te}$  powder, the  $\text{Cu}_2\text{Te}$ -deposited SPF was fabricated and its saturable absorption properties was investigated. The  $\text{Cu}_2\text{Te}$  sheets were transferred onto the polished region of the commercial SPF (Phoenix Photonics) using the optical deposition method. The polished region area had a dimension of 17 mm in length and width of about  $125 \mu\text{m}$  (the diameter of the cladding). This area is about  $2 \mu\text{m}$  above the core-cladding interface. Slow evaporation at ambient temperature resulted in the  $\text{Cu}_2\text{Te}$ -deposited SPF, as shown in Fig. 2(a). The saturable absorption of the device was realized through the nonlinear interaction of the  $\text{Cu}_2\text{Te}$  sheets with the evanescent field of light in the SPF. The existence of the evanescent field in the SPF can be observed by injecting the visible light, as shown in the inset of Fig. 2(b). The nonlinear saturable absorption of the  $\text{Cu}_2\text{Te}$ -deposited SPF was investigated based on a balanced twin-detector measurement method. The measurement was carried out using a commercial ELMO femtosecond laser (central wavelength: 1560 nm, repetition rate: 100 MHz, pulse duration:  $\sim 150$  fs). The intensity dependent absorption is fitted by  $\alpha(I) = \left[ \frac{\Delta\alpha}{1 + \left(\frac{I}{I_{sat}}\right)} \right] + \alpha_{ns}$ , where  $\Delta\alpha$ ,  $I_{sat}$  and  $\alpha_{ns}$  represent modulation depth, saturation intensity and non-saturable loss, respectively. The modulation depth of the SA was measured to be 25%, whereas the

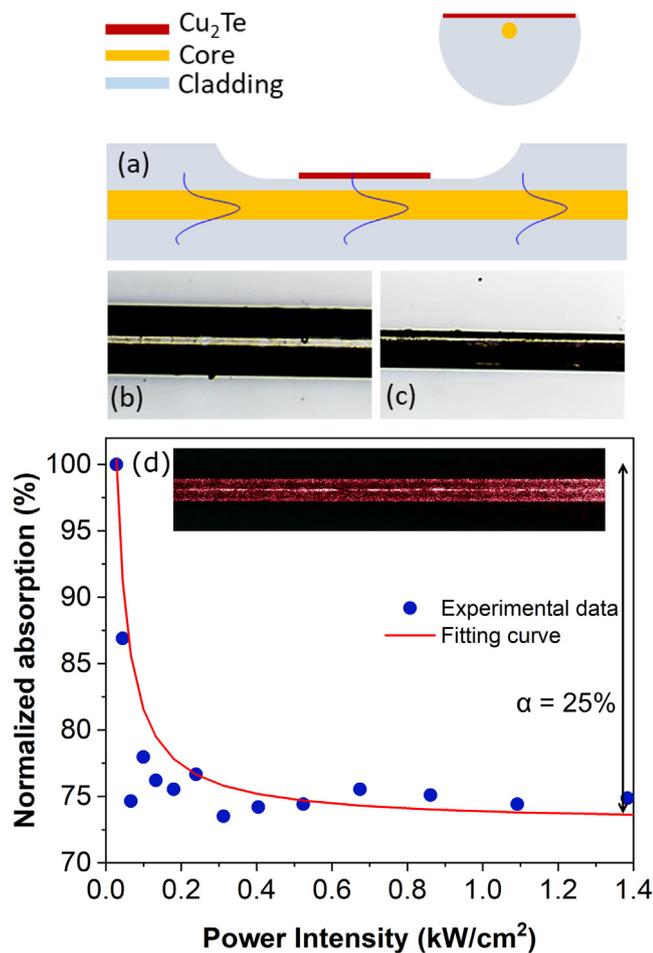


Fig. 2. (a) Illustration of the  $\text{Cu}_2\text{Te}$ -deposited SPF; (b) photomicrograph of SPF before and (c) after polished; and (d) nonlinear absorption curve. Inset shows the evanescent field from SPF after injecting with red light source. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

saturation intensity is  $0.08 \text{ kW/cm}^2$  which corresponded to the peak power of  $0.025 \text{ mW}$  and silica fiber cross-sectional area of  $3.24 \times 10^{-7} \text{ cm}^2$ .

