

Probing the magnetic phases in the Ni-V alloy close to the disordered ferromagnetic quantum critical point with μ SR

A Schroeder¹, R Wang¹, P J Baker², F L Pratt², S J Blundell³,
T Lancaster^{3,4}, I Franke³ and J S Möller³

¹ Department of Physics, Kent State University, Kent OH 44242, USA

² ISIS Facility, STFC Rutherford Appleton Laboratory, Harwell Oxford, OX11 0QX, UK

³ Oxford University Department of Physics, Clarendon Laboratory, Parks Road, Oxford, OX1 3PU, UK

E-mail: aschroe2@kent.edu

Abstract. Zero (ZF) and longitudinal field (LF) muon spin relaxation data of the d -metal alloy $\text{Ni}_{1-x}\text{V}_x$ are presented at several vanadium concentrations x below and above the critical $x_c \approx 11\%$ where long-range ferromagnetic (FM) order is suppressed. The clear single precession frequency observed for Ni, as expected for a homogeneous FM, changes to rather damped oscillation with small V substitution at $x = 4\%$, confirming magnetic inhomogeneities caused by the less magnetic V environments in the magnetic Ni matrix. Furthermore, local fields and spatial field distributions can be estimated to characterize different inhomogeneous regimes developing with x in the FM phase of $\text{Ni}_{1-x}\text{V}_x$. In the regime of $x = 7 - 10\%$ a Kubo Toyabe function well describes the low temperature ZF and LF asymmetry data supporting a static Gaussian field distribution. Closer to the quantum critical concentration a single scale static Kubo Toyabe function with one field distribution is not sufficient to describe the muon relaxation. These data indicate that further changes in spatial distributions and dynamics are evolving as expected within the critical regime of a disordered quantum critical point.

1. Introduction

The study of quantum phase transitions (QPT) is a promising route to understand the origin of unconventional properties and in finding new phases in correlated many body systems. Introduction of “disorder”, through e.g. impurities or sample imperfection can significantly change the nature of the QPT. In itinerant system with Heisenberg symmetry a new quantum critical point is predicted with exotic properties such as observable quantum Griffiths effects [1]. Examples for disordered systems are not rare, but the degree and significance of the “disorder” is often not clear. The binary alloy $\text{Ni}_{1-x}\text{V}_x$ shows the obvious signs of such a disordered QPT from bulk investigations [2, 3] and is therefore an ideal example for further microscopic studies. μ SR is ideally suited as spectroscopic and local probe [4] to reveal the first insights into these disordered magnetic phases as it has been successfully employed for spin glasses [5], reduced moment correlated electron systems like heavy fermions [6, 7, 8] and phase separated systems close to first order transitions like transition metal compounds [9].

⁴ Present address: Durham University, Centre for Materials Physics, South Road, Durham DH1 3LE, UK



The ferromagnetic (FM) order of Ni is rapidly suppressed replacing Ni by V. The critical temperature T_c decreases approximately linearly and vanishes towards a V-concentration of $x_c \approx 11.4\%$ [2] as shown in Fig. 1(a). Magnetization measurements show the typical weak itinerant FM response in the FM phase for $x \leq 11\%$, the “non FM ordered”, paramagnetic (PM) phase beyond x_c is not compatible with independent magnetic moments. Instead, the magnetization displays power laws in temperature (T) and magnetic field with non universal exponents which vary with x in a wide regime $x \geq x_c$ [2, 3]. This marks the quantum Griffiths phase (GP) (see e.g. in [10]), characterized by a distribution of magnetic clusters with different fluctuation rates. This phase is distinctly different than in clean QPT, where universal critical exponents are ideally observed only at x_c . We suspect that $\text{Ni}_{1-x}\text{V}_x$ is a disordered magnet but direct evidence is lacking. Friedel [11] already proposed that the drastic average magnetic moment reduction (see e.g. m_s in Fig 2(a)) in Ni introducing V is not homogeneous, as V (with 5 d-electron less than Ni) might reduce severely the spin of all the Ni-neighbors creating large voids in the magnetization density. Early neutron measurements detected a large magnetic moment disturbance in Ni-V alloys [12]. Our bulk magnetization measurements only show (a mean field behavior of) the order parameter of a FM and no clear signs of the disorder. Only close to x_c towards very low temperatures deviations from a long range FM become apparent, cluster glass freezing is observed [2] which does not show on this linear T-scale and is not discussed here. Here we use μSR to display the “disorder” in $\text{Ni}_{1-x}\text{V}_x$, probing the local fields in the FM phase. We present how to characterize the inhomogeneities in the FM and how they evolve upon dilution approaching a disordered QCP.

