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# Computing eigenvalues and Hermite interpolation for Dirac systems with eigenparameter in boundary conditions

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## Abstract

Eigenvalue problems with eigenparameter appearing in the boundary conditions usually have complicated characteristic determinant where zeros cannot be explicitly computed. In this paper we use the derivative sampling theorem 'Hermite interpolations' to compute approximate values of the eigenvalues of Dirac systems with eigenvalue parameter in one or two boundary conditions. We use recently derived estimates for the truncation and amplitude errors to compute error bounds. Using computable error bounds, we obtain eigenvalue enclosures. Examples with tables and illustrative figures are given. Also numerical examples, which are given at the end of the paper, give comparisons with the classical sinc-method in Annaby and Tharwat (BIT Numer. Math. 47:699-713, 2007) and explain that the Hermite interpolations method gives remarkably better results.

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## 1 Introduction

Let  $\sigma > 0$  and  $PW_{\sigma}^2$  be the Paley-Wiener space of all  $L^2(\mathbb{R})$ -entire functions of exponential type  $\sigma$ . Assume that  $f(t) \in PW_{\sigma}^2 \subset PW_{2\sigma}^2$ . Then  $f(t)$  can be reconstructed via the Hermite-type sampling series

$$f(t) = \sum_{n=-\infty}^{\infty} \left[ f\left(\frac{n\pi}{\sigma}\right) S_n^2(t) + f'\left(\frac{n\pi}{\sigma}\right) \frac{\sin(\sigma t - n\pi)}{\sigma} S_n(t) \right], \quad (1.1)$$

where  $S_n(t)$  is the sequences of sinc functions

$$S_n(t) := \begin{cases} \frac{\sin(\sigma t - n\pi)}{(\sigma t - n\pi)}, & t \neq \frac{n\pi}{\sigma}, \\ 1, & t = \frac{n\pi}{\sigma}. \end{cases} \quad (1.2)$$

Series (1.1) converges absolutely and uniformly on  $\mathbb{R}$ , cf. [1-4]. Sometimes, series (1.1) is called the derivative sampling theorem. Our task is to use formula (1.1) to compute eigenvalues of Dirac systems numerically. This approach is a fully new technique that uses the recently obtained estimates for the truncation and amplitude errors associated with (1.1),

cf. [5]. Both types of errors normally appear in numerical techniques that use interpolation procedures. In the following we summarize these estimates. The truncation error associated with (1.1) is defined to be

$$R_N(f)(t) := f(t) - f_N(t), \quad N \in \mathbb{Z}^+, t \in \mathbb{R}, \quad (1.3)$$

where  $f_N(t)$  is the truncated series

$$f_N(t) = \sum_{|n| \leq N} \left[ f\left(\frac{n\pi}{\sigma}\right) S_n^2(t) + f'\left(\frac{n\pi}{\sigma}\right) \frac{\sin(\sigma t - n\pi)}{\sigma} S_n(t) \right]. \quad (1.4)$$

It is proved in [5] that if  $f(t) \in PW_\sigma^2$  and  $f(t)$  is sufficiently smooth in the sense that there exists  $k \in \mathbb{Z}^+$  such that  $t^k f(t) \in L^2(\mathbb{R})$ , then, for  $t \in \mathbb{R}$ ,  $|t| < N\pi/\sigma$ , we have

$$\begin{aligned} |R_N(f)(t)| &\leq T_{N,k,\sigma}(t) \\ &:= \frac{\xi_{k,\sigma} E_k |\sin \sigma t|^2}{\sqrt{3}(N+1)^k} \left( \frac{1}{(N\pi - \sigma t)^{3/2}} + \frac{1}{(N\pi + \sigma t)^{3/2}} \right) \\ &\quad + \frac{\xi_{k,\sigma} (\sigma E_k + k E_{k-1}) |\sin \sigma t|^2}{\sigma (N+1)^k} \left( \frac{1}{\sqrt{N\pi - \sigma t}} + \frac{1}{\sqrt{N\pi + \sigma t}} \right), \end{aligned} \quad (1.5)$$

where the constants  $E_k$  and  $\xi_{k,\sigma}$  are given by

$$E_k := \sqrt{\int_{-\infty}^{\infty} |t^k f(t)|^2 dt}, \quad \xi_{k,\sigma} := \frac{\sigma^{k+1/2}}{\pi^{k+1} \sqrt{1 - 4^{-k}}}. \quad (1.6)$$

The amplitude error occurs when approximate samples are used instead of the exact ones, which we cannot compute. It is defined to be

$$\begin{aligned} \mathcal{A}(\varepsilon, f)(t) &= \sum_{n=-\infty}^{\infty} \left[ \left\{ f\left(\frac{n\pi}{\sigma}\right) - \tilde{f}\left(\frac{n\pi}{\sigma}\right) \right\} S_n^2(t) \right. \\ &\quad \left. + \left\{ f'\left(\frac{n\pi}{\sigma}\right) - \tilde{f}'\left(\frac{n\pi}{\sigma}\right) \right\} \frac{\sin(\sigma t - n\pi)}{\sigma} S_n(t) \right], \quad t \in \mathbb{R}, \end{aligned} \quad (1.7)$$

where  $\tilde{f}(\frac{n\pi}{\sigma})$  and  $\tilde{f}'(\frac{n\pi}{\sigma})$  are approximate samples of  $f(\frac{n\pi}{\sigma})$  and  $f'(\frac{n\pi}{\sigma})$ , respectively. Let us assume that the differences  $\varepsilon_n := f(\frac{n\pi}{\sigma}) - \tilde{f}(\frac{n\pi}{\sigma})$ ,  $\varepsilon'_n := f'(\frac{n\pi}{\sigma}) - \tilde{f}'(\frac{n\pi}{\sigma})$ ,  $n \in \mathbb{Z}$ , are bounded by a positive number  $\varepsilon$ , i.e.,  $|\varepsilon_n|, |\varepsilon'_n| \leq \varepsilon$ . If  $f(t) \in PW_\sigma^2$  satisfies the natural decay conditions

$$|\varepsilon_n| \leq \left| f\left(\frac{n\pi}{\sigma}\right) \right|, \quad |\varepsilon'_n| \leq \left| f'\left(\frac{n\pi}{\sigma}\right) \right|, \quad (1.8)$$

$$|f(t)| \leq \frac{M_f}{|t|^{\alpha+1}}, \quad t \in \mathbb{R} - \{0\}, \quad (1.9)$$

$0 < \omega \leq 1$ , then for  $0 < \varepsilon \leq \min\{\pi/\sigma, \sigma/\pi, 1/\sqrt{e}\}$ , we have, [5],

$$\begin{aligned} \|\mathcal{A}(\varepsilon, f)\|_\infty &\leq \frac{4e^{1/4}}{\sigma(\omega+1)} \left\{ \sqrt{3}e(1+\sigma) + ((\pi/\sigma)A + M_f)\rho(\varepsilon) \right. \\ &\quad \left. + (\sigma + 2 + \log(2))M_f \right\} \varepsilon \log(1/\varepsilon), \end{aligned} \quad (1.10)$$

where

$$A := \frac{3\sigma}{\pi} \left( |f(0)| + M_f \left( \frac{\sigma}{\pi} \right)^\omega \right), \quad \rho(\varepsilon) := \gamma + 10 \log(1/\varepsilon), \quad (1.11)$$

and  $\gamma := \lim_{n \rightarrow \infty} [\sum_{k=1}^n \frac{1}{k} - \log n] \cong 0.577216$  is the Euler-Mascheroni constant.

The classical [6] sampling theorem of Whittaker, Kotel'nikov and Shannon (WKS) for  $f \in PW_\sigma^2$  is the series representation

$$f(t) = \sum_{n=-\infty}^{\infty} f\left(\frac{n\pi}{\sigma}\right) S_n(t), \quad t \in \mathbb{R}, \quad (1.12)$$

where the convergence is absolute and uniform on  $\mathbb{R}$  and it is uniform on compact sets of  $\mathbb{C}$ , cf. [6–8]. Series (1.12), which is of Lagrange interpolation type, has been used to compute eigenvalues of second-order eigenvalue problems; see, e.g., [9–15]. The use of (1.12) in numerical analysis is known as the sinc-method established by Stenger, cf. [16–18]. In [10, 12], the authors applied (1.12) and the regularized sinc-method to compute eigenvalues of Dirac systems with a derivation of the error estimates as given by [19, 20]. In [12] the Dirac system has an eigenparameter appearing in the boundary conditions. The aim of this paper is to investigate the possibilities of using Hermite interpolations rather than Lagrange interpolations, to compute the eigenvalues numerically. Notice that, due to Paley-Wiener's theorem [21],  $f \in PW_\sigma^2$  if and only if there is  $g(\cdot) \in L^2(-\sigma, \sigma)$  such that

$$f(t) = \frac{1}{\sqrt{2\pi}} \int_{-\sigma}^{\sigma} g(x) e^{ixt} dx. \quad (1.13)$$

Therefore  $f'(t) \in PW_\sigma^2$ , i.e.,  $f'(t)$  also has an expansion of the form (1.12). However,  $f'(t)$  can be also obtained by the term-by-term differentiation formula of (1.12)

$$f'(t) = \sum_{n=-\infty}^{\infty} f\left(\frac{n\pi}{\sigma}\right) S'_n(t), \quad (1.14)$$

see [6, p.52] for convergence. Thus the use of Hermite interpolations will not cost any additional computational efforts since the samples  $f(\frac{n\pi}{\sigma})$  will be used to compute both  $f(t)$  and  $f'(t)$  according to (1.12) and (1.14), respectively.

