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Boundary value behaviors for solutions of the equilibrium equations with angular velocity

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Abstract

This work is concerned with a mixed boundary value problems for the slow equilibrium equations with prescribed angular velocity. As an application, we find sufficient conditions for the existence and uniqueness of blow-up solutions under weaker conditions.

Keywords: axisymmetric; equilibrium equation; blow-up solution

1 Introduction

In 4-D space, the equilibrium equations for a self-gravitating fluid rotating about the x_4 axis with prescribed velocity $\Omega(r)$ can be written as

$$\begin{cases} \nabla P = \rho \nabla (-\Phi + \int_0^r s \Omega^2(s) ds), \\ \Delta \Phi = 4\pi g \rho. \end{cases} \quad (1.1)$$

Here ρ , g , and Φ denote the density, gravitational constant, and gravitational potential, respectively, P is the pressure of the fluid at a point $x \in \mathbb{R}^4$, and $r = \sqrt{x_1^2 + x_2^2}$. We want to find axisymmetric equilibria and therefore always assume that $\rho(x) = \rho(r, x_4)$.

For the density ρ , from (1.1)₂ we can obtain the induced potential

$$\Phi_\rho(x) = -g \int \frac{\rho(y)}{|x-y|} dy, \quad (1.2)$$

Obviously, Φ_ρ is decreasing when ρ is increasing.

In the study of this model, Auchmuty [1] proved the existence of an equilibrium solution if the angular velocity satisfied certain decay conditions. For a constant angular velocity, Miyamoto [2] has proved that there exists an equilibrium solution if the angular velocity is less than certain constant and that there is no equilibrium for large velocity. Pang *et al.* [3] talked about the exact numbers of the stationary solutions. For many other interesting results, see references [4–6].

Under more general conditions than in [2], we prove that there exists an equilibrium solution under the following constraint set

$$\mathcal{A}_M := \left\{ \rho \mid \rho \geq 0, \rho \text{ is axisymmetric}, \int \rho \, dx = M \right\}. \quad (1.3)$$

A standard method to obtain steady states is prescribing the minimizer of the stellar energy functional. The main problem is to show that the steady state has finite mass and compact support. To approach this problem, we define the energy functional

$$F(\rho) := \int Q(\rho) \, dx - \int \rho J(r) \, dx - \frac{g}{2} \iint \frac{\rho(x)\rho(y)}{|x-y|} \, dy \, dx. \quad (1.4)$$

Here

$$Q(\rho) = \frac{1}{\gamma-1} P, \quad J(r) = \int_0^r s \Omega^2(s) \, ds, \quad (1.5)$$

In this paper, we assume that $J(r)$ is nonnegative, continuous, and bounded on $[0, +\infty)$ and P is nonnegative, continuous, and strictly increasing for $s > 0$ and satisfies:

$$P: \lim_{\rho \rightarrow 0} P(\rho) \rho^{-1} = 0, \quad \lim_{\rho \rightarrow +\infty} P(\rho) \rho^{-\frac{4}{3}} = +\infty.$$

In Section 2, first we prove the existence of a minimizer of the energy functional F in \mathcal{A}_M . Then we give the properties of minimizers; they are stationary solutions of equation (1.1) with finite mass and compact support. The main difficulty in the proof is the loss of compactness due to the unboundedness of \mathbb{R}^4 . To prevent the mass from running off to spatial infinity along a minimizing sequence, our variational approach is related to the concentration-compactness principle due to Fang and Li [4]. For many other interesting results, see references [6–8].