2. Experimentals

Polycrystalline samples of the alloy $\text{Ni}_{1-x}\text{V}_x$ were prepared as spherical pellets in the concentration range $x = 0$ to $x = 12.3\%$ as before [2]. The fcc structure remains unchanged with x , the lattice constant increases by 1% for $x = 15\%$ [2]. The muon data were collected using the Dolly instrument at PSI, using several pellets of each composition wrapped in silver foil, at low temperatures, in zero (ZF) and longitudinal magnetic fields (LF).

3. Evolution of disorder in the ferromagnetic state in $\text{Ni}_{1-x}\text{V}_x$

Fig. 2 shows an overview of the muon asymmetry for several $\text{Ni}_{1-x}\text{V}_x$ samples with different V concentration x at low temperatures T in zero field. In pure Ni a nearly undisturbed muon precession is observed as expected for a clean FM with a single frequency of $f = 20.2$ MHz or $\omega = 127$ rad/ μs corresponding to a local field of $B_0 = \omega/\gamma = 0.15$ T as seen before [13] (using the known gyromagnetic ratio of the muon as $\gamma = 2\pi \times 135.5$ MHz/T). Upon doping with V these oscillations appear more “damped”, shift to larger time scales and finally disappear beyond the critical concentration in the PM phase. Already for $x \geq 4\%$ only the first minimum remains visible due to a distribution of fields revealing an inhomogeneous FM. Some subtle qualitative changes can already be seen in Fig. 2: Note that $x = 4\%$ still shows a second minor maximum or “nose”, and that the minimum for $x = 11\%$ is rather shallow.

In order to model the total asymmetry $A(t)$ for all temperatures we use different polarization functions $P(t)$ for the FM component, a simple exponential function for the PM component (approximating effectively electronic spins of clusters and the diluted V nuclear moments) and a small constant background A_{BG} due to the Ag.

$$A(t) = A_{FM}P(t) + A_{PM}\exp(-\lambda t) + A_{BG} \quad (1)$$

$$P_{osci}(t; \omega, \Gamma, \lambda_L) = \frac{1}{3}\exp(-\lambda_L t) + \frac{2}{3}\cos(\omega t + \phi)\exp(-\Gamma t). \quad (2)$$

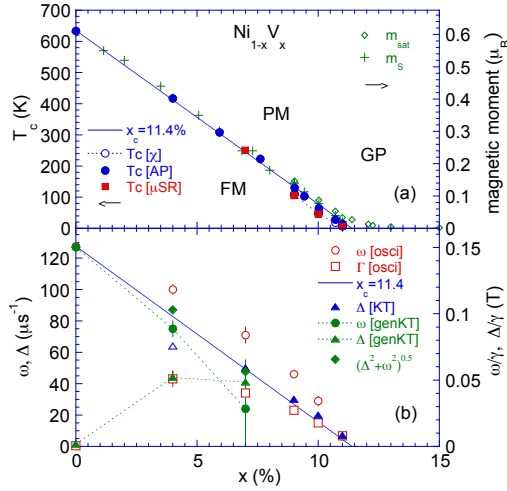


Figure 1. (a) Phase diagram for $\text{Ni}_{1-x}\text{V}_x$: critical temperature T_c vs. V-concentration x , separating ferromagnetic (FM), paramagnetic (PM) and quantum Griffiths phases (GP). T_c from magnetization (AP, χ) [2] and μSR . Magnetic moment determined at high (m_{sat}) and low fields (m_s). Straight line indicates $x_c = 11.4\%$. (b) Angular frequency or local field scales characterizing the center (ω) and distribution width (Δ) from various $P(t)$ models vs. x .