Consider the Dirac system which consists of the system of differential equations

$$u'_2(x) - r_1(x)u_1(x) = \lambda u_1(x), \quad u'_1(x) + r_2(x)u_2(x) = -\lambda u_2(x), \quad x \in [0, 1] \quad (1.15)$$

and the boundary conditions

$$\alpha_1 u_1(0) - \alpha_2 u_2(0) = -\lambda (\alpha'_1 u_1(0) - \alpha'_2 u_2(0)), \quad (1.16)$$

$$\beta_1 u_1(1) - \beta_2 u_2(1) = -\lambda (\beta'_1 u_1(1) - \beta'_2 u_2(1)), \quad (1.17)$$

where  $r_1(\cdot), r_2(\cdot) \in L^1(0, 1)$  and  $\alpha_i, \beta_i, \alpha'_i, \beta'_i \in \mathbb{R}$ ,  $i = 0, 1$ , satisfying

$$\begin{aligned} ((\alpha'_1, \alpha'_2) = (0, 0) \text{ or } \alpha'_1 \alpha_2 - \alpha_1 \alpha'_2 > 0) \quad \text{and} \\ ((\beta'_1, \beta'_2) = (0, 0) \text{ or } \beta_1 \beta'_2 - \beta'_1 \beta_2 > 0). \end{aligned} \quad (1.18)$$

The eigenvalue problem (1.15)-(1.17) will be denoted by  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  when  $(\alpha'_1, \alpha'_2) \neq (0, 0) \neq (\beta'_1, \beta'_2)$ . It is a Dirac system when the eigenparameter  $\lambda$  appears linearly in both boundary conditions. The classical problem when  $\alpha'_1 = \alpha'_2 = \beta'_1 = \beta'_2 = 0$ , which we denote by  $\Gamma(r, \alpha, \beta, 0, 0)$ , is studied in the monographs of Levitan and Sargsjan [22, 23]. Annaby and Tharwat [24] used Hermite-type sampling series (1.1) to compute the eigenvalues of problem  $\Gamma(r, \alpha, \beta, 0, 0)$  numerically. In [25], Kerimov proved that  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  has a denumerable set of real and simple eigenvalues with  $\pm\infty$  as the limit points. Similar results are established in [26] for the problem when the eigenparameter appears in one condition, *i.e.*, when  $\alpha'_1 = \alpha'_2 = 0, (\beta'_1, \beta'_2) \neq (0, 0)$  or equivalently when  $(\alpha'_1, \alpha'_2) \neq (0, 0)$  and  $\beta'_1 = \beta'_2 = 0$ , where also sampling theorems have been established. These problems will be denoted by  $\Gamma(r, \alpha, \beta, 0, \beta')$  and  $\Gamma(r, \alpha, \beta, \alpha', 0)$ , respectively. The aim of the present work is to compute the eigenvalues of  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  and  $\Gamma(r, \alpha, \beta, 0, \beta')$  numerically by the Hermite interpolations with an error analysis. This method is based on sampling theorem, Hermite interpolations, but applied to regularized functions hence avoiding any (multiple) integration and keeping the number of terms in the Cardinal series manageable. It has been demonstrated that the method is capable of delivering higher-order estimates of the eigenvalues at a very low cost; see [24]. In Sections 2 and 3, we derive the Hermite interpolation technique to compute the eigenvalues of Dirac systems with error estimates. We briefly derive some necessary asymptotics for Dirac systems' spectral quantities. The last section contains three worked examples with comparisons accompanied by figures and numerics with the Lagrange interpolation method.

## 2 Treatment of $\Gamma(r, \alpha, \beta, \alpha', \beta')$

In this section we derive approximate values of the eigenvalues of  $\Gamma(r, \alpha, \beta, \alpha', \beta')$ . Recall that  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  has a denumerable set of real and simple eigenvalues, *cf.* [25]. Let  $\varphi(\cdot, \lambda) = (\varphi_1(\cdot, \lambda), \varphi_2(\cdot, \lambda))^T$  be a solution of (1.15) satisfying the following initial:

$$\varphi_1(0, \lambda) = \alpha_2 + \lambda\alpha'_2, \quad \varphi_2(0, \lambda) = \alpha_1 + \lambda\alpha'_1. \quad (2.1)$$

Here  $A^T$  denotes the transpose of a matrix  $A$ . Since  $\varphi(\cdot, \lambda)$  satisfies (1.16), then the eigenvalues of the problem  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  are the zeros of the function

$$\Delta(\lambda) := (\beta_1 + \lambda\beta'_1)\varphi_1(1, \lambda) - (\beta_2 + \lambda\beta'_2)\varphi_2(1, \lambda). \quad (2.2)$$

Similarly to [22, p.220],  $\varphi_1(\cdot, \lambda)$  and  $\varphi_2(\cdot, \lambda)$  satisfy the system of integral equations

$$\varphi_1(x, \lambda) = -(\alpha_1 + \lambda\alpha'_1) \sin \lambda x + (\alpha_2 + \lambda\alpha'_2) \cos \lambda x + \mathcal{T}_1\varphi_1(x, \lambda) + \widetilde{\mathcal{T}}_2\varphi_2(x, \lambda), \quad (2.3)$$

$$\varphi_2(x, \lambda) = (\alpha_2 + \lambda\alpha'_2) \sin \lambda x + (\alpha_1 + \lambda\alpha'_1) \cos \lambda x - \widetilde{\mathcal{T}}_1\varphi_1(x, \lambda) + \mathcal{T}_2\varphi_2(x, \lambda), \quad (2.4)$$

where  $\mathcal{T}_i$  and  $\widetilde{\mathcal{T}}_i$ ,  $i = 1, 2$ , are the Volterra operators defined by

$$\begin{aligned} \mathcal{T}_i u(x, \lambda) &:= \int_0^x \sin \lambda(x-t) r_i(t) u(t, \lambda) dt, \\ \widetilde{\mathcal{T}}_i u(x, \lambda) &:= \int_0^x \cos \lambda(x-t) r_i(t) u(t, \lambda) dt, \quad i = 1, 2. \end{aligned} \quad (2.5)$$

For convenience, we define the constants

$$\begin{aligned} c_1 &:= \max\{|\alpha_1| + |\alpha_2|, |\alpha'_1| + |\alpha'_2|\}, & c_2 &:= \int_0^1 [|r_1(t)| + |r_2(t)|] dt, \\ c_3 &:= c_1 c_2, & c_4 &:= c_3 \exp(c_2), \\ c_5 &:= \max\{|\beta_1| + |\beta_2|, |\beta'_1| + |\beta'_2|\}, & c_6 &:= c_4 c_5. \end{aligned} \quad (2.6)$$

Define  $h_1(\cdot, \lambda)$  and  $h_2(\cdot, \lambda)$  to be

$$h_1(x, \lambda) := \mathcal{T}_1 \varphi_1(x, \lambda) + \tilde{\mathcal{T}}_2 \varphi_2(x, \lambda), \quad h_2(x, \lambda) := -\tilde{\mathcal{T}}_1 \varphi_1(x, \lambda) + \mathcal{T}_2 \varphi_2(x, \lambda). \quad (2.7)$$

As in [12] we split  $\Delta(\lambda)$  into two parts via

$$\Delta(\lambda) := \mathcal{G}(\lambda) + \mathcal{S}(\lambda), \quad (2.8)$$

where  $\mathcal{G}(\lambda)$  is the known part

$$\begin{aligned} \mathcal{G}(\lambda) &:= (\beta_1 + \lambda \beta'_1) (-(\alpha_1 + \lambda \alpha'_1) \sin \lambda + (\alpha_2 + \lambda \alpha'_2) \cos \lambda) \\ &\quad - (\beta_2 + \lambda \beta'_2) ((\alpha_2 + \lambda \alpha'_2) \sin \lambda + (\alpha_1 + \lambda \alpha'_1) \cos \lambda) \end{aligned} \quad (2.9)$$

and  $\mathcal{S}(\lambda)$  is the unknown one

$$\mathcal{S}(\lambda) := (\beta_1 + \lambda \beta'_1) h_1(1, \lambda) - (\beta_2 + \lambda \beta'_2) h_2(1, \lambda). \quad (2.10)$$

Then the function  $\mathcal{S}(\lambda)$  is entire in  $\lambda$  for each  $x \in [0, 1]$  for which, cf. [12],

$$|\mathcal{S}(\lambda)| \leq c_6 (1 + |\lambda|)^2 e^{|\Im \lambda|}, \quad \lambda \in \mathbb{C}. \quad (2.11)$$

The analyticity of  $\mathcal{S}(\lambda)$  as well as estimate (2.11) are not adequate to prove that  $\mathcal{S}(\lambda)$  lies in a Paley-Wiener space. To solve this problem, we will multiply  $\mathcal{S}(\lambda)$  by a regularization factor. Let  $\theta > 0$  and  $m \in \mathbb{Z}^+$ ,  $m \geq 4$ , be fixed. Let  $\mathcal{F}_{\theta, m}(\lambda)$  be the function

$$\mathcal{F}_{\theta, m}(\lambda) := \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^m \mathcal{S}(\lambda), \quad \lambda \in \mathbb{C}. \quad (2.12)$$

We choose  $\theta$  sufficiently small for which  $|\theta \lambda| < \pi$ . More specifications on  $m, \theta$  will be given latter on. Then  $\mathcal{F}_{\theta, m}(\lambda)$ , see [12], is an entire function of  $\lambda$  which satisfies the estimate

$$|\mathcal{F}_{\theta, m}(\lambda)| \leq \frac{c_0^m c_6 (1 + |\lambda|)^2}{(1 + \theta |\lambda|)^m} e^{|\Im \lambda| (1 + m\theta)}, \quad \lambda \in \mathbb{C}. \quad (2.13)$$

Moreover,  $\lambda^{m-3} \mathcal{F}_{\theta, m}(\lambda) \in L^2(\mathbb{R})$  and

$$E_{m-3}(\mathcal{F}_{\theta, m}) = \sqrt{\int_{-\infty}^{\infty} |\lambda^{m-3} \mathcal{F}_{\theta, m}(\lambda)|^2 d\lambda} \leq \sqrt{2} c_0^m c_6 \xi_0, \quad (2.14)$$

where

$$\xi_0 := \sqrt{\frac{1}{\theta^{2m-1}} \left( \frac{3 + 2m^2 - 6\theta + 6\theta^2 + 4m\theta - 5m}{4m^3 - 12m^2 + 11m - 3} + \frac{6\theta^3(\theta + 2m - 5)}{(4m^3 - 12m^2 + 11m - 3)(m - 2)(2m - 5)} \right)}.$$

