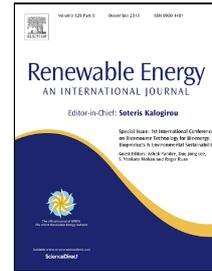


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Improving spectral modification for applications in solar cells: a review

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Abstract

The spectral mismatch between solar cells and incident radiation is a fundamental factor limiting their efficiencies. There exist materials and luminescent processes which can modify the incident sunlight's properties to better suit the cell's optimal absorption regions. This makes for an interesting area of research and promising technique for enhancing the efficiency of solar cells which is important for environmental reasons. It is intended for this review to provide the reader with historical and up-to-date developments of the application of spectral modification to solar cells and contribute to growing its impact on real-world PV devices. We concisely outline the underlying principles of three spectral modification processes: upconversion (UC), downconversion (DC) and luminescent downshifting (LDS). For each section we present up to date experimental results for applications to a range of solar PV technologies and discuss their drawbacks. With particular focus on UC, we then review how nanostructures or integrated optics might overcome these problems. Finally, we discuss practical challenges associated with advancing this approach for commercialisation and opportunities spectral modification presents; namely where future research should focus and via a cost analysis with a simple formula that can be used to determine financial viability for the deployment of this technology.

Keywords

Downconversion, Efficiency, Luminescent Downshifting, Nanostructures, Solar Cells, Upconversion.

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76 1. Introduction

77 The maximum theoretical efficiency of a 1.1 eV band gap silicon solar cell was calculated
 78 through a detailed balance model by Shockley and Queisser in 1961 to be approximately 30%
 79 [1]. However, this limit can be theoretically surpassed by the construction of solar cells in
 80 which the design and assumptions in this model need not hold true. Today there are a range
 81 of solar technologies which use new materials, going beyond traditional (or so-called 1st
 82 generation) silicon wafer based modules. These include 2nd generation thin film and 3rd
 83 generation nanostructured devices whose efficiencies and cost per watt are outlined in Table
 84 1. Despite promising progress, if the efficiencies could be further enhanced in a non-
 85 expensive manner, then the price of solar electricity could be further reduced which would
 86 ease the transition to a sustainable global economy.

87

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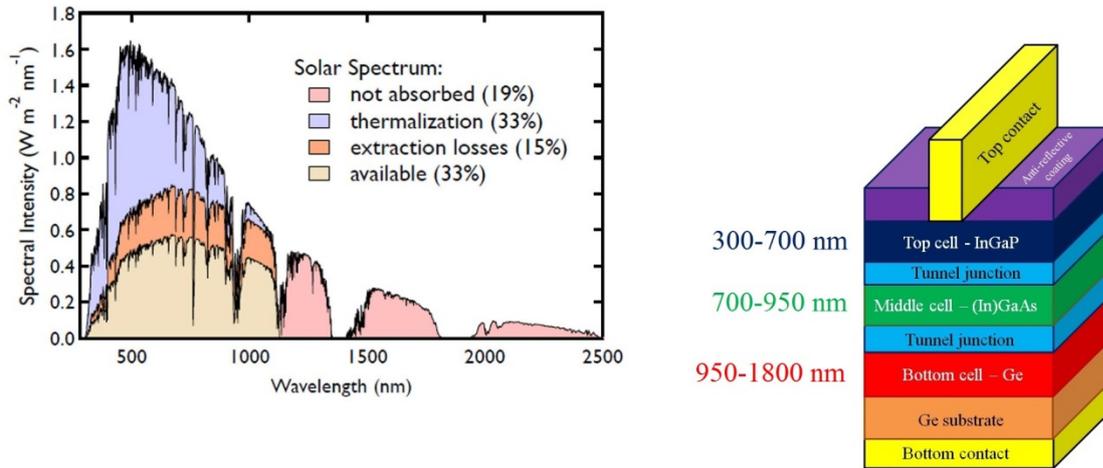
89 **Table 1 –Various solar technologies with their maximum efficiency and cost per peak watts as of**
 90 **2016. Table data from [2], [3] and [4].**

91

92 Solar cell efficiency is fundamentally hampered by the following parameters [5]:

- 93 • Non-absorption of low-energy photons
- 94 • Thermalization losses from the absorption of high-energy photons
- 95 • Extraction losses due to unavoidable charge carrier recombination

96 In practice, optical losses [6] also occur from incomplete absorption, reflection and shading,
 97 in addition to electrical losses from parasitic series and shunt resistances [7] which further
 98 reduce the system efficiency from its theoretical maximum. However, as seen in Fig. 1, the
 99 greatest efficiency losses arise from non-absorption of high wavelength light and
 100 thermalization of short wavelength light. Collectively this problem is referred to as the
 101 spectral mismatch between the incident solar spectrum and absorption profile of the cell; it
 102 stems from a defined energy called the band gap energy, E_G , which is present in all materials
 103 that can harness solar power. Photons with $E \geq E_G$ have the potential to be absorbed and
 104 generate photocurrent; whereas photons with $E < E_G$ will be transmitted through the cell
 105 without adding to the output load. Furthermore photons of $E > E_G$ contribute to
 106 thermalization losses. This is because when absorbed they possess excess energy than
 107 required to generate photocurrent, so electrons are promoted into the conduction band with
 108 additional kinetic energy which is wasted as heat. Therefore the spectral mismatch is the key
 109 detriment to a cell's efficiency. Multi-junction or tandem cells seek to rectify this problem by
 110 combining a series of cells with varying bandgaps which are connected together either
 111 mechanically or monolithically into a single structure [8]. Through this technique, it is
 112 possible to harness a wider range of the solar spectrum, attaining higher efficiencies. In 1980
 113 De Vos used a detailed balance model again to calculate a maximum efficiency of 68% was
 114 theoretically attainable for an infinite stack of varying band gap cells [9].



115

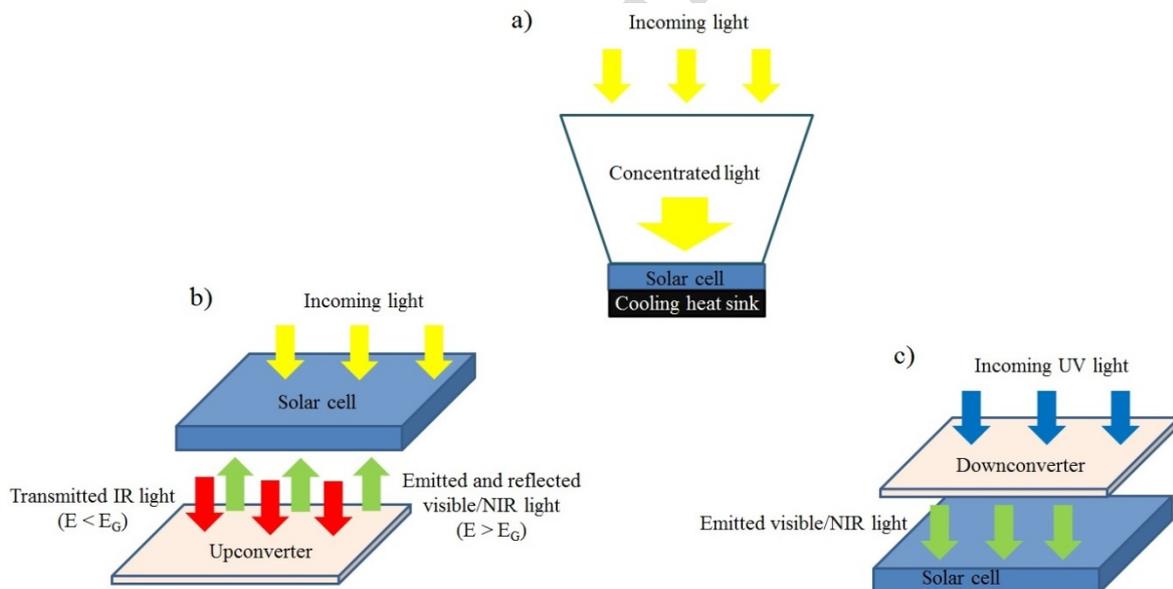
116 **Fig. 1. The percentages of the solar spectrum represented as the energy losses that occur in a**
 117 **silicon cell. The spectral mismatch is responsible for approximately 50% of these through non-**
 118 **absorption and thermalization. The structure of an epitaxial triple junction cell and regions of**
 119 **the spectrum it can absorb are shown. Figures modified from [5] and [10].**

120

121 In the Shockley-Queisser model it is assumed the cell only undergoes illumination by the
 122 ambient solar spectrum. From their equations, short-circuit current density increases linearly
 123 and open-circuit voltage logarithmically with the incident light concentration, hence the
 124 overall efficiency would rise under greater irradiance. According to De Vos' model a
 125 maximum efficiency of 86% is hypothetically achievable if using infinite junctions under
 126 concentrated sunlight, approaching the separately derived thermodynamic limit [11]. In
 127 reality, this means if lenses or parabolic mirrors [12] are utilized to achieve great light
 128 concentrations then higher efficiencies become possible. Also because optical materials are
 129 cheaper than semiconductor structures, a smaller area of photovoltaic surface is required
 130 which is advantageous for reducing production costs. As a result, there has been much
 131 research into and commercial deployment of concentrating solar PV technology in recent
 132 years [13]; a major challenge is maintaining a low cell temperature to maintain performance
 133 [14].

134 Through a combination of concentrating sunlight and multi-junction structure, remarkably
 135 efficient laboratory manufactured cells have been produced in recent years (46% in 2014
 136 [15]). Unfortunately, these ultrahigh efficiency panels are not commercially viable for global
 137 deployment and are mainly limited to regions of the world with high direct normalized
 138 irradiance (DNI) or for use in satellites where a low weight-electrical power ratio is essential
 139 [16], so it is worth exploring alternative methods of PV cell design which surpass the
 140 Shockley-Queisser limit. One such innovative technique is to challenge the assumption that
 141 only photons with energy equal or greater than the band-gap create an electron-hole pair to
 142 generate photocurrent; so rather than modifying the materials to fit the spectrum as with
 143 tandem cells, we modify the spectrum to suit the material.

144 It is estimated 20% of the total photon energy in the standardised AM1.5G solar spectrum is
 145 beneath the bandgap of a monocrystalline silicon solar cell [17]. If the energy of these NIR
 146 photons could be harnessed then the amount of sunlight being converted to electricity, and
 147 hence the efficiency of the solar cell would increase. This was proposed by Wolf as early as
 148 1960 [18] and later explored by Trupke et al [19] who calculated a maximum theoretical
 149 efficiency of 47.6% for un-concentrated sunlight and 63.2% under maximal terrestrial
 150 concentration (46,200 suns). In reality, a single photon with energy less than the band-gap
 151 *cannot* be absorbed but if the energies of two or more are combined into a single photon,
 152 through a process known as upconversion (UC), then this new quanta *can* excite electrons
 153 into the conduction band and lead to more current being produced as predicted. In practice
 154 this involves a UC layer placed beneath the solar cell to absorb transmitted radiation and re-
 155 emit it towards the cell at useful wavelengths. On the other hand, as shown by Fig. 1, an even
 156 greater portion of spectral energy is wasted via thermalization. It is also apparent solar cells
 157 give the best photovoltaic response at wavelengths closer to the band-gap of the material.
 158 Therefore, spectral modification in the opposing direction would also aid efficiency. There
 159 are two ways of producing lower energy photons from ones of higher energy;
 160 downconversion (DC) (also referred to as energy transfer by quantum cutting [20] (QC)) and
 161 luminescent downshifting (LDS). Unlike UC, it is logical to place the spectral modification
 162 material above the solar cell so a low transmittance to already optimal wavelengths is
 163 essential to result in a net gain in efficiency [21].



164

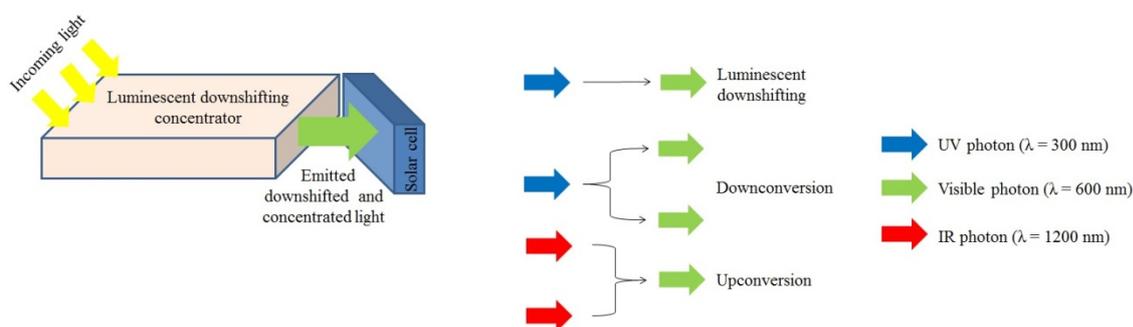
165 **Fig. 2. The principles of: a) Concentrating PV; a lens concentrates incoming light which**
 166 **improves solar cell performance and reduces material demand, while a cooling sink prevents**
 167 **the cell from overheating. b) An upconversion layer-solar cell system. A fraction of the**
 168 **transmitted sub-band gap radiation undergoes upconversion and is redirected toward the cell**
 169 **where it is utilized. c) A downconversion layer atop a solar cell which absorbs short wavelength**
 170 **radiation and emits at photon energies better suited to the cell, increasing overall efficiency.**

171

172

173 DC can be regarded as the inverse of UC. In this process a single high energy photon splits its
 174 energy between two longer wavelength photons. In 2002, Trupke et al [22] calculated a 1.1
 175 eV band-gap solar cell's maximum theoretical efficiency could increase from 30.9% to
 176 38.6%, following the addition of an optimal DC layer. Conversely, LDS is a process which
 177 Stokes shifts the energy of a short wavelength photon to a lower energy which is particularly
 178 useful for multi-crystalline silicon or CdTe cells with their poor responses to UV or blue light
 179 [23]. However, unlike UC and DC, LDS will not be able to supersede the Shockley-Queisser
 180 limit alone since no additional photons are generating electron-hole pairs [24]. However
 181 combining LDS materials with a concentrating geometry is an attractive route to improving
 182 cell efficiency as it can improve both the spectral response and the intensity of the incident
 183 light [25].

184 Many materials have been investigated for their UC, DC and LDS properties with studies
 185 predominantly focused on rare earth [26] (RE) doped lattices but also emerging alternatives
 186 such as quantum dots [27] (QDs) and organo-metallic dyes [28]. This review will present up
 187 to date investigations into spectral conversion using these materials. The limitations of these
 188 studies and possible solutions such as concentrating optics and utilising nanoscale
 189 phenomena will be discussed. Finally, we suggest where research should continue if spectral
 190 conversion is to emerge as a serious competitor to existing multijunction technology.



