

Use of the Lorentz-Boosted Frame Transformation to Simulate Free-Electron Laser Amplifier Physics

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Abstract.

Recently [1] it has been pointed out that numerical simulation of some systems containing charged particles with highly relativistic directed motion can be speeded up by orders of magnitude by choice of the proper Lorentz boosted frame. A particularly good example is that of short wavelength free-electron lasers (FELs) in which a high energy ($E_0 \geq 250$ MeV) electron beam interacts with a static magnetic undulator. In the optimal boost frame with Lorentz factor γ_F , the red-shifted FEL radiation and blue shifted undulator have identical wavelengths and the number of required time-steps (presuming the Courant condition applies) decreases by a factor of γ_F^2 for fully electromagnetic simulation.

We have adapted the WARP code [2] to apply this method to several FEL problems including coherent spontaneous emission (CSE) from pre-bunched e-beams, and strong exponential gain in a single pass amplifier configuration. We discuss our results and compare with those from the "standard" FEL simulation approach which adopts the eikonal approximation for propagation of the radiation field.

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INTRODUCTION

It is well known that in general, explicit, fully electromagnetic simulation will have its time step Δt limited by the Courant condition corresponding to the numerical grid spacing and/or that necessary to achieve sufficient temporal resolution of the highest frequencies important to the physics of the particular situation. For problems in which a highly relativistic charged particle beam is present, the overall system time and/or length scale L_{sim} can be large and the ratio of scale lengths $L_{sim}/c\Delta t$ can become enormous. Recently, Vay [1] pointed out that for some of these problems performing the simulation in a Lorentz-boosted frame offers potentially orders of magnitude speed-up in computation time.

A natural candidate for boosted frame calculations is a short wavelength free-electron laser (FEL). Here a sample problem could have the resonant radiation wavelength $\lambda_R = 10$ nm, an undulator wavelength $\lambda_u = 25$ mm, and a system length $L_{sim} \geq 10$ m. Performing this simulation in the laboratory frame requires $\sim 4 \times 10^9$ time steps or greater. Presuming a "moving window" type simulation centered about the electron beam of length $l_b \approx 100 \mu\text{m}$ and radius $r_b \approx 50 \mu\text{m}$, the number of grid points for an 2D axisymmetric (or slab) model exceeds 10^8 (and likely 3-10 times greater to properly model transverse diffraction effects).

Even today with massively parallel systems, doing just one fully electromagnetic simulation in the lab frame with these parameters would be an enormous undertaking, and not one that scales easily to multi-parameter investigations. Instead, realizing for most problems that the normalized gain bandwidth is small and that the radiation gain length $L_G \gg \lambda_R$, for a quarter century or more FEL simulationists have applied the eikonal approximation (also known as the slowly-varying envelope approximation). This changes the E&M field equations from hyperbolic to parabolic (*e.g.*, diffusive) in nature, thereby permitting Δz and Δt to approach $\sim L_G/8$ and $\sim 0.1\lambda_R^2/c\Delta\lambda_{BW}$, respectively. Numerically, this leads to the grid spacings Δz and Δr being of order centimeters and microns, respectively, rather than nanometers, resulting in the total number of grid points being 10^6 or less and time steps being 10^3 or less with an overall speedup of 10^8 or greater over full electromagnetic simulation.

The natural boosted frame for FEL computations is the so-called “ponderomotive” frame in which the e-beam longitudinal speed (when in the undulator) is zero on average. In this frame the red-shifted FEL resonant wavelength $\lambda'_R = 2\gamma_F\lambda_R$ is equal to the blue-shifted undulator wavelength $\lambda'_u = \lambda_u/\gamma_F$. Here $\gamma_F^2 \equiv \gamma_0^2/(1+a_w^2)$ with a_w being the normalized, RMS undulator strength. The Lorentz transformation to the boosted frame shrinks the undulator by a factor γ_F and increases the radiation wavelength by the same factor times two, resulting in an overall decrease of the needed number of time steps by a factor $\approx 2\gamma_F^2$. Likewise, from the point of view of the Courant condition, the increase in Δt permits (in general) a similar increase in the spatial grid zone size so that the savings in 2- and 3-D simulations can be immense. However, in cases where the electron beam is much longer than the so-called slippage length ($l_{slip} \equiv N_u\lambda_R$) in the lab frame, this length exceeds the undulator length in the boosted frame and one factor of γ_F is lost (*i.e.*, the ratio of l_b/λ_R remains constant independently of γ_F).