Throughout this paper, for simplicity of presentation, we use \int to denote $\int_{\mathbb{R}^4}$ and use $\|\cdot\|_p$ to denote $\|\cdot\|_{L^p(\mathbb{R}^4)}$. Define

$$\begin{aligned} B_R(x) &:= \{y \in \mathbb{R}^4 \mid |y-x| \leq R\}, & B_{R,K}(x) &:= \{y \in \mathbb{R}^4 \mid R \leq |y-x| \leq K\}, \\ F_{\text{pot}}(\rho) &:= -\frac{g}{2} \iint \frac{\rho(x)\rho(y)}{|x-y|} \, dy \, dx = -\frac{1}{8\pi g} \int |\nabla \Phi_\rho|^2 \, dx < 0. \end{aligned} \quad (1.6)$$

We denote by C a generic positive constant and by χ the indicator function.

2 Minimizer of the energy

In this section, we present some properties of the functional F and prove the existence of a minimizer. It is easy to verify that the function F is invariant under any vertical shift, that is, if $\rho \in \mathcal{A}_M$, then $T\rho(x) := \rho(x + ae_3) \in \mathcal{A}_M$ and $F(T\rho) = F(\rho)$ for any $a \in \mathbb{R}$. Here $e_3 = (0, 0, 1)$. Therefore, if (ρ_n) is a minimizing sequence of F in \mathcal{A}_M , then $(T\rho_n)$ is a minimizing sequence of F in \mathcal{A}_M too. First, we give some estimates.

Lemma 2.1 *Let $\rho \in L^1 \cap L^\gamma(\mathbb{R}^4)$. If $1 \leq \gamma \leq \frac{3}{2}$, then $\Phi \in L^r(\mathbb{R}^4)$ for $3 < r < \frac{3\gamma}{3-2\gamma}$, and*

$$\|\Phi\|_r^\rho \leq C(\|\rho\|_1^\alpha \|\rho\|_\gamma^{1-\alpha} + \|\rho\|_1^\beta \|\rho\|_\gamma^{1-\beta}) \quad (2.1)$$

for $0 < \alpha, \beta < 1$. If $\gamma > \frac{3}{2}$, then Φ is bounded and continuous and satisfies (2.1) with $r = +\infty$.

Proof The proof can be found in [1]. \square

Lemma 2.2 For $\rho \in L^1 \cap L^{\frac{4}{3}}(\mathbb{R}^4)$, we have $\nabla \Phi \in L^2(\mathbb{R}^4)$.

Proof Interpolation inequality [9] implies

$$\|\rho\|_{\frac{6}{5}} \leq \|\rho\|_1^{1/3} \|\rho\|_{4/3}^{2/3}.$$

By Sobolev's theorem, $\|\Phi\|_6 \leq C\|\rho\|_{\frac{6}{5}}$. So

$$\|\nabla \Phi\|_2^2 = 4\pi g \|\rho \Phi\|_1 \leq C\|\rho\|_{\frac{6}{5}} \|\Phi\|_6 \leq C\|\rho\|_{\frac{6}{5}}^2.$$

From the above estimates we can complete our proof. \square

Lemma 2.3 Assume that P_1 holds. Then there exists a nonnegative constant C , depending only on $\frac{1}{|x|}$, M , and $J(r)$, such that $F \geq -C$.

Proof For $\rho \in \mathcal{A}_M$, since P_1 holds, similarly to [2], we know that there exists a constant $S_1 > 0$ such that

$$\begin{aligned} F(\rho) &\geq \int_{\rho < S_1} Q(\rho) + \int_{\rho \geq S_1} Q(\rho) - M\|J\|_{\infty} - CM^{2/3} \int \rho^{4/3} \\ &\geq \int_{\rho < S_1} Q(\rho) + \frac{1}{2} \int_{\rho \geq S_1} Q(\rho) - M\|J\|_{\infty} - CM^{2/3} \int_{\rho < S_1} \rho^{4/3} \\ &\geq \frac{1}{2} \int Q(\rho) - M\|J\|_{\infty} - CM^{5/3} S_1^{1/3}. \end{aligned}$$

So $F \geq -C_1$ with $C_1 = M\|J\|_{\infty} - CM^{5/3} S_1^{1/3}$. \square

Let $h_M = \inf_{\mathcal{A}_M} F$. A simple scaling argument shows that $h_M < 0$: let $\bar{\rho}(x) = \varepsilon^3 \rho(\varepsilon x)$, then $\int \bar{\rho} = \int \rho$. Since $\lim_{\rho \rightarrow 0} Q(\rho) \rho^{-1} = 0$, it is easy to see that for ε small enough, $\int Q(\bar{\rho}) = \int \varepsilon^{-3} Q(\varepsilon^3 \rho) \rightarrow 0$. Therefore, $h_M < 0$.