191

192 **Fig. 3. Design of a luminescent concentrating PV system which both concentrates light intensity**
 193 **and shifts it to favourable wavelengths whilst directing it to the cell, along with a visual**
 194 **representation of 3 typical spectral conversion processes: UC, DC and LDS. Figure modified**
 195 **from [34].**

196

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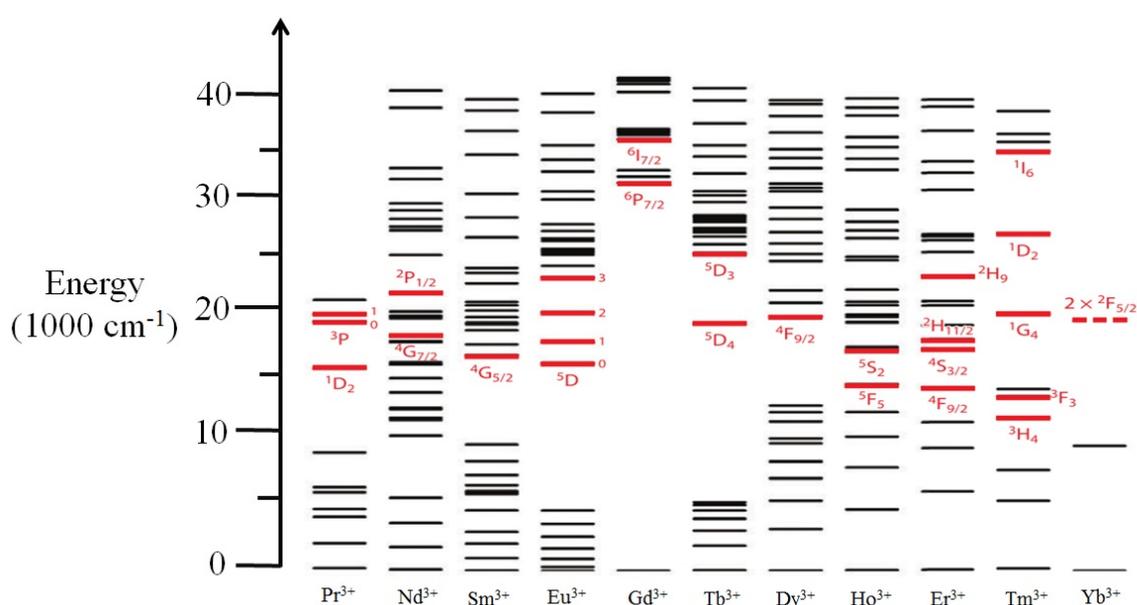
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200 2. Working principles of upconversion

201 2.1 Upconversion via rare earth ions

202 Bloembergen first suggested energy levels in transition and rare earth metal compounds
 203 could absorb IR radiation then reemit at a shorter wavelength, acting as *quantum counters*
 204 [29]. This process came to be known as *upconversion* and has been studied extensively,
 205 predominantly for applications in laser physics [30] but more recently for applications in bio
 206 imaging [31], and the focus of this review, photovoltaics. The rare earth metals in their
 207 trivalent ion state are excellent materials for UC due to their ladder-like energy structure [32]
 208 and their application to UC is the predominant scope of this review. For clarity, the term rare
 209 earth metal usually refers to elements 57-71 (the lanthanides) in addition to scandium and
 210 yttrium, but for the purposes of this paper we will use *rare earth metals* to refer exclusively
 211 to the lanthanides. Specifically, the lanthanide 3+ ions all of which have the electron
 212 configuration [Xe] 4fⁿ with n ranging from 0 to 14 [33]. It is the abundance of possible states
 213 in the 4f-orbital that allow for such a diverse range of emissions and absorptions. Although
 214 such transitions between states are parity forbidden upon doping into a host lattice,
 215 interaction with the crystal field leads to symmetry breaking and allows these transitions to
 216 take place [34]. Dieke [35] gives an excellent characterization of these states under
 217 spectroscopic notation in the so-called Dieke energy level diagrams.



218

219 **Fig. 4. A Dieke style diagram showing the energy levels of the lanthanide ions. The states**
 220 **highlighted in red tend to emit upconverted radiation when excited. Figure modified from [80].**

221

222 There are three predominant mechanisms as described by Shalav et al which allow for
 223 photon UC in a rare earth doped compound at realistic solar concentrations [36].

224

225 • Ground State Absorption/Excited State Absorption (GSA/ESA) is a single ion
226 process whereby an electron is promoted from the ground state to an excited
227 state by absorbing a photon. While in this excited state a second photon is
228 absorbed leading to an even higher energy state. When this electron relaxes
229 radiatively to the ground state it will emit a photon of energy equal to the sum
230 of the absorbed photons.

231

232 • Energy Transfer Upconversion (ETU) is a two ion process when nearby ions
233 both absorb from the ground state into the excited state. One ion (the
234 sensitizer) then transfers its energy to the other (the activator), causing the
235 sensitizer to return to the ground state and the activator to emit a higher energy
236 photon.

237

238 • Photon Avalanche (PA) is an unconventional UC process discovered in 1979
239 [37], in which energy mismatch between the incoming photon and the
240 intermediate/ultimate excitation states leads to an exponential population of
241 intermediate states via cross relaxation resonance.

242

243 Furthermore, in recent years two additional mechanics have been identified:

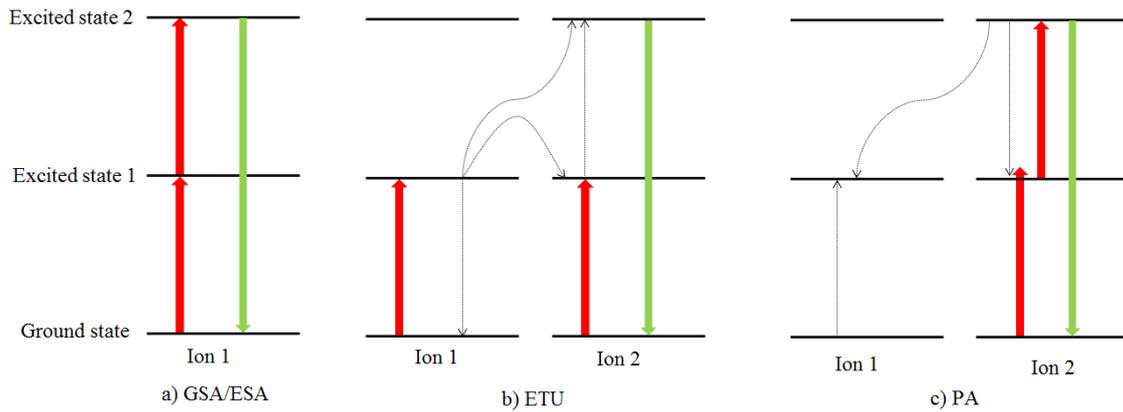
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245 • Co-operative Energy Transfer (CET) is a three step process during which two
246 sensitizer ions undergo excitation and transfer their energy to a single ground
247 state activator.

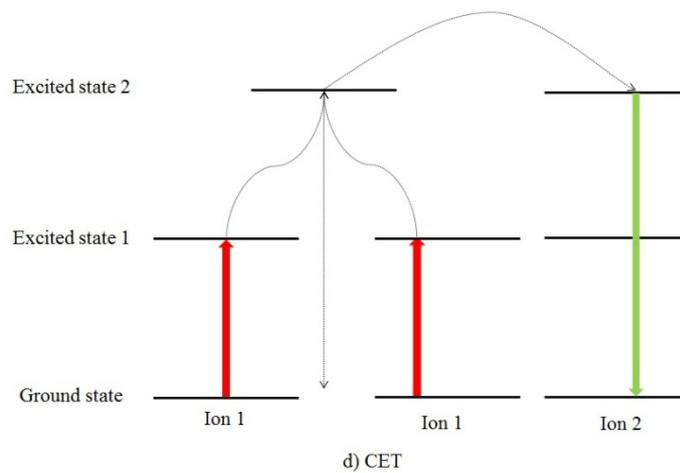
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249 • Energy Migration Upconversion (EMU) is a four step process recently
250 proposed in 2011 by Liu et al [38] which takes place in a core-shell
251 nanostructure and involves the interaction of 4 ions: a sensitizer, accumulator,
252 migrator and activator.

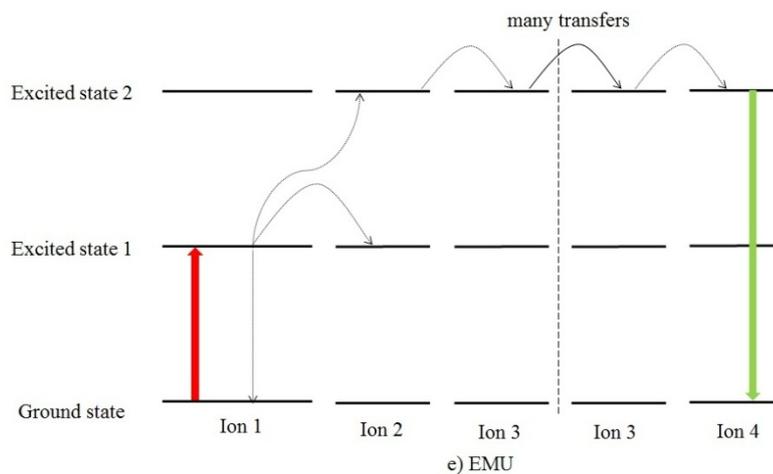
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254



255



256

257 **Fig. 5. Energy diagram representation of the upconversion mechanisms in rare earth ions: a)**
 258 **excited state absorption (ESA), b) energy transfer upconversion (ETU), c) photon avalanche, d)**
 259 **co-operative energy transfer and e) energy migration upconversion. Figure modified from [83].**

260

261 Of these ETU is the most efficient [39] so is aimed to be the dominant process in
262 experimental set-ups. The mechanism for a given UC material may be determined by
263 analysing the fluorescence time decay or the pump power-luminescent intensity relation. An
264 example of the former is demonstrated by Yang et al [40] where ESA and ETU processes
265 were differentiated by the profile of the time decay; ESA took place during the excitation
266 pulse and exhibited exponential behaviour (akin to direct excitation), while ETU had a longer
267 persistence and non-exponential characteristics relating to the lifetime of the level providing
268 energy transfer. On the other hand, Pollnau et al [41] showed UC luminescent intensity to
269 have a dependence on absorbed pump power and in certain cases ETU and ESA could be
270 distinguished by the measured slopes on a logarithmic plot; when n sequential photons are
271 absorbed in a mechanism excited by pump power P , the emission intensity will show a P^n
272 behaviour. Their theoretical results were verified with experimental investigations on four
273 rare earth ion containing compounds. Although anomalies were found in this relationship for
274 high-power pumping regimes which had a logarithmic slope gradient of 1, regardless of the
275 number of steps in the mechanism [42].

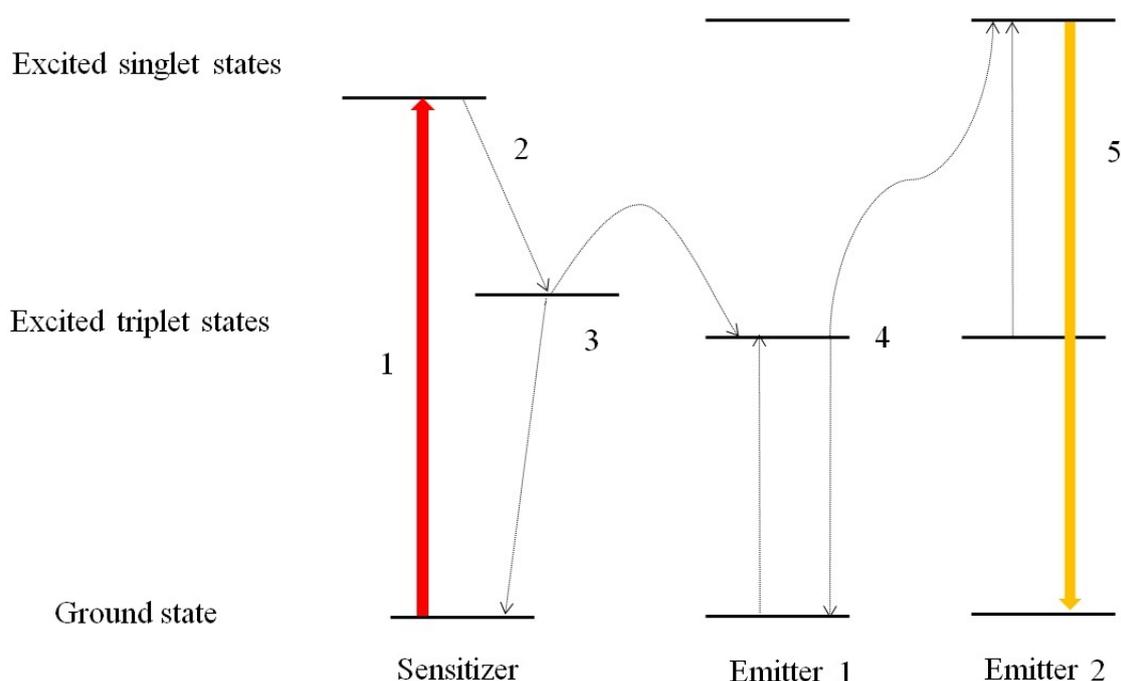
276 Detailed state by state mechanisms presented on energy level diagrams have been proposed
277 for specific ions, examples of which will be seen in section 3.3. However, since 2014 there
278 have been debates in the literature about discrepancies in the frequently cited diagrams for
279 popular UC materials and alternatives put forward. [43]. Furthermore, there have been
280 investigations into the dependence on temperature of the ratio of UC emission from different
281 energy states [44].

282

283

284 2.2 Upconversion via TTA

285 An alternative UC approach, utilizing organometallic compounds [45] has been suggested:
286 triplet-triplet annihilation (TTA). The mechanism differs from the ones discussed thus far as
287 it involves varying spin states. It could have advantages over rare earth ionic based UC due to
288 the low excitation energies required (of the order mW/cm^2 , comparable to ambient solar
289 illumination) and the readily tuneable excitation and emission wavelengths of the sensitizers
290 and acceptors [46]. This potential was identified by Balushev et al [47] in 2006 following its
291 demonstration under non-coherent solar radiation.



292

293

294 **Fig. 6. The states and steps involved in photon upconversion via TTA. The numbered processes**
 295 **are explained in the bulk text. Figure modified from [48].**

296

297 As shown in Fig. 11 this process occurs via 5 steps and is summarised by Cheng et al [48] as
 298 follows.

- 299 1. A low energy photon is absorbed by a sensitizer in the ground state.
- 300 2. Intersystem crossing (ISC) introduces the first triplet state.
- 301 3. The energy from this triplet state is transferred via a (Dexter [49]) triplet energy
 302 transfer (TET) process to a ground state emitter molecule which is excited to a triplet
 303 state.
- 304 4. TTA then occurs when two emitters in the excited triplet state undergo a *collisional*
 305 *complex*, yielding one in a (higher) excited singlet state and one in the ground state.
- 306 5. The excited singlet state decays, emitting a higher energy photon.

307

308 Similarly to UC via rare earth ions, this process can be evidenced by the fluorescent
 309 intensity's quadratic dependence on incident light power since TTA involves an interaction
 310 between two sensitized triplet acceptor molecules [50]. Kinetic studies by Cheng et al [51]
 311 disproved earlier assumptions in efficiency limits from spin-statistics and found the process

312 efficiency could reach 40%. Since the scope of this review tends towards rare earth ions, the
313 reader is referred to [52] for a recent and in-depth report on the developments in this field.

