

Magnetoconductivity of icosahedral and amorphous Al–Pd–Re films — a comparison

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Abstract

The magnetoconductivity of quasicrystals is often discussed in the frame of quantum corrections, namely weak (anti-) localisation and electron–electron interaction. A premise for both effects is a strong elastic scattering of conduction electrons. Amorphous and icosahedral phases are discussed as Hume–Rothery alloys with an electronically induced structural peak at the diameter of the Fermi sphere. Therefore, both should exhibit quantum corrections. The preparation of quasicrystalline films via the amorphous route offers the possibility to compare the magnetoconductivity on samples of identical composition but different structure. We report on magnetoconductivity measurements at temperatures between 0.2 and 22 K and for magnetic fields up to 16 T. With the exception of the electronic diffusion constant, amorphous as well as icosahedral Al–Pd–Re films can be described by nearly the same set of parameters if the samples are well on the metallic side of the metal–insulator transition. © 2000 Elsevier Science B.V. All rights reserved.

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1. Introduction

Quantum corrections to the resistivity are well known for amorphous systems in the two-dimensional (2D) [1] or three-dimensional (3D) cases [2,3]. The condition for the occurrence of quantum corrections is a strong elastic scattering which generates an extremely short mean free path in the range of only several atomic distances. However, the coherence length of the electronic wave function is much larger, since it is determined by inelastic scattering processes. In amorphous systems, the large elastic scattering is brought about by scattering from the short or medium range order. This fact can be expressed in the condition

$$K_{pe} = 2k_F \quad (1)$$

with k_F the Fermi vector and K_{pe} the position of an electronically induced peak in the structure factor $S(K)$ [4].

In the quasicrystalline state, it is a priori not clear where a large elastic scattering should come from, as the structure is well ordered. However, it has been shown for Al–Cu–Fe that the most intense diffraction peaks of the icosahedral structure (18, 29) and (20, 32) are at the position of the

electronically induced peak of the amorphous system and also coincide with the diameter of the Fermi sphere [5]. So, Eq. (1) is fulfilled also in the quasicrystalline phase, indicating a matching of the electronic wavelength at the Fermi energy E_F with a characteristic length of the quasicrystalline structure. This may cause a resonant-like elastic scattering [6] of the conduction electrons which leads to the occurrence of quantum corrections and to transport anomalies in the case of a moderate scattering. For a extremely strong scattering, as in some samples of *i*-Al–Pd–Re, this effect may result in a metal–insulator transition (MIT) [7]. Since the electronic states near E_F are no longer extended but localised, the Bergmann concept of *weak* localisation based on the back-scattering due to a huge number of elastic scattering processes breaks down in the vicinity of the MIT.

It has been shown that the low temperature behaviour of barely metallic quasicrystals like *i*-Al–Cu–Fe and *i*-Al–Pd–Mn can be explained in terms of weak localisation (WL) and electron–electron interaction (EEI) theories [8–10]. In the case of (probably) insulating *i*-Al–Pd–Re, Poon et al. [7] mentioned that the magnetoconductivity (MC) could not be explained in the framework WL and EEI theories, while Rodmar and co-workers [11,12] postulated their applicability for samples at least up to a conductivity ratio $\sigma_{295\text{K}}/\sigma_{4.2\text{K}}=10$. Therefore, they were forced to use a

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strongly sample dependent spin-orbit scattering time, which to our opinion is unlikely.

Here we want to compare the conductivity as a function of temperature T and magnetic field B for amorphous and quasicrystalline 3D films on the metallic side of the MIT and discuss them in the framework of WL and EEI theories.

2. Film preparation and characterisation

Al–Pd–Re films were prepared by co-sputtering with two magnetron sources, one with a Re target and the other with a sectional target of Al and Pd. Quartz glass was used as substrate, held at room temperature during the deposition of the amorphous films with a thickness of 220 nm. Due to the positions of the two sources in respect to the substrate, the defined composition gradient was achieved along the substrate holder. With this technique, in one preparation process, a set of 20 amorphous samples was produced, with a composition, slightly and systematically changing from one sample to the next, cutting the ternary phase diagram at, or close to the optimal composition of the quasicrystalline phase. The difference in Re-content between the insulating (f_{in}) and the metallic film (f_m) is 0.8 at.% while the composition of the insulating film is located near $Al_{72}Pd_{20.5}Re_{7.5}$ as estimated by elastic recoil detection and Rutherford back-scattering.