### 3. Experimental setup

A ring TDFL cavity was constructed to generate mode-locked pulses in  $2.0 \mu\text{m}$  region as shown in Fig. 3. In order to achieve this, the  $\text{Cu}_2\text{Te}$ -deposited SPF was incorporated in the TDFL cavity. The thulium-doped fiber amplifier consists of a piece of 4 m TmDF200 thulium-doped fiber (TDF) with a group velocity dispersion (GVD) parameter of  $-21.9 \text{ ps}^2/\text{km}$  at  $1940 \text{ nm}$  and peak absorption of  $20 \text{ dB/m}$  at  $1550 \text{ nm}$ . The TDF was dual-pumped by  $1550 \text{ nm}$  laser diodes (LDs) through the  $1550 \text{ nm}$  ports of two  $1550/2000 \text{ nm}$  wavelength division multiplexers (WDMs). In addition to the fiber amplifier, the cavity consists of an isolator (ISO) to force unidirectional operation of the ring cavity, polarization controller (PC) to adjust the polarization state of the circulating light. A 90:10 fiber-pigtailed optical coupler (OC) was added to the cavity for temporal and optical analysis. All of the fibers in the cavity, with the exception of the TDF, constitute of single mode fibers (SMF-28s) with a GVD of  $-71.4 \text{ ps}^2/\text{km}$  at  $1940 \text{ nm}$ . The total cavity length was  $25.5 \text{ m}$  and net cavity dispersion is estimated as  $-1.62 \text{ ps}^2$  (anomalous dispersion). The performance of the laser was examined using a Yokogawa AQ6375 optical spectrum analyser

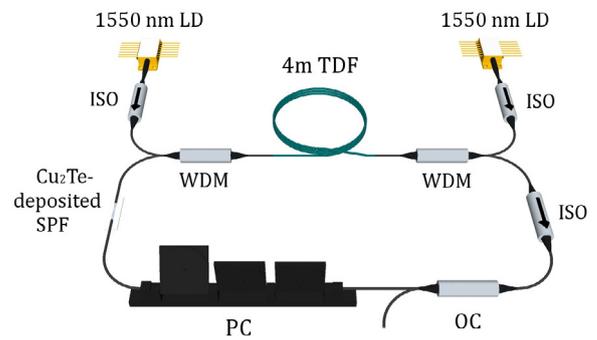


Fig. 3. Experimental setup of the passively mode-locked TDFL.

(OSA), a  $5 \text{ GSa/s}$  Rohde & Schwarz RTM3002 oscilloscope (OSC) and a  $7.8 \text{ GHz}$  Anritsu MS2683A radio frequency spectrum analyser (RFSA) coupled with a  $10 \text{ GHz}$  photodetector and an optical autocorrelator (APE PulseCheck).

### 4. Results and discussion

By carefully adjusting the PC in the laser cavity, stable mode-locked pulse train was observed at threshold pump power of  $87 \text{ mW}$ . The pulse train as shown in Fig. 4(a) was observed at pulse repetition frequency of  $8.1 \text{ MHz}$ , which corresponds well to the cavity round trip time of  $124 \text{ ns}$ . This verifies the fundamental operation of the mode-locked laser. Fig. 4(b) shows the enlarged optical spectrum of the generated pulses, centered at  $1951 \text{ nm}$  with a full-width half maximum (FWHM) of  $2.6 \text{ nm}$ . As seen from the inset of Fig. 4(b), Kelly sidebands were observed which are characteristic of an all-anomalous dispersion or solitonic mode-locked laser. The autocorrelation trace as shown in Fig. 4(c) is well-fitted with a  $\text{Sech}^2$  profile, showing a deconvolved pulse duration of  $1.58 \text{ ps}$ .