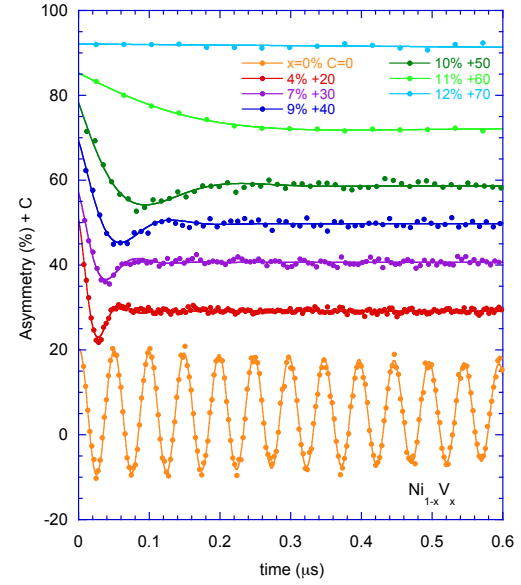


Figure 2. Asymmetry of $\text{Ni}_{1-x}\text{V}_x$ vs time for different V-concentrations x , ranging from $x = 0$ to $x = 12.3\%$ at low temperatures ($1.5\text{K} < T < T_c/4$). Each is shifted by a constant C for clarity. The line represents a fit to the oscillator model with P_{osci} .

First, we apply an oscillator model P_{osci} (see lines in Fig. 2), a simplified “damped” expansion of a single precession frequency as expected in a perfect clean FM polycrystal like Ni with one preferred stopping site. This form is often used and also provides here a parametrization for all $x \leq 11\%$. For low T , A_{PM} and λ are essentially zero, later we need to set $\lambda_L = \lambda$ for all T . Also we keep $\phi = 0$ for higher x to produce consistent parameters. The main frequency ω decreases with x as shown in Fig. 1(b). The damping constant Γ or transverse field depolarization rate is increased for $x = 4\%$ then decreases further with x . Γ can be seen as a measure of field distribution, as ω is indicating some local field, but comparison to other models like P_{genKT} (see eq.(4) below) show that in the “damped” cases ω is not exactly the mean field nor the field distribution. Γ has half the value of ω at lowest T for all samples with $7\% < x < 10\%$, which marks the regime with a Gaussian field distribution as discussed below.

More revealing simple models can be applied, but are restricted to certain x -regimes. In effect, for $x = 7 - 10\%$ a static Kubo-Toyabe (KT) function P_{KT} describes the data well for low T , indicating a Gaussian distribution of local fields [14] in the FM state.

$$P_{KT}(t; \Delta, \nu = 0) = \frac{1}{3} + \frac{2}{3}(1 - (\Delta t)^2) \exp(-\frac{1}{2}(\Delta t)^2), \quad (3)$$

The single parameter, the width Δ of the distribution decreases with x towards x_c as shown in Fig. 1(b). This field distribution follows the magnetic moment m_s on the same straight line in Fig.1(a,b). It is remarkable that inside a FM a “random” magnetic field distribution is found as originally shown for randomly oriented frozen moments [14]. We suspect here that the distribution of different local fields indicates a spatial variation of magnetic moments in

a long range ordered FM state. These muon data provide evidence that this system is indeed inhomogeneous or “disordered” as required to show this disordered QCP with distinct properties.

A Kubo Toyabe fit P_{KT} is possible for $x = 4\%$ but can be improved by a generalized form P_{genKT} [15] allowing a finite center field $H_0 = \omega/\gamma \neq 0$ from large domains or crystallites which is the center of a Gaussian distribution with width $\Delta H = \Delta/\gamma$.

$$P_{genKT}(t; \Delta, \omega, \nu = 0) = \frac{1}{3} + \frac{2}{3} \exp(-\frac{1}{2}(\Delta t)^2) [\cos(\omega t) - \frac{1}{\omega}(\Delta t)^2 \sin(\omega t)]. \quad (4)$$

The two parameters ω and Δ are shown in Fig. 1(b). Δ is similar for both $x = 4\%$ and $x = 7\%$, overall the maximum field distribution appears at about $\frac{1}{2}x_c$. The characteristic field H_0 in $x = 4\%$ is still high due to intact magnetic domains, while H_0 is rather close to zero in $x = 7\%$, the field distribution ΔH dominates. The transition from dominating FM domains to random distribution occurs in between these concentrations at about $\frac{1}{2}x_c$. That matches the concentration $x = 5.6\%$ where the number of reduced moment Ni-atoms (due to one or more V neighbors) exceeds the number of undisturbed “magnetic” Ni-atoms (with only Ni neighbors) in a fcc lattice. The local muon probe distinguishes the different degrees of inhomogeneities in this FM introduced by V substitution. The model P_{genKT} provides even a direct quantitative assessment of an internal field distribution width Δ/γ and a center local field ω/γ . An effective field $\sqrt{\Delta^2 + \omega^2}/\gamma$ finally scales with the mean magnetic moment $m_s(x)$.

