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New insights on weighted pseudo almost periodic solutions in a Lasota–Ważewska system

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Abstract

We study the weighted pseudo almost periodic solutions of a Lasota–Ważewska system. With the aid of fixed point theory and differential inequality strategies, we give a set of new sufficient criteria that guarantee the existence and global exponential stability of weighted pseudo almost periodic solutions to a Lasota–Ważewska system. The obtained results of this manuscript are completely innovative and complement the work of Shao (Appl. Math. Lett. 43:90–95, 2015) to some degree. So far, no scholars have investigated this aspect.

MSC: 34C25; 34C13; 34C20; 34K25; 92B20

Keywords: Weighted pseudo almost periodic solution; Lasota–Ważewska system; Global exponentially stability; Time delay

1 Introduction

In 1988, Ważewska-Czyżewska and Lasota [2] proposed the delayed differential model

$$\dot{w}(t) = -\varrho(t)w(t) + \sum_{k=1}^p \theta_k(t)e^{-\eta_k(t)w(t-\rho_k(t))} \quad (1.1)$$

to describe the survival of red cells in an animal [1]. In this model, p is a positive integer, $w(t)$ stands for the number of red blood cells at time t , $\varrho(t)$ stands for the death rate of the red blood cell, $\theta_k(t)$ and $\eta_k(t)$ are related to the production of red blood cells per unit time, and $\rho_k(t)$ represents the time to produce a red blood cell. For details, see [2].

We know that the death rate or harvesting rate usually change under different seasonal fluctuations. In addition, the actual living environment of species have weighted pseudo almost periodic nature due to the effect of human activities and industrial production, for example, the exhaust emission and reconstruction of rivers. Based on this viewpoint, we think that it is reasonable to suppose that the coefficients in model (1.1) are weighted pseudo almost periodic functions, which can be expressed as an almost periodic component plus the weighted ergodic perturbation. Thus a key problem arises: seek the new suf-

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ficient conditions to ensure the existence and exponential stability of the weighted pseudo almost periodic solution for model (1.1).

Unfortunately, until now, no scholars have considered the weighted pseudo almost periodic solutions for model (1.1). To make up for this deficiency and inspired by the previous discussion, in this work, we concentrate on the weighted pseudo almost periodic solutions for model (1.1).

Let $\mathcal{BC}(R, R)$ denote the set of all bounded continuous functions from R to R , and let $\bar{g} = \sup_{t \in R} g(t)$ and $\underline{g} = \inf_{t \in R} g(t)$ for a bounded continuous function $g(t)$. The initial condition of system (1.1) is given by

$$w(s) = \psi(s), \quad s \in [-\bar{\rho}, 0], \tag{1.2}$$

where $\psi \in \mathcal{BC}([-\bar{\rho}, 0], R^+)$, $\rho = \max_{k=1,2,\dots,p} \bar{\rho}_k$, and R^+ denotes the nonnegative real numbers.

We plan the paper as follows. In Sect. 2, we present some preliminary knowledge on weighted pseudo almost periodic solution. In Sect. 3, we investigate the existence and global exponential stability of weighted pseudo almost periodic solution to model (1.1). In Sect. 4, numerical simulations are put into effect. The conclusion is given to end this paper in Sect. 5.

2 Basic knowledge

Throughout this manuscript, let \mathcal{U} denote the collection of functions (weights) $v : R \rightarrow (0, +\infty)$, which are locally integrable on R and satisfy

$$\lim_{\gamma \rightarrow +\infty} v([- \gamma, \gamma]) = +\infty, \quad \text{where } v([- \gamma, \gamma]) = \int_{-\gamma}^{\gamma} v(s) ds \quad (\gamma > 0).$$

Let $\|w\| = |w|$, $g^+ = \sup_{t \in R} |g(t)|$, $g^- = \inf_{t \in R} |g(t)|$, $\mathcal{U}_\infty = \{v | v \in \mathcal{U}, \inf_{t \in R} v(w) = v_0 > 0\}$, $\mathcal{U}_\infty^+ = \{v | v \in \mathcal{U}, \lim_{|w| \rightarrow +\infty} \sup \frac{v(aw)}{v(w)} < +\infty, \lim_{\gamma \rightarrow +\infty} \sup \frac{v([-a\gamma, a\gamma])}{v([- \gamma, \gamma])} < +\infty \text{ for } a \in (0, +\infty)\}$. Denote the set of bounded continuous function from R to R by $\mathcal{BC}(R, R)$. So $(\mathcal{BC}(R, R), \|\cdot\|_\infty)$ is a Banach space with norm $\|g\|_\infty = \sup_{t \in R} \|g(t)\|$. Denote the set of the almost periodic functions from R to R by $\mathcal{AP}(R, R)$. Define