314 2.3 Upconversion via quantum dots

315 A QD is a semiconductor nanocrystal with its excitons confined in three spatial dimensions.
316 A fascinating attribute of QDs that make them ideal for applications in solar PV is the fact
317 their absorption properties depend on size, so it is possible by nanoengineering to fine tune a
318 QD to absorb radiation at specified wavelengths [53]. With these abilities, it is no surprise
319 they have come under investigation as spectral modification materials [54-5]. Shang et al [56]
320 in their 2015 article reviewed how QDs could be utilized for UC. Two QDs of different band-
321 gaps are connected by a semiconductor rod. Photons are absorbed in the lower energy QD,
322 creating an electron-hole pair. This is followed by a second photon absorption which allows
323 the hole to tunnel to the other QD where it recombines to emit a higher energy photon. There
324 are two ways of this occurring as shown in Fig. 12, they described two mechanisms: a) direct
325 intraband hole absorption and b) an Auger [57] mediated process. However, unlike TTA-UC
326 these QDs would require higher solar concentrations (by a factor of 10^5 compared to the
327 ambient spectrum) in order to activate this kind of UC.

328

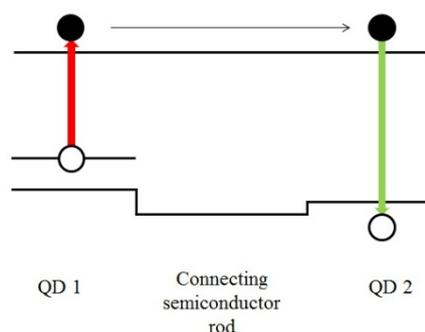
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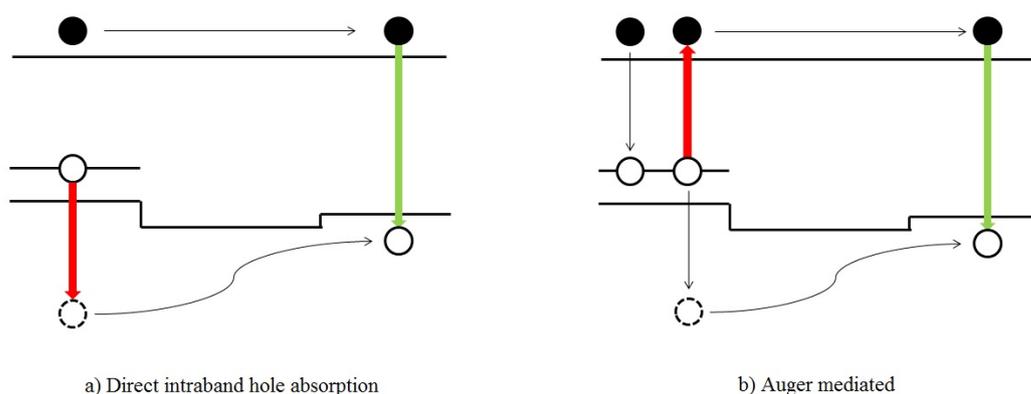
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338 **Fig. 7. Schematic of the possible upconversion processes in a quantum dot nanostructure. The**
 339 **overall process for QDs 1 and 2 connected by a semiconductor rod is shown foremost and can**
 340 **occur via two mechanisms: a) direct intraband hole absorption and b) an Auger mediated**
 341 **process. Figure modified from [54].**

342

343

344 2.4 Upconversion via thermal radiation

345 A radically different approach to UC has been demonstrated in recent years. The underlying
 346 principle of UC via thermal radiation is to use light to raise a substance's temperature,
 347 thereby shifting its black body emission spectrum to shorter wavelengths [58] as
 348 demonstrated in Fig. 13. These emitted photons could then be utilized by a solar cell to
 349 increase efficiency. The mechanism still involves a rare-earth ion but instead of its states
 350 decaying and emitting photons, the energy is absorbed into the lattice producing heat.

351

352 3. Upconversion experiments

353 3.1 Early experiments

354 Initially, Impurity Photovoltaic (IPV) devices were envisaged that would dope metal ions
 355 directly into the cell substrate and through their energy levels allow IR light to undergo UC
 356 and generate additional electron-hole pairs [59]. Yet in 1996 Gibart et al [60] found the
 357 efficiency of a GaAs cell to be 2.5% under 891 nm illumination (less photon energy than the
 358 cell GaAs bandgap), when a 100 μm thick vitroc ceramic containing Er^{3+} and Yb^{3+} ions was
 359 applied to the rear of the cell. This first ever proof-of principle experiment showed UC could
 360 utilise sub band gap photons in a solar cell. It has since been accepted separate up converting

361 layers have more potential over IPV as they can be materially optimised independently from
 362 the cell [61] and in early experiments on silicon, the additional layer outperformed IPV cells
 363 [62].

364 Rare earth metals were identified as the leading up converting materials in this field and
 365 Auzel gives an excellent review of the mechanisms responsible [63] in more detail than
 366 outlined in the previous section. In the years following, there have been many more proof-of
 367 principle experiments demonstrating UC to be beneficial to the efficiency of many types of
 368 solar cell (see Appendix 1). However, despite the growing interest and performances in UC
 369 devices, the area faces many challenges. Even the best results have only led to modest
 370 increases in efficiency. This is primarily due to the low quantum yield of the UC process and
 371 narrow absorption bands of the rare earth metals [64]. Thus, we present in the later sections
 372 of this review paper, innovative techniques, experimental data and where to go next if we are
 373 to make the process more efficient, which is necessary if it is to become a viable technology
 374 for commercial solar cells.

375 3.2 Upconversion efficiency and its limitations

376 When quantifying the efficiency of UC processes, several values are quoted. Frequently
 377 presented are the internal or external photoluminescent quantum yields (*iPLQYs/ePLQYs*),
 378 which are sometimes referred to as UCQY. These are defined as the fraction of absorbed or
 379 incident photons emitted at the UC wavelength, so that:

$$380 \quad iPLQY = \frac{\text{number of photons emitted}}{\text{number of photons absorbed}} \quad (\text{Equation 1})$$

$$381 \quad ePLQY = \frac{\text{number of photons emitted}}{\text{number of incident photons}} \quad (\text{Equation 2})$$

382 Therefore, an *ePLQY* of 50% would imply a perfect absorption and two low-energy photons
 383 to one high energy photon process. These values can be measured experimentally by using an
 384 integrating sphere [65]. If an upconverter is combined with a solar cell the overall external
 385 quantum efficiency (*EQE*) of the system is given by:

$$386 \quad EQE = \frac{\text{number of electron hole pairs generated}}{\text{number of incident photons}} \quad (\text{Equation 3})$$

387 Studies had until recently predominantly focused on monochromatic illumination of UC
 388 materials with speculation for PV applications. However, as the field has developed, there
 389 have been more reported physical PV-UC systems having their power conversion efficiency
 390 (PCE) and short-circuit current density (J_{SC}) measured under a solar simulator. Furthermore,
 391 in phosphor luminescence studies it is common to quote the normalized EQE (*NEQE*)
 392 measured in cm^2/W and defined as:

$$393 \quad NEQE = \frac{EQE}{\text{incident light power density}} \quad (\text{Equation 4})$$

394 Although one should be careful when contrasting these values, since UC is a non-linear
 395 process [66] so quantum efficiency is not independent from incident light concentration. The
 396 iPLQY follows a power law:

$$397 \quad iPLQY \propto P_{in}^n \quad (\text{Equation 5})$$

398 P_{in} is the incident power and n the number of steps in the UC mechanism. On the other hand,
 399 the EQE of single junction solar cells has a linear relationship with increasing intensity.
 400 Combining the two processes for a UC-PV system we see the overall EQE in response to sub-
 401 band gap illumination scales:

$$402 \quad EQE \propto \frac{P_{in}^n}{P_{in}} \propto P_{in}^{(n-1)} \quad (\text{Equation 6})$$

403 So for a predominantly ETU process which has two steps, the EQE is expected to scale
 404 linearly with increased radiation. This law holds at low irradiance but at higher irradiance
 405 deviates from theoretical value due to competing processes [67]. Therefore, if a higher
 406 concentration of light on the upconverter can be achieved the PV efficiency rise will be
 407 greater.

408 Shalav et al identified three dominant losses that depopulate the excited states [68].
 409 Collectively these processes are known as quenching. Firstly, multi-phonon transitions are a
 410 type of non-radiative relaxation to lower energy states; if more than six phonons are required
 411 then radiative relaxation will dominate [69]. Hence it is essential to use low phonon energy
 412 hosts such as halides [70] when designing compounds for UC experiments. Secondly,
 413 radiative relaxation can occur, resulting in the emission of a photon of lower energy than
 414 desired. And finally, energy transfer called cross-relaxation can take place between ions,
 415 leading to unwanted low energy photon emission. This is particularly prevalent for high rare-
 416 earth doping concentrations [71], so the ratios should be carefully selected.

417 Since the UC process is the transfer of energy between states in a material, it can be modelled
 418 by simulations. The fundamental dynamics of a luminescent body are underpinned by the
 419 following rate equation [72]:

$$420 \quad \frac{dN}{dt} = - (A_r + A_{nr})N \quad (\text{Equation 7})$$

421 Where N is the population of excited states and A_r and A_{nr} are the radiative and non-radiative
 422 decay rates. From this equation one can see in the absence of illumination the excited
 423 population decays exponentially with time constant [73]:

$$424 \quad \tau = \frac{1}{A_r + A_{nr}} \quad (\text{Equation 8})$$

425

426 The quantum efficiency of the photoluminescence is simply the ratio of radiative decay to
 427 total decay [74]:

428

429

$$\eta = \frac{A_r}{A_r + A_{nr}} = \frac{\tau_r}{\tau} \quad (\text{Equation 9})$$

431

432 So the quantum efficiency can be calculated by comparing observed lifetimes to τ_r which can
 433 be estimated from Judd-Ofelt theory [75]. Furthermore, empirical observations into non-
 434 radiative decay rates have shown the so called “energy gap law” [76]:

$$A_{nr} = A_{nr}(0)e^{-\alpha(\Delta E)} \quad (\text{Equation 10})$$

436 In this law, $A_{nr}(0)$ and α are constants that depend on the host lattice and ΔE the gap in
 437 energy between the states. This re-iterates the need for low phonon-energy host lattices since
 438 a large energy gap will require more phonons to bridge, hence the exponentially decaying
 439 decay rate. With these equations (and more complex quantum mechanics if required) it is
 440 possible to model up conversion processes in rare earth doped lattices with a range of energy
 441 levels. These simulations can then be compared to experimental data. In 2012 Fischer et al
 442 [77] demonstrated this by modelling the transition probabilities in 20% Er^{3+} doped NaYF_4
 443 and consistency with observed data meant using rate equations to simulate behaviour was
 444 physically reasonable. Shyichuk et al [78] extensively analysed ETU dynamics in
 445 $\text{YVO}_4:\text{Yb}^{3+}, \text{Er}^{3+}$. A set of rate equations was compiled and solved numerically using the
 446 Runge-Kutta method. Agreement with experiment showed they were successfully able to
 447 predict the macroscopic properties and dynamics of the system from microscopic
 448 considerations of the Er^{3+} ions’ behaviour. Their new methodology is advantageous to prior
 449 theoretical models as it can be generalised to any crystal matrix with known structure and
 450 various experimental illumination set ups. These studies increase our understanding of the
 451 fundamental behaviour in a UC material, so its development is a necessary task if we are to
 452 produce a highly efficient upconverter.

453

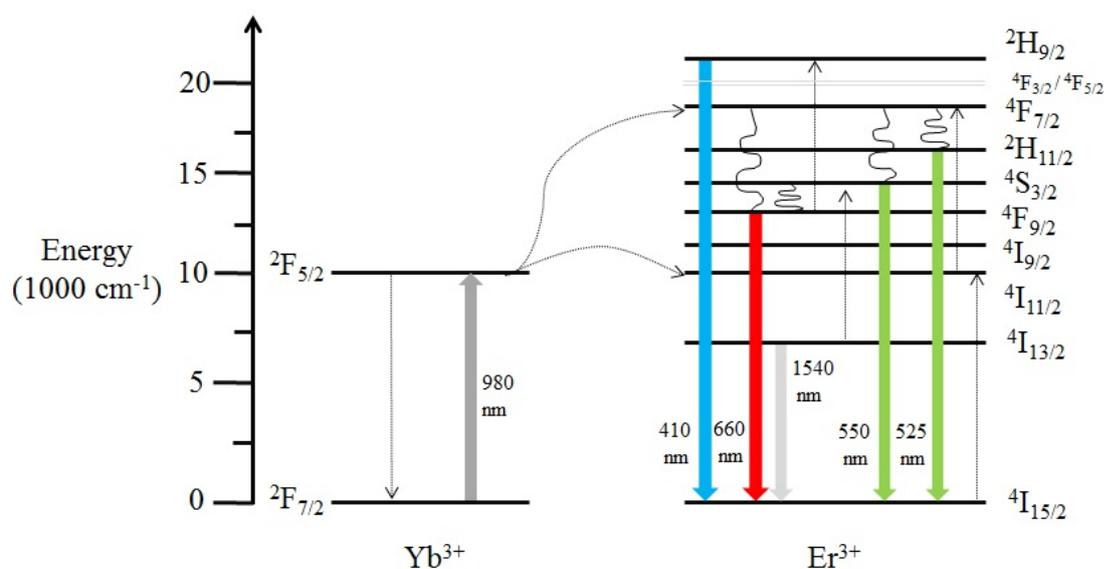
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456 3.3 Rare earth ions used for upconversion

457 Of the rare earth ions, Er^{3+} has stood out as an excellent candidate and has been widely
 458 tested for UC. This is because it has a suitable absorption spectrum centred at 1523 nm, with
 459 emission lines predominantly at 980nm, resulting from the $I_{11/2}$ to $I_{15/2}$ transition [79]. These
 460 values are remarkable for PV applications since 1523nm is an intense region of IR radiation
 461 in the sea level solar spectrum ($25\text{W}/\text{m}^2$) and a photon of wavelength 980nm has energy close
 462 to the band gap of a crystalline silicon solar cell (c-Si). $\text{Yb}^{3+} - \text{Er}^{3+}$ co-doped lattices have too
 463 shown excellent potential. Yb^{3+} has a larger cross-section than most ions ($9.11 \times 10^{-21} \text{cm}^2$ at

464 976 nm for the ${}^2F_{7/2}$ to ${}^2F_{5/2}$ transition) [80] and can act as an excellent sensitizer to another
 465 rare earth ion. These UC materials have higher UCQYs with approximately 50% of absorbed
 466 NIR photons emitted in the visible range [81]. The absorption and emission spectra of Yb^{3+} -
 467 Er^{3+} co-doped compounds are better suited for applications in wide band-gap PV, such as
 468 amorphous silicon or perovskite cells. The absorption peak for Yb^{3+} - Er^{3+} is located at 980
 469 nm, with emission across the 400-700nm range [82]. Additional rare earth metals have
 470 potential across the required wavelengths. In their 2013 review Huang et al [83] classify all
 471 possible transitions with relevance for spectral conversion in solar PV. However, Er^{3+} and
 472 Yb^{3+} - Er^{3+} aside, their frequency in the literature of being applied to solar cells is sparse.