What we have just proved is that  $\mathcal{F}_{\theta,m}(\lambda)$  belongs to the Paley-Wiener space  $PW_{\sigma}^2$  with  $\sigma = 1 + m\theta$ . Since  $\mathcal{F}_{\theta,m}(\lambda) \in PW_{\sigma}^2 \subset PW_{2\sigma}^2$ , then we can reconstruct the functions  $\mathcal{F}_{\theta,m}(\lambda)$  via the following sampling formula:

$$\mathcal{F}_{\theta,m}(\lambda) = \sum_{n=-\infty}^{\infty} \left[ \mathcal{F}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) S_n^2(\lambda) + \mathcal{F}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \frac{\sin(\sigma\lambda - n\pi)}{\sigma} S_n(\lambda) \right]. \quad (2.15)$$

Let  $N \in \mathbb{Z}^+$ ,  $N > m$ , and approximate  $\mathcal{F}_{\theta,m}(\lambda)$  by its truncated series  $\mathcal{F}_{\theta,m,N}(\lambda)$ , where

$$\mathcal{F}_{\theta,m,N}(\lambda) := \sum_{n=-N}^N \left[ \mathcal{F}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) S_n^2(\lambda) + \mathcal{F}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \frac{\sin(\sigma\lambda - n\pi)}{\sigma} S_n(\lambda) \right]. \quad (2.16)$$

Since all eigenvalues are real, then from now on we restrict ourselves to  $\lambda \in \mathbb{R}$ . Since  $\lambda^{m-3} \mathcal{F}_{\theta,m}(\lambda) \in L^2(\mathbb{R})$ , the truncation error, cf. (1.5), is given for  $|\lambda| < \frac{N\pi}{\sigma}$  by

$$|\mathcal{F}_{\theta,m}(\lambda) - \mathcal{F}_{\theta,m,N}(\lambda)| \leq T_{N,m-3,\sigma}(\lambda), \quad (2.17)$$

where

$$\begin{aligned} T_{N,m-3,\sigma}(\lambda) &:= \frac{\xi_{m-3,\sigma} E_{m-3} |\sin \sigma \lambda|^2}{\sqrt{3}(N+1)^{m-3}} \left( \frac{1}{(N\pi - \sigma\lambda)^{3/2}} + \frac{1}{(N\pi + \sigma\lambda)^{3/2}} \right) \\ &\quad + \frac{\xi_{m-3,\sigma} (\sigma E_{m-3} + (m-3)E_{m-4}) |\sin \sigma \lambda|^2}{\sigma(N+1)^{m-3}} \left( \frac{1}{\sqrt{N\pi - \sigma\lambda}} + \frac{1}{\sqrt{N\pi + \sigma\lambda}} \right). \end{aligned} \quad (2.18)$$

The samples  $\{\mathcal{F}_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  and  $\{\mathcal{F}'_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$ , in general, are not known explicitly. So, we approximate them by solving numerically  $8N + 4$  initial value problems at the nodes  $\{\frac{n\pi}{\sigma}\}_{n=-N}^N$ . Let  $\{\tilde{\mathcal{F}}_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  and  $\{\tilde{\mathcal{F}}'_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  be the approximations of the samples of  $\{\mathcal{F}_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  and  $\{\mathcal{F}'_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$ , respectively. Now we define  $\tilde{\mathcal{F}}_{\theta,m,N}(\lambda)$ , which approximates  $\mathcal{F}_{\theta,m,N}(\lambda)$

$$\begin{aligned} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda) &:= \sum_{n=-N}^N \left[ \tilde{\mathcal{F}}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) S_n^2(\lambda) + \tilde{\mathcal{F}}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \frac{\sin(\sigma\lambda - n\pi)}{\sigma} S_n(\lambda) \right], \quad N > m. \end{aligned} \quad (2.19)$$

Using standard methods for solving initial problems, we may assume that for  $|n| < N$ ,

$$\left| \mathcal{F}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) - \tilde{\mathcal{F}}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \right| < \varepsilon, \quad \left| \mathcal{F}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) - \tilde{\mathcal{F}}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \right| < \varepsilon \quad (2.20)$$

for a sufficiently small  $\varepsilon$ . From (2.13) we can see that  $\mathcal{F}_{\theta,m}(\lambda)$  satisfies the condition (1.9) when  $m \geq 4$  and therefore whenever  $0 < \varepsilon \leq \min\{\pi/\sigma, \sigma/\pi, 1/\sqrt{e}\}$ , we have

$$|\mathcal{F}_{\theta,m,N}(\lambda) - \tilde{\mathcal{F}}_{\theta,m,N}(\lambda)| \leq \mathcal{A}(\varepsilon), \quad \lambda \in \mathbb{R}, \quad (2.21)$$

where there is a positive constant  $M_{\mathcal{F}_{\theta,m}}$  for which, cf. (1.10),

$$\begin{aligned} \mathcal{A}(\varepsilon) := & \frac{2e^{1/4}}{\sigma} \left\{ \sqrt{3}e(1+\sigma) + \left( \frac{\pi}{\sigma}A + M_{\mathcal{F}_{\theta,m}} \right) \rho(\varepsilon) \right. \\ & \left. + (\sigma + 2 + \log(2))M_{\mathcal{F}_{\theta,m}} \right\} \varepsilon \log(1/\varepsilon). \end{aligned} \quad (2.22)$$

Here

$$A := \frac{3\sigma}{\pi} \left( |\mathcal{F}_{\theta,m}(0)| + \frac{\sigma}{\pi} M_{\mathcal{F}_{\theta,m}} \right), \quad \rho(\varepsilon) := \gamma + 10 \log(1/\varepsilon).$$

In the following, we use the technique of [27], where only the truncation error analysis is considered, to determine enclosure intervals for the eigenvalues; see also [24, 28]. Let  $\lambda^*$  be an eigenvalue with  $|\theta\lambda^*| < \pi$ , that is,

$$\Delta(\lambda^*) = \mathcal{G}(\lambda^*) + \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \mathcal{F}_{\theta,m}(\lambda^*) = 0.$$

Then it follows that

$$\begin{aligned} & \mathcal{G}(\lambda^*) + \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda^*) \\ &= \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda^*) - \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \mathcal{F}_{\theta,m}(\lambda^*) \\ &= \left[ \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda^*) - \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \mathcal{F}_{\theta,m,N}(\lambda^*) \right] \\ & \quad + \left[ \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \mathcal{F}_{\theta,m,N}(\lambda^*) - \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \mathcal{F}_{\theta,m}(\lambda^*) \right] \end{aligned}$$

and so

$$\left| \mathcal{G}(\lambda^*) + \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda^*) \right| \leq \left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N,m-3,\sigma}(\lambda^*) + \mathcal{A}(\varepsilon)).$$

Since  $\mathcal{G}(\lambda^*) + \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda^*)$  is given and  $\left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N,m-3,\sigma}(\lambda^*) + \mathcal{A}(\varepsilon))$  has computable upper bound, we can define an enclosure for  $\lambda^*$  by solving the following system of inequalities:

$$\begin{aligned} - \left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N,m-3,\sigma}(\lambda^*) + \mathcal{A}(\varepsilon)) & \leq \mathcal{G}(\lambda^*) + \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda^*) \\ & \leq \left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N,m-3,\sigma}(\lambda^*) + \mathcal{A}(\varepsilon)). \end{aligned} \quad (2.23)$$

Its solution is an interval containing  $\lambda^*$ , and over which the graph

$$\mathcal{G}(\lambda^*) + \left( \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right)^{-m} \tilde{\mathcal{F}}_{\theta, m, N}(\lambda^*)$$

is squeezed between the graphs

$$-\left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N, m-3, \sigma}(\lambda^*) + \mathcal{A}(\varepsilon)) \quad (2.24)$$

and

$$\left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N, m-3, \sigma}(\lambda^*) + \mathcal{A}(\varepsilon)). \quad (2.25)$$

Using the fact that

$$\tilde{\mathcal{F}}_{\theta, m, N}(\lambda) \rightarrow \mathcal{F}_{\theta, m}(\lambda)$$

uniformly over any compact set, and since  $\lambda^*$  is a simple root, we obtain, for large  $N$  and sufficiently small  $\varepsilon$ ,

$$\frac{\partial}{\partial \lambda} \left( \mathcal{G}(\lambda) + \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^{-m} \tilde{\mathcal{F}}_{\theta, m, N}(\lambda) \right) \neq 0$$

in a neighborhood of  $\lambda^*$ . Hence the graph of  $\mathcal{G}(\lambda) + \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^{-m} \tilde{\mathcal{F}}_{\theta, m, N}(\lambda)$  intersects the graphs  $-\left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N, m-3, \sigma}(\lambda) + \mathcal{A}(\varepsilon))$  and  $\left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N, m-3, \sigma}(\lambda) + \mathcal{A}(\varepsilon))$  at two points with abscissae  $a_-(\lambda^*, N, \varepsilon) \leq a_+(\lambda^*, N, \varepsilon)$  and the solution of the system of inequalities (2.23) is the interval

$$I_{\varepsilon, N} := [a_-(\lambda^*, N, \varepsilon), a_+(\lambda^*, N, \varepsilon)]$$

and in particular  $\lambda^* \in I_{\varepsilon, N}$ . Summarizing the above discussion, we arrive at the following lemma which is similar to that of [27] for Sturm-Liouville problems.