In order to induce a transition from the amorphous to the quasicrystalline state, samples were annealed at 950 K in vacuum ($p_{950K} < 5 \times 10^{-8}$ mbar) for about 10 h. As resistance was measured during the transition, the annealing was stopped when the resistance saturated. X-ray diffraction, SEM and TEM showed [13] that the resulting films considered here are single phased icosahedral with grain sizes up to 1 μ m.

3. Results and discussion

Films may be classified electronically as being either insulating or (poorly) metallic. Insulating 3D films exhibit infinite resistivity or zero conductivity σ at $T=0$. In contrast, 3D films are called poorly metallic if they display non-zero positive conductivity at $T=0$. A useful technique to identify the MIT was previously introduced [14]. The mathematical function $w(T)$

$$w(T) = \frac{d \ln \sigma}{d \ln T} = \frac{(T/\sigma)d\sigma}{dT}, \quad (2)$$

exhibits distinctively different temperature behaviours for insulating and metallic films.

For *strongly insulating* films exhibiting variable-range hopping, with

$$\sigma(T) = \sigma_0 \left[\exp \left(-\frac{T_0}{T} \right) \right]^y, \quad (3)$$

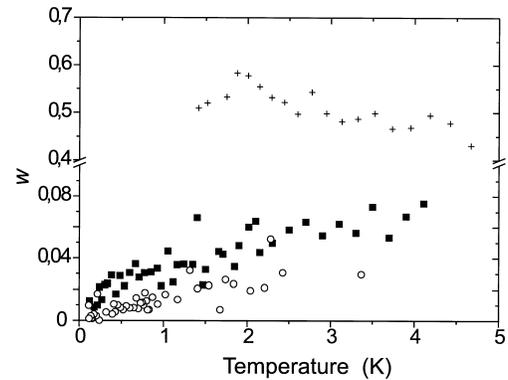


Fig. 1. The $w = d \ln \sigma / d \ln T$ dependence upon temperature for the icosahedral Al–Pd–Re films: f_{in} (+), f_{bi} (■) and f_m (○).

where σ_0 is the prefactor, T_0 the characteristic temperature, and y is an exponent. Inserting Eq. (3) into Eq. (2) yields

$$w(T) = y \left(\frac{T_0}{T} \right)^y. \quad (4)$$

Notice that $w(T)$ increases to infinity as the temperature approaches absolute zero.

In contrast, the conductivity of a 3D metallic film at sufficiently low T can be described by a power law

$$\sigma(T) = \sigma(0) + CT^z, \quad (5)$$

where $\sigma(0)$ is the positive conductivity at $T=0$ and C is the prefactor of the power law term. Eq. (5) might approximate the conductivity contribution of critical states or from the 3D EEI and WL theories. Inserting Eq. (5) into Eq. (2) yields

$$w(T) = \frac{zCT^z}{[\sigma(0) + CT^z]} = \frac{zCT^z}{\sigma(T)}. \quad (6)$$

If the film is *metallic* and exhibits a finite positive conductivity $\sigma(T=0)$, then $w(T)$ should extrapolate to zero for $T \rightarrow 0$ K.

Fig. 1 shows the $w(T)$ behaviour for three different icosahedral films (f_{in} , f_m , f_{bi}). Film f_m exhibits w -values which tend to zero for $T \rightarrow 0$ K and is clearly metallic. In contrast, the w behaviour exhibited by film f_{in} suggests that the w -values extrapolate to a *finite value* or even increase down to smaller temperatures. Hence, this film appears to be insulating. Film f_{bi} seems to be in the vicinity of a MIT as its w -values tend to a very small but probably finite value for $T \rightarrow 0$.

Low- T MC data are plotted in Figs. 2–5 for the three different icosahedral and an amorphous film f_{am} . From a qualitative point of view all *metallic* samples regardless if amorphous or icosahedral exhibit exclusively negative MC for all measured temperatures. This is as expected for films containing high Z elements like Pd, and especially, Re due to strong spin-orbit scattering. The positive MC at $B < 3$ T and $T = 4.2$ K for film f_{in} (Fig. 2) is typical for *i*-Al–Pd–Re samples on the insulating side of the MIT and cannot be interpreted using the theory of WL.

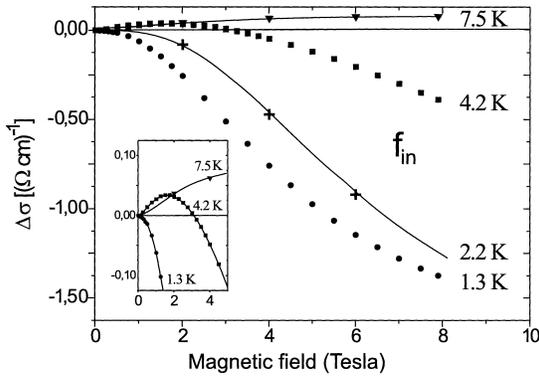


Fig. 2. Magnetoconductivity of the probably insulating film f_{in} as a function of magnetic field.