473

474

475 **Fig. 8. Energy level diagram showing the possible emissions from the Yb^{3+} - Er^{3+} ion pair. The**
 476 **dashed lines represent energy transfer, the solid lines radiative decay and the curly lines non-**
 477 **radiative decay. Figure modified from [81].**

478

479

480 Dy^{3+} has a similar energy level structure to Er^{3+} , so has theoretical potential for UC in solar
 481 cells. Er^{3+} - Dy^{3+} co-doped lattices can absorb at 1300nm and emit at visible wavelengths, but
 482 investigations showed improved UC efficiency was needed to match the emission intensity as
 483 compared to 1.5 μm excitation [84]. Likewise, Tm^{3+} ions absorb at 1220nm and emit at
 484 808nm. Rodríguez-Rodríguez et al [85] analysed UC in Tm^{3+} doped glasses for applications
 485 in PV yet they estimated only an additional EQE of 0.1% was attainable at high solar
 486 concentrations. This negligible improvement meant the authors concluded its use unlikely to
 487 be justified in solar cells. However, the Ho^{3+} ion does appear more promising. It has a wide
 488 absorption in the 1150-1225nm range [86]; this part of the ambient solar spectrum is
 489 approximately twice as intense as at the 1530nm centred Er^{3+} absorption. Ho^{3+} doped
 490 fluorindate glasses have been successfully tested for efficient UC and excitingly they are

491 transparent at 1540 nm. This is significant because it means there is potential for a multi-
492 wavelength UC device containing a Ho^{3+} layer atop an Er^{3+} layer, which allows for photon
493 UC occurring simultaneously in the two biggest regions of the sub band-gap spectrum. This
494 design was conceptualised by Lahoz et al [87] in 2011 but to this author's knowledge, no
495 such device has been assembled and tested on a solar cell as of December 2016.

496 **3.4 Synthesis of upconversion materials**

497 When choosing UC materials, a host lattice with a low phonon energy and high chemical
498 stability is of most importance, with Sun et al [88] identifying ideal compounds as REF_3 ,
499 REOF or MREF_n ($M = \text{Li, Na, K, Ba}$; $n = 4, 5$). Thus unsurprisingly one of the most popular
500 host lattices for doping Er^{3+} has been NaYF_4 with widely reproduced syntheses originating
501 from Kramer's group in 2004 [89]. There are two occurring phase structures of NaYF_4 ; cubic
502 α - NaYF_4 and hexagonal β - NaYF_4 . Of these the hexagonal β -phase has a much greater
503 potential for UC devices as its emission has been shown to be comparably greater than the α -
504 phase [90]. Hence it is crucial in the synthesis to ensure cubic impurities are kept to a
505 minimum. Also it is important to carefully select the doping concentration of the rare earth
506 ion. Too much causes cross-relaxation processes to dominate, leading to low UC efficiency;
507 too little and large distances between the rare earth ions results in a low energy transfer rate
508 and similarly poor yield [91]. It's therefore desirable to strike a balance between these
509 extremes. A further practical consideration for applications to PV cells is that the rare earth
510 doped compounds exist in powder form, so a chemical etchant is required to act as a host
511 matrix and allow for construction of a real-world system. Polyfluorocyclobutane (PFCB) has
512 been a popular choice for its transparency to NIR radiation and similar refractive index to
513 NaYF_4 in order to reduce scattering [92]. These properties are vital to not incur further optical
514 losses. Thus, when preparing UC materials, it is important to use the ideal doping
515 concentrations and ratio of UC compound to host matrix. Ivaturi et al [93] explored
516 thoroughly the effects of varying Er^{3+} doping concentrations in a NaYF_4 lattice and
517 phosphor-PFCB matrix mass ratio. They found the optimum Er^{3+} doping level to be 25% and
518 phosphor concentration in the matrix to be 84.9 w/w %. Despite these fundamental losses,
519 there are other sources of inefficiency that arise when building a practical UC device.
520 Boccolini et al [94] have argued that losses occur due to re-absorption in a UC layer, so
521 optimizing the thickness is essential for an efficient process. Using simulations, they were
522 able to show the thicknesses that yielded the greater intensities of up converted light.

523 Alternative materials seek to improve the efficiency increase by obtaining a higher up
524 conversion quantum yield. Fluoride glasses such as ZBLAN (which contains fluorides of Zr,
525 Ba, Ln Al and Na) have been doped with Er^{3+} , reaching an impressive 12.7% absolute up
526 conversion efficiency [95]. Other emerging host lattices for this Er^{3+} transition, have included
527 30% doped BaY_2F_8 and 10% doped $\text{Gd}_2\text{O}_2\text{S}$, which reached eUCQYs of 9.5% and 12%
528 respectively [96-97]. The use broadband as opposed to monochromatic illumination also
529 provides higher quantum yields and is more like the ambient solar spectrum which is relevant
530 for real world systems. Goldschmidt et al speculated that the higher yields arise from a
531 broader incident spectrum containing more photons at resonant wavelengths for the
532 intermediate transitions in a UC process [98]. MacDougall et al [99] first investigated the

533 effect of broadband illumination on the UCQY of 10% Er³⁺ doped NaYF₄. They found the
534 iUCQY to increase from 8.7% to 16.2% upon widening the bandwidth from 12 to 61nm.
535 Then Arnaoutakis et al [100] prepared a UC structure to the specifications recommended by
536 Ivaturi et al [101]. Using a quartz tungsten halogen lamp and relevant filters, they illuminated
537 the sample under various broadband excitations. Again it was demonstrated that the wider
538 incoming spectrum (1400-1600nm) produced the highest iUCQY of 4.27%, compared to
539 1.61% at 1500-1550nm.

540 **3.5 Upconversion applications for silicon PV**

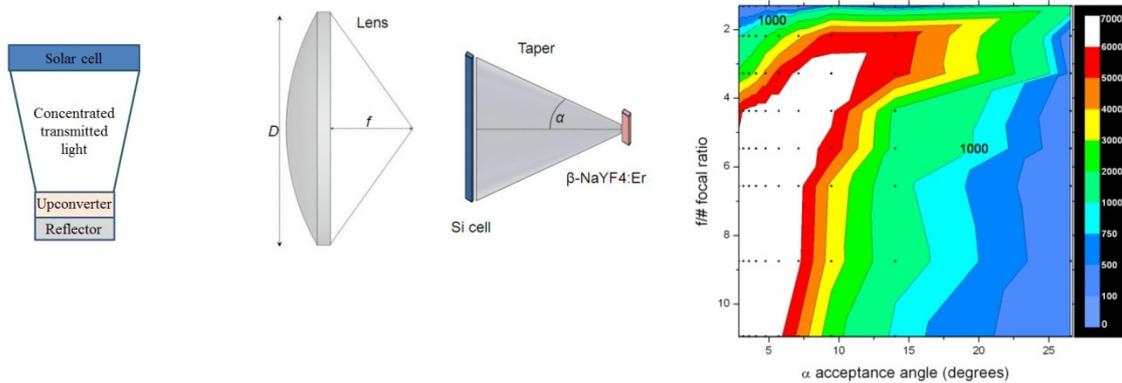
541 In an early significant study, Shalav-et al [102] placed a planar bi-facial c-Si cell above a
542 20% Er³⁺ doped NaYF₄ UC layer. This set up was illuminated by a 6mW 1523 nm laser and a
543 maximum EQE of 3.4% was measured, however the cell performed poorly under lower
544 intensity illumination (only a few times better than an IPV Er³⁺ doped cell). The group had
545 also prepared Tm³⁺ and Ho³⁺ doped NaYF₄ phosphors but unfortunately were unable to
546 characterise them due to the intense light required. The authors calculated an overall
547 maximum theoretical efficiency increase of 2.4% absolute and recommended a much broader
548 absorption spectrum was needed if this approach was to have an impact on commercial solar
549 cell production. Following this group's work, more UC layers have been applied to solar cells
550 and had their efficiencies measured, including under broadband excitation. In 2011
551 Goldschmidt et al [103] applied 20% Er³⁺ doped NaYF₄ to a silicon solar cell and under
552 1460-1600 nm radiation, the UC efficiency was 1.03% which under ambient solar incidence
553 would lead to a relative efficiency increase of 0.014% (the cell's regular efficiency was
554 16.7%). Again it was concluded this was too low to have an immediate significant impact.
555 However, by 2013, a potential relative efficiency increase of 0.16% was found via short
556 circuit current density measurements; 4.03 mA/cm² on a silicon cell that had been applied
557 with 25% Er³⁺ doped NaYF₄ under 1450-1600nm illumination [104]. The group had also
558 found a novel way to calculate relative solar concentrations for broadband spectral
559 mismatches at given wavelengths, useful for UC applied to concentrating PV Results from
560 2015 show further progress; 30% Er³⁺ doped BaY₂F₈ achieved an extra short-circuit current
561 density of (17.2 ± 3.0) mA/cm². This represents a relative enhancement of the short circuit
562 current of (0.55 ± 0.14) %, 35-times higher than values published in 2011 [105].

563 Other types of silicon cells have too been tested for UC. A significant increase in efficiency
564 of a multi-crystalline silicon solar cell (mc-Si) was reported in 2011 by Chen et al [106].
565 La₂Mo₂O₉ was co-doped with Yb³⁺ - Er³⁺ or Yb³⁺ - Ho³⁺. Against precedent, they placed the
566 UC layers on top of the cell, but this was able to reduce reflectance losses and in fact
567 increased the cell efficiency by 1.5-2.7% relatively despite a low UC process performance. A
568 widely studied class of thin film solar cells are those based on amorphous silicon (often
569 abbreviated to a-Si or a-Si:H if undergone hydrogenation). These cells have the advantage
570 that they require far less material than traditional wafer-based crystalline cells which leads to
571 a lower energy payback time [107]. Also UC devices were predicted by Atre and Dionne to
572 be of more use to thin film technologies due to their wider band gap (meaning a greater
573 portion of the solar spectrum's photons lie below the band gap) [108]. In 2010 de Wild et al
574 [109] reported a 0.01 mA/cm² short-circuit current density on an amorphous silicon solar cell

575 under 980nm illumination that had had β -NaYF₄ doped with 18% Yb³⁺ and 2% Er³⁺ applied
576 to the rear. This represented only a slight increase from the background short-circuit current
577 of 5 μ A/cm². However, sources of loss were discussed and efficiency improvement
578 suggestions made. A major drawback was identified as the poor response of the cell to rear-
579 side illumination following UC. In a recent study Qu et al [110] prepared rare earth co-doped
580 β -NaYF₄ phosphors with different combinations: 20% Yb³⁺ - 1 % Er³⁺, 25% Yb³⁺ - 1 % Ho³⁺
581 and 60% Yb³⁺ - 0.5 % Tm³⁺. The Er³⁺ doped performed best increasing overall efficiency by
582 7.3% relative whilst also providing an improved scattering function. Furthermore in 2015,
583 Chen et al [111] observed a remarkable improvement over NaYF₄ by studying 28% Er³⁺
584 doped BaCl₂. Unlike previous studies these ions were excited at two distinct wavelengths
585 (980 and 808 nm). They found upon broadband excitation >800nm its emission at 410-667
586 nm was substantially more intense with great potential for enhancing the efficiency of a-Si: H
587 solar cells. In a later publication [112] the same group stated under non-coherent NIR
588 simulated sunlight the photo response of an a-Si:H cell greatly enhanced to 0.6 mA/cm².
589 Moreover tri-wavelength simultaneous excitation at 1530, 980 and 808 nm has been observed
590 through an Er³⁺ doped LaF₃ embedded oxyfluoride glass ceramic applied to an a-Si:H cell
591 [113].

592 **3.5.1 Enhanced upconversion via concentrating PV**

593 Recalling the non-linear nature of the UC process, a conceptual design for a spectral and
594 geometric concentrator up converting system was proposed by Goldschmidt et al in 2008.
595 They predicted their specifications could raise the conversion efficiency of a bifacial solar
596 cell from 23% to 25% [114]. Fischer et al used a Fresnel lens under a solar simulator to
597 concentrate light onto an Er³⁺ doped NaYF₄ layer beneath a bifacial solar cell. At a level of
598 210 suns the short-circuit current increased by 13.1 mA/cm², corresponding to a relative
599 efficiency increase of 0.19% [115]. A further complication arises because the optimal solar
600 concentration for the cell will differ to the ideal conditions for the spectral conversion layer.
601 Some groups have investigated this imbalance and constructed integrated optics-solar PV
602 systems to compensate for the difference in concentration levels. Optical fibres combined
603 with dielectric tapers [116] have been investigated as components that can increase light
604 concentration to 2000 suns at specific wavelengths for PV applications. This integrated optics
605 approach has been tested to look for an increase in photo response in a UC-PV device under
606 sub-band gap illumination [117] and results showed the concentration on the UC layer can be
607 independently optimized regardless of the intensity at the cell's surface (an important step
608 forward for UC's potential in CPV). Arnaoutakis et al [118] demonstrated this by designing
609 such a device which placed a compound parabolic concentrator between a bifacial silicon
610 solar cell and up conversion layer as shown in Fig. 15. The group observed an increase in the
611 cell's EQE of 1.80% with the integrated system as opposed to 1.33% with the non-integrated
612 up conversion layer. Fischer et al [119] have investigated the relation between excitation
613 power density and Er³⁺ doping level in the successful lattices NaYF₄ and Gd₂O₂S. Using
614 simple models, they determined the optimum doping level decreases as the illumination
615 increases. This is something for groups to bear in mind when preparing UC materials for use
616 in set-ups of varying concentration.