The time-bandwidth product is  $0.323$ , which is just slightly higher than the theoretical limit of  $0.315$ , indicating that the laser cavity has a low chirp. Fig. 4(d) depicts the radio frequency spectrum measured around fundamental frequency with a resolution bandwidth (RBW) of  $300 \text{ Hz}$ . A sharp and clean peak is observed at  $8.1 \text{ MHz}$  with SNR of  $>60 \text{ dB}$ , demonstrating good mode-locking operation. RF spectrum, on a wider range span of  $1 \text{ GHz}$  as shown in the inset is observed free from spectral modulation, indicating no Q-switching mode-locking. Fig. 5 shows the average output power as a function of pump power. At the maximum pump power of  $462 \text{ mW}$ , the laser average output power was measured around  $3.2 \text{ mW}$ , with a peak power of  $250 \text{ W}$  and a pulse energy of  $0.39 \text{ nJ}$ . The pulse energy as measured in this work is high with no pulse splitting observed. This could be explained using the soliton area theorem [40], where mode-locked fiber lasers operating at  $2 \mu\text{m}$  can typically support high energy soliton pulse due to the low nonlinearity and large dispersion of silica fiber. Besides, the use of evanescent-wave-type SA such as SPF can also be important factor to enhance the soliton pulse energy [12,41,42]. A slope efficiency of  $0.8\%$  was obtained from this laser. The low power efficiency could arise from the high non-saturable loss of the  $\text{Cu}_2\text{Te}$ -deposited SPF, possibly due to the large polishing depth and roughness of the polished surface which would result in severe scattering losses. On top of this, there are also losses from the components used in the cavity, which will also reduce the efficiency of the laser.

The long-term stability of the mode-locking operation was further evaluated by observing the output spectrum for 4 h. The output spectrum was recorded every 30 mins over a period of 4 h under fixed experimental condition such as pump power of  $192 \text{ mW}$ , as shown in Fig. 6(a). Neither the center wavelength drifted nor a new wavelength component was generated during the measurement period. Moreover, no significant degradation of the optical spectra is observed,

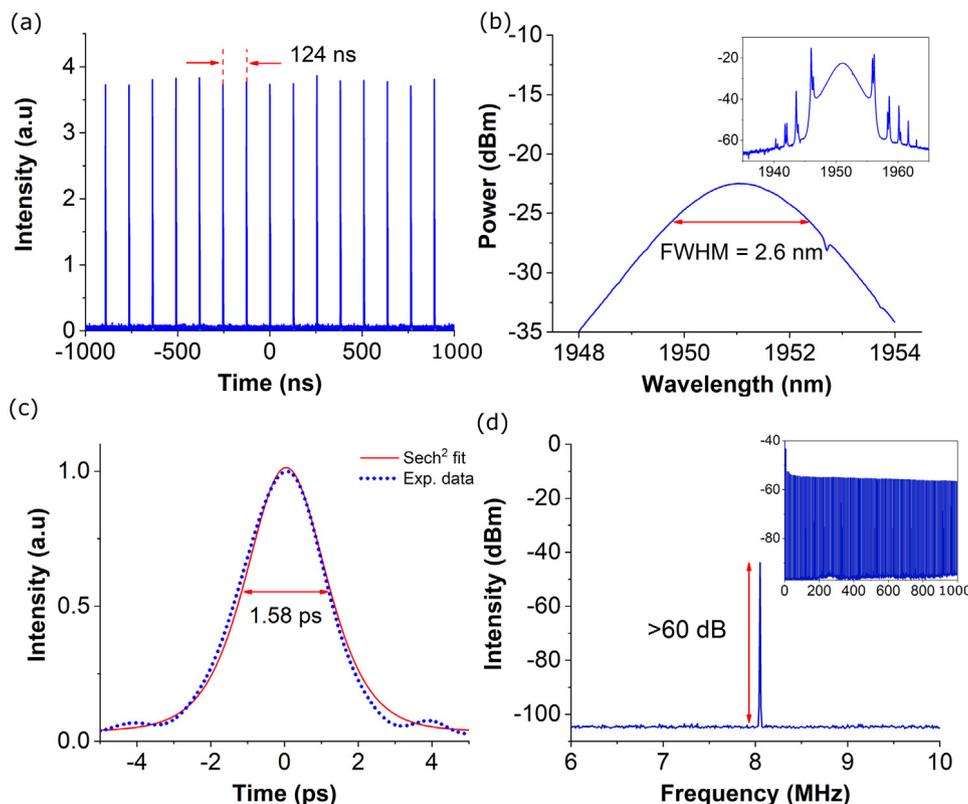


Fig. 4. (a) Mode-locked pulse train; (b) enlarged optical spectrum with FWHM of 2.6 nm. Inset shows full soliton spectrum; (c) autocorrelation trace with a Sech<sup>2</sup> fit; and (d) radio frequency spectrum, on a 10 MHz range span with a 300 Hz RBW measured around fundamental frequency of 8.1 MHz. Inset shows spectrum on wider span range of 1 GHz.