Lemma 2.4 Assume that P_1 holds. Then for every $0 < \tilde{M} \leq M$, we have $h_{\tilde{M}} \geq (\frac{\tilde{M}}{M})^{\frac{5}{3}} h_M$.

Proof Let $\tilde{\rho}(x) = \rho(ax)$ and $\tilde{J}(r) = J(ax)$, where $a = (M/\tilde{M})^{1/3} \geq 1$. So, for any $\rho \in \mathcal{A}_M$ and $\tilde{\rho} \in \mathcal{A}_{\tilde{M}}$, we have

$$F(\tilde{\rho}) = \int Q(\tilde{\rho}) - \int \tilde{\rho} \tilde{J} + F_{\text{pot}}(\tilde{\rho}) \geq b^{-3} F(\rho). \quad (2.2)$$

The mappings $\mathcal{A}_M \rightarrow \mathcal{A}_{\tilde{M}}$, $\rho \rightarrow \tilde{\rho}$, $J \rightarrow \tilde{J}$ are all one-to-one and onto, which completes our proof. \square

From Lemma 2.3 we immediately obtain that any minimizing sequence $(\rho_n) \in \mathcal{A}_M$ of F satisfies

$$\int \rho_n^{4/3} = \int_{\rho_n < S_1} \rho_n^{4/3} + \int_{\rho_n \geq S_1} \rho_n^{4/3} < MS_1^{1/3} + \int cQ(\rho_n) < 2cF(\rho_n) + C + MS_1^{1/3}.$$

Lemma 2.5 Let (ρ_n) be bounded in $L^{4/3}(\mathbb{R}^4)$ and $\rho_n \rightharpoonup \rho_0$ weakly in $L^{4/3}(\mathbb{R}^4)$. Then, for any $R > 0$,

$$\int |\nabla \Phi_{\chi_{B_R} \rho_n}|^2 dx \rightarrow \int |\nabla \Phi_{\chi_{B_R} \rho_0}|^2 dx.$$

Proof By Sobolev theorem and Lemma 2.1 we can complete the proof. \square

Lemma 2.6 Assume that P_1 holds. Let $(\rho_n)_{n=1}^\infty \subset \mathcal{A}_M$ be a minimizing sequence of $F(\rho)$. Then there exist a sequence $(a_n)_{n=1}^\infty \subset \mathbb{R}^4$ and $\delta_0 > 0$, $R_0 > 0$ such that

$$\int_{a_n + B_R} \rho_n(x) dx \geq \delta_0, \quad R \geq R_0,$$

for all sufficiently large $n \in \mathbb{N}$.

Proof Split the potential energy:

$$\begin{aligned} -\frac{2}{g} F_{\text{pot}} &:= \int \int_{|x-y| \leq 1/R} \frac{\rho_n(x) \rho_n(y)}{|x-y|} dy dx + \int \int_{1/R < |x-y| < R} \dots + \int \int_{|x-y| \geq R} \dots \\ &:= I_1 + I_2 + I_3. \end{aligned}$$

From Lemma 2.2 we easily see that $I_1 \leq \frac{C}{R}$. The estimates for I_2 and I_3 are straightforward:

$$\begin{aligned} I_2 &\leq R \int \int_{|x-y| < R} \rho_n(x) \rho_n(y) dx dy \leq MR \sup_{a \in \mathbb{R}^4} \int_{a+B_R} \rho_n(x) dx; \\ I_3 &= \int \int_{|x-y| \geq R} \frac{\rho(x) \rho(y)}{|x-y|} dy dx \leq \frac{M^2}{R}. \end{aligned}$$

