

Uncertainty inequalities for the Heisenberg group

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Abstract. We establish the Heisenberg–Pauli–Weyl uncertainty inequalities for Fourier transform and the continuous wavelet transform on the Heisenberg group.

Keywords. Heisenberg group; Fourier transform; wavelet transform; uncertainty inequality; heat kernel.

1. Introduction

The classical Heisenberg–Pauli–Weyl uncertainty inequality for the Fourier transform on \mathbb{R}^n is

$$\|f\|_2^4 \leq c_n \int_{\mathbb{R}^n} |x|^2 |f(x)|^2 dx \int_{\mathbb{R}^n} |\xi|^2 |\hat{f}(\xi)|^2 d\xi, \quad (1.1)$$

which can also be written in the form:

$$\|f\|_2^2 \leq c_n \| |x| f \|_2 \| (-\Delta)^{1/2} f \|_2,$$

where Δ denotes the Laplacian.

By means of the classical one-dimensional uncertainty inequality (1.1), Singer [13] obtained an uncertainty inequality for the continuous wavelet transform on a real line:

$$k_\phi \|f\|_2^2 \leq \|\tau W_\phi(f)(a, \tau)\|_{L^2(d\tau da/a^2)} \|\xi \hat{f}(\xi)\|_2, \quad (1.2)$$

where ϕ is an admissible wavelet and k_ϕ is an appropriate positive constant. For more on the history and the relevance of the uncertainty inequality, we refer the readers to the survey [5], the books [6,8], and the papers [2,10,11].

For the Heisenberg group H^n , Thangavelu [16] proved the following theorem.

Theorem 1.1. For a normalized L^2 function f on H^n , one has

$$\sqrt{n} \left(\frac{\pi}{2}\right)^{\frac{n+1}{2}} \leq \left(\int_{H^n} |z|^2 |f(z, t)|^2 dz dt \right)^{1/2} \left(\int_{H^n} |\mathcal{L}^{1/2} f(z, t)|^2 dz dt \right)^{1/2},$$

where \mathcal{L} is the Heisenberg sub-Laplacian.

Later, Sitaram *et al.* [14] obtained a generalized theorem.

Theorem 1.2. For $f \in L^2(H^n)$, $0 \leq \gamma < n + 1$, one has

$$\|f\|_2^2 \leq K \left(\int_{H^n} |(z, t)|^{2\gamma} |f(z, t)|^2 dz dt \right)^{1/2} \left(\int_{H^n} |\mathcal{L}^{\gamma/2} f(z, t)|^2 dz dt \right)^{1/2},$$

where K is a constant.

In this paper, we will prove the following two extensions:

Theorem 1.3. For $f \in L^2(H^n)$, $a, b > 0$, one has

$$\|f\|_2^2 \leq C \left(\int_{H^n} |(z, t)|^{2a} |f(z, t)|^2 dz dt \right)^{\frac{b}{a+b}} \left(\int_{H^n} |\mathcal{L}^{b/2} f(z, t)|^2 dz dt \right)^{\frac{a}{a+b}}, \quad (1.3)$$

where C is a constant.

Theorem 1.4. For $f \in H_\alpha^\sigma$, $\phi \in AW_\alpha^\sigma$ and $a, b > 0$, one has

$$\|f\|_2^2 \leq C \left(\int_0^\infty \int_{H^n} |(z, t)|^{2a} |W_\phi f(z, t)|^2 \frac{dz dt d\rho}{\rho^{n+2}} \right)^{\frac{b}{a+b}} \times \left(\int_{H^n} |\mathcal{L}^{b/2} f(z, t)|^2 dz dt \right)^{\frac{a}{a+b}},$$

where C is a constant.

The idea that we modify Theorem 1.2 to Theorem 1.3 originates from [1], where the analogue of inequality (1.3) is established for a locally compact space. The main approach of the proof of Theorem 1.3 is based on the estimate of the heat kernel together with the relation between the sub-Laplacian and the group Fourier transform. Our method is different from that in [1]. Furthermore, as an application of (1.3) we deduce an uncertainty inequality for the continuous wavelet transform on the Heisenberg group, which is an extension of (1.2).

