

# EFFECT OF NOZZLE LENGTH ON GAS ENTRAINMENT CHARACTERISTICS OF VERTICAL LIQUID JET

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**Key Words:** Vertical Liquid Jet, Nozzle Length, Bubble Penetration Depth, Gas Entrainment Rate, Jet Surface Change

For a vertical plunging liquid jet system using nozzles which had various values of the length-to-diameter ratio  $L_n/D_n$ , characteristics such as the bubble penetration depth  $Z$  and the gas entrainment rate  $Q_g$  were studied experimentally in an air-water system. When nozzles of an  $L_n/D_n$  ratio exceeding 15 were employed, the values of  $Z$  and  $Q_g$  were confirmed to be almost independent of  $L_n/D_n$ . Empirical correlations for predicting  $Z$  and  $Q_g$  are also presented.

## Introduction

An aeration system based on gas entrainment by a liquid jet is attractive compared to conventional gas-sparging systems because of its simplicity in construction and operation.<sup>2)</sup> The earlier researchers adopted four operating variables which affect the gas entrainment characteristics of a liquid jet. They are: nozzle diameter, nozzle velocity, nozzle height, and nozzle angle. However, it is insufficient to investigate or discuss the characteristics of the liquid jet by these four factors only. An important factor that cannot be overlooked is differences in the nozzle design, especially the nozzle length-to-diameter ratio  $L_n/D_n$ .<sup>4,5,13,21)</sup> Reviewing the existing studies on inclined or vertical plunging water jets, most of these works were carried out in long nozzles which had a value of  $L_n/D_n$  exceeding 50.<sup>8,9,11,17-21)</sup> In such systems, useful information is available on the gas entrainment characteristics. In contrast, there are few studies in which the characteristics of the liquid jet were investigated over a wide range of  $L_n/D_n$ . Considering practical application of the jet aeration system, it is assumed that there are cases where a long nozzle with a large  $L_n/D_n$  ratio cannot be used due to the structural limitation of the reactors or difficulties in the piping and the liquid feeding, etc. For example, consider the case of using a nozzle of 0.02 m diameter: the nozzle must be longer than 1.0 m to satisfy the condition of  $L_n/D_n$  greater than 50. Long nozzles of a large  $L_n/D_n$  seems to be of no practical use. That suggests the necessity of a study to clarify the range of effective  $L_n/D_n$  ratio of the nozzle to be used. The bubble penetration depth  $Z$  and the gas entrainment

rate  $Q_g$  in the vertical jet system using short nozzles of  $L_n/D_n=5$  were reported in the previous work.<sup>15)</sup> In the present study, the relationship between  $Z$  or  $Q_g$  and the operating conditions in vertical jet systems was investigated experimentally by varying widely the  $L_n/D_n$  ratio. Also, to obtain a clearer view of the suggestions made by some authors that when the liquid jet is discharged from the nozzle into a gas phase a change of the jet occurs<sup>13)</sup> and that the jet surface change affects the gas entrainment,<sup>4,7,12,14)</sup> the change of the jet with variation of operating conditions was observed, and some studies were made of the relation between these changes and the gas entrainment behavior.

## 1. Experimental

The experimental apparatus was almost the same as that employed in the previous work.<sup>15)</sup> The apparatus consisted of a liquid pump, a straight cylindrical nozzle (PVC) with inside diameter  $D_n$  between 0.007 and 0.017 m, and a bath of height 1.0 m and section 1.0 m by 0.65 m. The jet velocity  $V_n$  was varied from 2.0 to 13.5 m/s, depending on  $D_n$ . The range of  $L_n/D_n$  ratio variation was between 5 and 70. The nozzle height  $H_n$  from the bath surface to the nozzle exit ranged from 0.025 to 0.75 m. The value of  $Z$  was measured by a scale fitted to the bath wall. In measuring the gas entrainment rate  $Q_g$ , the same volume of liquid as that of the entrained gas into the length-variable cylinder having a cap was introduced into a replacement vessel. By measuring the volume of liquid replaced by the entrained gas, the volumetric rate  $Q_g$  of entrained gas was determined. Additional details of the apparatus and procedures were shown in the previous paper.<sup>15)</sup> The change of the jet surface was observed by taking photographs. All the

Received October 29, 1986. Correspondence concerning this article should be addressed to A. Ohkawa.

experiments were carried out at 293 K with an air–water system.

