

# A QFT Approach to Data Streaming in Natural and Artificial Neural Networks <sup>†</sup>

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**Abstract:** In the actual panorama of machine learning (ML) algorithms, the issue of the real-time information extraction/classification/manipulation/analysis of data streams (DS) is acquiring an ever-growing relevance. They arrive generally at high speed and always require an unsupervised real-time analysis for individuating long-range and higher order correlations among data that are continuously changing over time (phase transitions). This emphasizes the infinitary character of the issue, i.e., the continuous change of the signifying number of degrees of freedom characterizing the statistical representation function, challenging the classical ML algorithms, both in their classical and quantum versions, as far as all are based on the (stochastic) search for the global minimum of some cost/energy function. The physical analogue must be studied in the realm of quantum field theory (QFT) for dissipative systems as biological and neural systems, which are able to map between different phases of quantum fields, using the formalism of the Bogoliubov transform (BT). By applying the BT in a reversed way, on the system-thermal bath energetically balanced states, it is possible to define the powerful computational tool of the “doubling of the degrees of freedom” (DDF), making the choice of the signifying finite number of the degrees of freedom dynamic and then automatic, so to suggest a different class of unsupervised ML algorithms for solving the DS issue.



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## 1. Introduction: The Infinitely Many Degrees of Freedom in Data Streaming

During the last twenty years, considerable research has been conducted on the development of probabilistic machine learning algorithms, especially in the field of the artificial neural networks (ANN), for dealing with the problem of *data streaming* classification and, more generally, for the *real-time* information extraction/manipulation/analysis of (infinite) data streams (see [1–4] for updated reviews on this topic). For instance, sensor networks, healthcare monitoring, social networks, financial markets, etc., are among the main sources of data streams, often arriving at high speed and always requiring a real-time analysis, above all for individuating *long-range and higher order correlations* among data that are continuously *changing over time*.

Indeed, the standard statistical machine learning algorithms in ANN models, starting from their progenitor, the so-called *back-propagation* (BP) algorithm [5], were developed for *static* and huge bases of data (“big data”), but are systematically inadequate and unadaptable for the analysis of DS. That is, for the analysis and processing of *dynamic* bases of data, characterized by sudden changes in the correlation length among the variables (phase transitions), and by the unpredictable variation of the number of the signifying degrees of freedom of the probability distributions. We will refer to the characterization

of the system deriving from the dynamically evolving set of infinitely many degrees of freedom as the *infinitary* character of the system.

From the computational standpoint, the solution to the *infinitary* character of the DS problem is in principle unreachable by a Turing Machine (TM), either classical or quantum (QTM). Indeed, for dealing with the DS infinitary challenge, the increasing of the computational speed, derived by the usage of quantum machine learning algorithms, is not very helpful, either using “quantum gates” (QTM), or using “quantum annealing” (*quantum Boltzmann Machine* (QBM)); both objects of intensive research during the last few years (see [6–8] for updated reviews).

Generally, in the case of ANNs, the improvement provided by the *Boltzmann-Machine* (BM) learning algorithm to the stochastic gradient descent (GD) algorithm, for the weight refresh of the BP learning algorithm, is that the BM uses “thermal fluctuations” for jumping out of the local minima of the cost function (*simulated annealing*), to avoid the main limitation of the GD algorithm in machine learning [9].

In this framework, the advantage of *quantum annealing* in a QBM is that it uses the “quantum vacuum fluctuations” instead of the thermal fluctuations of the classical annealing for bringing the system out of swallow (local) minima, by resorting to the “quantum tunnelling” effect [10]. This outperforms the thermal annealing, especially where the potential energy (cost) landscape consists of high but thin barriers surrounding shallow local minima [11,12]. However, despite the improvement that, at least in some specific cases, QBM can provide to the procedure for finding the absolute minimum size/length/cost/distance among a very large, even though *finite* set of possible solutions, the problem of DS remains because in this case this “finitary” supposition does not hold.

## 2. The Analogy with the Infinitary Character of QFT Dynamics in Brains

As the analogy with the coarse-graining problem in statistical physics emphasizes very well, the search for the global minimum of the energy function makes sense *after* the system performed a phase transition. That is, after that a sudden change in the correlation length among the variables of the system, generally under the action of an *external field*, determined a new way by which the variables are aggregated for defining the signifying number of the degrees of freedom  $N$ , characterizing the system statistics after the transition.

In other terms, the infinitary challenge implicit in DS is related to *phase transitions* so that, from QFT standpoint, this is the same phenomenon of the infinite number of degrees of freedom of the Haag Theorem [13], characterizing the quantum superposition in QFT systems in *conditions far from equilibrium*. This requires the extension of the QFT formalism to dissipative systems, inaugurated by the pioneering works of N. Bogoliubov [14,15] and H. Umezawa [16–18]. The Bogoliubov transform allows mapping between the different phases of the boson and the fermion quantum fields, making the dissipative QFT—differently from QM and from QFT for closed system—able to describe systems continuously undergoing phase transitions.