The KT function P_{KT} does not work as well towards the more diluted regime approaching x_c , the minimum becomes very shallow, for e.g. 11% , questioning the long range order of this compound. Beyond x_c , for $x = 12.3\%$, a small depolarization rate ($\lambda = 0.07\mu s^{-1}$ for 2K) describes the data in ZF. Extended time, more temperature and magnetic field data and analysis are required to reveal more details about the fluctuation rates beyond the FM regime.

4. Contrasting the field distributions in $Ni_{1-x}V_x$ with $x = 9\%$ and $x = 11\%$

The asymmetry spectra for $x = 9\%$ can be modeled well by a KT function. For higher temperature we used a dynamic KT approximation with finite fluctuation rates ν [14, 16] for the FM contribution (as presented in Fig. 3(a)). Application of a longitudinal field (LF) at low T confirms the main parameter choices of a KT description (see Fig.3(b)). The straight lines shown are a KT fit with unchanged parameters from ZF with internal field = external LF, better fit results are achieved with reduced internal fields in these FM spherical samples. Fig.4 (a-d) presents the T-dependence of the fit-results of the exploratory oscillator model and the KT model. Panel (a) shows the main field distribution Δ together with ω and Γ . $\Delta(T)$ follows nearly a mean field T-dependence like the magnetic moment $m_s(T)$ vanishing at T_c . The fixed ratio ω/Γ of 2 (observed for several x at low T) changes towards higher T towards 1, before ω vanishes. The ferromagnetic amplitude ratio, the ferromagnetic contribution divided through the total contribution, $A_{FM}/(A_{FM} + A_{PM})$, increases rapidly below T_c and then saturates to about 1 independent of the model (see Fig.4(b)). The depolarization rate λ is quite small in the PM regime and also very small, close to the detection limit, in the FM for $T < 0.8T_c$ (as λ_L) in P_{osci} . This rate peaks towards T_c , indicating some critical dynamics. The dynamic KT model reveals directly critical fluctuations increasing rapidly at T_c (see Fig.4(d)). We see that below about $0.8T_c$ magnetic fluctuations become irrelevant. All quantities as A_{AF} , Δ , ν indicate a critical temperature of $T_c = 108$ K consistent with low field susceptibility χ measurements [3], a bit lower than the usual high field Arrott plot (AP) determination of $T_c = 130$ K.

Magnetization measurements indicate that $x = 11\%$ is still a ferromagnet displaying Arrott plots and no indication of a cluster glass transition is found in magnetic susceptibility. A strong ZF depolarization is observed, but the minimum is too shallow for the static KT form even at lowest T (see Fig. 3(d)). The ZF data can be modeled quite well applying a *dynamic* KT function ($P(t, \Delta, \nu > \Delta)$), using the strongly fluctuating approximation of a Keren function

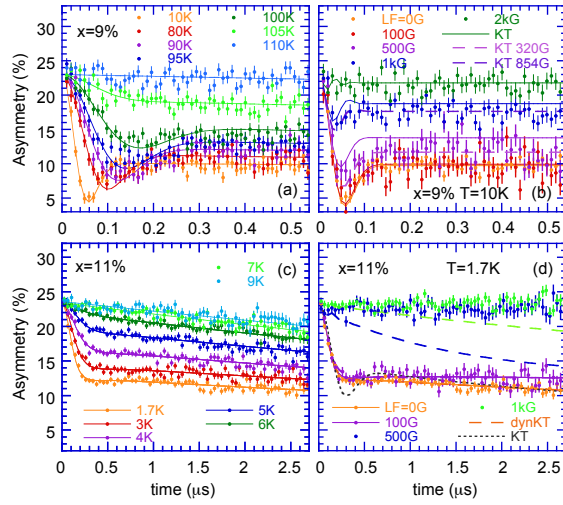


Figure 3. (a) ZF spectra for $\text{Ni}_{1-x}\text{V}_x$ with $x = 9\%$. Lines from P_{KT} . (b) LF spectra for $x = 9\%$ at $T = 10\text{ K}$. Lines from P_{KT} : field = external LF (solid), adjusted internal field (dashed). (c) ZF spectra for $x = 11\%$. Lines from P_{GBG} . (d) LF spectra for $x = 11\%$. Dotted line: static P_{KT} . Dashed lines: dynamic P_{KT} . Solid lines: P_{GBG} .