## 2. Results and Discussion

### 2.1 Bubble penetration depth $Z$

In our previous work,<sup>15)</sup> the choice of non-dimensional parameters to be employed in plotting the data of  $Z$  was made by referring to a theoretical analysis of the momentum equation for the axisymmetric submerged jet<sup>1)</sup> analogous to the vertical plunging jet system. And the relationship between  $Z$  and the operating conditions when nozzles having a constant  $L_n/D_n$  were employed was derived as an expression of functional form:

$$Z/D_n = \text{func.}(F_r, H_n/D_n) \quad (1)$$

The effect of  $L_n/D_n$  ratio on the change of the dimensionless bubble penetration depth  $Z/D_n$  was further investigated, along with the effects of the Froude number  $F_r$  ( $= V_n/\sqrt{gD_n}$ ) and the ratio of nozzle height to diameter,  $H_n/D_n$ . **Figure 1** shows the typical relationship between  $Z/D_n$  and  $L_n/D_n$ . As can be seen from these results,  $Z/D_n$  decreases with  $L_n/D_n$ , if the  $L_n/D_n$  ratio is smaller than 15, but is almost constant above this range. **Figure 2** shows the effect of  $H_n/D_n$  on  $Z/D_n$ . The effect of  $H_n/D_n$  was the same as that observed for the nozzles of  $L_n/D_n = 5$ .<sup>15)</sup> That is, regardless of  $L_n/D_n$ ,  $Z/D_n$  at  $H_n/D_n$  below 20 tended to decrease with  $H_n/D_n$  and be almost constant above that range. Moreover, in experiments using short nozzles of  $D_n$  from 0.001 to 0.004 m but of unknown values of  $L_n/D_n$ , Suci and Smigelschi<sup>16)</sup> observed that  $Z/D_n$  at  $H_n/D_n$  above 20 becomes almost independent of  $H_n/D_n$ , similarly to our results.

### 2.2 Correlation of $Z$

The bubble penetration depth  $Z$  when the short nozzles of  $L_n/D_n = 5$  were employed was shown to be represented by the following form.<sup>15)</sup>

$$Z/D_n = (H_n/D_n)^A 10^{f(F_r)} \quad (2)$$

where  $A$  is constant and  $f(F_r)$  is a function of  $F_r$ . A correlation of  $Z/D_n$  in which the effect of  $L_n/D_n$  was further taken into consideration on the right-hand side of Eq. (2) was carried out, by utilizing such results as shown in Figs. 1 and 2. As plotted in **Fig. 3**, the values of  $Z/D_n$  were correlated by the following form within 15% error.

$$Z/D_n = (L_n/D_n)^\alpha (H_n/D_n)^\beta 10^{f(F_r)} \quad (3)$$

$$f(F_r) = \gamma + \delta(\text{Log } F_r) + \xi(\text{Log } F_r)^2 \quad (4)$$

The values of empirical constants  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\delta$  and  $\xi$ , which change depending on the range of  $L_n/D_n$  and  $H_n/D_n$ , are summarized in **Table 1**. Equations (3) and (4) are applicable when  $L_n/D_n$  is 5–70,  $H_n/D_n$  is 1.4–108, and  $F_r$  is 4.8–51.5.

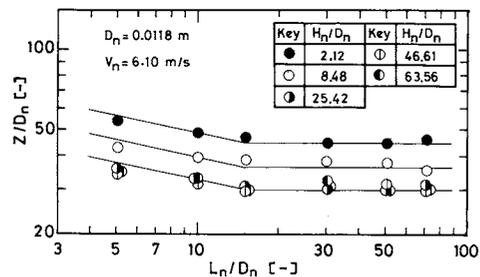


Fig. 1. Effect of  $L_n/D_n$  on  $Z/D_n$ .

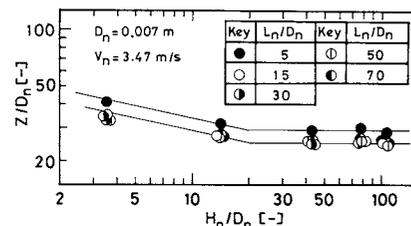


Fig. 2. Effect of  $H_n/D_n$  on  $Z/D_n$ .