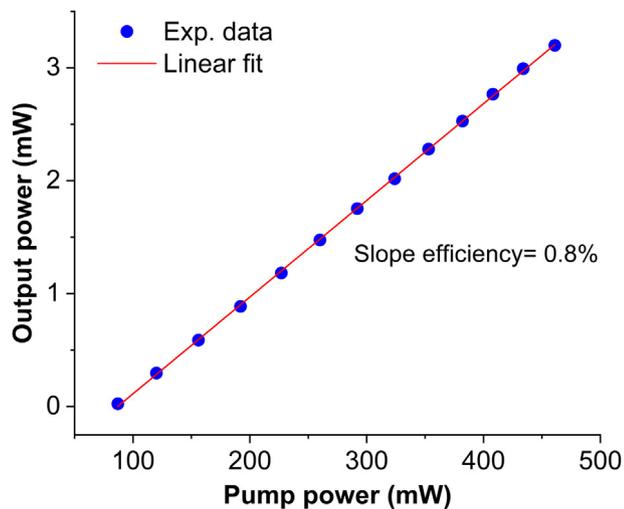


Fig. 5. Average output power against pump power.

as shown in Fig. 6(b). This verifies that the mode-locking operation possesses excellent repeatability and long-term stability. In order to ensure the mode-locking operation in the TDFL was induced by the Cu<sub>2</sub>Te-deposited SPF, a polarization dependent loss (PDL) measurement was performed. The PDL of the SA device was measured around 0.2 dB, which was low enough to exclude nonlinear polarization rotation (NPR) effect [43]. This was further confirmed by removing the SA and placing an SPF with no material coating. In this case no mode-locking was observed. The generation of stable soliton pulses using Cu<sub>2</sub>Te-deposited

SPF in 2 μm region has clearly indicated the performance of the copper chalcogenide as an effective SA.

### 5. Conclusions

A passively mode-locked TDFL based on copper chalcogenide was demonstrated for the first time. The mode-locked pulses were obtained based on the evanescent field interaction, where the Cu<sub>2</sub>Te was deposited onto a SPF. The pulsed laser had a repetition rate of 8.1 MHz and a pulse width of 1.58 ps, as well as a pulse energy of 0.39 nJ and a peak power of 250 W. The RF spectrum of the laser indicated good stability with an SNR of more than 60 dB. The laser was also stable when tested for a long-term stability, showing no sign of instability for a period of 4 h. The mode-locked TDFL operating in the 2 μm wavelength region would have useful applications in the medical field and as sensors.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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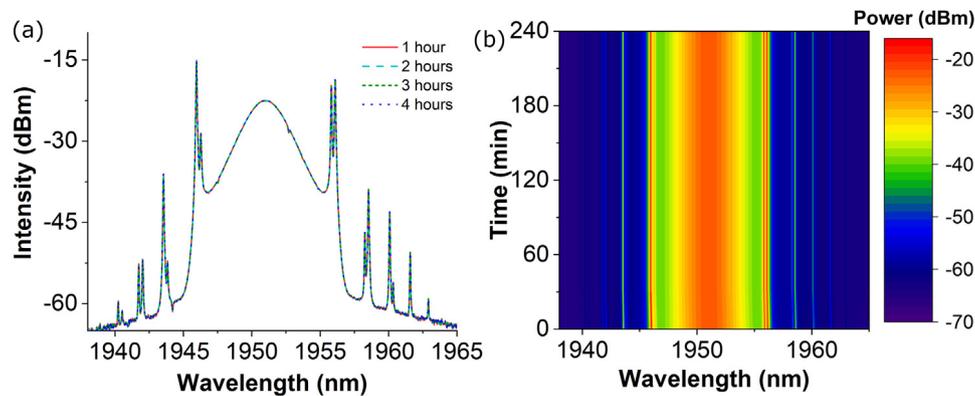


Fig. 6. (a) Spectra acquired after 1-h (red line), 2 h (cyan dash line), 3 h (green short dash line) and 4 h (blue dot line); and (b) 3D spectrum of mode-locking operation over 240 min. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

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