Indeed, inspired by the modeling of natural brains as many-body systems, the QFT dissipative formalism has been used to model ANNs [19–21]. The mathematical formalism of QFT (details in [22]) requires that for *open* (dissipative) systems, such as the brain, which is in a permanent “trade” or “dialogue” with its environment, the degrees of freedom of the system (the brain), say  $A$  need to be “doubled” by introducing the degrees of freedom  $\tilde{A}$  describing the environment, according to the coalgebraic scheme:  $A \rightarrow A \times \tilde{A}$ . Indeed, Hopf coproducts (sums) are generally used in quantum physics to calculate the *total energy* of a superposition quantum state. In the case of a dissipative system, the coproducts represent the total energy of a state balanced between the system and its thermal bath. In this case, because the two terms of the coproduct are not mutually interchangeable such as in the case of QM closed systems, we are led to consider the *non-commutative  $q$ -deformed* Hopf (co)algebras, out of which the *Bogoliubov transformations* involving the  $A$ ,  $\tilde{A}$  modes are derived, and where the  $q$ -deformation parameter is a thermal parameter, strictly related with the Bogoliubov transform.

These transformations induce *phase transitions*, i.e., transitions through physically distinct spaces of the states describing different dynamical regimes in which the system can sit. The brain is thus, continuously undergoing phase transitions (*criticalities*), under the action of the inputs from the environment ( $\tilde{A}$  modes). The brain activity is, therefore, the result of a continual balancing of fluxes of energy (in all its forms), exchanged with the environment. The balancing is controlled by the minimization of the free energy at each step of time evolution. Since fluxes “in” for the brain ( $A$  modes) are fluxes “out” for the environment ( $\tilde{A}$  modes), and vice-versa, the  $\tilde{A}$  modes are the *time-reversed images* of the  $A$  modes; they represent the *Double* of the system [23]. In such a way, by the doubling of the algebras—of the state spaces, and of the Hilbert spaces—, and thus by inserting the degrees of freedom of the environment (thermal bath), the Hamiltonian canonical representation of a (closed) dynamic system can be recovered also in the case of a dissipative system.

### 3. From Natural to Artificial Quantum Neural Net Dynamics

From the theoretical computer science (TCS) standpoint, this means that the system satisfies the notion of a particular type of automaton, the *Labelled State Transition Machine* (LTM), i.e., the so-called *infinite-state* LTM, coalgebraically interpreted, and used in TCS for modelling *infinite streams of data* [21,24]. However, the *doubling of the degrees of freedom* (DDF)  $\{A, \tilde{A}\}$  just illustrated, and characterizing a dissipative QFT system, acts as a *dynamic, unsupervised* selection criterion of admissible because balanced states (minimum of the free energy). Effectively, it acts as a mechanism of “phase locking” between the data flow (environment) and the system dynamics.

Moreover, each system-environment entangled (doubled) state is *univocally* characterized by a *dynamically generated code*, or *dynamic labelling* (memory addresses). In our model, an input triggers the *spontaneous breakdown of the symmetry* (SBS) of the system dynamical equations. As a result of SBS, massless modes, called Nambu–Goldstone (NG) modes, are dynamically generated [25,26]. They are boson quanta of *long-range correlations* among the system elementary components, and their *coherent condensation value*  $\mathcal{N}$  in the system ground state (the least energy state or vacuum state  $|0\rangle$ ), that in our dissipative case is a *balanced*, or 0-sum energy state with  $T > 0$ ) describes the recording of the information carried by that input, *indexed univocally (labeled) in*  $\mathcal{N}$  [21].

*Coherence* denotes that the long-range correlations are not destructively interfering in the system ground state. The memory state turns out to be, therefore, a *squeezed coherent state*:  $|0(\theta)\rangle_{\mathcal{N}} = \sum_j w_j(\theta) |w_j\rangle_{\mathcal{N}}$  to which Glauber information entropy measure  $Q$  directly applies [27], with  $|w_j\rangle$  denoting (the statistical weights of the) states of  $A$  and  $\tilde{A}$  pairs, and  $\theta$  is the time- and temperature-dependent Bogoliubov transformation parameter.  $|0(\theta)\rangle_{\mathcal{N}}$  is, therefore, a time-dependent state at finite temperature  $T > 0$ . It is effectively an entangled (system-environment) state of the modes  $A$  and  $\tilde{A}$ , which provides the mathematical description of the *unavoidable* interdependence between the brain and its environment. Coherence and entanglement imply that quantities relative to the  $A$  modes depend on the corresponding ones of the  $\tilde{A}$  modes.

### 4. Conclusion: A Possible Quantum Optics Implementation of the DDF Algorithm

To conclude, the more natural implementation of such a quantum computational architecture for unsupervised data streaming machine learning, based on the DDF principle, is provided by an optical ANN, using the tools of *optical interferometry*, just as in the applications discussed in [28,29]. The fully programmable architecture of this optical chip allows us “to depict” over coherent light waves as many interference figures as we like, and overall, to maintain their *phase coherences* stable in time, so to allow the implementation of quantum computing architectures (either quantum gates, or squeezed coherent states), working at *room temperature*. In our application for data streaming analysis, the DDF principle can be applied in a recursive way, by using the *mutual information* as a measure of *phase distance*, such as an optimization tool of error minimization of the input-output phase

mismatch. In this architecture, indeed, the input of the net is not on the initial conditions of the net dynamics, such as in the ANN architectures based on statistical mechanics, but on the *boundary conditions* (thermal bath) of the system, to implement the architecture of a net in *continuous learning*, as required by the data streaming challenge.

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