2. Preliminaries

The Heisenberg group, denoted by H^n , is a nilpotent Lie group of step two whose underlying manifold is $\mathbb{C}^n \times \mathbb{R}$ equipped with the group law

$$(z, t)(z', t') = \left(z + z', t + t' + \frac{1}{2} \operatorname{Im} z \bar{z}' \right).$$

For $(z, t) \in H^n$, the homogeneous norm of (z, t) is given by $|(z, t)| = (|z|^4 + |t|^2)^{1/4}$. Then the ball of radius r centered at (z, t) is defined by $B_r(z, t) = \{(z', t') \in H^n : |(z, t)^{-1}(z', t')| < r\}$. Let S^n be the unit sphere in H^n and suppose f is a measurable function, one has (see [3])

$$\int_{H^n} f(z, t) dz dt = \int_{S^n} \int_0^\infty f(r\zeta) r^{2n+1} dr d\zeta. \quad (2.1)$$

Let $\pi_\lambda(z, t)$ ($z = x + iy$, $\lambda \in \mathbb{R}^* = \mathbb{R} \setminus \{0\}$) be the Schrödinger representations acted on $\varphi \in L^2(\mathbb{R}^n)$ by

$$\pi_\lambda(z, t)\varphi(\xi) = e^{i\lambda t} e^{i\lambda(x \cdot \xi + \frac{1}{2}x \cdot y)}\varphi(\xi + y).$$

Given a function $f \in L^1(H^n)$, its group Fourier transform \hat{f} is defined to be the operator-valued function and

$$\hat{f}(\lambda) = \int_{H^n} f(z, t)\pi_\lambda(z, t)dzdt.$$

Let $d\mu(\lambda) = (2\pi)^{-n-1}|\lambda|^n d\lambda$. Then one has the inversion of Fourier transform

$$f(z, t) = \int_{-\infty}^{\infty} \text{tr}(\pi_\lambda^*(z, t)\hat{f}(\lambda))d\mu(\lambda)$$

and the Plancherel formula

$$\|f\|_2^2 = \int_{\mathbb{R}^*} \|\hat{f}(\lambda)\|_{HS}^2 d\mu(\lambda).$$

Suppose f and g are measurable functions on H^n , then their convolution is defined by

$$f * g(z, t) = \int_{H^n} f((z, t)(-w, -s))g(w, s)dwds.$$

It follows from the definition of the Fourier transform that $\widehat{f * g}(\lambda) = \hat{f}(\lambda)\hat{g}(\lambda)$. In addition, one has the generalized Yong inequality:

$$\|f * g\|_r \leq \|f\|_p \|g\|_q,$$

where $\frac{1}{r} = \frac{1}{p} + \frac{1}{q} - 1$.

Now consider the Heisenberg sub-Laplacian

$$\mathcal{L} = - \sum_{j=1}^n (X_j^2 + Y_j^2),$$

where $X_j = \frac{\partial}{\partial x_j} - \frac{1}{2}y_j \frac{\partial}{\partial t}$, $Y_j = \frac{\partial}{\partial y_j} + \frac{1}{2}x_j \frac{\partial}{\partial t}$. For the Schrödinger representations π_λ one has

$$\pi_\lambda^*(X_j) = i\lambda\xi_j, \quad \pi_\lambda^*(Y_j) = \frac{\partial}{\partial \xi_j}.$$

So that $\pi_\lambda^*(\mathcal{L}) = -\Delta + \lambda^2|\xi|^2 = H(\lambda)$ is the Hermite operator. Let Φ_α ($\alpha \in \mathbb{N}^n$) stand for the normalized Hermite functions on \mathbb{R}^n . For $\lambda \in \mathbb{R}^*$, define $\Phi_\alpha^\lambda(\xi) = |\lambda|^{\frac{n}{4}}\Phi_\alpha(|\lambda|^{\frac{1}{2}}\xi)$. Then one has

$$H(\lambda)\Phi_\alpha^\lambda = (2|\alpha| + n)|\lambda|\Phi_\alpha^\lambda.$$

One important relevance of sub-Laplacian \mathcal{L} is the heat semigroup defined by

$$(e^{-s\mathcal{L}} f)(z, t) = q_s * f(z, t),$$

where q_s is the heat kernel given by

$$q_s(z, t) = c_n \int_{-\infty}^{\infty} e^{-i\lambda t} \left(\frac{\lambda}{\sinh \lambda s} \right)^n e^{-\frac{1}{4}(\lambda \coth \lambda s)|z|^2} d\lambda$$

with the positive constant c_n . Note that the heat kernel is a C^∞ function on $H^n \times (0, \infty)$ and its Fourier transform is (see p. 86 of [18])

$$\hat{q}_s(\lambda) = e^{-sH(\lambda)}.$$

More details about the sub-Laplacian and the heat kernel on Heisenberg group can be found in [4,15,17].