$$PAP_0^v(R, R) = \left\{ \phi \in \mathcal{BC}(R, R) \mid \lim_{\gamma \rightarrow +\infty} \frac{1}{v([- \gamma, \gamma])} \int_{-\gamma}^{\gamma} |\phi(t)| v(t) dt = 0 \right\}.$$

We say that a function $g \in \mathcal{BC}(R, R)$ is weighted pseudo almost periodic if $g = f + \phi$, where $f \in \mathcal{AP}(R, R)$ and $\phi \in PAP_0^v(R, R)$. Here f is called the almost periodic component, and ϕ is called the weighted ergodic perturbation. We denote by $PAP^v(R, R)$ the space of all weighted almost periodic functions. For more detail, see [3–15].

Lemma 2.1 ([2]) *Let $\bar{\varrho} : R \rightarrow (0, +\infty)$ be a bounded continuous function, and let $\alpha^M, \alpha^m, \beta^M, \beta^m, \sigma, \chi$ be positive constants such that*

$$\begin{aligned} \alpha^m e^{-\int_s^t \bar{\varrho}(v) dv} &\leq e^{-\int_s^t \varrho(v) dv} \leq \alpha^M e^{-\int_s^t \bar{\varrho}(v) dv}, \quad \forall t, s \in R \text{ and } t - s \geq 0, \\ -\beta^M &= \left(-\bar{\varrho}(t)\chi + \alpha^M \sum_{k=1}^r \theta_k(t) \right)^+, \quad \chi > \sigma, \\ -\beta^m &= \left(-\varrho^*(t)\sigma + \alpha^m \sum_{k=1}^r \theta_k(t) e^{-\eta_k(t)\chi} \right)^-. \end{aligned}$$

Then there exists t_0 such that the solution $w(t)$ of system (1.1) with initial condition (1.2) satisfies $\sigma < w(t) < \chi$ for $t > t_0$.

We introduce the following assumptions:

(A₁) For $k = 1, 2, \dots, p, \theta_k, \eta_k, \rho_k \in \mathcal{PAP}^v(R, R), \varrho \in \mathcal{AP}(R, R),$

$$M[\varrho] = \lim_{\mathcal{T} \rightarrow +\infty} \frac{1}{\mathcal{T}} \int_t^{t+\mathcal{T}} \varrho(s) ds > 0.$$

(A₂) $v \in \mathcal{U}_\infty^+,$ and there exist constants $\epsilon > 0$ and $\zeta > 0$ such that

$$\sup_{t \in R} \left\{ -\bar{\varrho}(t) + \zeta^{-1} \alpha^M \sum_{k=1}^p |\theta_k(s)| |\eta_k(s)| e^{-\eta_k^- \sigma} \right\} < -\epsilon.$$

3 Main results

Lemma 3.1 Define the operator

$$\Gamma(\psi)(t) = \int_{-\infty}^t e^{-\int_s^t \varrho(v) dv} \left[\zeta^{-1} \sum_{k=1}^p \theta_k(t) e^{-\zeta \eta_k(t) \psi(t-\rho_k(t))} \right] ds, \quad \psi \in \mathcal{PAP}^v(R, R).$$

Then Γ maps $\mathcal{PAP}^v(R, R)$ into itself.

Proof In view of (A₁) and Lemmas 2.1 and 2.3 of [16], we have that

$$\sum_{k=1}^p \theta_k(t) e^{-\zeta \eta_k(t) \psi(t-\rho_k(t))} \in \mathcal{PAP}^v(R, R).$$

According to (A₁)–(A₂) and applying the proof of Lemma 2.1 of [17], we can see that $\Gamma(\psi) \in \mathcal{BC}(R, R).$ Then there exist $\mathcal{E}(t) \in \mathcal{AP}(R, R)$ and $\mathcal{F}(t) \in \mathcal{PAP}_0^v(R, R)$ such that

$$\sum_{k=1}^p \theta_k(t) e^{-\zeta \eta_k(t) \psi(t-\rho_k(t))} = \mathcal{E}(t) + \mathcal{F}(t).$$