**Lemma 2.1** *For any eigenvalue  $\lambda^*$ , we can find  $N_0 \in \mathbb{Z}^+$  and sufficiently small  $\varepsilon$  such that  $\lambda^* \in I_{\varepsilon, N}$  for  $N > N_0$ . Moreover,*

$$[a_-(\lambda^*, N, \varepsilon), a_+(\lambda^*, N, \varepsilon)] \rightarrow \{\lambda^*\} \quad \text{as } N \rightarrow \infty \text{ and } \varepsilon \rightarrow 0. \quad (2.26)$$

*Proof* Since all eigenvalues of  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  are simple, then for large  $N$  and sufficiently small  $\varepsilon$ , we have  $\frac{\partial}{\partial \lambda} (\mathcal{G}(\lambda) + \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^{-m} \tilde{\mathcal{F}}_{\theta, m, N}(\lambda)) > 0$  in a neighborhood of  $\lambda^*$ . Choose  $N_0$  such that

$$\mathcal{G}(\lambda) + \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^{-m} \tilde{\mathcal{F}}_{\theta, m, N_0}(\lambda) = \pm \left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N_0, m-3, \sigma}(\lambda) + \mathcal{A}(\varepsilon))$$

has two distinct solutions which we denote by  $a_-(\lambda^*, N_0, \varepsilon) \leq a_+(\lambda^*, N_0, \varepsilon)$ . The decay of  $T_{N, m-3, \sigma}(\lambda) \rightarrow 0$  as  $N \rightarrow \infty$  and  $\mathcal{A}(\varepsilon) \rightarrow 0$  as  $\varepsilon \rightarrow 0$  will ensure the existence of the solutions  $a_-(\lambda^*, N, \varepsilon)$  and  $a_+(\lambda^*, N, \varepsilon)$  as  $N \rightarrow \infty$  and  $\varepsilon \rightarrow 0$ . For the second point, we recall



that  $\tilde{\mathcal{F}}_{\theta,m,N}(\lambda) \rightarrow \mathcal{F}_{\theta,m}(\lambda)$  as  $N \rightarrow \infty$  and as  $\varepsilon \rightarrow 0$ . Hence, by taking the limit, we obtain

$$\begin{aligned}\mathcal{G}(a_+(\lambda^*, \infty, 0)) + \left(\frac{\sin \theta \lambda^*}{\theta \lambda^*}\right)^{-m} \mathcal{F}_{\theta,m}(a_+(\lambda^*, \infty, 0)) &= 0, \\ \mathcal{G}(a_-(\lambda^*, \infty, 0)) + \left(\frac{\sin \theta \lambda^*}{\theta \lambda^*}\right)^{-m} \mathcal{F}_{\theta,m}(a_-(\lambda^*, \infty, 0)) &= 0,\end{aligned}$$

that is,  $\Delta(a_+) = \Delta(a_-) = 0$ . This leads us to conclude that  $a_+ = a_- = \lambda^*$  since  $\lambda^*$  is a simple root.

Let  $\tilde{\Delta}_N(\lambda) := \mathcal{G}(\lambda) + \left(\frac{\sin \theta \lambda}{\theta \lambda}\right)^{-m} \tilde{\mathcal{F}}_{\theta,m,N}(\lambda)$ . Then (2.17) and (2.21) imply

$$|\Delta(\lambda) - \tilde{\Delta}_N(\lambda)| \leq \left|\frac{\sin \theta \lambda}{\theta \lambda}\right|^{-m} (T_{N,m-3,\sigma}(\lambda) + \mathcal{A}(\varepsilon)), \quad |\lambda| < \frac{N\pi}{\sigma}. \quad (2.27)$$

Therefore  $\theta, m$  must be chosen so that for  $|\lambda| < \frac{N\pi}{\sigma}$

$$m \geq 4, \quad |\theta \lambda| < \pi.$$

Let  $\lambda^*$  be an eigenvalue and  $\lambda_N$  be its approximation. Thus  $\Delta(\lambda^*) = 0$  and  $\tilde{\Delta}_N(\lambda_N) = 0$ . From (2.27) we have  $|\tilde{\Delta}_N(\lambda^*)| \leq \left|\frac{\sin \theta \lambda^*}{\theta \lambda^*}\right|^{-m} (T_{N,m-3,\sigma}(\lambda^*) + \mathcal{A}(\varepsilon))$ . Now we estimate the error  $|\lambda^* - \lambda_N|$  for an eigenvalue  $\lambda^*$ .  $\square$

**Theorem 2.2** *Let  $\lambda^*$  be an eigenvalue of  $\Gamma(r, \alpha, \beta, \alpha', \beta')$ . For sufficient large  $N$ , we have the following estimate:*

$$|\lambda^* - \lambda_N| < \left|\frac{\sin \theta \lambda_N}{\theta \lambda_N}\right|^{-m} \frac{T_{N,m-3,\sigma}(\lambda_N) + \mathcal{A}(\varepsilon)}{\inf_{\zeta \in I_{\varepsilon,N}} |\Delta'(\zeta)|}. \quad (2.28)$$

Moreover,  $|\lambda^* - \lambda_N| \rightarrow 0$  when  $N \rightarrow \infty$  and  $\varepsilon \rightarrow 0$ .

*Proof* Since  $\Delta(\lambda_N) - \tilde{\Delta}_N(\lambda_N) = \Delta(\lambda_N) - \Delta(\lambda^*)$ , then from (2.27) and after replacing  $\lambda$  by  $\lambda_N$ , we obtain

$$|\Delta(\lambda_N) - \Delta(\lambda^*)| \leq \left|\frac{\sin \theta \lambda_N}{\theta \lambda_N}\right|^{-m} (T_{N,m-3,\sigma}(\lambda_N) + \mathcal{A}(\varepsilon)). \quad (2.29)$$

Using the mean value theorem yields that for some  $\zeta \in J_{\varepsilon,N} := [\min(\lambda^*, \lambda_N), \max(\lambda^*, \lambda_N)]$ ,

$$|(\lambda^* - \lambda_N) \Delta'(\zeta)| \leq \left|\frac{\sin \theta \lambda_N}{\theta \lambda_N}\right|^{-m} (T_{N,m-3,\sigma}(\lambda_N) + \mathcal{A}(\varepsilon)), \quad \zeta \in J_{\varepsilon,N} \subset I_{\varepsilon,N}. \quad (2.30)$$

Since the eigenvalues are simple, then for sufficiently large  $N$   $\inf_{\zeta \in I_{\varepsilon,N}} |\Delta'(\zeta)| > 0$  and we get (2.28). The rest of the proof follows from the fact that  $\Delta_N(\lambda)$  converges uniformly to  $\Delta(\lambda)$  in  $\mathbb{R}$  and  $\mathcal{A}(\varepsilon) \rightarrow 0$  when  $\varepsilon \rightarrow 0$ .  $\square$

### 3 The case of $\Gamma(r, \alpha, \beta, 0, \beta')$

This section includes briefly a treatment similar to that of the previous section for the eigenvalue problem  $\Gamma(r, \alpha, \beta, 0, \beta')$  introduced in Section 1 above. Notice that the condition (1.18) implies that the analysis of problem  $\Gamma(r, \alpha, \beta, 0, \beta')$  is not included in that of

$\Gamma(r, \alpha, \beta, \alpha', \beta')$ . Let  $\psi(\cdot, \lambda) = (\psi_1(\cdot, \lambda), \psi_2(\cdot, \lambda))^T$  be a solution of (1.15) satisfying the following initial:

$$\psi_1(0, \lambda) = \alpha_2, \quad \psi_2(0, \lambda) = \alpha_1. \quad (3.1)$$

Therefore, the eigenvalues of the problem in question are the zeros of the function

$$\Omega(\lambda) := (\beta_1 + \lambda\beta'_1)\psi_1(1, \lambda) - (\beta_2 + \lambda\beta'_2)\psi_2(1, \lambda). \quad (3.2)$$

Similarly to [22, p.220],  $\psi(\cdot, \lambda)$  satisfies the system of integral equations

$$\psi_1(x, \lambda) = \alpha_2 \cos \lambda x - \alpha_1 \sin \lambda x + \mathcal{T}_1 \psi_1(x, \lambda) + \tilde{\mathcal{T}}_2 \psi_2(x, \lambda), \quad (3.3)$$

$$\psi_2(x, \lambda) = \alpha_1 \cos \lambda x + \alpha_2 \sin \lambda x - \tilde{\mathcal{T}}_1 \psi_1(x, \lambda) + \mathcal{T}_2 \psi_2(x, \lambda), \quad (3.4)$$

where  $\mathcal{T}_i$  and  $\tilde{\mathcal{T}}_i$ ,  $i = 1, 2$ , are the Volterra operators defined in (2.5) above. Define  $g_1(\cdot, \lambda)$  and  $g_2(\cdot, \lambda)$  to be

$$g_1(x, \lambda) := \mathcal{T}_1 \psi_1(x, \lambda) + \tilde{\mathcal{T}}_2 \psi_2(x, \lambda), \quad g_2(x, \lambda) := -\tilde{\mathcal{T}}_1 \psi_1(x, \lambda) + \mathcal{T}_2 \psi_2(x, \lambda). \quad (3.5)$$

As in [12] we split  $\Omega(\lambda)$  into

$$\Omega(\lambda) := \mathcal{K}(\lambda) + \mathcal{U}(\lambda), \quad (3.6)$$

where  $\mathcal{K}(\lambda)$  is the known part

$$\mathcal{K}(\lambda) := (\beta_1 + \lambda\beta'_1)(\alpha_2 \cos \lambda - \alpha_1 \sin \lambda) - (\beta_2 + \lambda\beta'_2)(\alpha_1 \cos \lambda + \alpha_2 \sin \lambda) \quad (3.7)$$

and  $\mathcal{U}(\lambda)$  is the unknown one

$$\mathcal{U}(\lambda) := (\beta_1 + \lambda\beta'_1)g_1(1, \lambda) - (\beta_2 + \lambda\beta'_2)g_2(1, \lambda). \quad (3.8)$$

Then  $\mathcal{U}(\lambda)$  is entire in  $\lambda$  for each  $x \in [0, 1]$  for which, see [12],

$$|\mathcal{U}(\lambda)| \leq c_6(1 + |\lambda|)e^{|\Im \lambda|}, \quad \lambda \in \mathbb{C}. \quad (3.9)$$