Now we are going to recall some facts about the continuous wavelet transform on the Heisenberg group. Let $P_\alpha (\alpha \in \mathbb{N}^n)$ be the projection from $L^2(\mathbb{R}^n)$ to one-dimensional subspace spanned by Φ_α . For $\sigma = +$ or $-$, define

$$H_\alpha^\sigma = \{f \in L^2(H^n) : \hat{f}(\lambda) = \hat{f}(\lambda)P_\alpha \quad \text{and} \quad \hat{f}(\lambda) = 0 \text{ if } \lambda \notin \mathbb{R}^\sigma\}.$$

By Theorem 1 of [9] one has

$$L^2(H^n) = \bigoplus_{\alpha \in \mathbb{N}^n} (H_\alpha^+ \oplus H_\alpha^-).$$

For a nonzero function $\phi \in H_\alpha^\sigma$, if it satisfies

$$C_\phi = \int_{\mathbb{R}^*} \frac{\|\hat{\phi}(\lambda)\|_{HS}^2}{|\lambda|} d\lambda < +\infty,$$

we then call ϕ an admissible wavelet and write $\phi \in AW_\alpha^\sigma$. Given a function $f \in H_\alpha^\sigma$, the continuous wavelet transform of f with respect to ϕ is defined by

$$W_\phi f(z, t, \rho) = \langle f, U(z, t, \rho)\phi \rangle,$$

where

$$U(z, t, \rho)\phi(z', t') = \rho^{-(n+1)/2} \phi\left(\frac{z' - z}{\sqrt{\rho}}, \frac{t' - t - \frac{1}{2}\text{Im} z z'}{\rho}\right).$$

Note the following facts: for $f \in H_\alpha^\sigma$ and $\phi \in AW_\alpha^\sigma$, one has

$$\int_{H^n} W_\phi f(z, t, \rho) \pi_\lambda(z, t) dz dt = \rho^{(n+1)/2} \hat{f}(\lambda) \hat{\phi}(\rho\lambda)^* \tag{2.2}$$

and

$$\int_0^\infty \int_{H^n} |W_\phi f(z, t, \rho)|^2 \frac{dz dt d\rho}{\rho^{n+2}} = C_\phi \|f\|_2^2. \tag{2.3}$$

For further details of wavelet transform on H^n , we refer to [7,9,12].

3. Proof of Theorems 1.3 and 1.4

In order to prove our main theorems, we first need some lemmas for preparation. Throughout the paper, we will use C to denote the positive constant, which is not necessarily same at each occurrence.

Lemma 3.1. The heat kernel $q_s(z, t)$ satisfies the estimate

$$\|q_s\|_2^2 \leq Cs^{-n-1}$$

with some positive constants C .

Proof. By Proposition 2.8.2 of [18] one has that

$$|q_s(z, t)| \leq Cs^{-n-1}e^{-\frac{A}{s}(|z|^2+|t|)}$$

holds for some positive constants C and A . Since $(|z|^4 + |t|^2)^{\frac{1}{2}} \leq |z|^2 + |t|$, we have

$$|q_s(z, t)| \leq Cs^{-n-1}e^{-\frac{A}{s}|(z,t)|^2}.$$

Thus by (2.1) we get

$$\begin{aligned} \|q_s\|_2^2 &\leq C \int_{H^n} s^{-2n-2} e^{-\frac{2A}{s}|(z,t)|^2} dz dt \\ &= C \int_0^\infty s^{-2n-2} e^{-\frac{2A}{s}r^2} r^{2n+1} dr \\ &= Cs^{-n-1}, \end{aligned}$$

which proves this lemma. □

Lemma 3.2. Let $f \in L^2(H^n)$. Then for $0 \leq \gamma < n + 1$, one has

$$\|f * q_s\|_2 \leq Cs^{-\gamma/2} \left(\int_{H^n} |(z, t)|^{2\gamma} |f(z, t)|^2 dz dt \right)^{1/2}$$

with some positive constants C .