Since

$$M[\varrho] = \lim_{\mathcal{T} \rightarrow +\infty} \frac{1}{\mathcal{T}} \int_t^{t+\mathcal{T}} \varrho(s) ds > 0,$$

applying the theory of exponential dichotomy of [18], we get that

$$\int_{-\infty}^t e^{-\int_s^t \varrho(v) dv} \mathcal{E}(s) ds \in \mathcal{AP}(R, R) \tag{3.1}$$

is a solution of the following equation

$$\dot{u}(t) = -\varrho(t)u(t) + \mathcal{E}(t). \tag{3.2}$$

By the proof of Lemma 2.3 of [19] we have that

$$\int_{-\infty}^t e^{-\int_s^t \varrho(v) dv} \mathcal{F}(s) ds \in \mathcal{PAP}_0^v(R, R). \tag{3.3}$$

In view of (3.1) and (3.3), we conclude that

$$\Gamma(\psi)(t) = \int_{-\infty}^t e^{-\int_s^t \varrho(v) dv} \left[\zeta^{-1} \sum_{k=1}^p \theta_k(t) e^{-\zeta \eta_k(t) \psi(t-\rho_k(t))} \right] ds \in \mathcal{PAP}^\nu(R, R)$$

and Γ maps $\mathcal{PAP}^\nu(R, R)$ into itself. This completes the proof. \square

Theorem 3.1 *In addition to the condition of Lemma 2.1, suppose that (\mathcal{A}_1) – (\mathcal{A}_2) are fulfilled. Then model (1.1) has a unique weighted pseudo almost periodic solution, which is globally exponentially stable.*

Proof Let $u(t) = \zeta^{-1}w(t)$. Then system (1.1) becomes

$$\dot{u}(t) = -\varrho(t)u(t) + \zeta^{-1} \sum_{k=1}^p \theta_k(t) e^{-\zeta \eta_k(t) u(t-\rho_k(t))}. \tag{3.4}$$

For $\psi, \phi \in \mathcal{PAP}^\nu(R, R)$, in view of (\mathcal{A}_1) and (\mathcal{A}_2) , we have

$$\begin{aligned} & |(\Gamma\psi)(t) - (\Gamma\phi)(t)| \\ &= \left| \int_{-\infty}^t e^{-\int_s^t \varrho(v) dv} \left[\zeta^{-1} \sum_{k=1}^p \theta_k(s) \left[e^{-\zeta \eta_k(s) \psi(s-\rho_k(s))} - e^{-\zeta \eta_k(s) \phi(s-\rho_k(s))} \right] \right] ds \right| \\ &\leq \int_{-\infty}^t e^{-\int_s^t \bar{\varrho}(v) dv} \zeta^{-1} \alpha^M \sum_{k=1}^p \theta_k(s) \eta_k(s) e^{-\eta_k \sigma} |\psi(s - \rho_k(s)) - \phi(s - \rho_k(s))| ds \\ &\leq \sup_{t \in R} \int_{-\infty}^t e^{-\int_s^t \bar{\varrho}(v) dv} \zeta^{-1} \alpha^M \sum_{k=1}^p |\theta_k(s)| |\eta_k(s)| e^{-\eta_k \sigma} ds \|\psi(t) - \phi(t)\|_\infty \\ &\leq \sup_{t \in R} \int_{-\infty}^t e^{-\int_s^t \bar{\varrho}(v) dv} [\bar{\varrho}(s) - \epsilon] ds \|\psi(t) - \phi(t)\|_\infty \\ &\leq \sup_{t \in R} \int_{-\infty}^t e^{-\int_s^t \bar{\varrho}(v) dv} d\left(-\int_s^t \bar{\varrho}(v) dv\right) - \frac{\epsilon}{2} \int_{-\infty}^t e^{-\int_s^t \bar{\varrho}(v) dv} ds \|\psi(t) - \phi(t)\|_\infty \\ &\leq \left(1 - \frac{\epsilon}{2\bar{\varrho}^+}\right) \|\psi(t) - \phi(t)\|_\infty \end{aligned} \tag{3.5}$$

for $l = 1, 2, \dots, r$. By (\mathcal{A}_2) we easily see that $(1 - \frac{\epsilon}{2\bar{\varrho}^+}) \in (0, 1)$, and thus Γ is a contraction mapping. Thus Γ has a unique fixed point $u^* \in \mathcal{PAP}^\nu(R, R)$ and satisfies $\Gamma u^* = u^*$. Based on this analysis, we conclude that model (1.1) has a unique weighted pseudo almost periodic solution $w^* = \zeta u^* \in \mathcal{PAP}^\nu(R, R)$.