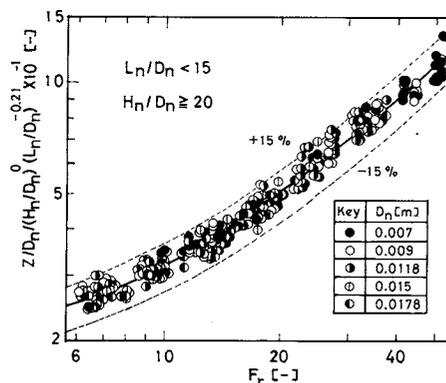


Fig. 3. Empirical correlation of  $Z/D_n$ .

Table 1. Values of  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\delta$  and  $\xi$  in Eqs. (3) and (4)

$L_n/D_n$ [—]	$H_n/D_n$ [—]	$\alpha$ [—]	$\beta$ [—]	$\gamma$ [—]	$\delta$ [—]	$\xi$ [—]
5–15	1.4–20	–0.21	–0.21	1.55	–0.15	0.37
	20–107	–0.21	0	1.27	–0.09	0.32
15–70	1.4–20	0	–0.21	1.26	–0.06	0.34
	20–107	0	0	1.09	–0.14	0.35

Some reports have dealt with  $Z$  in vertical jet systems.<sup>6,16,20)</sup> However, values of  $Z$  obtained in the system using nozzles for which the  $L_n/D_n$  ratio was specified have been reported only by Van de Sande and Smith.<sup>20)</sup> They employed vertical nozzles of  $D_n$  from 0.0039 to 0.012 m which had a value of  $L_n/D_n$  exceeding 50 and measured  $Z$  at  $H_n = 0.2$  m and at  $V_n$  below 14 m/s. **Figure 4** shows a comparison of the results calculated with those observed.<sup>20)</sup> Good coincidence can be seen.

### 2.3 Gas entrainment rate $Q_g$

In the previous paper,<sup>15)</sup> the dimensionless gas

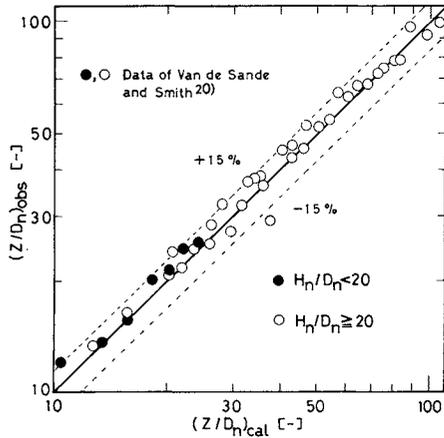


Fig. 4. Comparison of  $Z/D_n$  calculated from Eqs. (3) and (4) with values obtained by Van de Sande and Smith.<sup>20)</sup>

entrainment rate  $Q_g/Q_L$  in nozzles having a constant  $L_n/D_n$  was shown as a function of  $F_r$ ,  $H_n/D_n$  and the Ohnesorge number  $O_h (= \mu\sqrt{\rho\sigma D_n})$ .

$$Q_g/Q_L = \text{func.}(F_r, O_h, H_n/D_n) \quad (5)$$

The effect of  $L_n/D_n$  on  $Q_g/Q_L$  was further investigated. **Figure 5** shows the typical relationship between  $Q_g/Q_L$  and  $F_r$ .  $Q_g/Q_L$  is found to increase with  $F_r$  and  $H_n/D_n$ , but their rate of increase differs in the ranges of  $F_r$  below and above 20 at which  $V_n$  is about 5 m/s. For the effect of  $L_n/D_n$  and  $H_n/D_n$ , on the other hand,  $Q_g/Q_L$  when the nozzle of a large  $L_n/D_n$  is employed at a large  $H_n/D_n$  tended to increase linearly with  $F_r$  (Figs. 5b and 5c). **Figure 6** shows the typical relationship between  $Q_g/Q_L$  and  $L_n/D_n$ . These results indicate that  $Q_g/Q_L$  increases with increase of  $L_n/D_n$ , if the  $L_n/D_n$  ratio is smaller than 15, but is almost constant above this range.