Proof. We assume the non-trivial case that the integral of the right-hand side is finite. Let $B_r = B_r(0, 0)$ and set $f_r = f \chi_{B_r}$, $f^r = f \chi_{H^n \setminus B_r}$. Note that

$$\|\hat{f}(\lambda)\|_{HS}^2 = \sum_\alpha \|\hat{f}(\lambda)\Phi_\alpha^\lambda\|_2^2.$$

Then by the Plancherel formula we have

$$\begin{aligned} \|f^r * q_s\|_2^2 &= \int_{\mathbb{R}^*} \sum_\alpha \|\hat{f}^r(\lambda)e^{-sH(\lambda)}\Phi_\alpha^\lambda\|_2^2 d\mu(\lambda) \\ &= \int_{\mathbb{R}^*} \sum_\alpha \|\hat{f}^r(\lambda)e^{-s(2|\alpha|+n)|\lambda|}\Phi_\alpha^\lambda\|_2^2 d\mu(\lambda) \\ &\leq \int_{\mathbb{R}^*} \sum_\alpha \|\hat{f}^r(\lambda)\Phi_\alpha^\lambda\|_2^2 d\mu(\lambda) \\ &= \int_{H^n} |f^r|^2 dz dt \\ &\leq \frac{1}{r^{2\gamma}} \int_{H^n \setminus B_r} |(z, t)|^{2\gamma} |f(z, t)|^2 dz dt. \end{aligned}$$

On the other hand, by the generalized Yong inequality and Lemma 3.1 we have

$$\begin{aligned} \|f_r * q_s\|_2 &= \|q_s\|_2 \|f_r\|_1 \\ &\leq C s^{-\frac{n+1}{2}} \left(\int_{H^n} |(z,t)^{2\gamma}| f(z,t)^2 dz dt \right)^{1/2} \left(\int_{B_r} |(z,t)^{-2\gamma}| dz dt \right)^{1/2}. \end{aligned}$$

Note that the second integral of the right-hand side is controlled by $C r^{-\gamma+n+1}$ since $0 \leq \gamma < n + 1$. Therefore,

$$\begin{aligned} \|f * q_s\|_2 &\leq \|f^r * q_s\|_2 + \|f_r * q_s\|_2 \\ &\leq C r^{-\gamma} (1 + s^{-(n+1)/2} r^{n+1}) \left(\int_{H^n} |(z,t)^{2\gamma}| f(z,t)^2 dz dt \right)^{1/2}. \end{aligned}$$

Choosing $r = s^{1/2}$ then gives the desired estimate. □

Lemma 3.3. For $f \in L^2(H^n)$, if $0 < a < n + 1, 0 < b \leq 2$, one has

$$\begin{aligned} \|f\|_2^2 &\leq C \left(\int_{H^n} |(z,t)^{2a}| f(z,t)^2 dz dt \right)^{\frac{b}{a+b}} \\ &\quad \times \left(\int_{H^n} |\mathcal{L}^{b/2} f(z,t)|^2 dz dt \right)^{\frac{a}{a+b}}, \end{aligned} \tag{3.1}$$

where C is a constant.

Proof. By Lemma 3.2, we have

$$\begin{aligned} \|f\|_2 &\leq \|f * q_s\|_2 + \|f - f * q_s\|_2 \\ &\leq C s^{-a/2} \left(\int_{H^n} |(z,t)^{2a}| f(z,t)^2 dz dt \right)^{1/2} \\ &\quad + \left(\int_{\mathbb{R}^*} \sum_{\alpha} \|\hat{f}(\lambda) (1 - e^{-s(2|\alpha|+n)|\lambda|}) \Phi_{\alpha}^{\lambda}\|_2^2 d\mu(\lambda) \right)^{1/2} \\ &\leq C s^{-a/2} \left(\int_{H^n} |(z,t)^{2a}| f(z,t)^2 dz dt \right)^{1/2} \\ &\quad + \left(\int_{\mathbb{R}^*} \sum_{\alpha} \|\hat{f}(\lambda) \frac{1 - e^{-s(2|\alpha|+n)|\lambda|}}{(s(2|\alpha|+n)|\lambda|)^{b/2}} \right. \\ &\quad \left. \times (s(2|\alpha|+n)|\lambda|)^{b/2} \Phi_{\alpha}^{\lambda}\|_2^2 d\mu(\lambda) \right)^{1/2}. \end{aligned}$$

Note that if $s \geq 1$, then for $b \geq 0$,

$$\frac{1 - e^{-s}}{s^{b/2}} \leq 1;$$

if $0 < s < 1$, then for $0 \leq b \leq 2$,

$$\frac{1 - e^{-s}}{s^{b/2}} \leq C e^{-s} s^{1-b/2} \leq C,$$

since $\frac{1-e^{-s}}{e^{-s}s^{b/2}} \rightarrow 1 (s \rightarrow 0^+)$. Thus we obtain

$$\begin{aligned} \|f\|_2 &\leq C s^{-a/2} \left(\int_{H^n} |(z, t)|^{2a} |f(z, t)|^2 dz dt \right)^{1/2} \\ &\quad + C s^{b/2} \left(\int_{\mathbb{R}^*} \sum_{\alpha} \|\widehat{\mathcal{L}^{b/2} f}(\lambda) \Phi_{\alpha}^{\lambda}\|_2^2 d\mu(\lambda) \right)^{1/2}, \end{aligned} \tag{3.2}$$

where we have applied the fact $\widehat{f}(\lambda)H(\lambda)^{b/2} = \widehat{\mathcal{L}^{b/2} f}(\lambda)$ to get the last term. Minimizing the right-hand side of (3.2) then gives the desired result. \square