Next, we will prove the exponential stability of the solution w^* . Let $w(t)$ be an arbitrary solution of (1.1) with initial condition $\psi(t)$ satisfying (1.2). Let $v(t) = \zeta^{-1}(w(t) - w^*(t))$. Then we have

$$\dot{v}(t) = -\varrho(t)v(t) + \zeta^{-1} \sum_{k=1}^p \theta_k(t) \left[e^{-\eta_k(t) w(t-\rho_k(t))} - e^{-\eta_k(t) w^*(t-\rho_k(t))} \right]. \tag{3.6}$$

Let

$$\Phi(\varphi) = \sup_{t \in R} \left\{ \varphi - \bar{\varrho}(t) + \zeta^{-1} \alpha^M \sum_{k=1}^p |\theta_k(s)| |\eta_k(s)| e^{-\eta_k^- \sigma} e^{\varphi \rho_k(t)} \right\}, \quad \varphi \in [0, \bar{\varrho}^-]. \tag{3.7}$$

Then

$$\Phi(0) = \sup_{t \in R} \left\{ -\bar{\varrho}(t) + \zeta^{-1} \alpha^M \sum_{k=1}^p |\theta_k(s)| |\eta_k(s)| e^{-\eta_k^- \sigma} \right\} \leq -\epsilon < 0. \tag{3.8}$$

Since $\Phi(\varphi)$ is continuous, we can choose a constant $\mu \in (0, \bar{\varrho}^-)$ such that

$$\Phi(\mu) = \sup_{t \in R} \left\{ \mu - \bar{\varrho}(t) + \zeta^{-1} \alpha^M \sum_{k=1}^p |\theta_k(s)| |\eta_k(s)| e^{-\eta_k^- \sigma} \right\}, \quad \varphi \in [0, \bar{\varrho}^-]. \tag{3.9}$$

Set

$$\|\psi - w^*\|_{\zeta} = \sup_{t \in (-\bar{\rho}, 0]} \zeta^{-1} |\psi(t) - w(t)| \tag{3.10}$$

and choose $\mathcal{N} > \alpha^M + 1$. Then we obtain that

$$\|v(t)\| \leq (\|\psi - w^*\|_{\zeta} + \epsilon) e^{-\mu t} < (\mathcal{N}(\|\psi - w^*\|_{\zeta} + \epsilon) e^{-\mu t}), \quad t \in (-\bar{\rho}, 0]. \tag{3.11}$$

Next, we will prove that

$$\|v(t)\| \leq (\|\psi - w^*\|_{\zeta} + \epsilon) e^{-\mu t} < (\mathcal{N}(\|\psi - w^*\|_{\zeta} + \epsilon) e^{-\mu t}), \quad t > 0. \tag{3.12}$$

If (3.12) were not true, then there would exist $\xi > 0$ such that

$$\begin{cases} |v(\xi)| = \mathcal{N}(\|\psi - w^*\|_{\zeta} + \epsilon) e^{-\mu \xi}, \\ |v(t)| = \mathcal{N}(\|\psi - w^*\|_{\zeta} + \epsilon) e^{-\mu t}, \quad t \in (-\bar{\rho}, \xi). \end{cases} \tag{3.13}$$

In view of

$$\begin{aligned} & \dot{v}(s) + \varrho(s)v(s) \\ &= \zeta^{-1} \sum_{k=1}^p \theta_k(s) [e^{-\eta_k(s)\zeta w(s-\rho_k(s))} - e^{-\eta_k(s)\zeta w^*(s-\rho_k(s))}], \quad s \in [0, t], t \in [0, \xi], \end{aligned} \tag{3.14}$$

we have

$$\begin{aligned} v(s) &= v(0) e^{-\int_0^t \varrho(v) dv} + \int_0^t e^{-\int_s^t \varrho(v) dv} \\ &\quad \times \left\{ \zeta^{-1} \sum_{k=1}^p \theta_k(s) [e^{-\eta_k(s)\zeta w(s-\rho_k(s))} - e^{-\eta_k(s)\zeta w^*(s-\rho_k(s))}] \right\} ds, \quad t \in [0, \xi]. \end{aligned} \tag{3.15}$$