#### 2.4 Correlation of $Q_g/Q_L$

From Fig. 5, the following functional form can be derived.

$$Q_g/Q_L \propto F_r^{f(H_n/D_n)} \quad (6)$$

Also, from the results for the range of  $L_n/D_n$  below 15 in Fig. 6, the relationships of Eqs. (7) and (8) which depend on the range of  $V_n$  are obtained.

$$Q_g/Q_L \propto (L_n/D_n)^{0.49} (V_n < 5 \text{ m/s}) \quad (7)$$

$$Q_g/Q_L \propto (L_n/D_n)^{0.31} (V_n \geq 5 \text{ m/s}) \quad (8)$$

Considering further the effects of  $H_n/D_n$  and  $O_h$ , the final correlation was carried out. As seen in an example of Fig. 7,  $Q_g/Q_L$  could be correlated by the following form within 30% error.

$$Q_g/Q_L = a F_r^{f(H_n/D_n)} O_h^b (H_n/D_n)^c (L_n/D_n)^d \quad (9)$$

$$f(H_n/D_n) = e (H_n/D_n)^f \quad (10)$$

The values of empirical constants  $a$ ,  $b$ ,  $c$ ,  $d$ ,  $e$  and  $f$  in Eqs. (9) and (10) which change depending on the

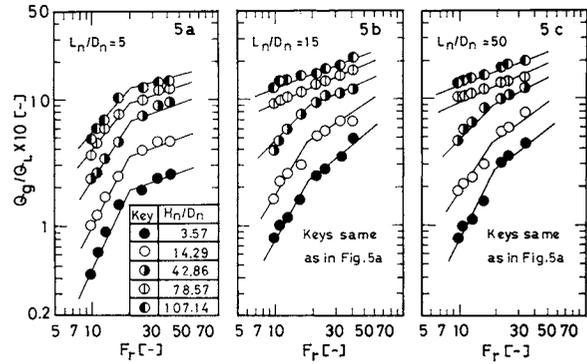


Fig. 5. Relationship between  $Q_g/Q_L$  and  $F_r$  for nozzles of  $D_n = 0.007 \text{ m}$ .

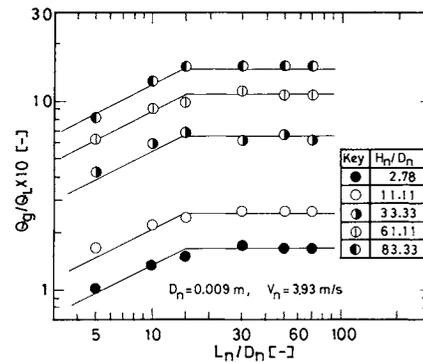


Fig. 6. Effect of  $L_n/D_n$  on  $Q_g/Q_L$ .

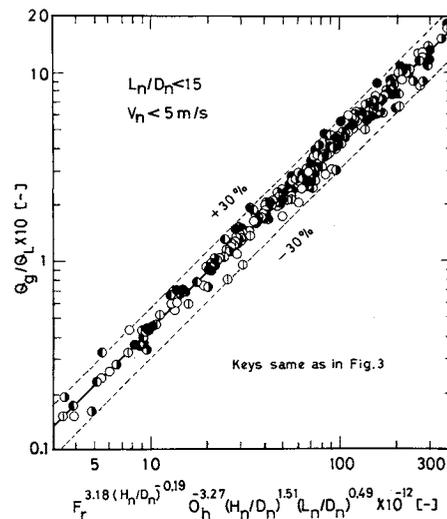


Fig. 7. Empirical correlation of  $Q_g/Q_L$ .

range of  $L_n/D_n$  and  $V_n$  are summarized in Table 2. Equations (2) and (10) are applicable when  $F_r$  is 4.8–51.5;  $H_n/D_n$  is 1.4–108;  $O_h$  is  $8.82 \times 10^{-4}$ – $1.41 \times 10^{-3}$ ; and  $L_n/D_n$  is 5–70.