To study various standard FEL problems, we used the WARP simulation code [2] with its standard full E&M solver in slab-mode geometry (*e.g.*, $x-z$ or $y-z$). A special Python script implemented the linearly-polarized undulator fields in the boosted frame. For amplifier runs, a second script that initialized the input radiation at the upstream longitudinal boundary (in the boosted frame) also was applied. Special diagnostics were added which measured the forward-moving radiation intensity through a transverse plane fixed in the lab frame (*e.g.*, at a fixed z relative to the undulator entrance). In order to avoid requiring initialization of the E - and B -fields associated with a beam pulse with a net current and charge, we “added” a positron beam with the exact same charge and current distribution at $t = 0$ in the simulation; this doubles the effective current for radiation production. When these beams enter the transverse fields of the undulator, they are separated in the wiggle plane, creating an attractive (but unphysical) electric field. However, for optical and shorter wavelength FEL’s operating in the Compton regime (*e.g.*, $I_B \leq 10$ kA) we expect these additional forces to be negligible in terms of radiation production (as was spot-checked by removing the faux positron beam).

We now present example results for some short wavelength test cases of stimulated emission from pre-bunched beams and a high-gain amplifier. Due to space limitations, a more detailed exposition will be given elsewhere.

SIMULATIONS OF COHERENT STIMULATED EMISSION (CSE) FROM A PRE-BUNCHED BEAM

Emission by a Single Ultrashort Bunch

A number of simulations were done examining the emission from an ultrashort bunch ($l_b \ll \lambda_R$) traveling through a moderate length undulator. Theoretically, one expects the total emission to scale as $Q_b^2 N_u a_w^2 / (1 + a_w^2)$. This scaling was confirmed in the WARP runs in the boosted frame, as was the relative growth of harmonic emission as a_w approached 1. As expected, even harmonic emission peaked off-axis; however, it was not exactly zero on-axis because of the presence of 2nd and other even harmonic components in the undulator field associated with the ramp in strength at undulator entrance and exit. As N_u increased, this on-axis, even harmonic emission became proportionately less and less relative to the fundamental and odd harmonics.

One advantage of a full E&M code over an eikonal code is that it can study situations where there is broadband emission and/or conditions where the slowly-varying envelope approximation are no longer strongly valid. Such a case occurs in the entrance and exit ramp regions of the undulator where the local B_{peak} (and thus the local value of λ_R) changes significantly over a few undulator periods or less. We studied the role of the ramp length for the flux from a single, low charge micropulse with optical phase extent $\pi/4$ propagating in an undulator of total length 0.5 m. With a functional matching ramp dependence of $B_{\perp} = B_0 (z/L_M) \sin(2\pi z/\lambda_u)$ for $0 \leq z \leq L_M$, we did a series of simulations varying L_M from 2 to 8 λ_u with $\lambda_u = 25$ mm. The other FEL parameters were $\lambda_R = 200$ nm, $a_w = 0.2$, $\gamma_0 = 255$, and $Q_b = 1.33$ fC. As shown in the left plot of Fig. 1, diagnostics of integrated flux reveal a linear increase with z for $z \geq L_M$. Extrapolating the linear portion back toward $z = 0$ indicates a y-intercept equal to nearly exactly $L_M/2$. The significance of this result is that for eikonal FEL simulation codes, the adiabatic matching sections at the entrance and exit to undulator segments should to lowest order be treated as being composed of a drift section and a full strength undulator section, each of length $L_M/2$.

Emission by a Periodic Series of Identical Ultrashort Pulses

In order to compare the predicted results from full E&M simulation in the boosted frame with that of a standard eikonal code, we performed a simulation of a very low current ($\langle I_b \rangle = 1$ A), high energy ($\gamma_0 = 354$) long pulse ($l_b = 10$ mm $\gg l_{slip} = 2.4$ mm) composed of a periodic series of ultrashort pulses, each with duration of one-eighth the resonant wavelength of 200 nm. The energy spread and normalized emittance were sufficiently low (values of 0 and 0.1 mm-mrad respectively) that essentially no debunching occurred over the duration of the undulator ($L_u = 0.3$ m, $\lambda_u = 25$ mm, $a_w = 1.0$). The right graph of Fig. 1 shows the radiation flux through a transverse plane 50 mm downstream of the undulator; the rapid oscillations correspond to half the radiation period as expected. The average power of 240 W measured over the interior portion of the radiation pulse (where slippage effects would disappear) is very close to the 280 W value for