Hence

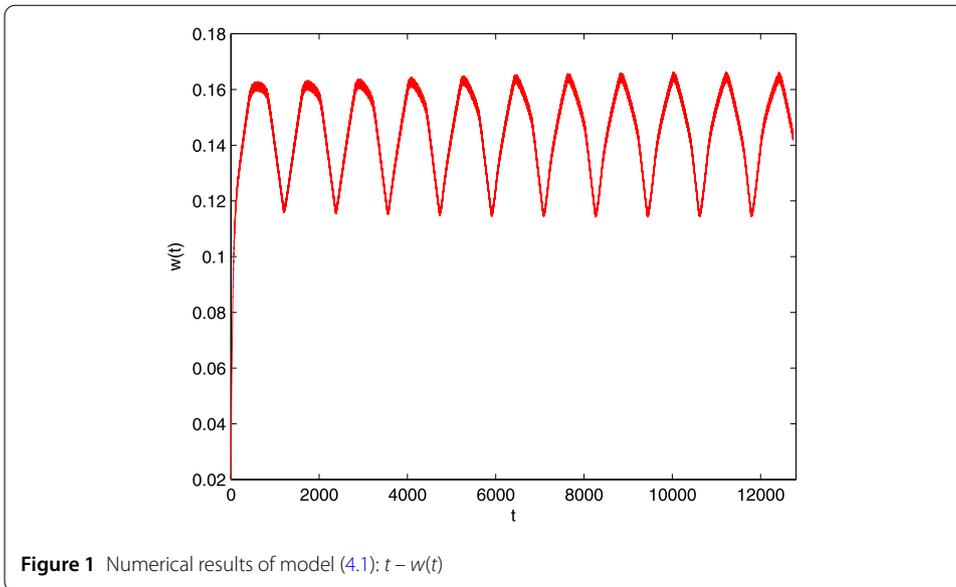
$$\begin{aligned}
 |v(s)| &= \left| v(0)e^{-\int_0^t \varrho(v) dv} + \int_0^t e^{-\int_s^t \varrho(v) dv} \right. \\
 &\quad \times \left. \left\{ \zeta^{-1} \sum_{k=1}^p \theta_k(s) \left[e^{-\eta_k(s)\zeta w(s-\rho_k(s))} - e^{-\eta_k(s)\zeta w^*(s-\rho_k(s))} \right] \right\} ds \right| \\
 &\leq \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) \alpha^M e^{-\int_0^\xi \bar{\varrho}(v) dv} + \int_0^\xi e^{-\int_s^\xi \bar{\varrho}(v) dv} \alpha^M \\
 &\quad \times \left[\zeta^{-1} \sum_{k=1}^p |\theta_k(s)| e^{-\eta_k \bar{\sigma}} |v(s - \rho_k(s))| \right] ds \\
 &\leq \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) \alpha^M e^{-\int_0^\xi \bar{\varrho}(v) dv} + \int_0^\xi e^{-\int_s^\xi \bar{\varrho}(v) dv} \alpha^M \\
 &\quad \times \left[\zeta^{-1} \sum_{k=1}^p |\theta_k(s)| e^{-\eta_k \bar{\sigma}} \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu(s-\rho_k(s))} \right] ds \\
 &\leq \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu\xi} \alpha^M e^{-\int_0^\xi (\bar{\varrho}(v)-\mu) dv} + \int_0^\xi e^{-\int_s^\xi (\bar{\varrho}(v)-\mu) dv} \alpha^M \\
 &\quad \times \left[\zeta^{-1} \sum_{k=1}^p |\theta_k(s)| e^{-\eta_k \bar{\sigma}} \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{\mu\rho_k(s)} \right] ds \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu\xi} \\
 &\leq \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu\xi} \alpha^M e^{-\int_0^\xi (\bar{\varrho}(v)-\mu) dv} \\
 &\quad + \int_0^\xi e^{-\int_s^\xi (\bar{\varrho}(v)-\mu) dv} (\bar{\varrho}(v) - \mu) ds \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu\xi} \\
 &\leq \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu\xi} \left[1 - \left(1 - \frac{\alpha^M}{\mathcal{N}} \right) e^{-\int_0^\xi (\bar{\varrho}(v)-\mu) dv} \right] \\
 &\leq \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu\xi}. \tag{3.16}
 \end{aligned}$$