Recently, Van de Sande<sup>19)</sup> proposed Eq. (11) for predicting  $Q_g/Q_L$  in the vertical water jet system, by utilizing the data of inclined jets using nozzles of  $L_n/D_n \geq 50$ .

$$Q_g/Q_L = 0.04 F_r^{0.56} (H_n/D_n)^{0.4} \quad (11)$$

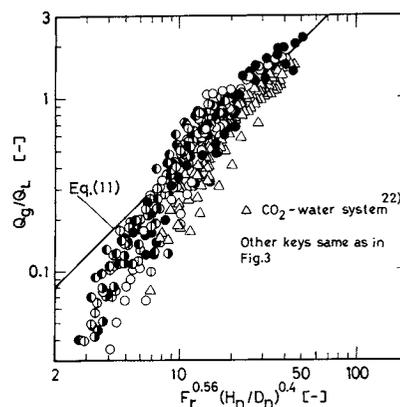
**Table 2.** Values of a, b, c, d, e and f in Eqs. (9) and (10)

$L_n/D_n$ [—]	$V_n$ [m/s]	a [—]	b [—]	c [—]	d [—]	e [—]	f [—]
5-15	2-5	$4.33 \times 10^{-15}$	-3.27	1.51	0.49	3.18	-0.19
	5-13.5	$2.57 \times 10^{-4}$	-0.55	0.84	0.31	$5.98 \times 10^{-1}$	-0.25
15-70	2-5	$1.09 \times 10^{-9}$	-1.96	1.51	0	2.22	-0.34
	5-13.5	$1.47 \times 10^{-3}$	-0.30	0.84	0	$9.37 \times 10^{-1}$	-0.21

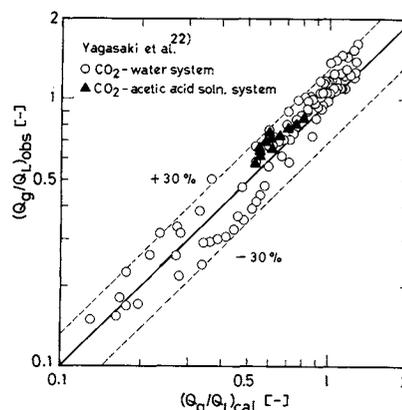
Bin<sup>3)</sup> also has discussed the applicability of Eq. (11) to the results obtained from other available data on inclined jets. **Figure 8** shows a comparison between  $Q_g/Q_L$  values calculated from Eq. (11) and those observed in the present work and in a CO<sub>2</sub>-water system ( $D_n=0.0028-0.0073$  m,  $H_n=0.02-0.25$  m and  $V_n=2.0-12.4$ ).<sup>22)</sup> As can be seen from Fig. 8, it is clear that Eq. (11) cannot be applied to all regions of gas entrainment. This is considered to be mainly due to the situation that Eq. (11) was based not on data actually measured but on data obtained by extrapolating those of inclined jets to the angle of 90°. **Figure 9** then shows a comparison of the observed  $Q_g/Q_L$  values<sup>22)</sup> with those calculated from the present correlations. The data ( $D_n=0.0057$  m,  $H_n=0.1$  m and  $V_n=5-10$  m/s)<sup>22)</sup> in an aqueous CH<sub>3</sub>COOH solution (0.1-10.0 vol%) system are also plotted. According to these results, Eqs. (9) and (10) seem to be valid for predicting not only the entrainment rate of CO<sub>2</sub> in water but also that in an aqueous CH<sub>3</sub>COOH solution when vertical nozzles with various  $L_n/D_n$  ratios are employed.

### 2.5 Change of jet surface

**Figure 10** shows typical photographs of the jet. The jet type observed within the present experimental conditions was similar to the types (A), (B), (C) and (D) shown in Fig. 10. Moreover, no break-up of liquid jets occurred. **Figure 11** shows some typical patterns of jet surface obtained by photographic observations. For the change of jet surface with  $H_n/D_n$ , the jet of types (A) and (B) was found to be produced at  $H_n/D_n$  below 20 and that of types (C) and (D) above 20. This value of 20 corresponds to the transition  $H_n/D_n$  value on the change of  $Z/D_n$  as shown in Fig. 2. For the effect of  $V_n$ , as seen in the change of (A)→(B) or (C)→(D), the surface disturbance or the sinuosity, etc. became large gradually with increasing  $V_n$ . But, their way of changing differed slightly depending on the range of  $V_n$ . That is, the degree of their surface changes at  $V_n$  above about 5 m/s tended to be on the whole smaller than that below 5 m/s. In Fig. 5a, the increasing tendency of  $Q_g/Q_L$  at  $V_n$  exceeding 5 m/s was shown to be small compared to that below this range. One of the reasons for this seems to be related to the differences in jet surface change as observed above. On the other



**Fig. 8.** Comparison of  $Q_g/Q_L$  calculated from Eq. (11) with values obtained in this work and in a CO<sub>2</sub>-water system.<sup>22)</sup>