Proof of Theorem 1.3. If $a \geq n + 1$, let $a' < n + 1$, then we have for all $\epsilon > 0$,

$$\frac{|(z, t)|^{a'}}{\epsilon^{a'}} \leq 1 + \frac{|(z, t)|^a}{\epsilon^a},$$

which implies that

$$\| |(z, t)|^{a'} f \|_2 \leq \epsilon^{a'} \|f\|_2 + \epsilon^{a'-a} \| |(z, t)|^a f \|_2.$$

Optimizing in ϵ then gives the Landau–Kolmogorov inequality:

$$\| |(z, t)|^{a'} f \|_2 \leq C \|f\|_2^{1-a'/a} \| |(z, t)|^a f \|_2^{a'/a}. \tag{3.3}$$

Now for $b > 2$, we set $b' \leq 2$. Analogues to the above case we get

$$\| \mathcal{L}^{b'/2} f \|_2 \leq C \|f\|_2^{1-b'/b} \| \mathcal{L}^{b/2} f \|_2^{b'/b}. \tag{3.4}$$

Plugging (3.3) and (3.4) into (3.1) with a replaced by a' , b replaced by b' then suggests the desired result. \square

Proof of Theorem 1.4. First, by the Plancherel formula and (2.2) we have

$$\begin{aligned} &\int_0^\infty \int_{H^n} | \mathcal{L}^{b/2} W_{\phi} f(z, t, \rho) |^2 \frac{dz dt d\rho}{\rho^{n+2}} \\ &= \int_0^\infty \int_{\mathbb{R}^*} \| \widehat{f}(\lambda) \widehat{\phi}^*(\rho\lambda) H(\lambda)^{b/2} \|_{HS}^2 d\mu(\lambda) \frac{d\rho}{\rho} \\ &\leq \int_0^\infty \int_{\mathbb{R}^*} \| \widehat{f}(\lambda) H(\lambda)^{b/2} \|_{HS}^2 \| \widehat{\phi}^*(\rho\lambda) \|_{HS}^2 d\mu(\lambda) \frac{d\rho}{\rho} \\ &= C_{\phi} \int_{H^n} | \mathcal{L}^{b/2} f(z, t) |^2 dz dt. \end{aligned}$$

Next from Theorem 1.3, we get

$$\begin{aligned} \int_{H^n} | W_{\phi} f(z, t, \rho) |^2 dz dt &\leq C \left(\int_{H^n} |(z, t)|^{2a} | W_{\phi} f(z, t, \rho) |^2 dz dt \right)^{\frac{b}{a+b}} \\ &\quad \times \left(\int_{H^n} | \mathcal{L}^{b/2} W_{\phi} f(z, t, \rho) |^2 dz dt \right)^{\frac{a}{a+b}}. \end{aligned}$$

Integrating with respect to $d\rho/\rho^{n+2}$ we obtain

$$\int_0^\infty \int_{H^n} |W_\phi f(z, t, \rho)|^2 \frac{dz dt d\rho}{\rho^{n+2}} \leq C \int_0^\infty \left(\int_{H^n} |(z, t)|^{2a} |W_\phi f(z, t, \rho)|^2 dz dt \right)^{\frac{b}{a+b}} \times \left(\int_{H^n} |\mathcal{L}^{b/2} W_\phi f(z, t, \rho)|^2 dz dt \right)^{\frac{a}{a+b}} \frac{d\rho}{\rho^{n+2}}.$$

Then by Hölder's inequality and (2.3) we have

$$\|f\|_2^2 \leq C \left(\int_0^\infty \int_{H^n} |(z, t)|^{2a} |W_\phi f(z, t, \rho)|^2 \frac{dz dt d\rho}{\rho^{n+2}} \right)^{\frac{b}{a+b}} \times \left(\int_{H^n} |\mathcal{L}^{b/2} f(z, t)|^2 dz dt \right)^{\frac{a}{a+b}}.$$

This completes the proof of this theorem. \square

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