By (3.13) we get that (3.12) holds. Let $\epsilon \rightarrow 0^+$. Then

$$\|v(t)\| < \mathcal{N}(\|\psi - w^*\|_\zeta + \epsilon) e^{-\mu t}, \quad t > 0. \tag{3.17}$$

Therefore we conclude that the pseudo almost periodic solution of model (1.1) is globally exponentially stable. The proof is finished. \square

Remark 3.1 Shao [1] studied the pseudo almost periodic solution of model (1.1). In this paper, we considered the weighted pseudo almost periodic solution. All the assumptions in this paper are different from those in [1]. All the derived results in [1] cannot be applicable to system (1.1) to establish the existence and globally exponential stability of weighted pseudo almost periodic solution for model (1.1). Based on this viewpoint, our results on the existence and globally exponential stability of weighted pseudo almost periodic solution for model (1.1) are essentially new and complement earlier works to a certain extent.



4 Computer simulations

Consider the system

$$\dot{w}(t) = -\varrho(t)w(t) + \sum_{k=1}^2 \theta_k(t)e^{-\eta_k(t)w(t-\rho_k(t))}, \tag{4.1}$$

where $\varrho(t) = 17 + 10 \cos 1000t$, $\theta_1(t) = 2(1 + 0.5|\sin 1.7t|)$, $\theta_2(t) = 2(0.8 + 0.2|\sin 1.7t|)$, $\eta_1(t) = -1 + 0.2|\sin 5t|$, $\eta_2(t) = -0.9 + 0.5|\sin 30t|$, $\rho_1(t) = e^{0.2|\sin t|}$, $\rho_2(t) = e^{0.3|\sin t|}$. Let $\bar{\varrho}(t) = 17$, $M[\varrho] = 17$, $\alpha^M = \frac{1}{200}$, $\alpha^m = -\frac{1}{200}$, $\chi = 0.8$, $\sigma = 0.2$, $\zeta = 2$, $\epsilon = 3$. Then

$$\alpha^m e^{-\int_s^t \bar{\varrho}(v) dv} \leq e^{-\int_s^t \varrho(v) dv} \leq \alpha^M e^{-\int_s^t \bar{\varrho}(v) dv} \quad \text{for } t, s \in \mathbb{R} \text{ such that } t - s \geq 0$$

and

$$\sup_{t \in \mathbb{R}} \left\{ -\bar{\varrho}(t) + \zeta^{-1} \alpha^M \sum_{k=1}^p |\theta_k(s)| |\eta_k(s)| e^{-\eta_k^- \sigma} \right\} \approx -3.087 < -3 = -\epsilon.$$

Thus all the hypotheses of Theorem 3.1 are satisfied, and so system (4.1) has a unique weighted pseudo almost periodic solution, which is globally exponentially stable. Figure 1 reveals this fact.

5 Conclusions

In this paper, we have discussed the existence and globally exponential stability of weighted pseudo almost periodic solutions for a Lasota–Ważewska system. Using the fixed point theory and differential inequalities, we establish new sufficient criteria ensuring the existence and globally exponential stability of weighted pseudo almost periodic solutions for the Lasota–Ważewska model. The derived results complement some earlier publications to some extent. Up to now, to the best of our knowledge, it is the first time to deal with this aspect. In the near future, we will investigate the pseudo almost automorphic solutions for the Lasota–Ważewska model.

Acknowledgements

The authors would like to thank the referees and the editor for helpful suggestions incorporated into this paper.

Funding

The work is supported by National Natural Science Foundation of China (No. 61673008), Project of High-level Innovative Talents of Guizhou Province ([2016]5651), Major Research Project of The Innovation Group of The Education Department of Guizhou Province ([2017]039), Innovative Exploration Project of Guizhou University of Finance and Economics ([2017]5736-015), Project of Key Laboratory of Guizhou Province with Financial and Physical Features ([2017]004), Hunan Provincial Key Laboratory of Mathematical Modeling and Analysis in Engineering, Changsha University of Science & Technology (2018MMAEZD21), University Science and Technology Top Talents Project of Guizhou Province (KY[2018]047), Guizhou University of Finance and Economics(2018XZD01), and Foundation of Science and Technology of Guizhou Province ([2019]1051).

Availability of data and materials

Data sharing not applicable to this paper as no datasets were generated or analyzed during the current study.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

All authors have read and approved the final manuscript.

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Received: 10 March 2020 Accepted: 17 April 2020 Published online: 18 May 2020

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