Therefore,

$$\sup_{a \in \mathbb{R}^4} \int_{a+B_R} \rho_n(x) dx \geq \frac{1}{MR} \left(-\frac{2}{g} F_{\text{pot}} - \frac{M^2}{R} - \frac{C}{R} \right). \quad (2.3)$$

We know that $F_{\text{pot}}(\rho_n) < 0$ from (1.6). Thus, when R large enough, $-F_{\text{pot}} > 0$ dominates the sign of (2.3), so that there exist $\delta_0 > 0$, $R_0 > 0$ as required. \square

We are now ready to show the existence of a minimizer of h_M , provided that P_1 holds.

Theorem 2.1 Assume that P_1 holds. Let $(\rho_n) \in \mathcal{A}_M$ be a minimizing sequence of F . Then there exist a subsequence, still denoted by (ρ_n) , and a sequence of translations $T\rho_n := \rho_n(\cdot + a_n e_3)$ with constant a_n and $e_3 = (0, 0, 1)$ such that

$$F(\rho_0) = \inf_{\mathcal{A}_M} F(\rho_n) = h_M$$

and $T\rho_n \rightharpoonup \rho_0$ weakly in $L^{4/3}(\mathbb{R}^4)$. For the induced potentials, we have $\nabla \Phi_{T\rho_n} \rightarrow \nabla \Phi_{\rho_0}$ strongly in $L^2(\mathbb{R}^4)$.

Remark 2.1 Without admitting the spatial shifts, the assertion of the theorem is false: Given a minimizer ρ_0 and a sequence of shift vectors $(a_n e_3) \in \mathbb{R}^4$, the functional F is translation invariant, that is, $F(T\rho) = F(\rho)$. But if $|a_n e_3| \rightarrow \infty$, then this minimizing sequence converges weakly to zero, which is not in \mathcal{A}_M .

Proof Split $\rho \in \mathcal{A}_M$ into three different parts:

$$\rho = \chi_{B_{R_1}} \rho + \chi_{B_{R_1, R_2}} \rho + \chi_{B_{R_2, \infty}} \rho := \rho_1 + \rho_2 + \rho_3$$

with

$$I_{lm} := \int \int \frac{\rho_l(x) \rho_m(y)}{|x-y|} dy dx, \quad l, m = 1, 2, 3.$$

Thus,

$$F(\rho) := F(\rho_1) + F(\rho_2) + F(\rho_3) - I_{12} - I_{13} - I_{23}.$$

If we choose $R_2 > 2R_1$, then

$$I_{13} \leq 2 \int_{B_{R_1}} \rho(x) dx \int_{B_{R_2, \infty}} |y|^{-1} \rho(y) dy \leq \frac{C_1}{R_2}.$$

Next we estimate I_{12} and I_{23} :

$$\begin{aligned} I_{12} + I_{23} &= - \int \rho_1 \Phi_2 dx - \int \rho_2 \Phi_3 dx = \frac{1}{4\pi g} \int \nabla(\Phi_1 + \Phi_3) \cdot \nabla \Phi_2 dx \\ &\leq C_2 \|\rho_1 + \rho_3\|_{\frac{6}{5}} \|\nabla \Phi_2\|_2 \leq C_3 \|\nabla \Phi_2\|_2, \end{aligned}$$

where $\Phi_l = \Phi_{\rho_l}$.