Define  $\mathcal{R}_{m,\theta}(\lambda)$  to be

$$\mathcal{R}_{m,\theta}(\lambda) = \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^m \mathcal{U}(\lambda), \quad \lambda \in \mathbb{C}, \quad (3.10)$$

where  $\theta$  is sufficiently small, for which  $|\theta \lambda| < \pi$  and  $m$  are as in the previous section, but  $m \geq 3$ . Hence

$$|\mathcal{R}_{m,\theta}(\lambda)| \leq \frac{c_0^m c_6 (1 + |\lambda|)}{(1 + \theta |\lambda|)^m} e^{|\Im \lambda| (1 + m\theta)}, \quad \lambda \in \mathbb{C} \quad (3.11)$$

and  $\lambda^{m-2}\mathcal{R}_{m,\theta}(\lambda) \in L^2(\mathbb{R})$  with

$$E_{m-2}(\mathcal{R}_{m,\theta}) = \sqrt{\int_{-\infty}^{\infty} |\lambda^{m-2}\mathcal{R}_{m,\theta}(\lambda)|^2 d\lambda} \leq c_0^m c_6 \omega_0, \quad (3.12)$$

where

$$\omega_0 := \sqrt{\frac{2(3-5m+2m^2-3\theta+2m\theta+\theta^2)}{\theta^{2m-1}(-3+11m-12m^2+4m^3)}}.$$

Thus,  $\mathcal{R}_{m,\theta}(\lambda)$  belongs to the Paley-Wiener space  $PW_{\sigma}^2$  with  $\sigma = 1 + m\theta$ . Since  $\mathcal{R}_{\theta,m}(\lambda) \in PW_{\sigma}^2 \subset PW_{2\sigma}^2$ , then we can reconstruct the functions  $\mathcal{R}_{\theta,m}(\lambda)$  via the following sampling formula:

$$\mathcal{R}_{\theta,m}(\lambda) = \sum_{n=-\infty}^{\infty} \left[ \mathcal{R}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) S_n^2(\lambda) + \mathcal{R}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \frac{\sin(\sigma\lambda - n\pi)}{\sigma} S_n(\lambda) \right]. \quad (3.13)$$

Let  $N \in \mathbb{Z}^+$ ,  $N > m$ , and approximate  $\mathcal{R}_{\theta,m}(\lambda)$  by its truncated series  $\mathcal{R}_{\theta,m,N}(\lambda)$ , where

$$\mathcal{R}_{\theta,m,N}(\lambda) := \sum_{n=-N}^N \left[ \mathcal{R}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) S_n^2(\lambda) + \mathcal{R}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \frac{\sin(\sigma\lambda - n\pi)}{\sigma} S_n(\lambda) \right]. \quad (3.14)$$

Since all eigenvalues are real, then from now on we restrict ourselves to  $\lambda \in \mathbb{R}$ . Since  $\lambda^{m-2}\mathcal{R}_{\theta,m}(\lambda) \in L^2(\mathbb{R})$ , the truncation error, cf. (1.5), is given for  $|\lambda| < \frac{N\pi}{\sigma}$  by

$$|\mathcal{R}_{\theta,m}(\lambda) - \mathcal{R}_{\theta,m,N}(\lambda)| \leq T_{N,m-2,\sigma}(\lambda), \quad (3.15)$$

where

$$\begin{aligned} T_{N,m-2,\sigma}(\lambda) &:= \frac{\xi_{m-2,\sigma} E_{m-2} |\sin \sigma \lambda|^2}{\sqrt{3}(N+1)^{m-2}} \left( \frac{1}{(N\pi - \sigma\lambda)^{3/2}} + \frac{1}{(N\pi + \sigma\lambda)^{3/2}} \right) \\ &\quad + \frac{\xi_{m-2,\sigma} (\sigma E_{m-2} + (m-2)E_{m-3}) |\sin \sigma \lambda|^2}{\sigma(N+1)^{m-2}} \left( \frac{1}{\sqrt{N\pi - \sigma\lambda}} + \frac{1}{\sqrt{N\pi + \sigma\lambda}} \right). \end{aligned} \quad (3.16)$$

The samples  $\{\mathcal{R}_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  and  $\{\mathcal{R}'_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$ , in general, are not known explicitly. So, we approximate them by solving numerically  $8N+4$  initial value problems at the nodes  $\{\frac{n\pi}{\sigma}\}_{n=-N}^N$ . Let  $\{\tilde{\mathcal{R}}_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  and  $\{\tilde{\mathcal{R}}'_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  be the approximations of the samples of  $\{\mathcal{R}_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$  and  $\{\mathcal{R}'_{\theta,m}(\frac{n\pi}{\sigma})\}_{n=-N}^N$ , respectively. Now we define  $\tilde{\mathcal{R}}_{\theta,m,N}(\lambda)$ , which approximates  $\mathcal{R}_{\theta,m,N}(\lambda)$

$$\begin{aligned} \tilde{\mathcal{R}}_{\theta,m,N}(\lambda) &:= \sum_{n=-N}^N \left[ \tilde{\mathcal{R}}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) S_n^2(\lambda) + \tilde{\mathcal{R}}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \frac{\sin(\sigma\lambda - n\pi)}{\sigma} S_n(\lambda) \right], \quad N > m. \end{aligned} \quad (3.17)$$

Using standard methods for solving initial problems, we may assume that for  $|n| < N$ ,

$$\left| \mathcal{R}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) - \tilde{\mathcal{R}}_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \right| < \varepsilon, \quad \left| \mathcal{R}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) - \tilde{\mathcal{R}}'_{\theta,m} \left( \frac{n\pi}{\sigma} \right) \right| < \varepsilon \quad (3.18)$$

for a sufficiently small  $\varepsilon$ . From (2.13) we can see that  $\mathcal{R}_{\theta,m}(\lambda)$  satisfies the condition (1.9) when  $m \geq 3$  and therefore whenever  $0 < \varepsilon \leq \min\{\pi/\sigma, \sigma/\pi, 1/\sqrt{e}\}$ , we have

$$|\mathcal{R}_{\theta,m,N}(\lambda) - \tilde{\mathcal{R}}_{\theta,m,N}(\lambda)| \leq \mathcal{A}(\varepsilon), \quad \lambda \in \mathbb{R}, \quad (3.19)$$

where there is a positive constant  $M_{\mathcal{R}_{\theta,m}}$  for which, cf. (1.10),

$$\begin{aligned} \mathcal{A}(\varepsilon) := & \frac{2e^{1/4}}{\sigma} \left\{ \sqrt{3}e(1+\sigma) + \left( \frac{\pi}{\sigma}A + M_{\mathcal{R}_{\theta,m}} \right) \rho(\varepsilon) \right. \\ & \left. + (\sigma + 2 + \log(2))M_{\mathcal{R}_{\theta,m}} \right\} \varepsilon \log(1/\varepsilon). \end{aligned} \quad (3.20)$$

Here

$$A := \frac{3\sigma}{\pi} \left( |\mathcal{R}_{\theta,m}(0)| + \frac{\sigma}{\pi} M_{\mathcal{R}_{\theta,m}} \right), \quad \rho(\varepsilon) := \gamma + 10 \log(1/\varepsilon).$$

As in the above section, we have the following lemma.

**Lemma 3.1** *For any eigenvalue  $\lambda^*$  of the problem  $\Gamma(r, \alpha, \beta, 0, \beta')$ , we can find  $N_0 \in \mathbb{Z}^+$  and sufficiently small  $\varepsilon$  such that  $\lambda^* \in \mathcal{I}_{\varepsilon,N}$  for  $N > N_0$ , where*

$$\mathcal{I}_{\varepsilon,N} := [b_-(\lambda^*, N, \varepsilon), b_+(\lambda^*, N, \varepsilon)],$$

$b_-, b_+$  are the solutions of the inequalities

$$-\left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N,m-2,\sigma}(\lambda) + \mathcal{A}(\varepsilon)) \leq \tilde{\Omega}_N(\lambda) \leq \left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N,m-2,\sigma}(\lambda) + \mathcal{A}(\varepsilon)). \quad (3.21)$$

Moreover,

$$[b_-(\lambda^*, N, \varepsilon), b_+(\lambda^*, N, \varepsilon)] \rightarrow \{\lambda^*\} \quad \text{as } N \rightarrow \infty \text{ and } \varepsilon \rightarrow 0. \quad (3.22)$$

Let  $\tilde{\Omega}_N(\lambda) := \mathcal{K}(\lambda) + \left( \frac{\sin \theta \lambda}{\theta \lambda} \right)^{-m} \tilde{\mathcal{R}}_{\theta,m,N}(\lambda)$ . Then (3.15) and (3.19) imply

$$|\Omega(\lambda) - \tilde{\Omega}_N(\lambda)| \leq \left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N,m-2,\sigma}(\lambda) + \mathcal{A}(\varepsilon)), \quad |\lambda| < \frac{N\pi}{\sigma}. \quad (3.23)$$

Therefore,  $\theta, m$  must be chosen so that for  $|\lambda| < \frac{N\pi}{\sigma}$ ,

$$m \geq 3, \quad |\theta \lambda| < \pi.$$

Let  $\lambda^*$  be an eigenvalue and  $\lambda_N$  be its approximation. Thus  $\Omega(\lambda^*) = 0$  and  $\tilde{\Omega}_N(\lambda_N) = 0$ . From (3.23) we have  $|\tilde{\Omega}_N(\lambda^*)| \leq \left| \frac{\sin \theta \lambda^*}{\theta \lambda^*} \right|^{-m} (T_{N,m-2,\sigma}(\lambda^*) + \mathcal{A}(\varepsilon))$ . Now we estimate the error  $|\lambda^* - \lambda_N|$  for an eigenvalue  $\lambda^*$ . Finally, we have the following estimate.

**Theorem 3.2** Let  $\lambda^*$  be an eigenvalue of the problem  $\Gamma(r, \alpha, \beta, 0, \beta')$ . For sufficient large  $N$ , we have the following estimate:

$$|\lambda^* - \lambda_N| < \left| \frac{\sin \theta \lambda_N}{\theta \lambda_N} \right|^{-m} \frac{T_{N,m-2,\sigma}(\lambda_N) + \mathcal{A}(\varepsilon)}{\inf_{\zeta \in \mathcal{I}_{\varepsilon,N}} |\Omega'(\zeta)|}. \quad (3.24)$$

Moreover,  $|\lambda^* - \lambda_N| \rightarrow 0$  when  $N \rightarrow \infty$  and  $\varepsilon \rightarrow 0$ .