**Fig. 9.** Comparison of  $Q_g/Q_L$  calculated from Eqs. (9) and (10) with values obtained by Yagasaki and Kuzuoka.<sup>22)</sup>

hand, as for the change of jet surface with variation of  $L_n/D_n$ , the surface disturbance, the degree of sinuosity, etc. were found to become large gradually with increasing  $L_n/D_n$  if the  $L_n/D_n$  ratio is smaller than 15 (this value corresponds to the transition  $L_n/D_n$  value on the change of  $Z/D_n$  as shown in Fig. 1). But, their changes became practically indistinguishable above this range, as seen in the change between 2 and 3 in each jet form. Most of the air entrained by the jets as shown in the form of (A) to (D) is considered to be air entrapped in the surface of the jet<sup>4)</sup> or air carried away by the sinuous undulations of the jet<sup>7,14)</sup> and so on. One of the causes of almost constant  $Q_g/Q_L$  in the nozzles of  $L_n/D_n \geq 15$  seems to be

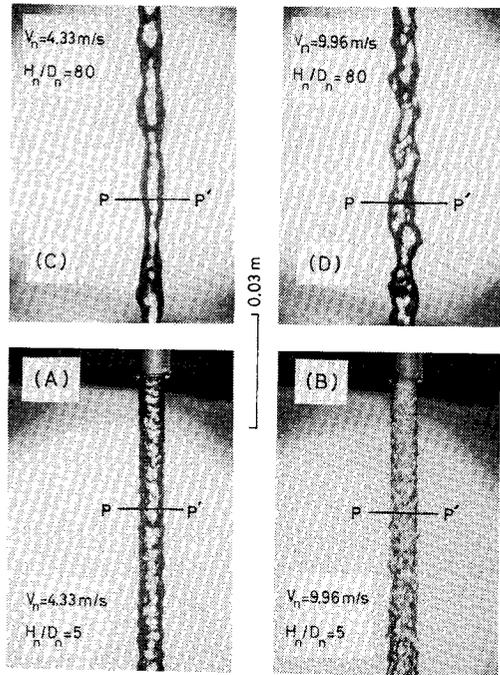


Fig. 10. Photographs of the jet for a nozzle of  $D_n=0.007$  m having a value of  $L_n/D_n=5$  (the plane  $P-P'$  in each jet represents the position corresponding to the  $H_n/D_n$  ratio designated).

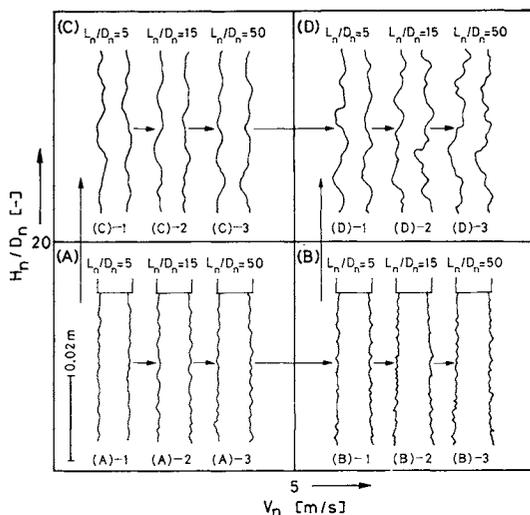


Fig. 11. Schematic representation of jet surface disturbances observed for nozzles of  $D_n=0.007$  m having various  $L_n/D_n$  ratios under varied conditions of  $H_n/D_n$  and  $V_n$  [ $V_n=4.33$  m/s for (A) and (C) and  $V_n=9.96$  m/s for (B) and (D)].

related to the situation that the way of changing of jet surface is almost the same from one to another when such a nozzle was employed at fixed conditions of  $H_n/D_n$  and  $V_n$ .

Moreover, emphasis is placed here on the fact that the above discussions are based only on surface observations of the jet before plunging. An important factor that must be further taken into consideration is the difference in flow states at the plunging point on

the bath surface, as has been reported in a study using inclined water jets.<sup>10)</sup> Further studies are necessary, in order to investigate such flow states around the plunging point and to clarify the relationship between their changes and those of the surface of the vertical jet.