617

618 **Fig. 9. A hypothetical cross-section of a UC-PV system with an array of compound parabolic**
 619 **concentrators between the layers, along with the schematic of a system modelled by changing**
 620 **parameters α and focal ratio (f/D). High concentrations $>1000X$ may be achieved. Figure**
 621 **modified from [120] and taken from [79].**

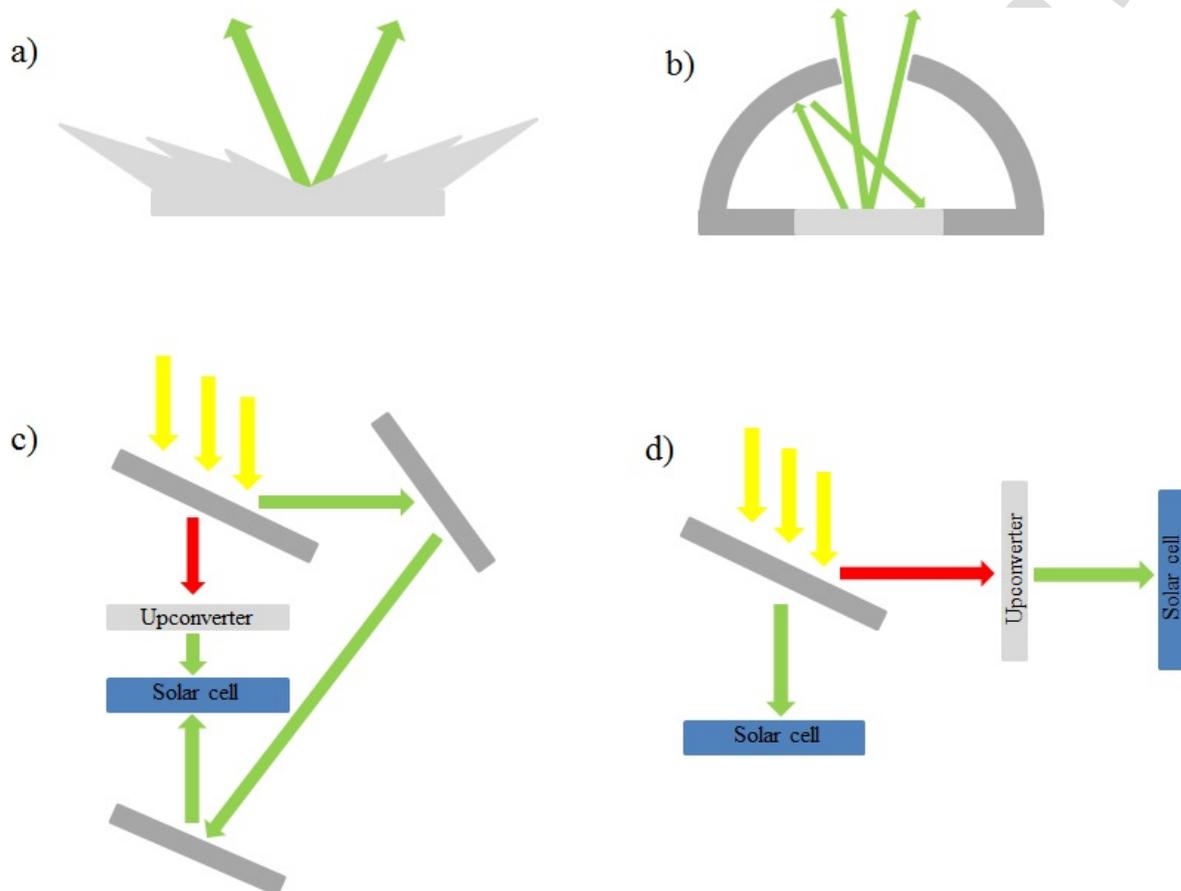
622 Geometric optical modelling has also been applied to concentrating UC systems to optimise
 623 set up for maximum concentration of light. Ray-tracing optics and Monte Carlo [120]
 624 simulations often form the bases of these investigations. Arnaoutakis et al [121] modelled the
 625 optical system shown in Fig. 17 to determine the geometric concentration of light at the UC
 626 sample. It was found to reach 6000X for acceptance angles below 10° and $f/\#$ greater than 3.
 627 If such an intense beam could be configured it would result in a much higher quantum yield
 628 of the UC layer.

629

630 3.5.2 Enhanced upconversion via thermal radiation

631 In 2014 Wang et al [122] showed UC with an efficiency of 16% could be achieved by
 632 thermal radiation. Yb^{3+} doped ZrO_2 was the material of choice for its high melting point, high
 633 thermal conductivity and strong infrared absorption. In juxtaposition to conventional UC host
 634 lattices, multi-phonon relaxation becomes a *desirable* process since incoming radiation needs
 635 to be dissipated as heat to the crystal. The oxide powder was heated to 2850K under a 976
 636 nm, 570 W/cm^2 laser and Wien's displacement law and conduction loss considerations used
 637 to give the efficiency. Further investigations used concentrated white-light to test if this
 638 glowing could be harnessed to improve solar cell efficiency. It was also compared to the
 639 widely used Er^{3+} doped NaYF_4 and found to perform worse at low excitation powers but
 640 better at high concentrations. Practical considerations when designing such a system must be
 641 to thermally isolate the solar cell from the radiation source, as from solar cell and
 642 semiconductor theory allowing the temperature to increase would have a detriment on its
 643 performance. Geometrical arrangements are briefly discussed whereby it is assumed a lens
 644 beneath a bifacial solar cell would concentrate transmitted photons, focus them onto a
 645 thermal UC material then redirect its radiation back towards the cell. Boriskina and Chen

646 [123] further developed a conceptual thermal UC-PV platform. Angular selectivity could be
 647 realised by nanomaterials or a reflective cavity with small aperture. It was envisaged to
 648 achieve a 73% maximum up conversion efficiency and a limiting silicon PV cell efficiency of
 649 45%, surpassing the Shockley-Queisser limit. In 2015 Chen et al [124] tested thermal UC in a
 650 CeVO_4 powder over the range 808-1064 nm in a system with an amorphous silicon solar cell.
 651 Upon sub band-gap illumination photocurrents of order 1 mA/cm^2 were observed,
 652 demonstrating “colossal” potential for wide band-gap solar cells.



653

654 **Fig. 10. Schematics of possible physical UC via thermal radiation systems: a) angular selectivity**
 655 **by surface nano-patterning, b) enclosing the upconverter into a reflective cavity with small**
 656 **aperture, c) geometrical and spectral splitting of sunlight onto a cell d) same concept as c) but**
 657 **with two cells. Figure modified from [124].**

658 3.6 Upconversion applications for emerging PV

659 3.6.1 Dye Sensitized Solar Cells

660 Dye Sensitized Solar Cells (DSSCs) were proposed in the 1980s and first realised by Grätzel
 661 in 1991 [125]. These cells also have the advantage of being lightweight and cheap to produce
 662 so even if efficiencies are relatively low, they may economically compete with well-
 663 established technologies if improvements in stability can be implemented [126]. The most
 664 commonly used dyes are ruthenium-based containing carboxyl groups (such as N-719 [127])
 665 that absorb poorly beyond 700nm, so a UC enhancement would allow more of the spectrum
 666 to be utilized by the dye, generating more current and higher cell efficiency.

667 Shan and Demopoulos [128] first attempted to increase DSSC efficiency this way adding
668 Yb^{3+} - Er^{3+} co-doped LaF_3 into the TiO_2 . Unfortunately, although the NIR response increased,
669 the overall efficiency of the cell decreased due to higher charge recombination at the
670 interfaces. In a more successful study in 2010, Xie et al [129] constructed DSSCs with Yb^{3+} -
671 Er^{3+} co-doped into the TiO_2 thin film layer. This converted 980nm unabsorbed light into
672 visible wavelengths 510-700nm which could be used by the cell. Under a solar simulator the
673 cell achieved an efficiency of 7.28% compared to 6.41% for the un-doped DSSC,
674 representing a substantial relative efficiency increase of 13.6%. In synthesis optimization
675 studies it was found the DSSC with a TiO_2 : RE ratio of 1:3 had the greatest efficiency. A
676 beneficial secondary effect of directly inserting rare earths into the oxide is that it acts a p-
677 type dopant, raising the Fermi level of the electrons in the oxide and the redox potential in the
678 electrolyte, therefore increasing the photo voltage of the DSSC. In 2013 Ramasamy and Kim
679 [130] took a different approach in applying UC to DSSCs. They produced a novel rear
680 structure with silver reflectors and Fe^{3+} doped $\beta\text{-NaGdF}_4$: Yb^{3+} - Er^{3+} UC nanoparticles
681 (UCNPs). Iron had been found to increase the UC emission by a factor of 30 and the silver
682 particles further enhanced the process by 2 possible effects – increased scattering or surface-
683 coupled plasmon emission. The latter is an area of great interest to improving UC in solar
684 cells since overlap between a plasmonic resonance frequency and UC emission spectra can
685 lead to faster radiative decay rates in the UCNPs, increasing the intensity of the
686 luminescence. A 2017 study by Han et al [131] used hollow Er^{3+} - Yb^{3+} CeO_2 spheres as
687 UCNPs and reported a 27% relative increase in PCE for a DSSC.

688 3.6.2 Organic Solar Cells

689 Organic Solar Cells (OSCs) are polymer based photocurrent generating structures and like
690 DSSCs they are lightweight and cheap to produce. Common compounds used in OSC
691 fabrication include conjugated polymers such as polyphenylenevinylene (PPV) derivatives or
692 poly [3-hexylthiophene] (P3HT) with a fullerene derivative (PCBM) [132]. As with all solar
693 cells, OSCs suffer IR losses from a limited absorption range which stops at 800nm. Therefore
694 UC has potential to improve their efficiency.

695 In 2012 studies by Wang et al [133] added a composite layer of Yb^{3+} - Er^{3+} co-doped MoO_3
696 beneath the photoactive P3HT/PCBM region of the OSC. A modest increase in short-circuit
697 current of < 1% was observed under 1 sun. If greater improvements were to be made there
698 would need to be developments in UCQY, employment of light trapping/management and
699 thicker up converting films. Over the same period, Damitha Adikaari et al [134] instead used
700 Ho^{3+} - Yb^{3+} co-doped Y_2BaZnO_5 on organic PV devices to obtain an EQE of 0.0052% under
701 illumination by a 986 nm laser. Despite this small value they showed a wider excitation
702 spectrum and additional modifications such as transparent electrodes could lead to values
703 more like those achieved in amorphous silicon cells and an overall conversion efficiency rise
704 of 0.45%. In 2014 Chen et al [135] employed Yb^{3+} - Er^{3+} doped $\beta\text{-NaYF}_4$ UC particles in a
705 polymethylmethacrylate (PMMA) matrix and a silver reflector to an OSC. They observed a
706 22.2% enhancement in the short-circuit current density; however, a large portion of this was
707 due to visible light being reflected. Despite demonstrating technical feasibility, they attribute
708 only 1.6% of the photocurrent enhancement to the UC process and poor performance under

709 ambient solar conditions mean it is unlikely to be practically applied unless further
710 improvements are made. Zhao, Ji and Guo reviewed a range of organo-transition metal
711 complexes that could be used as sensitizing and accepting molecules in such a system [136].
712 Schulze et al [137] tested this through theoretical models and experimentally on three solar
713 cells: two heterojunction OSCs (P3HT: ICBA and PCDTBT: PC71BM) and one a-Si. They
714 showed proof-of-principle but again the measured short-circuit current enhancements were
715 negligible, only 0.2%, 2-3 orders of magnitude below a device relevant value. However,
716 TTA still may be promising over conventional rare earth ETU because of higher quantum
717 efficiency. Furthermore the models showed sensitizer concentration was the key parameter to
718 be optimized but was difficult to achieve the desired levels in the organic solvent. They
719 suggested incorporating the molecules into solid state compounds such as nano-particles or
720 polymer films as a way to get round this.

721 3.6.3 Quantum Dot Solar Cells

722 QD solar cells (QDSCs) are an emerging technology that have attracted considerable interest
723 as yet another low cost alternative to traditional silicon cells [138]. Despite tuneable
724 absorption spectra, there will always be wavelengths that a manufactured QD cannot absorb
725 beyond, so one would expect there is again potential for their use in conjunction with UC
726 devices. Until recently there were few examples of UC being applied to QDSCs. However, in
727 a recent study by Wang et al [139] in 2014, a Tm^{3+} - Yb^{3+} - Er^{3+} tri-doped $NaYF_4$ phosphor
728 was synthesised and annealed, then introduced to the photoanodes of a CdS/CdSe QDSC. A
729 maximum PCE of 4.37% was achieved, marking a 20% increase and 17.6% enhancement of
730 the photocurrent. The annealing pre-treatment is a contributing factor to this excellent result
731 as it eliminates defects which reduce charge recombination in the cell. Furthermore in 2017,
732 Ramachari et al [140] have explored the effects of incorporating Yb^{3+} - Er^{3+} co-doped ZrO_2
733 into a QDSC, leading to a 55% relative increase in the PCE.

734 3.6.4 Perovskite Solar Cells

735 An emerging PV technology known as perovskite solar cells (PSCs) are hybrid organic-
736 inorganic devices based on the crystal structure called perovskite [141]. Their low cost and
737 easy fabrication makes them a promising candidate in PV systems, and they demonstrate
738 reduced charge carrier losses and significantly improved stability as compared to DSSCs
739 [142]. Recent developments in their design has led to a rapid growth in power conversion
740 efficiency; 3.8% in 2009 to 6.5% in 2011 with a 20.1% system being verified by National
741 Renewable Energy Laboratory (NREL) by the end of 2015 [143]. These cells absorb between
742 300 and 800 nm so like other nanostructured and thin film technologies could benefit from a
743 UC (and/or DC) layer in order to use a greater portion of the solar spectrum. This was
744 envisaged by Wang [144] and in recent studies UC crystals or nanoparticles have been
745 successfully integrated with PSCs to give modest efficiency increases. There have been
746 different approaches in choosing where to put the UC materials. For example Chen et al [145]
747 applied a visibly transparent $LiYF_4: Yb^{3+}, Er^{3+}$ single crystal on top of a PSC to increase the
748 PCE by 7.9% under 7-8 sun illumination, whereas He et al [146] and Roh et al [147] have
749 added the UCNPs in situ by combining them with the mesoporous electrode, the latter

750 attaining a 13.7% PCE enhancement and latest results by Wang et al have shown a 20.8%
 751 improvement [148]. As with DSSCs there may be extra benefits to an internal addition of
 752 nanoparticles such as the aiding of quality perovskite crystal growth during the device
 753 fabrication.