In the following section, we have taken  $\theta = 1/(N - m)$ , where  $\sigma = 1 + m\theta$ , in order to avoid the first singularity of  $(\frac{\sin \theta \lambda_N}{\theta \lambda_N})^{-1}$ .

#### 4 Examples

This section includes three detailed worked examples illustrating the above technique accompanied by comparison with the sinc-method derived in [12]. It is clearly seen that the Hermite interpolations method gives remarkably better results. The first two examples are computed in [12] with the classical sinc-method where  $r_1(x) = r_2(x)$ . But in the last example, where eigenvalues cannot be computed concretely,  $r_1(x) \neq r_2(x)$ . By  $E_S$  and  $E_H$  we mean the absolute errors associated with the results of the classical sinc-method and our new method (Hermite interpolations), respectively. We indicate in these examples the effect of the amplitude error in the method by determining enclosure intervals for different values of  $\varepsilon$ . We also indicate the effect of the parameters  $m$  and  $\theta$  by several choices. Each example is exhibited via figures that accurately illustrate the procedure near to some of the approximated eigenvalues. More explanations are given below. Recall that  $a_{\pm}(\lambda)$  and  $b_{\pm}(\lambda)$  are defined by

$$a_{\pm}(\lambda) = \tilde{\Delta}_N(\lambda) \pm \left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N,m-3,\sigma}(\lambda) + \mathcal{A}(\varepsilon)), \quad |\lambda| < \frac{N\pi}{\sigma}, \quad (4.1)$$

$$b_{\pm}(\lambda) = \tilde{\Omega}_N(\lambda) \pm \left| \frac{\sin \theta \lambda}{\theta \lambda} \right|^{-m} (T_{N,m-2,\sigma}(\lambda) + \mathcal{A}(\varepsilon)), \quad |\lambda| < \frac{N\pi}{\sigma}, \quad (4.2)$$

respectively. Recall also that the enclosure intervals  $I_{\varepsilon,N} := [a_-, a_+]$  and  $\mathcal{I}_{\varepsilon,N} := [b_-, b_+]$  are determined by solving

$$a_{\pm}(\lambda) = 0, \quad |\lambda| < \frac{N\pi}{\sigma}, \quad (4.3)$$

$$b_{\pm}(\lambda) = 0, \quad |\lambda| < \frac{N\pi}{\sigma}, \quad (4.4)$$

respectively. We would like to mention that MATHEMATICA has been used to obtain the exact values for the three examples where eigenvalues cannot be computed concretely. MATHEMATICA is also used in rounding the exact eigenvalues, which are square roots.

**Example 1** The boundary value problem

$$u_2'(x) - x^2 u_1(x) = \lambda u_1(x), \quad u_1'(x) + x^2 u_2(x) = -\lambda u_2(x), \quad 0 \leq x \leq 1, \quad (4.5)$$

$$u_1(0) = -\lambda u_2(0), \quad u_1(1) = \lambda u_2(1) \quad (4.6)$$

is a special case of the problem  $\Gamma(r, \alpha, \beta, \alpha', \beta')$  when  $r_1(x) = r_2(x) = x^2$ ,  $\alpha'_1 = \alpha_2 = \beta'_1 = \beta_2 = 0$ ,  $\alpha_1 = -1$  and  $\alpha'_2 = \beta_1 = \beta'_2 = 1$ . Here the characteristic function is

**Table 1**  $N = 20, m = 10, \theta = 1/10$

$\lambda_k$	Sinc $\lambda_{k,N}$	Exact $\lambda_k$	Hermite $\lambda_{k,N}$	$E_S$	$E_H$
$\lambda_{-2}$	-1.505786875767961	-1.5057868758327264	-1.5057868758327246	$6.47653 \times 10^{-11}$	$1.77636 \times 10^{-15}$
$\lambda_{-1}$	-0.11141619186432938	-0.11141619146375636	-0.11141619146375908	$4.00573 \times 10^{-10}$	$2.72005 \times 10^{-15}$
$\lambda_0$	1.1223201536675476	1.1223201551741047	1.1223201551741295	$1.50656 \times 10^{-9}$	$2.4869 \times 10^{-14}$
$\lambda_1$	3.3830704087110752	3.383070408212596	3.3830704082125935	$4.98479 \times 10^{-10}$	$2.66454 \times 10^{-15}$

**Table 2**  $N = 20, m = 15, \theta = 1/5$

$\lambda_k$	Sinc $\lambda_{k,N}$	Exact $\lambda_k$	Hermite $\lambda_{k,N}$
$\lambda_{-2}$	-1.5057868758327237144550336	-1.5057868758327218561623117	-1.5057868758327237144491654
$\lambda_{-1}$	-0.1114161914637569327965574	-0.1114161914637563667829627	-0.1114161914637563668056111
$\lambda_0$	1.1223201551741129577354075	1.1223201551741041543767735	1.1223201551741041544693398
$\lambda_1$	3.3830704082126126090125379	3.3830704082125963004202471	3.3830704082125963003644934

**Table 3** Absolute error  $|\lambda_k - \lambda_{k,N}|$  for  $N = 20, m = 15, \theta = 1/5$

$\lambda_k$	$\lambda_{-2}$	$\lambda_{-1}$	$\lambda_0$	$\lambda_1$
$E_S$	$1.858 \times 10^{-15}$	$5.660 \times 10^{-16}$	$8.803 \times 10^{-15}$	$1.630 \times 10^{-14}$
$E_H$	$5.868 \times 10^{-21}$	$2.265 \times 10^{-20}$	$9.257 \times 10^{-20}$	$5.575 \times 10^{-20}$

**Table 4** For  $N = 20, m = 10$  and  $\theta = 1/10$ , the exact solutions  $\lambda_k$  are all inside the interval  $[a_-, a_+]$  for different values of  $\varepsilon$

$\lambda_k$	Exact $\lambda_k$	$[a_-, a_+], \varepsilon = 10^{-10}$	$[a_-, a_+], \varepsilon = 10^{-15}$
$\lambda_{-2}$	-1.5057868758327264	$[-1.650349, -1.220683]$	$[-1.508403, -1.502664]$
$\lambda_{-1}$	-0.11141619146375636	$[-0.179803, -0.071447]$	$[-0.130019, -0.100199]$
$\lambda_0$	1.1223201551741047	$[0.429491, 1.314588]$	$[0.884579, 1.206467]$
$\lambda_1$	3.383070408212596	$[3.314309, 3.464923]$	$[3.369349, 3.400197]$

$$E_7(\mathcal{F}_{\theta,m}) = 3.05294 \times 10^{11}, E_6(\mathcal{F}_{\theta,m}) = 2.53419 \times 10^9, \omega = 1, M_{\mathcal{F}_{\theta,m}} = 3.56048 \times 10^6.$$

**Table 5** With  $N = 20, m = 15$  and  $\theta = 1/5$ ,  $\lambda_k$  are all inside the interval  $[a_-, a_+]$  for different values of  $\varepsilon$

$\lambda_k$	Exact $\lambda_k$	$[a_-, a_+], \varepsilon = 10^{-10}$	$[a_-, a_+], \varepsilon = 10^{-15}$
$\lambda_{-2}$	-1.5057868758327218561623117	$[-1.652755, -1.334613]$	$[-1.505894, -1.505678]$
$\lambda_{-1}$	-0.1114161914637563667829627	$[-0.331996, 0.121731]$	$[-0.111834, -0.111003]$
$\lambda_0$	1.1223201551741041543767735	$[0.923906, 1.285003]$	$[1.120633, 1.124014]$
$\lambda_1$	3.3830704082125963004202471	$[3.241846, 3.533914]$	$[3.382059, 3.384093]$

$$E_{12}(\mathcal{F}_{\theta,m}) = 1.61064 \times 10^{13}, E_{11}(\mathcal{F}_{\theta,m}) = 1.71043 \times 10^{11}, \omega = 1, M_{\mathcal{F}_{\theta,m}} = 3.98665 \times 10^6.$$

$$\Delta(\lambda) := 2\lambda \cos\left(\frac{1}{3} + \lambda\right) - (\lambda^2 - 1) \sin\left(\frac{1}{3} + \lambda\right). \quad (4.7)$$

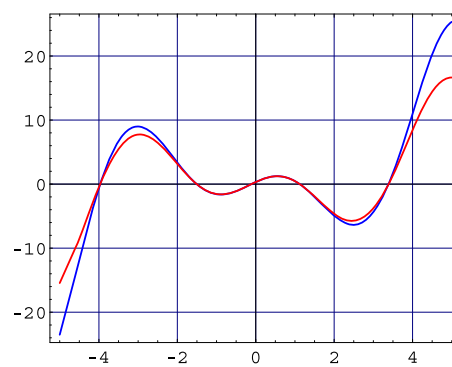
The function  $\mathcal{G}(\lambda)$  will be

$$\mathcal{G}(\lambda) := 2\lambda \cos \lambda + (1 - \lambda^2) \sin \lambda. \quad (4.8)$$

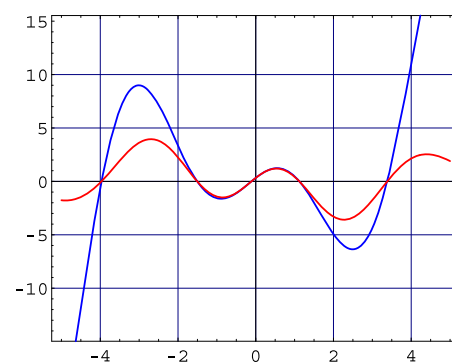
As is clearly seen, eigenvalues cannot be computed explicitly. Five tables indicate the application of our technique to this problem and the effect of  $\varepsilon$ ,  $\theta$  and  $m$  (Tables 1, 2, 3, 4 and 5). By exact, we mean the zeros of  $\Delta(\lambda)$  computed by Mathematica.