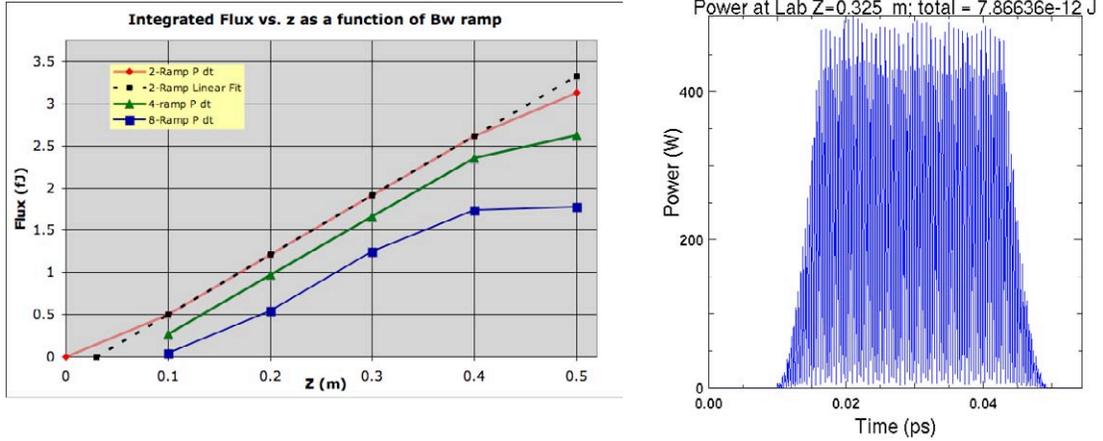


FIGURE 1. (a) WARP results for integrated 200-nm radiation flux (in the lab frame) through a transverse plane vs. z for various lengths matching ramps in $B_{\perp}(z)$. From top to bottom, the ramp lengths L_M were 5, 10, and 20 mm respectively with $\lambda_u = 25$ mm. (b) Lab-frame $P(t)$ at a plane 2.5 cm downstream of the undulator exit for a 200-nm FEL with 50 ultrashort microbunches. The time-averaged power over the central pulse is 240 W. A time-steady, slab-mode GINGER run with identical e-beam parameters predicts 280 W. Please see text for more details concerning FEL parameters used to produce both figures.

a time-steady, pre-bunched pulse predicted by the eikonal GINGER code [3] (run in slab mode geometry with the same e-beam and undulator parameters taking into account the matching ramp length correction described above).

Emission by a Long Pulse with Sub-harmonic Bunching

There is great current interest among many FEL groups in exploiting modulator-radiator configurations where an e-beam is first strongly bunched by an external laser at a long wavelength λ_M in the modulator and then radiates at a much shorter wavelength in the radiator, whose period and magnetic strength are tuned such that its resonant wavelength λ_R is a harmonic of λ_M . We studied such a configuration in the boosted frame by examining emission in a “radiator” undulator from an e-beam strongly pre-bunched at $\lambda_M = 6\lambda_R$. Here again we used a very low current to minimize any complications due to gain. As expected from numerous studies using eikonal codes, on axis there is strong emission at λ_R . However, due to the full E&M code being able to study emission at wavelengths much longer and shorter than λ_R , we can also examine sideband emission whose separation (in wavenumber) from λ_R are harmonics of λ_M . We find that such non-resonant sideband emission does exist but, unlike radiation at the resonant wavelength for a prebunched beam, it does not grow quadratically with z .

HIGH GAIN SINGLE PASS AMPLIFIER SIMULATION

In addition to the pre-bunched, ultra-low current beam simulations discussed above, we also performed a simulation of a high gain FEL amplifier. The critical e- and undulator parameters were $I_B = 1$ kA (for each of the e^- and e^+ beams), $\gamma_0 = 353.6$, $\epsilon_N = 1$ mm-mrad, $\sigma_E = 0$, $\lambda_R = 200$ nm, $a_w = 1.0$, $\lambda_u = 25$ mm, and $N_u = 40$. The input seed laser had $P_0 = 50$ MW, and a downstream transverse waist size $\omega_0 = 250$ μ m. The effective e-beam size in the (non-resolved) x - dimension was 250 μ m. From standard FEL theory, one expects a power gain length of about 0.25 m.

The WARP simulation showed a total gain P_F/P_0 of approximately 5.0 in the central region of the pulse where slippage effects should not apply. A time-steady, slab-mode GINGER simulation with L_u reduced to 0.95 m to take into account the adiabatic match sections showed a net gain of 5.6 . On the whole, this is good quantitative agreement given the strong exponential gain for these parameters and the uncertainty in getting exactly equivalent conditions in the e-beam match between the two simulations.

CONCLUSIONS

In this paper we have discussed the application of the boosted-frame formalism of Vay [1] to short wavelength FEL simulations. The results indicate that performing electromagnetic simulation in the ponderomotive frame saves orders of magnitude in CPU time over doing the same problem in the lab frame without any obvious loss of significant physics. However, for the great bulk of FEL problems of interest at sub-IR wavelengths, despite the speed-up relative to lab frame E&M simulation, the boosted frame approach is still orders of magnitude slower than the eikonal method. Consequently, we expect the eikonal method codes to still be used for the great bulk of FEL simulation. Nonetheless, in some instances, the boosted frame approach allows examination of effects such as matching ramps and sub-harmonic bunching that are difficult or impossible to study with eikonal codes. Further improvements (*e.g.*, multi-undulator-period averaging) may further increase the efficiency of the boosted frame method.

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