754

755 **Table 2 – Summary of UC investigations carried out on various types of solar cell.**

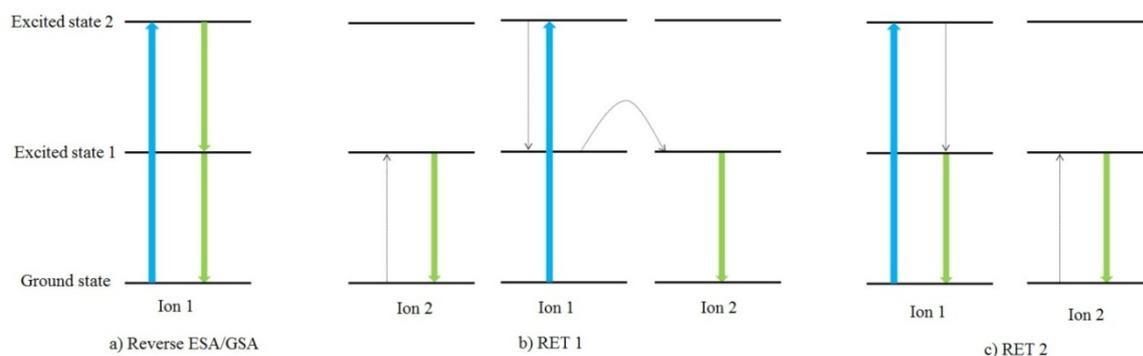
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757 4. Downconversion experiments

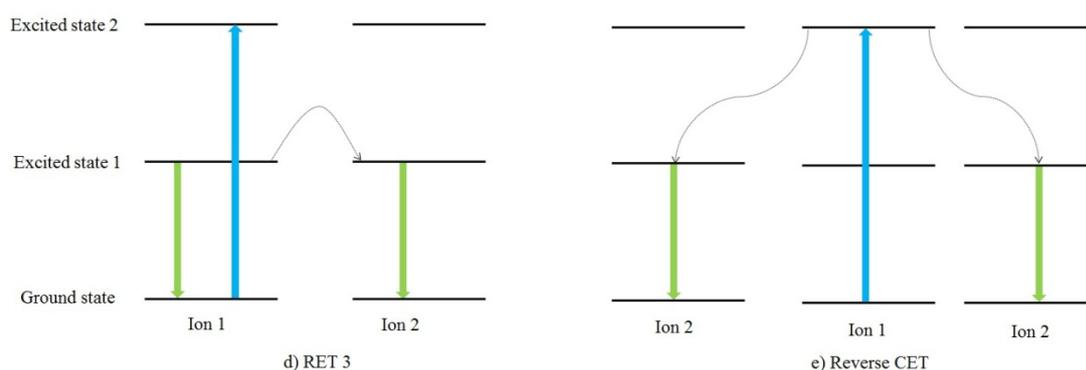
758 In addition to losses from sub-band gap photons it is estimated an even greater proportion of
 759 the solar spectral energy (recalling Fig. 1) lost from photons possessing too much energy. In
 760 silicon solar cells, a photon with energy greater than the band-gap will create an electron-hole
 761 pair but the once in the conduction band the charge carrier has excess energy which will
 762 mostly be lost as dissipated heat. Cells also have an optimal EQE that varies with wavelength
 763 (cf. Fig.6) so in theory getting as much of the spectrum as possible at that energy would be an
 764 excellent way to improve efficiency. Although contrary to UC, lower band gap cells or
 765 systems with a poor UV response stand to gain the most from an efficient DC layer.

766 4.1 Working principles of downconversion

767 Owing again to their diverse energy structure, rare earth materials have considerable potential
 768 for DC devices. Many of the responsible mechanisms for DC appear analogous to UC but in
 769 reverse. In their 2013 article Huang et al [149] describe five processes as typical. Firstly, a
 770 mechanism equivalent to GSA/ESA in reverse, a UV photon is absorbed in a single ion,
 771 promoting an electron from the ground state to a secondary excited state before decaying
 772 sequentially, emitting two NIR photons. There also exist three possible routes for DC via
 773 resonant energy transfer between two distinct ions as shown in the figure. Finally, a reverse
 774 co-operative energy transfer takes place when an ion absorbs a photon and transfers the
 775 energy to two additional ions promoting them to an excited state. These states then
 776 simultaneously decay emitting two low energy photons. Despite these similarities it has been
 777 shown unlike UC, DC tends to be a linear process independent of light intensity [150].



778



779
780 **Fig. 11. Diagram of the main mechanisms for downconversion: a) ESA in reverse b) c) and d)**
781 **resonant energy transfer between two ions and e) CET. Figure modified from [149].**

782

783 4.2 Downconversion materials

784 The DC process has been observed in several lanthanide compounds and from observing the
785 energy transfer efficiency in RE ions, theoretically attainable e-PLQYs greater than 150%
786 [151] (in cases approaching 200% [152]) have been consistently reached (cf. Equation 1, the
787 maximum possible quantum yield for a DC process will be 200%). Direct measurements of
788 e-PLQY are rarer in the literature but a value of over 100% has been reported [153]. Despite
789 this and Trupke's earlier promising analysis, in 2007 Strümpel et al [154] concluded DC to
790 have little practical promise in terrestrial PV due to the high excitation energies required
791 falling within UV light which is mostly blocked by the atmosphere. A further problem arises
792 from having to place the DC material on top of the solar cell which leads to parasitic
793 absorption and scattering losses, meaning an e-PLQY of 115% would be required just to
794 break even [155]. Nevertheless, research in this field continued with Huang et al cataloguing
795 many studies from their 2013 review of both single ion and paired ion lattices [156]. The
796 addition of other species can sensitize the DC process by increasing the absorption cross
797 section. These have often taken the form of divalent rare earth ions such as Eu^{2+} and Yb^{2+} or
798 transition metal ions such as Mn^{2+} . Tai et al [157] demonstrated Eu^{2+} acted as an efficient
799 energy sensitizer for Yb^{3+} in a SrAl_2O_4 lattice, reaching a quantum efficiency of 147.36%
800 under illumination by a 450 W xenon lamp. The broadband 250-450 nm light was shifted to
801 980 nm where it could be effectively utilised by a silicon solar cell. Most studies have
802 implied use for c-Si but from Fig. 6 it's clear other technologies could benefit too, such as Ge
803 cells with their very low band gap of 0.67eV (1850 nm). As with UC, it is important to
804 carefully select the dopant concentration and layer thickness to minimise self-absorption
805 losses. Boccolini et al have presented an optical model to determine the optimal thickness of
806 a DC layer which was validated by experimental analysis of a Ce^{3+} - Yb^{3+} co-doped borate
807 glass [158]. Li et al [159] have also demonstrated broadband DC in Yb^{3+} doped Na_2YMg_2
808 $(\text{VO}_4)_3$ which under UV illumination 240-400nm showed intense NIR emission at 974nm,
809 ideal for silicon cells. By analysing decay curves a mechanism based on combined energy

810 transfer was proposed. For a recent review focussing exclusively on the available DC
811 materials and their application to PV the reader is referred to [160].

812

813 **4.3 Downconversion applications for silicon PV**

814 Until recently there had been few physical DC-PV systems fabricated and their efficiency
815 enhancements recorded. However there have been an increasing number of studies since
816 2015. That year, Dumont et al [161] prepared Tb^{3+} - Yb^{3+} co-doped Silicon-Nitride matrix
817 DC layers via magnetron co-sputtering. By using novel optimization techniques and material
818 characterization, its yield was improved from prior studies and was deemed to be suitable for
819 application to a silicon solar cell, due to efficient of emission of 980 nm light following 325
820 nm illumination. It also had anti-reflective coating properties, a further reason for this being
821 an excellent candidate for being placed atop an actual silicon cell. Gonzales-Perez et al [162]
822 have further considered these design properties with their Eu^{3+} complex embedded in PMMA
823 which was applied to a reference silicon cell by spin coating or tape casting. The EQE of the
824 cell increased by 5.8% at 280 nm illumination due to a strong UV excitation band from 250-
825 360 nm with emission at 665 nm. The film also showed excellent thermal stability meaning it
826 could easily be applied in the module assembly of silicon cells and be used in CPV.
827 Additionally, Florêncio et al [163] reported a 7% enhancement in the efficiency of a
828 commercial silicon cell when Tb^{3+} - Yb^{3+} co-doped tellurite glass was placed on top as
829 compared to un-doped glass.

830

831 **4.4 Downconversion applications for emerging PV**

832

833 DC particles have also been applied to emerging PV cells. In a novel investigation Yao et al
834 [164] enhanced the conversion efficiency in a DSSC to 4.8% (an improvement of 245% from
835 pure TiO_2) using Eu^{3+} - Dy^{3+} co-doped ZnO applied to the TiO_2 photoanodes alongside
836 graphene loading to reduce recombination and interfacial resistance. Similarly in 2016, Hou's
837 group have reported a Eu^{3+} doped $ZnGa_2O_4$ nanophosphor in the porous TiO_2 of a PSC [165].
838 This led to an enhanced short-circuit photocurrent by a maximum of 4.08 mA/cm² and
839 relative power conversion efficiency increase of 34.4% due to improved light harvesting.

840

841

842 **Table 3 – Summary of DC investigations carried out on various types of solar cell.**

843

844

845

846 **5. Luminescent downshifting experiments**

847

848 **5.1 Luminescent downshifting materials**

849 Rare earth doped phosphors [166], organic dyes [167] and QDs [168] have all been
850 suggested to be utilised for achieving LDS structures for PV cells. The responsible
851 mechanisms will vary by material and do not appear to be as widely discussed in the
852 literature as per the rare earth ion energy level diagrams described for UC/DC.

853 However, Huang et al [169] outline five key characteristics for LDS materials:

- 854 1. Broadband absorption, particularly in regions of poor spectral response for the
855 relevant solar cell
- 856 2. High absorption coefficient and quantum efficiency
- 857 3. High transmittance and narrowband emission in region of high device response
- 858 4. Large Stokes shift to minimise self-absorption losses
- 859 5. Long term stability

860 Rothemund et al presented a simple, analytical optical model to calculate the potential EQE
861 gains for various solar cell technologies following the addition of an LDS layer. This could
862 be useful in future studies for evaluating experimental data, optimizing LDS layers and
863 screening potential LDS materials [170]. In another investigation Lipovšek et al [171] used 3-
864 D ray tracing and Mie scattering theory, to produce an experimentally verified model which
865 predicted OSCs could gain 6% short-circuit current from an LDS layer. The model accounted
866 for several parameters such as layer thickness, phosphor size distribution and volume
867 concentration. Furthermore, Lesyuk et al [172] have explored cadmium-free QDs as a non-
868 toxic alternative to the often used CdSe/ZnS QD.

869 **5.2 Luminescent downshifting applications for silicon PV**

870 Hung and Chen [173] prepared a Eu^{3+} doped gadolinium oxysulfide layer for a mc-Si solar
871 cell. Under ambient solar illumination conditions an enhancement in J_{SC} of 6.47 mA/cm^2 and
872 power conversion efficiency of 2.67% was observed. This material was simple and low cost
873 to produce and also acted as an anti-reflective coating, reducing losses from Fresnel
874 reflection. Similarly, Ho et al [174] used two Eu^{2+} doped phosphor species mixed with SiO_2
875 and spin coated the solution on a single silicon solar cell. The cell's J_{SC} increased 19.85% and
876 relative PCE by 15.97% due to a combination of broadband (512-610 nm) emission and
877 forward scattering. Tm^{3+} doped fluoride glasses have also been investigated for their LDS
878 properties and combined with a c-Si cell. Maalej et al [175] reported a modest 1.4% relative
879 PCE enhancement at a 1% Tm^{3+} doping concentration when compared to un-doped glass.
880 They concluded a greater absorption and collection of converted photons was needed to scale
881 up these glasses for commercial application.

882

883 **5.3 Luminescent downshifting applications for thin film PV**

884 2nd generation PV technologies such as CdTe or CIGS cells would stand to benefit a lot from
885 LDS due to their very poor short wavelength response (cf. Fig. 6). These are thin film
886 semiconductors intended to attain lower pay back times and less environmental impact than
887 1st generation c-Si modules [176]. A typical cell consists of a layer of CdTe or CIGS in
888 contact with CdS to form a p-n junction [177]. Uekert et al [178] recently prepared
889 nanostructured organosilicon luminophores embedded in ethyl vinyl acetate (EVA) or
890 polyvinyl butyral (PVB) to act as an LDS layer for a CIGS cell. The device showed a
891 remarkable improvement in EQE in the UV region from 1% to 55% at 360nm and a relative
892 PCE increase of up to 4.3%. Alternatively, Ross et al [179] applied organic dye-doped EVA
893 and fluorinated ethylene propylene (FEP) cover sheets to two CdTe modules which improved
894 the short-circuit current by 9.7% and 5.3% respectively. Meanwhile Hodgson et al [180]
895 have explored varying concentrations of semiconductor QDs encapsulated in a PMMA film
896 over a thin-film CdS/CdTe cell. They found an overall efficiency gain of 1.7% for a QD
897 concentration 48 $\mu\text{g}/\text{mm}^3$. In most recent studies by Song et al [181] Mn²⁺-Eu³⁺ co-doped
898 phosphate glasses were applied to CdS/CdTe cells for a relative increase in power conversion
899 efficiency of 7.14%. It was also noted how the thickness of the LDS layer is a key parameter
900 in determining overall efficiency gains; a thicker layer leads to more increased photocurrent
901 at short wavelengths but will also lead to more absorption of the desired wavelengths and
902 hence losses. Concentrating LDS has also been investigated as material that simultaneously
903 shifts wavelength and increases intensity would be most beneficial atop a solar cell. In one
904 such study Parel et al [182] tested different LDS structures on a CdTe/CdS solar cell. The
905 measured current output increased 10% by the addition of a concentrating LDS, compared to
906 5% without the light concentration. Furthermore, the addition of other species into the layer
907 improved the energy transfer and spectral absorption, resulting in an ultimate 20%
908 enhancement in the short-circuit current.

909

910 **Table 4 – Summary of LDS investigations carried out on various types of solar cell.**

911

912

913 **6. Enhancing spectral modification via nanostructures**

914 As advances in nanotechnology have progressed, it has been possible to engineer
915 nanostructures to the benefit of spectral conversion (particularly the UC process). One of the
916 key drawbacks of increasing solar cell efficiency via UC is the rare earth doped material's
917 limited absorption spectrum. When we compare their optimum absorption wavelengths to the
918 AM1.5G spectrum (cf. Fig 19), we see there are large parts of the sub-band gap region that do
919 not undergo UC; the 1050-1350 nm and 1580-1800 nm range is not utilised by either Yb³⁺-

920 Er³⁺ or Er³⁺ ions where they are attempting to improve the efficiency wide band gap or c-Si
921 PV respectively.

922

923 **6.1 Quantum dots**

924 QDs have been explored as a potential way to overcome this problem and enhance emission
925 from rare earth ion based UC layers [183]. In 2010 Pan et al [184] combined an Yb³⁺-Er³⁺
926 doped oxide UC layer with PbS QDs. The addition of the QDs effectively widened the
927 absorption spectrum by absorbing photons that would otherwise be transmitted and re-
928 emitting at wavelengths that could be utilized by the UC ions. When applied to a bifacial
929 silicon solar cell, a 60% increase in photocurrent was observed under IR illumination for the
930 PbS-UC layer than the UC layer alone.

931

932 **6.2 Photonic crystals**

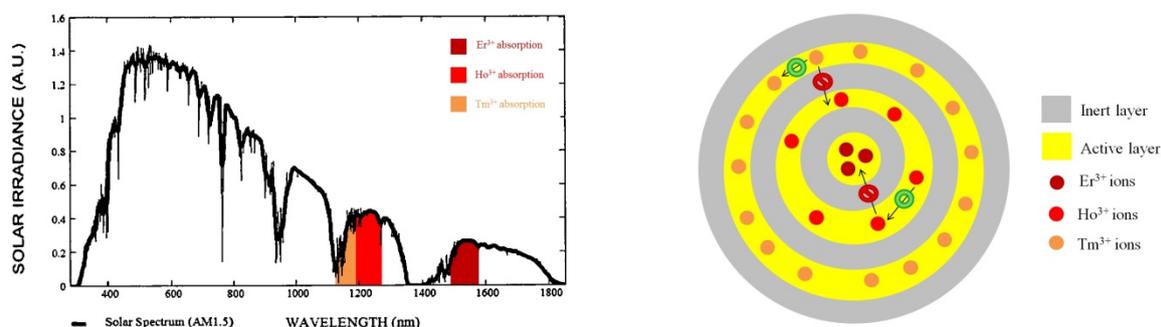
933 Further still improvements can be made by making use of another nanostructure, photonic
934 crystals. Photonic crystals are optical nanostructures that influence a photon's behaviour in a
935 solid [185]. Therefore they can be very useful for UC-PV, where direction and wavelength of
936 light undergoing interaction is critical for the process' efficiency. In 2015 Marques-Hueso et
937 al [186] coupled PbSe QDs with InP graphite lattice photonic crystals for a remarkable
938 improvement in UC properties. The photonic crystal acted as a wavelength modifier by
939 further shifting the QD emissions (reducing re-absorption) and concentrating the spectrum in
940 the region where the rare earth absorption is strongest. This led to a 158% increase in the
941 amount of QD emission matching the upconverter absorption. Furthermore, the vertical
942 emission grew by a factor of 7.8 which would be beneficial for a spectral concentrator located
943 beneath a bifacial cell as it significantly reduces scattering, meaning more photons reach the
944 cell.