Figures 1 and 2 illustrate the comparison between  $\Delta(\lambda)$  and  $\tilde{\Delta}_N(\lambda)$  for different values of  $m$  and  $\theta$ . Figures 3 and 4, for  $N = 20, m = 10$  and  $\theta = 1/10$ , illustrate the enclosure inter-

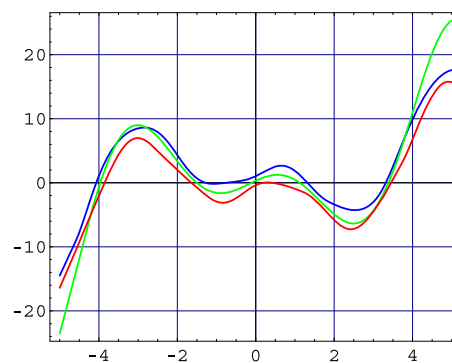
**Figure 1**  $\Delta(\lambda)$ ,  $\widetilde{\Delta}_N(\lambda)$  with  $N = 20$ ,  $m = 10$  and  $\theta = 1/10$ .



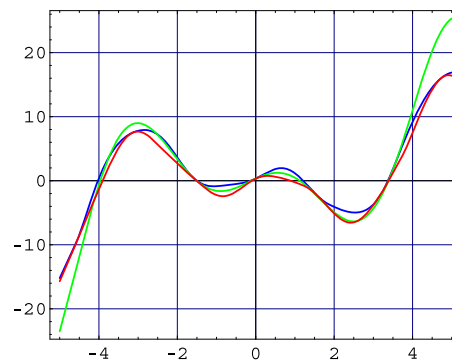
**Figure 2**  $\Delta(\lambda)$ ,  $\widetilde{\Delta}_N(\lambda)$  with  $N = 20$ ,  $m = 15$  and  $\theta = 1/5$ .



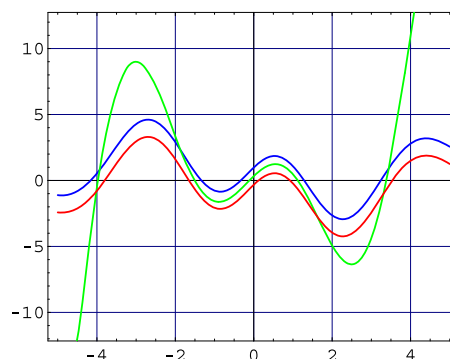
**Figure 3**  $a_+$ ,  $\Delta(\lambda)$ ,  $a_-$  with  $N = 20$ ,  $m = 10$ ,  $\theta = 1/10$  and  $\varepsilon = 10^{-10}$ .



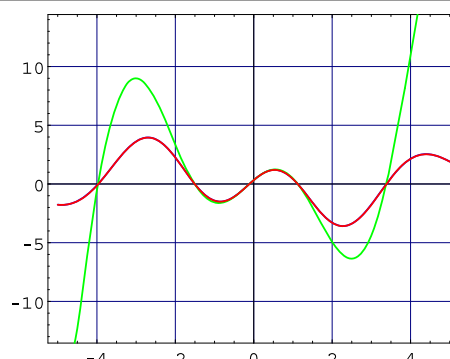
**Figure 4**  $a_+$ ,  $\Delta(\lambda)$ ,  $a_-$  with  $N = 20$ ,  $m = 10$ ,  $\theta = 1/10$  and  $\varepsilon = 10^{-15}$ .



**Figure 5**  $a_+$ ,  $\Delta(\lambda)$ ,  $a_-$  with  $N = 20$ ,  $m = 15$ ,  $\theta = 1/5$  and  $\varepsilon = 10^{-10}$ .



**Figure 6**  $a_+$ ,  $\Delta(\lambda)$ ,  $a_-$  with  $N = 20$ ,  $m = 15$ ,  $\theta = 1/5$  and  $\varepsilon = 10^{-15}$ .



vals for  $\varepsilon = 10^{-10}$  and  $\varepsilon = 10^{-15}$ , respectively. Also, Figures 5 and 6 illustrate the enclosure intervals for  $\varepsilon = 10^{-10}$  and  $\varepsilon = 10^{-15}$ , respectively, but for  $m = 15$ ,  $\theta = 1/5$ .

### Example 2 The Dirac system

$$u_2'(x) - x^2 u_1(x) = \lambda u_1(x), \quad u_1'(x) + x^2 u_2(x) = -\lambda u_2(x), \quad 0 \leq x \leq 1, \quad (4.9)$$

$$u_1(0) = 0, \quad u_2(1) = -\lambda u_1(1) \quad (4.10)$$

is a special case of the problem treated in the previous section with  $r_1(x) = r_2(x) = x^2$ ,  $\alpha_1 = \beta_1' = 1$ ,  $\alpha_2 = \beta_1 = \beta_2' = 0$  and  $\beta_2 = -1$ . The characteristic function is

$$\Omega(\lambda) := \cos\left(\frac{1}{3} + \lambda\right) - \lambda \sin\left(\frac{1}{3} + \lambda\right). \quad (4.11)$$

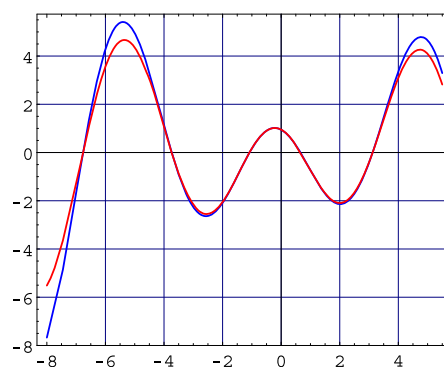
The function  $\mathcal{K}(\lambda)$  will be

$$\mathcal{K}(\lambda) := \cos \lambda - \lambda \sin \lambda. \quad (4.12)$$

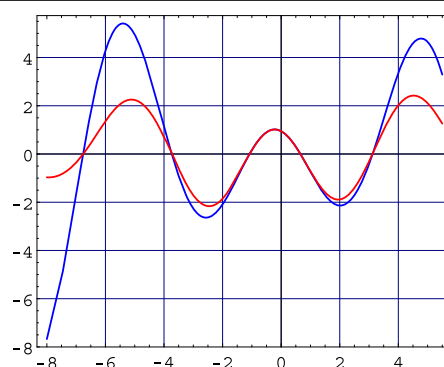
As in the previous example, Figures 7, 8, 9, 10, 11 and 12 illustrate the results of Tables 6, 7, 8, 9 and 10. Figures 7 and 8 illustrate the comparison between  $\Omega(\lambda)$  and  $\tilde{\Omega}_N(\lambda)$  for different values of  $m$  and  $\theta$ . Figures 9 and 10, for  $N = 20$ ,  $m = 6$  and  $\theta = 1/14$ , illustrate the enclosure intervals for  $\varepsilon = 10^{-10}$  and  $\varepsilon = 10^{-15}$ , respectively. Also, Figures 11 and 12 illustrate the enclosure intervals for  $\varepsilon = 10^{-10}$  and  $\varepsilon = 10^{-15}$ , respectively, but for  $m = 12$ ,  $\theta = 1/8$ .



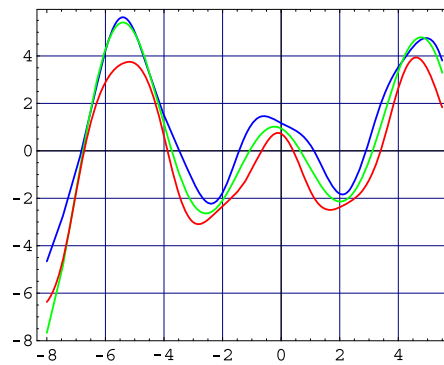
**Figure 7**  $\Omega(\lambda)$ ,  $\widetilde{\Omega}_N(\lambda)$  with  $N = 20$ ,  $m = 6$  and  $\theta = 1/14$ .



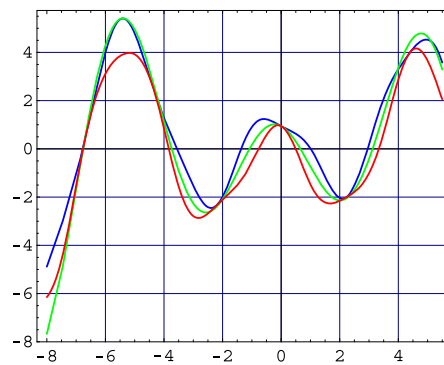
**Figure 8**  $\Omega(\lambda)$ ,  $\widetilde{\Omega}_N(\lambda)$  with  $N = 20$ ,  $m = 12$  and  $\theta = 1/8$ .



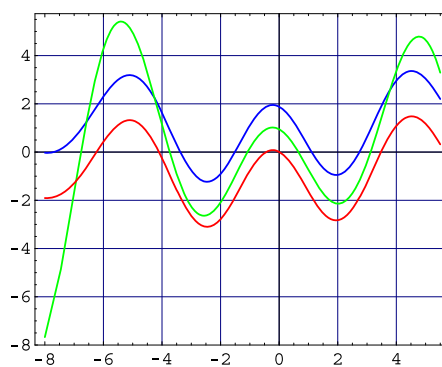
**Figure 9**  $b_+$ ,  $\Omega(\lambda)$ ,  $b_-$  with  $N = 20$ ,  $m = 6$ ,  $\theta = 1/14$  and  $\varepsilon = 10^{-10}$ .



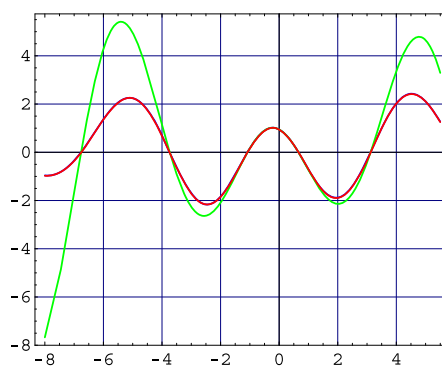
**Figure 10**  $b_+$ ,  $\Omega(\lambda)$ ,  $b_-$  with  $N = 20$ ,  $m = 6$ ,  $\theta = 1/14$  and  $\varepsilon = 10^{-15}$ .