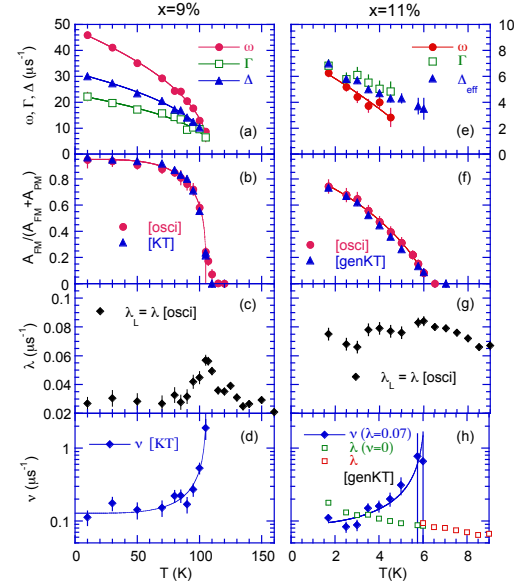


Figure 4. T-dependence of various fit parameters using P_{osci} , dynamic P_{KT} or P_{GBG} to model $A(t)$ as shown in Fig. 3 for $x = 9\%$ (a-d) and for $x = 11\%$ (e-h). Lines are guide to the eye.

[17, 16]. But the present LF data (e.g. in 500G) cannot be reproduced keeping the same fluctuation rate as in ZF (see dashed line in Fig. 3(d)). More intermediate fields are needed to study field dependent changes. We alternatively applied a *static* distribution function which describes multiple field distributions, narrow and wide ones to increase the weight towards very small field regions. Taking the response of many Gaussian distribution of fields with different Δ 's following a Gaussian distribution with width Δ_w and center Δ_0 gives the so called simple Gaussian broadened Gaussian function P_{GBG} [18, 19].

$$P_{GBG}(t; \Delta_{eff}, R, \nu) = \frac{1}{3} \exp(-\frac{2}{3}\nu t) + \frac{2}{3}((1 + R^2)/N)^{3/2} (1 - (\Delta t)^2/N) \exp(-\frac{1}{2N}(\Delta t)^2), \quad (5)$$

with $N = 1 + R^2 + R^2 \Delta^2 t^2$, where $R = \Delta_w/\Delta_0$ and $\Delta^2 = \Delta_{eff}^2 = \Delta_0^2 + \Delta_w^2$. This form describes the data well as shown in Fig. 3(a) (with $R=0.7$). Note that an additional PM contribution with a small depolarization rate is always present even for low T. The effective field distribution width Δ_{eff} is shown in Fig. 4(e) together with the P_{osci} parameters. They decrease towards higher T, a sharp critical temperature T_c cannot be determined. Also, the ratio of these parameters at $x = 11\%$ is different than those at $x = 9\%$ at low T, more similar to those at $x = 9\%$ close to T_c . The FM amplitude ratio increases only gradually towards low T (Fig. 4(f)), not as sharp as in $x = 9\%$. At the lowest T, only 75 % of the amplitude can be described by the FM component $P(t)$, while an extra (rest of 25 %) exponential “paramagnetic” contribution remains, independent of the model. The depolarization rate λ ($=\lambda_L$) shows barely a maximum at T_c and does not decrease towards low values as for $x = 9\%$ in the oscillator model (Fig.4(g)). Keeping P_{GBG} static ($\nu = 0$) the depolarization rate λ of the PM contribution increases further below T_c (Fig.4(h)). A small LF=100G suppresses these fluctuations (see Fig.3(d)), while the steep FM depolarization seems unaffected. Alternatively, allowing small fluctuations ($\nu > 0$) and keeping $\lambda = 0.07 \mu\text{s}^{-1}$ constant produces some critical dynamic towards T_c . Estimates of

$T_c \approx 6$ K from these muon data is in agreement with low field susceptibility χ measurements [3], much lower than Arrott plot (AP) extrapolations of $T_c = 14$ K [2].