945

946 **6.3 Core-shell nanostructures**

947 The small physical size of the UC ions or nanoparticles is also problematic for the process'
948 efficiency [187]. This results in a small cross section leads to low quantum yields even at
949 absorbing wavelengths. One possible route is the addition of antenna ligands to the UC ions,
950 in order to increase the absorption cross section. In one such study Zou et al [188] observed a
951 3300 enhancement in UC luminescence by using a cyanine dye (IR-806). Further
952 developments in the use of dye ligands to enhance UC are discussed in a recent letter by
953 Huang [189]. The creation of UCNPs in a core-shell design has also been proposed to
954 overcome this [190]. These UCNPs would typically consist of a core of rare earth ions
955 surrounded by a shell of phosphor and insulator (for example SiO₂) to achieve electrical
956 isolation of the up converting metal and reduce recombination. Furthermore, it is possible to
957 construct UCNPs with a multi-shell structure with each layer containing distinct RE ions,

958 allowing for greater broadband excitation. Investigations by Shao et al [191] have
 959 successfully shown these designs can absorb a wider part of the spectrum, 270 nm in their
 960 case. The use of inert layers helped the process by ensuring each ion acted in isolation and to
 961 minimise cross relaxation losses.



962

963 **Fig. 12. 2-D cross-section schematic of core-shell nanostructure with distinct ions in each**
 964 **region. The ions absorb at distinct wavelengths and are separated by an inert layer**
 965 **prevents the different elements from interacting with one another. The expanded absorption**
 966 **range from 1120-1260 nm is shown on the AM 1.5 solar spectrum. Figure modified from [191].**

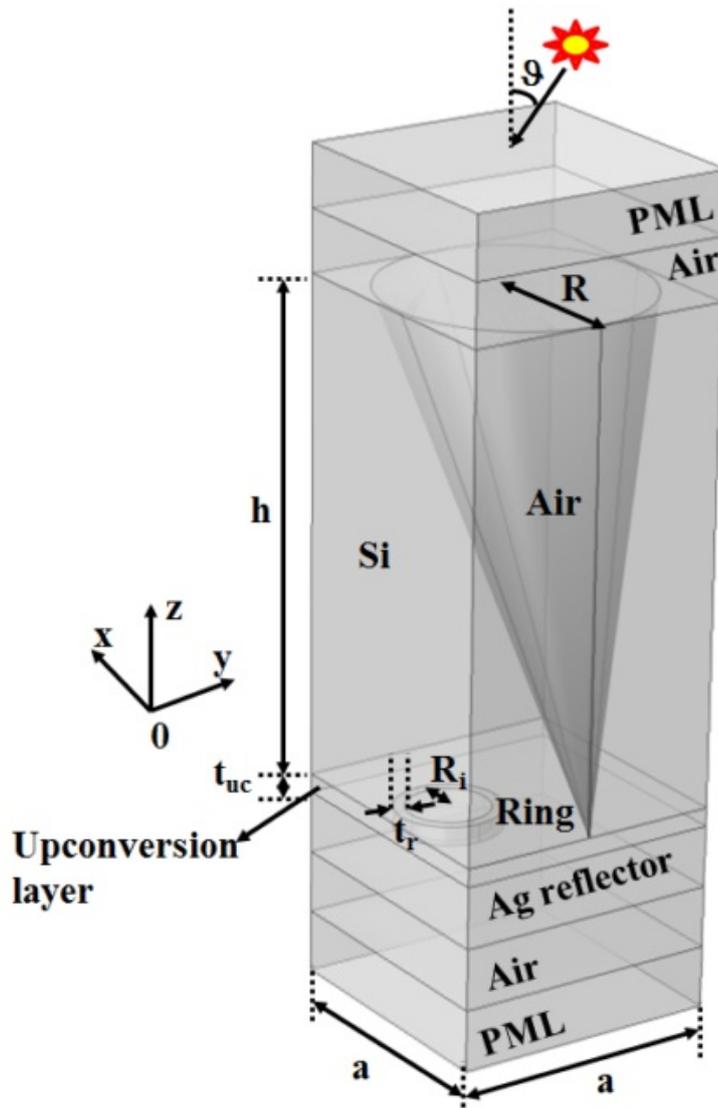
967

968

969 6.4 Plasmonic nanostructures

970 An additional nanoscale phenomenon that can benefit UC is plasmonic resonance [192]
 971 Plasmons are quasiparticles that are the quanta of vibrations in the electron density in metals.
 972 Surface plasmon resonance can occur at the boundary of two materials when the frequency of
 973 incoming radiation matches that of the oscillation of the electrons [193], leading to increased
 974 photon intensity being emitted from the surface. Such an effect could be useful for UC where
 975 the converted light has a low intensity. In practice these are often combined with core-shell
 976 design UCNPs. Atre et al [194] constructed an upconverter doped dielectric core encased in a
 977 crescent-shaped Ag shell. They calculated a 100 fold increase in the above band-gap power
 978 emission towards the cell. They concluded that these nanostructures would eventually enable
 979 low-cost single junction solar cells that exceed the Shockley Queisser limit. This concept has
 980 been successfully employed to a DSSC by Liu et al [195] who prepared Au-decorated core-
 981 shell β -NaYF₄:Er³⁺/Yb³⁺@SiO₂ UCNPs. The UC was able to induce localised surface
 982 plasmon resonance which increased the absorption of the dye sensitizer and hence the
 983 photocurrent response. Luoshan's group [196] also employed a similar method but their
 984 UCNPs had an additional @TiO₂ shell layer to provide a transport channel for generated
 985 electrons, enhance dye loading and isolate the Au from the electrolyte. These effects can also
 986 complement other nanostructures for even greater effect. Le and John [197] proposed
 987 synergistic plasmonic and photonic crystal trapping of light to enhance the UC process in thin
 988 film solar technologies. They modelled a slanted conical photonic crystal structure atop core-
 989 shell rings and a UC layer. It was found the IR light intensity in the UC region could reach

990 1400× the ambient solar spectrum, allowing for a much more efficient UC process. An
 991 additional beneficial effect for practical considerations was the improved light trapping and
 992 absorption spectrum for light at incident angles of up to 60°. Furthermore, Ahmed et al have
 993 for the first time demonstrated plasmonic LDS using QDs and Ag-UCNPs, resulting in a 60%
 994 increase in emission intensity and 22% rise in J_{SC} for both silicon and DSSCs over the 300-
 995 500nm illumination range [198].



996

997 **Fig. 13. The unit cell of the synergistic plasmonic and photon crystal structure with perfectly**
 998 **matched layers, denoted PML. Figure taken from [197].**

999 7. Discussion: challenges and opportunities

1000 As presented over the last decade there have been significant advances in UC; from a largely
 1001 conceptual field to the construction of efficiency enhanced physical systems. However
 1002 despite this progress the overall efficiency increases are modest at best and are yet to surpass
 1003 the Shockley-Queisser limit as predicted or be suitable for mass commercial deployment. We
 1004 must therefore ask what further progress must be made, if we are to defeat these challenges

1005 and utilise spectral conversion for solar cells in the world energy market as an alternative to
1006 fossil fuels?

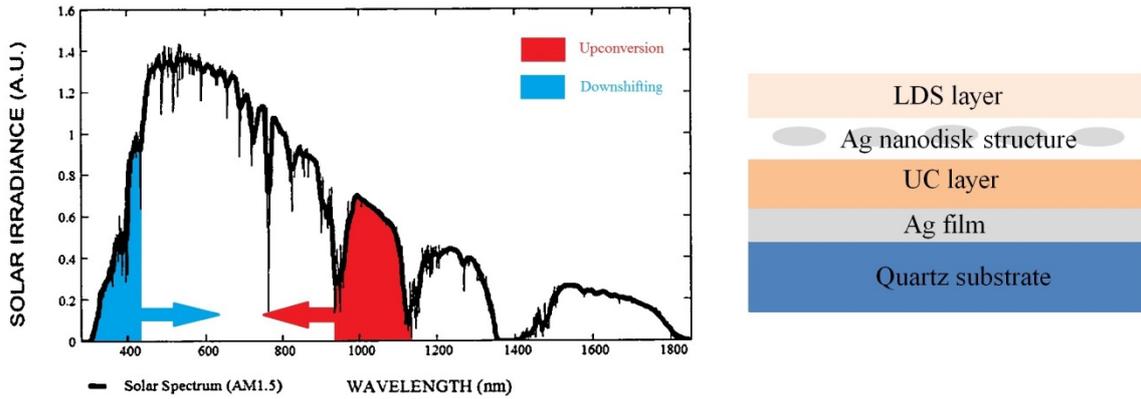
1007 The development of dedicated solar cells for UC-PV is a critical area of research to take this
1008 field forward. Rüdiger et al [199] present ways to practically implement these designs by
1009 identifying three criteria for a bifacial solar cell being used in an up converting system to
1010 fulfil: a very high transmittance to photons of wavelength $< 1200\text{nm}$ to illuminate the UC
1011 layer, a high EQE when illuminated by upconverted photons from the rear and a high
1012 efficiency itself to optimize the system. From UC data as of 2014 they found this could
1013 reasonably attain a 3% relative enhancement in overall efficiency and then fabricated a
1014 bifacial n-type silicon solar cell which achieved an EQE of 1.69% under 1508nm illumination
1015 (5 times higher NEQE than non-optimized cells). They concluded both sides being textured
1016 rather than planar improved performance and the double-layer ARC helped reduce optical
1017 losses. Further aspects for consideration are the external electronic circuit and ensuring
1018 minimal interference with the UC system and the ultimate design for physical deployment in
1019 panels of these structures. For example, how would they have the long term stability of on the
1020 market modules? There needs to follow a comprehensive classification of materials,
1021 including the ones presented in this review, for their effectiveness in increasing over all solar
1022 cell efficiency. It can then eventually be decided which materials are prime for adaptation in
1023 real systems; Er^{3+} / Er^{3+} - Yb^{3+} doped NaYF_4 or its competitors? The recent review on
1024 advances in NaYF_4 luminescence by Kumar et al [200] or a comparative analysis between
1025 fluorides by Favilla et al [201] are examples of this. A range of experimental synthesis
1026 techniques have been reported for the preparation of spectral modification materials in the
1027 studies referred to throughout this paper. Although discussing general optimization
1028 parameters such as doping concentration, polymer weight ratio and choice of materials, it is
1029 worth comparing the fabrication procedures as these will impact viability of research (since
1030 the reader may wish to know what can be constructed in their lab) and potentially
1031 commercialization. The costs and environmental impact of each method should be a matter
1032 for debate. Verification procedures such as x-ray diffraction, scanning electron microscopy
1033 and differential scanning calorimetry are widely used throughout the literature. These
1034 techniques should be utilized along with further modelling studies to maintain a decent
1035 theoretical understanding of the underlying mechanisms.

1036 Concentrator-UC systems, with components operating at distinct spectral intensities (cf. Fig.
1037 18) must be greater looked into and for a range of materials. Two-photon lithography
1038 allowing for the microscale construction of optics could be an exciting area to develop UC
1039 concentrators for thin film solar technologies [202]. The new fields of UC via thermal
1040 radiation and QDs can be contrasted for efficiency gains in these conditions as well,
1041 especially the latter since practical applications of that approach are lacking due to high
1042 concentrations required. The use of novel new materials like graphene may also aid the QD
1043 approach to UC being tested experimentally. Zhang et al [203] have reviewed the remarkable
1044 properties of graphene QDs for a range of applications including solar PV and in 2013 Lee et
1045 al [204] fabricated them and tested their effects on DSSCs.

1046 TTA-UC, again lacking applications to cells in the literature may prove to be an attractive
1047 route to increasing solar cell efficiency in the near future. Not only does its low required light
1048 intensities offer simpler deployment (perhaps eventually to areas of low-moderate DNI) but
1049 from a financial perspective its fundamental materials are not too expensive either. Schulze
1050 advocated zinc porphyrins to sensitise TTA-UC systems, which still give high efficiencies but
1051 at a fraction of the cost of heavier metals (since zinc is only approximately \$2/kg) [205].
1052 However as with many emerging photochemical technologies further research is needed into
1053 oxygen defence and encapsulation in order to prevent photo-degradation. To support this
1054 outlook, recently Wu's group have prepared solid state TTA-UC materials that utilize
1055 thermally activated delayed fluorescence molecules and noted the lack of heavy metals such
1056 as Pt or Pd could lower costs if the technology becomes widespread [206]. To further
1057 maximise potential enhancement to PV performance, multi-mode or broadband UC devices
1058 should be constructed and their benefits to solar cell efficiency compared with (or even
1059 combined with) materials that increase a given ions absorption spectrum. These steps of
1060 research could help to take UC experiments beyond their meagre increases in overall cell
1061 efficiency. Multimodal phosphors either in separate layers as proposed by Lahoz or in Shao's
1062 group's UCNPs designed in 6.3 should be tested for efficiency improvements and compared
1063 with what can be achieved through spectral conversion aided by the use of QDs. Since for
1064 spectral conversion to be a realistically viable approach it will have to capture as much of the
1065 spectrum as possible.