Denote $M_l = \int \rho_l$, $l = 1, 2, 3$. Then $M = M_1 + M_2 + M_3$. Using the above estimates and Lemma 2.4, we have

$$\begin{aligned} h_M - F(\rho) &\leq \left(1 - \left(\frac{M_1}{M}\right)^{5/3} - \left(\frac{M_2}{M}\right)^{5/3} - \left(\frac{M_3}{M}\right)^{5/3}\right) h_M + \frac{C_1}{R_2} + C_3 \|\nabla \Phi_2\|_2 \\ &\leq C_4 h_M M_1 M_3 + C_5 \left(\frac{1}{R_2} + \|\nabla \Phi_2\|_2\right), \end{aligned} \quad (2.4)$$

where C_4, C_5 are positive and depend on M but not on R_1 or R_2 . Let $(\rho_n) \in \mathcal{A}_M$ be a minimizing sequence and $(a_n e_3) \in \mathbb{R}^4$ such that Lemma 2.6 holds. Since F is translation invariant, the sequence $(T\rho_n)$ is a minimizing sequence too. So, $\|T\rho_n\|_1 \leq M$. Thus, there exists a subsequence, denoted by $(T\rho_n)$ again, such that $T\rho_n \rightharpoonup \rho_0$ weakly in $L^{\frac{4}{3}}(\mathbb{R}^4)$. By Mazur's lemma and Fatou's lemma,

$$\int Q(\rho_0) dx \leq \liminf_{n \rightarrow \infty} \int Q(T\rho_n) dx. \quad (2.5)$$

Now we want to show that

$$\nabla \Phi_{T\rho_n} \rightarrow \nabla \Phi_{\rho_0} \text{ strongly in } L^2(\mathbb{R}^4). \quad (2.6)$$

Due to Lemma 2.5, $\nabla \Phi_{T\rho_{n,1}+T\rho_{n,2}}$ converge strongly in $L^2(B_{R_2})$. Therefore, we only need to show that for any $\varepsilon > 0$,

$$\int |\nabla \Phi_{T\rho_{n,3}}|^2 dx < \varepsilon.$$

By Lemmas 2.1 and 2.2 it suffices to prove that

$$\int T\rho_{n,3} dx < \varepsilon. \quad (2.7)$$

Choosing $R_0 < R_1$, we obtain that $M_{n,1} \geq \delta_0$ for n large enough from Lemma 2.6. By (2.4) we have

$$\begin{aligned} -C_4 h_M \delta_0 M_{n,3} &\leq -C_4 h_M M_{n,1} M_{n,3} \\ &\leq \frac{C_5}{R_2} + C_5 \|\nabla \Phi_{0,2}\|_2 + C_5 \|\nabla \Phi_{n,2} - \nabla \Phi_{0,2}\|_2 + |F(T\rho_n) - h_M|, \end{aligned} \quad (2.8)$$

where $\Phi_{n,l}$ is the potential induced by $T\rho_{n,l}$, which in turn has mass $M_{n,l}$, $n \in \mathbb{N} \cup \{0\}$, and the index $l = 1, 2, 3$ refers to the splitting.

Given any $\varepsilon > 0$, by Lemma 2.6 we can increase $R_1 > R_0$ so that $C_5 \|\nabla \Phi_{0,2}\|_2 < \varepsilon/4$. Next, choose $R_2 > 2R_1$ such that the first term in (2.8) is less than $\varepsilon/4$. Now, since R_1 and R_2 are fixed, the third term converges to zero by Lemma 2.5. Since $(T\rho_n)$ is a minimizing sequence, we have $|F(T\rho_n) - h_M| < \varepsilon/4$ for suitable n . So, for n large enough,

$$-C_4 h_M \delta_0 M_{n,3} \leq \varepsilon, \quad \text{i.e., } M_{n,3} \leq \varepsilon;$$

thus, (2.7) holds, (2.6) follows, and

$$M \geq \int_{a_n + B_{R_2}} T\rho_n = M - M_{n,3} \geq M - \varepsilon.$$

Since $T\rho_n \rightharpoonup \rho_0$ weakly in $L^1(\mathbb{R}^N)$, it follows that for any $\varepsilon > 0$, there exists $R > 0$ such that

$$M \geq \int_{B_R} \rho_0 \geq M - \varepsilon;$$

thus,

$$\rho_0 \in L^1(\mathbb{R}^N) \text{ with } \int \rho_0 dx = M,$$

so that $\rho_0 \in \mathcal{A}_M$. Together with (2.5), we obtain

$$F(\rho_0) = \inf_{\mathcal{A}_M} F = h_M.$$

The proof is completed. \square

Next, we show that the minimizers obtained are steady states of equation (1.1).