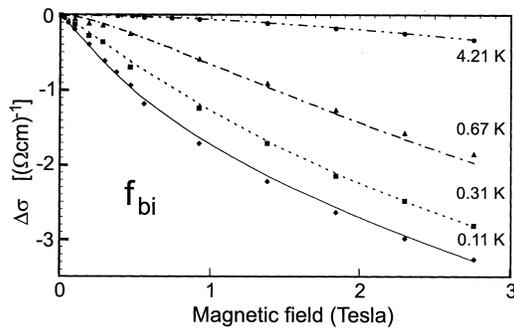


Fig. 3. Magnetoconductivity of the barely insulating film f_{bi} as a function of magnetic field; lines show a fit according to WL and EEI contributions.

As summarised in Fig. 6 for 20 different films, the absolute change of the conductivity $\Delta\sigma$ by a magnetic field $B=8\text{ T}$ at $T=1.3\text{ K}$ is very similar in size and depends only slightly from the conductivity. Especially, no significant change occurs at the transition from amorphous to icosahedral Al–Pd–Re films. Thereby, weak (anti-) localisation contributes about 80% to the negative MC magnitude. In the high field limit, it is proportional to $D^{-1/2}$ (D : electronic diffusion const.) and a function of the spin-orbit τ_{so} and inelastic scattering times $\tau_{ie}(T)$. As both scattering times probably do not show a considerable change with

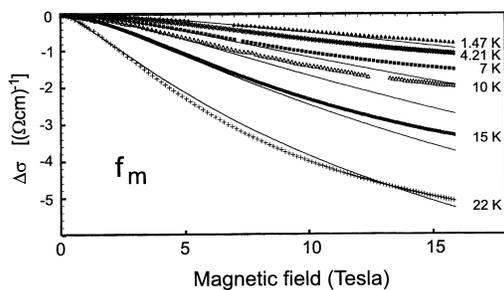


Fig. 4. Magnetoconductivity of the 'metallic' film f_m as a function of magnetic field; lines show a fit according to WL and EEI contributions.

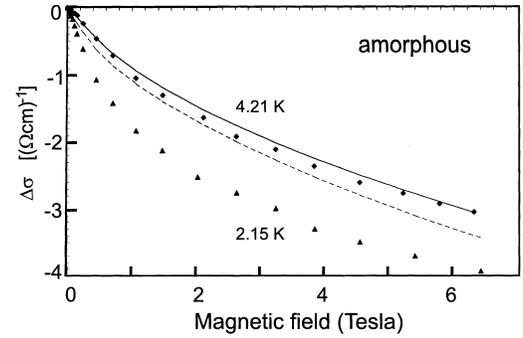


Fig. 5. Magnetoconductivity of an amorphous film f_{am} as a function of magnetic field; lines show a fit according to WL and EEI contributions.

structure, Fig. 6 suggests that also the diffusion constant is nearly constant crossing the transition from amorphous to icosahedral structure. This may be astonishing as the conductivity changes by a factor of 10^2 but is in accordance with measurements of the Hall mobility [15] which also shows a comparably small change.

For a quantitative analysis, the MC data for the metallic icosahedral film f_m were fitted using the contributions arising from EEI in the particle-hole [16] and the 3D WL including spin-orbit scattering and Zeeman splitting [17]. According to the lack of any other better formalism, we use the WL theory close to the MIT for film f_m ; the WL theory generally applies to highly metallic films.

A magnitude for the spin-orbit scattering time [18,19] was estimated from the expression $\tau_{so} \approx \tau_o(137/Z)^4$ to be $\tau_{so} \approx 10^{-13}\text{ s}$, where Z is the atomic number and $\tau_o \approx 10^{-15}\text{ s}$ is the elastic scattering time. The resulting fitting parameters are listed in Table 1 and are quite similar to the values reported in [11,12] extracted from the *metallic ones* of the bulky *i*-samples. The inelastic scattering times yielded from the MC fit at different T could be fitted to a simple power law $\tau_{ie}(T) = \tau_{ie0} T^p$, where $p \approx -1$. The MC fits describe the two metallic of the icosahedral films quite well while the MC of the amorphous film at $T=2.15\text{ K}$ is strongly influenced