## Conclusions

The relationship between the bubble penetration depth  $Z$  or the gas entrainment rate  $Q_g$  and operating conditions such as  $V_n$ ,  $H_n$  and  $D_n$  in a vertical plunging jet system where the nozzle length-to-diameter ratio  $L_n/D_n$  was varied over a wide range was investigated in an air-water system. When nozzles of  $L_n/D_n \geq 15$  were employed the values of  $Z$  and  $Q_g$  changed according to  $V_n$ ,  $H_n$  and  $D_n$ , but these values were almost independent of  $L_n/D_n$ . Based on the results obtained, empirical correlations useful for predicting  $Z$  and  $Q_g$  under any operating conditions including  $L_n/D_n$  were also presented.

## Acknowledgment

The authors wish to thank Prof. K. Endoh and Associate Prof. H. Imai of Hokkaido Univ. for their helpful suggestions. Thanks are also due to Prof. M. Fujii of Niigata Univ. for his kind advice.

## Nomenclature

$D_n$	= nozzle diameter	[m]
$F_r$	= Froude number	[—]
$g$	= acceleration of gravity	[m/s <sup>2</sup> ]
$H_n$	= nozzle height	[m]
$L_n$	= nozzle length	[m]
$O_h$	= Ohnesorge number ( $= \mu / \sqrt{\rho \sigma D_n}$ )	[—]
$Q_g$	= volumetric flow rate of entrained gas	[m <sup>3</sup> /s]
$Q_L$	= volumetric flow rate of liquid fed through the nozzle	[m <sup>3</sup> /s]
$V_n$	= jet velocity at nozzle exit	[m/s]
$Z$	= bubble penetration depth	[m]
$\rho$	= density	[kg/m <sup>3</sup> ]
$\mu$	= viscosity	[Pa · s]
$\sigma$	= surface tension	[N/m]

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## MATHEMATICAL MODELING OF TRANSDERMAL DRUG DELIVERY

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**Key Words:** Medical Chemical Engineering, Transdermal Drug Delivery, Percutaneous Absorption, Pharmacokinetics, Mathematical Model, Prodrug, Enzymatic Reaction

A dynamic mathematical model for transdermal drug delivery is developed on the basis of the bi-layer skin/two-compartment body model. The effects of metabolism reaction in the viable skin, the drug binding and reservoir function in the stratum corneum, and the solubility and diffusivity of the drug in the skin on the permeation rate-time profile are extensively simulated. The effects of the pharmacokinetic parameters on the plasma concentration profile are also analyzed. The present model is useful not only for analyzing the rate of skin permeation but also for predicting the plasma concentration after transdermal drug delivery.

### Introduction

For many years, drugs have been administered percutaneously to achieve localized pharmacological action. Recently, however, the skin has increasingly been used as a portal of entry for systemically active agents. In 1981, ALZA Corporation first introduced a product of transdermal therapeutic systems, Transderm-Scop<sup>®</sup>, for medication of motion sickness. In 1982, three transdermal delivery systems for nitroglycerin, Nitro-Dur<sup>®</sup>, Transderm-Nitro<sup>®</sup> and Nitrodisc, in the treatment of angina pectoris were introduced from Key Pharmaceuticals, Ciba Geigy and G.D. Searle, respectively. In 1985, the transdermal delivery system for clonidine, Catapres-TTS<sup>®</sup>, was introduced in antihypertensive treatment by Boehringer Ingelheim. Other transdermal drug delivery systems are also in various stages of develop-

ment and testing. We can expect a significant force of transdermal drug delivery in the near future.

In spite of the great impact of transdermal drug delivery on pharmaceutical society, the mechanism of drug permeation across the skin is not clear due to the complicated structure of the skin. During the last decade, a number of researchers have extensively studied percutaneous absorption of drugs both experimentally and theoretically. They have made great efforts to understand the penetration of drugs through the skin. The permeation data, however, were frequently analyzed on simplistic assumptions such as a single layer skin, a constant diffusivity, a steady-state condition, a simplified boundary condition, a compartment model of skin, etc.<sup>1-3,6,8,9,12-14,18,21,23,37,43,44</sup> A more rigorous approach is needed, at this stage of transdermal drug delivery studies, to elucidate the mechanism of percutaneous absorption.

The main resistance to drug transport across the

Received November 4, 1986. Correspondence concerning this article should be addressed to K. Tojo.