1066 This author would propose further investigations drawing advice from this review, into a
1067 range of solar cells including the very latest emerging technologies which have only recently
1068 been investigated for UC/DC applications. DC/LDS should be applied to more solar cells,
1069 especially those made from c-Si and CdTe as despite recent progress testing for overall
1070 efficiency measurements is still less frequently reported than UC. Also it was felt DC/LDS
1071 lacked the comprehensive optimization studies and standard lattices prevalent in the UC
1072 literature, so this should be an area of research interest in the near future. Simultaneous up
1073 and DC is a concept that should also be explored, to *both* harness sub band-gap photons *and*
1074 reduce losses from photons with excess energy. This could in theory highly enhance photon
1075 energy or at least surpass the Shockley-Queisser limit while allowing for imperfections in
1076 each conversion layer. There is very little experimental research into this, partly because of a
1077 number of physical dilemmas it raises. The addition of a DC or LDS layer raises the issue of
1078 photons of wavelengths desirable to the solar cell below being absorbed into the conversion
1079 layer and it crucially must have very high transmission to photons of UC energies. One proof-
1080 of-principle experiment for this concept was carried out by Yao et al [207] in which rare earth
1081 doped zinc oxide up and DC layers were introduced into a DSSC. Er^{3+} and Yb^{3+} were used
1082 for the UC layer, Eu^{3+} and Tb^{3+} for the DC. The DC layer also acted to reduce charge
1083 recombination. A 70% enhancement in efficiency was observed in comparison to the
1084 standard DSSCs, one of the largest relative increases cited in this review. In another recent
1085 novel study, Lee et al [208] designed and fabricated a metal nanodisk-insulator-metal
1086 structure which simultaneously enhances UC and LDS luminescence via plasmonic
1087 resonance. It was created by a nanotransfer printing method which can easily replicate this
1088 structure over a large area and achieved outstanding enhancements of UC and LDS intensity

1089 by factors of 174 and 29 respectively. These improved modifications from both NIR and UV
 1090 to visible could have remarkable potential in capturing even more of the solar spectrum for a
 1091 wide band-gap cell. Materials that can operate as up or down converters may reduce costs in
 1092 making these types of systems. Same material, multi-modal UC and DC emission has been
 1093 observed in compounds such as Er^{3+} doped K_2YbF_5 [209], Tb^{3+} - Yb^{3+} co-doped lithium
 1094 borate glass [210] and Tb^{3+} - Yb^{3+} co-doped GdPO_4 [211] which could all be attractive for PV
 1095 applications. Additionally, efficient DC has been reported in rare earth doped $\alpha\text{-NaYF}_4$ [212]
 1096 which contains the same elements and could be synthesised alongside the popular UC
 1097 phosphor $\beta\text{-NaYF}_4$.



1098 **Fig. 14. Graphical representation of the promising simultaneous spectral conversion (UC and**
 1099 **LDS) and the schematic of the layers fabricated. Figure modified from [208].**
 1100

1101 Finally, a quantitative financial comparison must be made between spectral modification and
 1102 competing solutions to surpassing the Shockley-Queisser limit, building on the cost analysis
 1103 work which came from the EPHOCELL project [213]. Multijunction cells use many
 1104 expensive materials and are difficult to fabricate, resulting in a high cost per watt even with
 1105 the use of concentrators, so a cost-effective efficient spectral modification layer on a single
 1106 junction cell could still be a sensible solution over multi-junction CPV. To demonstrate this
 1107 point mathematically we derive the condition for an efficiency enhancing modification to
 1108 lower the cost per peak watt. First, consider a single junction cell of arbitrary area with a cost
 1109 C_{cell} , maximum power output P_{max} and cost per peak watt C_{cell} / P_{max} . Then if a modification
 1110 layer with cost $C_{modification}$ is added to give a relative increase in efficiency ΔPCE , the new
 1111 cost per peak watt will be:

1112

1113

1114

$$1115 \quad \text{New cost per watt peak} = \frac{C_{cell} + C_{modification}}{(1 + \Delta PCE)P_{max}} \quad (\text{Equation 5})$$

1116 This can be used to directly compare costs with other approaches. Furthermore, comparing
 1117 the above equation with the cost per watt peak of the original cell and dividing by C_{cell}/P_{max}
 1118 yields the condition for the modification to lower the cost per watt peak:

1119

$$1 + \frac{C_{modification}}{C_{cell}} < \frac{1}{1 + \Delta PCE} \quad \text{(Equation 6)}$$

1120

1121 By using equation 12, approximate rare earth prices (assuming they dominate compared to
 1122 the lattice material), solar cell costs from Table 1 and the results from state of the art studies
 1123 presented in this review, an important estimation can be determined: the maximum mass of
 1124 rare earth than can be added per W of cell whilst still reducing the cost per W. Similar
 1125 methods and cost analysis by Ross which calculated the cost per W of power generated [214]
 1126 are a useful exercise and should be considered in future spectral modification studies. Le
 1127 Donne et al carried out a simple cost analysis [215] from their LDS Eu^{3+} complex on silicon
 1128 experiment and estimated the price per watt peak could fall by €0.06 if a 0.65% enhancement
 1129 in the J_{SC} could be obtained and the EPHOCELL project predicted similarly promising
 1130 economics. Given the current data for efficiency gains this may be the only likely way for UC
 1131 or DC to emerge into the market; small efficiency gains at a small cost per unit area.

1132 **Table 5 – Approximations for the maximum mass of rare earth ions that can be added per W of**
 1133 **solar cell to still reduce the cost per W.**

1134 If these practical efficiency gains are to be realized then the nanostructures discussed should
 1135 be applied to more physical set ups, in order to see if their reported effects can translate to
 1136 greater PCE enhancements. Advances in simulations based on rate equations will allow
 1137 theoretical models to predict and aid their implementation in physical systems. However,
 1138 their development may prove expensive if they are to be implemented on a larger scale.
 1139 Furthermore, rare earth metals necessary for up converting ions are also costly components
 1140 [216], so we must consider their financial contribution to the construction of UC solar panels.
 1141 For DC and LDS, whether the ions are doped into encapsulate glasses or applied directly via
 1142 a thin film will influence cost. LDS glasses could be coloured to an aesthetically pleasing
 1143 look, increasing their potential for building integrated PV. These factors should be
 1144 investigated and a matter for debate, to justify if the relative efficiency gains are worth the
 1145 added costs. Notwithstanding this competition, it may be possible to combine spectral
 1146 conversion and tandem solar cell technologies. It has been proposed III-V silicon tandem
 1147 cells may benefit from UC since the silicon component produces a lower current so would
 1148 benefit more for a given UC performance [217]. It also may make sense to have a UC layer
 1149 for the lowest band-gap material, since despite the use of multi-junctions widening the
 1150 absorption spectrum they do not cover the whole solar spectrum. Although perhaps it is a
 1151 futile exercise, as there is minimal energy to be gained from up converting photons with less
 1152 energy than the narrowest solar cell band-gaps. Another solution to attain theoretical
 1153 efficiencies above the Shockley-Queisser limit is multiple excitation generation (MEG); a

1154 phenomena where an absorbed photon can excite more than one electron to the conduction
1155 band in a semiconductor. Investigations have led to peak EQE's of over 100% reported in a
1156 QDSC [218] and Shpaisman et al [219] concluded MEG and UC could be combined to give a
1157 PCE of 38% for a single junction cell, although at a very optimistic efficiency for processes
1158 (25% for UC and 70% for MEG).

1159

1160 **8. Conclusion**

1161 We have reviewed spectral modification as a means to achieve greater solar cell efficiencies
1162 which is an important research area in the context of environmental issues and current
1163 engineering limitations. After briefly overviewing the theoretical basis for UC, DC and LDS,
1164 a number of studies have been presented that successfully attempt to apply these processes to
1165 harnessing greater energy from the solar spectrum. From the investigations using rare earth
1166 doped compounds, it seems the cells most likely to benefit from UC are wide band gap
1167 technologies such as amorphous silicon and DSSCs. On the contrary DC and LDS stand to
1168 improve cells with a poor UV response such as c-Si, CdTe and CIGS. However, all cells have
1169 an optimal wavelength for conversion efficiency, so as a handful of studies have tried, we
1170 may see more UC working together with DC or LDS on the same cell to efficiently shift as
1171 much of the incident spectrum towards this value as possible. The now commonly used c-Si
1172 cell appears to benefit from both upward and downward spectral modification, so this may be
1173 a focus of research although this depends on if emerging solar technologies improve towards
1174 commercial deployment in the coming years. Unfortunately, limitations arising from poor
1175 absorption and hence spectral conversion efficiency have prevented the exceptional gain in
1176 solar cell efficiency that is theoretically possible. However, nanoengineered phenomena such
1177 as plasmonic resonance and photonic crystals, in addition to macroscopic augmentations
1178 through concentrating optics have the potential to mitigate these drawbacks. Furthermore,
1179 novel new approaches such to UC via TTA, QDs or thermal radiation may aid in overcoming
1180 these problems. In conclusion spectral modification may not be the holy grail of solar cell
1181 design it first appears and in terms of performance is severely lagging behind concentrating
1182 multijunction PV. However, if research continues (including proposals in Section 7) while
1183 costs are kept low, investigations carried out may yield exciting prospects for expansion of
1184 solar PV.

1185

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1189

1190 **References**

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- Solar cell efficiency enhanced up to 70% relative by converting incident spectrum
- Range of cell technologies attain greater performance at short and long wavelengths
- 35 fold increase in short circuit current density enhancement from 2011-15
- Further potential from optics and nanostructures which are independently optimized
- Cost per watt could be reduced by €0.06 and economic analysis of up-to-date results

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Technology	Maximum efficiency (%)	Cost per W_p (\$)
Silicon (monocrystalline)	25.0	1.10-1.40
Silicon (multi-crystalline)	21.3	0.80-0.90
Silicon (amorphous)	13.6	0.45-0.53
Copper Indium Gallium Selenide	22.1	0.55-0.65
Cadmium Telluride	22.1	0.50-0.60
Dye Sensitized	11.9	0.25-0.40
Perovskite	22.1	0.20-0.30
Concentrator Multijunction	46.0	1.50-2.30

UC material	$\lambda_{\text{excitation}}$ (nm)	$\lambda_{\text{emission}}$ (nm)	Solar cell	Device response	Ref.
Er ³⁺ doped NaYF ₄	1460-1600	980	c-Si	+0.69 mA/cm ² J _{SC} under 732 suns	103
Er ³⁺ doped NaYF ₄	1450-1600	980	c-Si	+4.03 mA/cm ² J _{SC} under 77 suns	104
Er ³⁺ doped NaYF ₄	1450-1600	980, 805, 655, 540	c-Si	+13.1 mA/cm ² J _{SC} under 210 suns	115
Er ³⁺ doped BaY ₂ F ₈	1493	1000	c-Si	+17.2 mA/cm ² J _{SC} under 95 suns	105
Er ³⁺ doped NaYF ₄	1523	980	c-Si	1.33% EQE (without optic) and 1.80% EQE (with optic) under 0.024 W/cm ² laser	118
Er ³⁺ doped BaCl ₂	808, 980	667, 546, 530, 496, 455, 410	a-Si:H	+0.6 mA/cm ² J _{SC} under bi- wavelength 70 and 100 mW lasers	112
Er ³⁺ doped LaF ₃ in oxyfluoride ceramic	808, 980, 1530	657, 542	a-Si:H	+1.7 mA/cm ² J _{SC} under tri- wavelength 1400, 1200 and 800 W/cm ² lasers	113
CeVO ₄ via thermal UC	1064, 980, 808	750-700	a-Si:H	+0.91 mA/cm ² J _{SC} under 831 W/cm ²	124
Yb ³⁺ - Er ³⁺ co-doped NaYF ₄	980	650, 536, 518	a-Si:H	+1.74 mA/cm ² J _{SC} and +0.41% PCE under AM 1.5G	110

Yb ³⁺ -Er ³⁺ co-doped TiO ₂	980	700-510	DSSC	+0.87% PCE under AM 1.5	129
Yb ³⁺ - Er ³⁺ co-doped CeO ₂	980	525-680	DSSC	+1.54% PCE under AM 1.5	131
Yb ³⁺ - Er ³⁺ co-doped NaYF ₄	975	655, 540	OSC	+0.46% PCE under AM 1.5 G	135
Tm ³⁺ - Yb ³⁺ - Er ³⁺ co-doped NaYF ₄	980	650, 540, 520, 475, 400	QDSC	+2.15 mA/cm ² J _{SC} +0.73% PCE (+17.6% and 20% relative increases) under AM 1.5 G	139
Yb ³⁺ -Er ³⁺ co-doped ZrO ₂	970	725-500	QDSC	+1.06% PCE (+55% relative increase) under 0.1 W/cm ²	140
Yb ³⁺ - Er ³⁺ co-doped LiYF ₄	1540, 980	700-500	PSC	+0.87% PCE (+7.9% relative increase) under 7-8 suns	145
Yb ³⁺ - Er ³⁺ co-doped NaYF ₄	980	655, 543, 523, 408	PSC	+1.9% PCE (+13.7% relative increase) under AM 1.5 G	147
Yb ³⁺ -Er ³⁺ co-doped TiO ₂ nanorod arrays	980	512-680	PSC	+2.6% PCE (+20.8% relative increase) under AM 1.5 G	148

DC material	$\lambda_{\text{excitation}}$ (nm)	$\lambda_{\text{emission}}$ (nm)	Solar cell	Device response	Ref.
Eu ³⁺ in silicon nitride matrix	325	980	c-Si	N/A	161
Eu ³⁺ containing aromatic complex in PMMA	250-360	665	c-Si	+5.8% EQE under 100W Xe lamp	162
Tb ³⁺ - Yb ³⁺ co-doped tellurite glass	355, 482	548, 980	c-Si	+7% relative PCE under 0.1W/cm ²	163
Eu ³⁺ - Dy ³⁺ co-doped ZnO (alongside graphene loading)	338, 394	394, 458, 483, 575, 593, 611	DSSC	+6.06 mA/cm ² and +3.18% PCE (212% and 245% relative increases) under 0.1W/cm ²	164
Eu ³⁺ doped ZnGa ₂ O ₄	380-500	450-720	PSC	+3.48 mA/cm ² and +3.67% PCE (17.2% and 34.4% relative increases) under AM 1.5 G	165

LDS material	$\lambda_{\text{excitation}}$ (nm)	$\lambda_{\text{emission}}$ (nm)	Solar cell	Device response	Ref.
Eu ³⁺ doped Gd ₂ O ₂ S	200-350	624	mc-Si	+6.47 mA/cm ² and +2.67% PCE under 1 sun	173
Eu ²⁺ doped phosphors in SiO ₂	250-360	512, 610	c-Si	+19.85% J _{SC} and +16% relative PCE under AM 1.5 G	174
Tm ³⁺ doped fluoride glass	355, 482	548, 980	c-Si	+1.4% relative PCE under AM 1.5 G	175
Organosilicon lumiphores in EVA or PMMA	375	575	CIGS	+54% EQE at 360nm and +4.3% relative PCE under AM 1.5 G	178
Organic dyes in EVA and FEP	375, 475, 500, 520	425, 525, 560, 575	CdTe	+5.3% and +9.7% J _{SC} under 1.5 AM G	179
QDs in PMMA	300-500	540	CdTe	+1.7% PCE under AM 1.5 G	180
Mn ²⁺ - Eu ³⁺ doped phosphate glass	392	615	CdTe	+7.14% relative PCE under 0.1 W/cm ²	181

Dyes in PMMA	425-475	500-650	CdTe	+5% relative J_{SC} (non-concentrating) and up to +20% relative J_{SC} (concentrating) under 300 W Xenon lamp	182
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Rare earth ion	Process	Δ PCE relative (%)	Solar cell	Maximum mass of ion to reduce cost per watt (mg/W)	Ref.
Eu ³⁺	UC	20.8	PSC	0.16	161
Eu ³⁺	DC	34.4	PSC	0.06	162
Tb ³⁺ - Yb ³⁺	LDS	16.0	c-Si	0.15	163
Eu ³⁺ - Dy ³⁺	UC/DC	70.0	DSSC	0.44	164