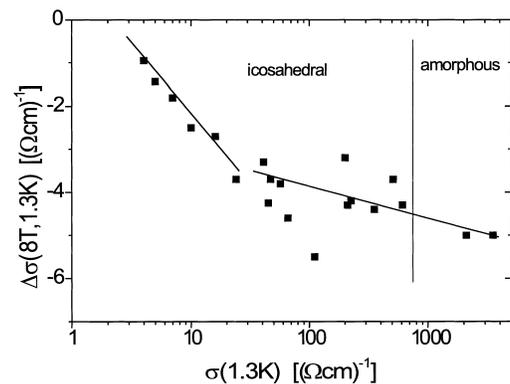


Fig. 6. Magnetoconductivity at $B=8\text{ T}$ and $T=1.3\text{ K}$ as a function of the conductivity at $T=1.3\text{ K}$ for icosahedral and amorphous Al–Pd–Re films.

Table 1
Fitting parameters to the MC data^a

Film	τ_{so} (s)	τ_{ie} (4.2 K) (s)	p	F	D (cm ² /s)
f_m	1.5e-13	3.6e-12	-1.03	0.2	0.75
f_{bi}	1.5e-13	1.8e-11	-1.03	0.2	0.15
f_{am}	1.5e-13	5.5e-11	-2	0.2	1.5

^a τ_{so} and τ_{ie} denote the spin-orbit and inelastic scattering times, respectively, F the electron screening parameter, D the electronic diffusion constant and p denotes the exponent of the power law for inelastic scattering times.

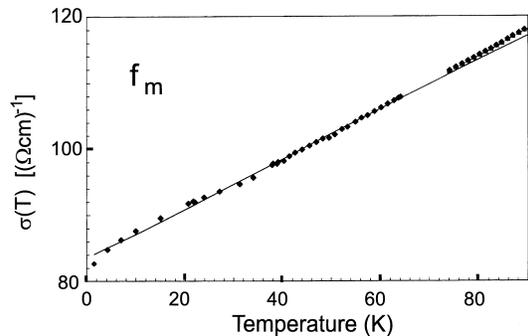


Fig. 7. Zero-field conductivity of the 'metallic' film f_m ; the line shows a fit according to WL and EEI contributions as obtained from the magnetoconductivity data.

from superconducting fluctuations which are a precursor to the superconducting transition at $T=0.55$ K.

The *zero-field* conductivity data of *i*-films f_m and f_{bi} are compared to data computed by using the parameters which were determined from the MC fits. The single parameter fit of $\sigma(T)=\sigma(T=0)+\Delta\sigma_{WL}+\Delta\sigma_{EEI}$ is impressively good for film f_m between 4 and 60 K with $\sigma(T=0)=18.7$ (Ohm cm)⁻¹ as shown in Fig. 7. In this case, the EEI expression contributes 70% to the total conductivity change with temperature, while the WL expression makes a smaller 30% contribution. Interestingly, the WL expression exhibits anti-localisation below 6 K.

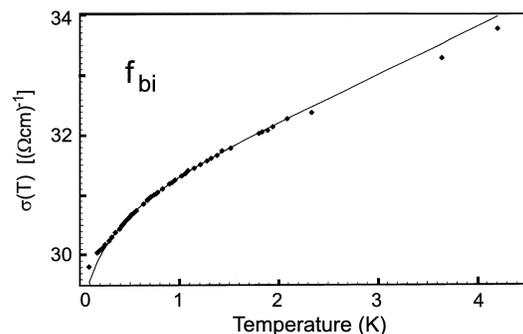


Fig. 8. Zero-field conductivity of the barely insulating film f_{bi} ; the line shows a fit according to WL and EEI contributions as obtained from the magnetoconductivity data.

For the barely insulating film f_{bi} (Fig. 8), initial inspection indicates agreement between data and theory. But the fit is unphysical, since a negative conductivity $\sigma(T=0)=-85$ (Ohm cm)⁻¹ had to be used. Thus, the WL contribution greatly overestimates the zero-field conductivity; and the WL theory breaks down already in the vicinity of the MIT. The above illustration marks the importance of fitting both MC and zero-field conductivity data only to metallic film and checking for consistency between all the fits.

4. Conclusions

Magnetoconductivity as well as zero-field conductivity data of icosahedral as well as amorphous Al-Pd-Re films can be described by WL and EEI theories if the films are well on the metallic side of the metal-insulator transition. Besides the electronic diffusion constant, amorphous and icosahedral films of very different conductivities can be described by nearly the same set of parameters.

Acknowledgements

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