While the inhomogeneous magnetism in $x = 9$ % could be well characterized by one scale, the KT distribution width Δ , multiple field distributions seemed better in $x = 11$ %. Furthermore, a PM contribution due to fluctuating clusters never vanishes at low T. These models are far too simple but indicate that fluctuating moments and the spatial inhomogeneities further evolve towards x_c . 11 % is not just a scaled down version of 9 % with a smaller field range.

5. Conclusion

$\text{Ni}_{1-x}\text{V}_x$ offers the opportunity to study a FM with different degrees of “spatial disorder” or inhomogeneities, introduced through V substitution from a pristine FM to a disordered QCP. Analyzing the muon asymmetry with a generalized KT model, P_{genKT} , the mean local field and the local field distribution could be determined for a large x regime. These parameters measure the degree of “disorder” and allow a distinction between different regimes within the FM state. For lower x intact domains with finite fields dominate, while for larger x the field distribution dominates. This distribution width is the dominating energy scale in a wide region $7\% \leq x \leq 10\%$, it follows the order parameter, m_s . Approaching the critical concentration, x_c , simple single scale fit functions cease to work. Multiple distributions with multiple time scales evolve at the edge of the Griffiths phase, dominated by clusters with different time scales. So far we only used simple models to probe local field distributions in the FM phase. To relate this information to the spatial variation of inhomogeneities and correlation length a microscopic model is needed. The simple fcc lattice of Ni-V looks promising as a good model system. The study of the complex scenario in the quantum critical regime and towards the quantum Griffiths phase is a further challenge in this d -metal alloy.

Acknowledgments

We thank S. Ubaid-Kassis for the original samples and W. Hayes for helpful discussions. This work was supported in part by OBR Research Challenge from KSU and part of this work was carried out at the Swiss Muon Source, Paul Scherrer Institute, CH.

References

- [1] Vojta T and Schmalian J 2005 *Phys. Rev. B* **72** 045438
- [2] Ubaid-Kassis S, Vojta T and Schroeder A 2010 *Phys. Rev. Lett.* **104** 066402 and refs therein
- [3] Schroeder A, Ubaid-Kassis S, and Vojta T 2011 *J. Phys.: Condens. Matter* **23** 094205
- [4] Yaouanc A and Dalman de Réotier P 2011 *Muon Spin Rotation, Relaxation, and Resonance* (Oxford: Oxford University Press) and refs therein
- [5] Uemura Y J, Yamazaki T, Harshman D R, Senba M and Ansaldo E J 1985 *Phys. Rev. B* **31** 546–563
- [6] Amato A 1997 *Rev. Mod. Phys.* **69** 1119–1180
- [7] MacLaughlin D E, Heffner R H, Bernal O O, Ishida K, Sonier J E, Nieuwenhuys G J, Maple M B and Stewart G R 2004 *Journal of Physics: Condensed Matter* **16** S4479
- [8] Adroja D T, Hillier A D, Park J G, Kockelmann W, McEwen K A, Rainford B D, Jang K H, Geibel C and Takabatake T 2008 *Phys. Rev. B* **78** 014412
- [9] Uemura Y J, Goko T, Gat-Malureanu I M et al 2007 *Nature Phys* **3** 29 – 35
- [10] Vojta T 2006 *J. Phys. A.: Math. Gen.* **39** R143–R205
- [11] Friedel J 1958 *Nuovo Cimento* **7**, 287–311
- [12] Collins M F and Low G G 1965 *Proceedings of the Physical Society* **86** 535
- [13] Foy M L G, Heiman N, Kossler W J and Stronach C E 1973 *Phys. Rev. Lett.* **30** 1064–1067
- [14] Hayano R S, Uemura Y J, Imazato J, Nishida N, Yamazaki T and Kubo R 1979 *Phys. Rev. B* **20** 850–859
- [15] Kornilov E I, Pomjakushin V Yu 1991 *Physics Letters A* **153** 364–367
- [16] Pratt F L 2000 *Physica B* **289-290** 710–714
- [17] Keren A 1994 *Phys. Rev. B* **50** 10039–10042
- [18] Noakes D R and Kalvius G M 1997 *Phys. Rev. B* **56** 2352–2355
- [19] Noakes D R 1999 *Journal of Physics: Condensed Matter* **11** 1589