Theorem 2.2 Let $\rho_0 \in \mathcal{A}_M$ be a minimizer of $F(\rho)$ with induced potential Φ_0 . Then

$$\Phi_0 + Q'(\rho_0) - J(r) = K_0 \text{ on the support of } \rho_0,$$

where K_0 is a constant. Furthermore, ρ_0 satisfies (1.1).

Proof We will derive the Euler-Lagrange equation for the variational problem. Let $\rho_0 \in \mathcal{A}_M$ be a minimizer with induced potential Φ_0 . For any $\epsilon > 0$, we define

$$V_\epsilon := \left\{ x \in \mathbb{R}^4 \mid \epsilon \leq \rho_0 \leq \frac{1}{\epsilon} \right\}.$$

For a test function $\omega \in L^\infty(\mathbb{R}^4)$ that has compact support and is nonnegative on V_ϵ^c , define

$$\rho_\tau := \rho_0 + \tau \omega - \tau \frac{\int \omega dy}{\text{meas}(V_\epsilon)} \chi_{V_\epsilon},$$

where $\tau \geq 0$ is small such that

$$\rho_\tau \geq 0, \quad \int \rho_\tau = \int \rho_0 = M.$$

Therefore, $\rho_\tau \in \mathcal{A}_M$. Since ρ_0 is a minimizer of $F(\rho)$, we have

$$\begin{aligned} 0 &\leq F(\rho_\tau) - F(\rho_0) \\ &= \int Q(\rho_\tau) - Q(\rho_0) dx - \int J(r)(\rho_\tau - \rho_0) + \frac{1}{2} \int (\rho_\tau \Phi_\tau - \rho_0 \Phi_0) dx \\ &\leq \int (Q'(\rho_0) - J(r))(\rho_\tau - \rho_0) dx + \int (\rho_\tau \Phi_0 - \rho_0 \Phi_0) dx + o(\tau) \\ &= \tau \int (Q'(\rho_0) - J(r) + \Phi_0) \left(\omega - \frac{\int \omega dy}{\text{meas}(V_\epsilon)} \chi_{V_\epsilon} \right) dx + o(\tau). \end{aligned}$$

Hence,

$$\int \left[Q'(\rho_0) - J(r) + \Phi_0 - \frac{1}{\text{meas}(V_\epsilon)} \left(\int_{V_\epsilon} Q'(\rho_0) - J(r) + \Phi_0 dy \right) \right] \omega dx \geq 0.$$

This holds for all test functions ω positive and negative on V_ϵ as specified above; hence, for all $\epsilon > 0$ small enough,

$$Q'(\rho_0) - J(r) + \Phi_0 = K_\epsilon \quad \text{on } V_\epsilon, \quad \text{and} \quad Q'(\rho_0) - J(r) + \Phi_0 \geq K_\epsilon \quad \text{on } V_\epsilon^c, \quad (2.9)$$

where K_ϵ is a constant. Taking the limit as $\epsilon \rightarrow 0$, we get

$$Q'(\rho_0) - J(r) + \Phi_0 = K_0 \text{ on the support of } \rho_0. \quad (2.10)$$

By taking the gradient of both sides of (2.10) we can prove that ρ_0 satisfies the equilibrium equation (1.1). \square

Competing interests

The authors declare that there is no conflict of interests regarding the publication of this article.

Authors' contributions

All authors contributed equally to the writing of this paper. All authors read and approved the final manuscript.

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