**Figure 11**  $b_+$ ,  $\Omega(\lambda)$ ,  $b_-$  with  $N = 20$ ,  $m = 12$ ,  $\theta = 1/8$  and  $\varepsilon = 10^{-10}$ .



**Figure 12**  $b_+$ ,  $\Omega(\lambda)$ ,  $b_-$  with  $N = 20$ ,  $m = 12$ ,  $\theta = 1/8$  and  $\varepsilon = 10^{-15}$ .



**Table 6**  $N = 20$ ,  $m = 6$ ,  $\theta = 1/14$

$\lambda_k$	Sinc $\lambda_{k,N}$	Exact $\lambda_k$	Hermite $\lambda_{k,N}$	$E_S$	$E_H$
$\lambda_{-2}$	-3.7364320716761927	-3.736432198331617	-3.7364321983463715	$1.26655 \times 10^{-7}$	$1.47544 \times 10^{-11}$
$\lambda_{-1}$	-1.0801974353048152	-1.0801976714797825	-1.0801976714531203	$2.36175 \times 10^{-7}$	$2.66622 \times 10^{-11}$
$\lambda_0$	0.6565189567613093	0.6565187872152198	0.6565187872029187	$1.69546 \times 10^{-7}$	$1.23012 \times 10^{-11}$
$\lambda_1$	3.118561532798614	3.1185614501648167	3.1185614501681216	$4.98479 \times 10^{-8}$	$3.30491 \times 10^{-12}$

**Table 7**  $N = 20$ ,  $m = 12$ ,  $\theta = 1/8$

$\lambda_k$	Sinc $\lambda_{k,N}$	Exact $\lambda_k$	Hermite $\lambda_{k,N}$
$\lambda_{-2}$	-3.736432198332202082929465	-3.736432198331617091212013	-3.736432198331617091189782
$\lambda_{-1}$	-1.080197671476027921290673	-1.080197671479782493157863	-1.080197671479782493947136
$\lambda_0$	0.6565187872242083579354743	0.6565187872152199183983102	0.6565187872152199230592640
$\lambda_1$	3.118561450158043898832776	3.118561450164816849643922	3.118561450164816845810261

**Table 8** Absolute error  $|\lambda_k - \lambda_{k,N}|$  for  $N = 20$ ,  $m = 12$ ,  $\theta = 1/8$

$\lambda_k$	$\lambda_{-2}$	$\lambda_{-1}$	$\lambda_0$	$\lambda_1$
$E_S$	$5.849 \times 10^{-13}$	$3.755 \times 10^{-12}$	$8.988 \times 10^{-12}$	$6.773 \times 10^{-12}$
$E_H$	$2.223 \times 10^{-20}$	$7.893 \times 10^{-19}$	$4.661 \times 10^{-18}$	$3.834 \times 10^{-18}$

**Table 9** For  $N = 20$ ,  $m = 6$  and  $\theta = 1/14$ , the exact solutions  $\lambda_k$  are all inside the interval  $[b_-, b_+]$  for different values of  $\varepsilon$

$\lambda_k$	Exact $\lambda_k$	$[b_-, b_+], \varepsilon = 10^{-10}$	$[b_-, b_+], \varepsilon = 10^{-15}$
$\lambda_{-2}$	-3.736432198331617091212013	$[-3.881037, -3.476447]$	$[-3.836682, -3.557513]$
$\lambda_{-1}$	-1.080197671479782493157863	$[-1.435432, -0.665868]$	$[-1.365324, -0.760935]$
$\lambda_0$	0.6565187872152199183983102	$[0.410872, 1.116247]$	$[0.492155, 1.004381]$
$\lambda_1$	3.118561450164816849643922	$[2.884061, 3.390359]$	$[2.940901, 3.331955]$

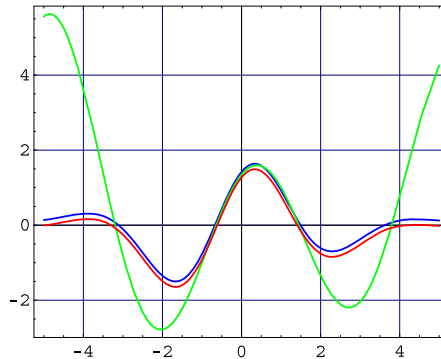
$$E_4(\mathcal{R}_{\theta,m}) = 2.9056 \times 10^7, E_3(\mathcal{R}_{\theta,m}) = 2.29859 \times 10^6, \omega = 1, M_{\mathcal{R}_{\theta,m}} = 98845.4.$$

**Table 10** With  $N = 20$ ,  $m = 12$  and  $\theta = 1/8$ ,  $\lambda_k$  are all inside the interval  $[b_-, b_+]$  for different values of  $\varepsilon$

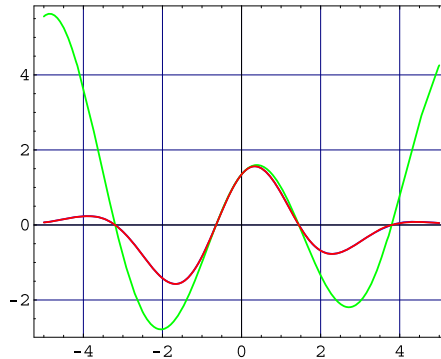
$\lambda_k$	Exact $\lambda_k$	$[b_-, b_+]$ , $\varepsilon = 10^{-10}$	$[b_-, b_+]$ , $\varepsilon = 10^{-15}$
$\lambda_{-2}$	-3.736432198331617	[-4.1011429, -3.3717065]	[-3.7364598, -3.7364045]
$\lambda_{-1}$	-1.0801976714797825	[-1.5078873, -0.4433678]	[-1.0808585, -1.07952734]
$\lambda_0$	0.6565187872152198	[0.0168549, 1.1086918]	[0.6528005, 0.6602210]
$\lambda_1$	3.1185614501648167	[2.7401391, 3.1185614]	[3.1157222, 3.1214041]

$$E_{10}(\mathcal{R}_{\theta,m}) = 6.2724 \times 10^{12}, E_9(\mathcal{R}_{\theta,m}) = 8.21004 \times 10^{11}, \omega = 1, M_{\mathcal{R}_{\theta,m}} = 501421.$$

**Figure 13**  $b_+$ ,  $\Omega(\lambda)$ ,  $b_-$  with  $N = 20$ ,  $m = 16$ ,  $\theta = 1/4$  and  $\varepsilon = 10^{-12}$ .



**Figure 14**  $b_+$ ,  $\Omega(\lambda)$ ,  $b_-$  with  $N = 20$ ,  $m = 16$ ,  $\theta = 1/4$  and  $\varepsilon = 10^{-15}$ .



**Table 11**  $N = 20$ ,  $m = 16$ ,  $\theta = 1/4$

$\lambda_k$	Exact $\lambda_k$	$\lambda_{k,N}$	$E_H$
$\lambda_{-2}$	-3.1976270593385675784857858037	-3.1976270593385675784857498452	$3.596 \times 10^{-23}$
$\lambda_{-1}$	-0.64351783872891518984316280760	-0.64351783872891518984316309998	$2.924 \times 10^{-25}$
$\lambda_0$	1.4487204290456776077365351429	1.4487204290456776077365176362	$1.751 \times 10^{-23}$
$\lambda_1$	3.8015200831700579923508826075	3.8015200831700579923509045951	$2.199 \times 10^{-23}$

**Table 12** With  $N = 20$ ,  $m = 16$  and  $\theta = 1/4$ ,  $\lambda_k$  are all inside the interval  $[b_-, b_+]$  for different values of  $\varepsilon$

$\lambda_k$	Exact $\lambda_k$	$[b_-, b_+]$ , $\varepsilon = 10^{-10}$	$[b_-, b_+]$ , $\varepsilon = 10^{-15}$
$\lambda_{-2}$	-3.1976270593385675784857858037	[-3.30255437, -3.11013060]	[-3.19791869, -3.19733846]
$\lambda_{-1}$	-0.64351783872891518984316280760	[-0.67219637, -0.61489406]	[-0.64356572, -0.64346999]
$\lambda_0$	1.4487204290456776077365351429	[1.40795687, 1.49107473]	[1.44812338, 1.44932224]
$\lambda_1$	3.8015200831700579923508826075	[3.60636554, 4.19453907]	[3.80103975, 3.80200804]

$$E_{14}(\mathcal{R}_{\theta,m}) = 2.16956 \times 10^{13}, E_{13}(\mathcal{R}_{\theta,m}) = 5.61116 \times 10^{12}, \omega = 1, M_{\mathcal{R}_{\theta,m}} = 3.15557 \times 10^6.$$

**Example 3** The boundary value problem

$$u_2'(x) - xu_1(x) = \lambda u_1(x), \quad u_1'(x) - u_2(x) = -\lambda u_2(x), \quad 0 \leq x \leq 1, \quad (4.13)$$

$$u_2(0) = 0, \quad u_1(1) = \lambda u_2(1) \quad (4.14)$$

is a special case of the problem  $\Gamma(r, \alpha, \beta, 0, \beta')$  when  $r_1(x) = x$ ,  $r_2(x) = -1$ ,  $\alpha_2 = \beta_1 = \beta_2' = 1$  and  $\alpha_1 = \beta_1' = \beta_2 = 0$ . Here the characteristic function is

$$\begin{aligned} \Omega(\lambda) := & 1/(\text{AiryAiPrime}[\lambda(1-\lambda)^{1/3}] \text{AiryBi}[\lambda(1-\lambda)^{1/3}] \\ & - \text{AiryAi}[\lambda(1-\lambda)^{1/3}] \text{AiryBiPrime}[\lambda(1-\lambda)^{1/3}]) \\ & \times [\lambda(1-\lambda)^{2/3} \text{AiryAi}[(1+\lambda)(1-\lambda)^{1/3}] \text{AiryBi}[\lambda(1-\lambda)^{1/3}] \\ & + \text{AiryAiPrime}[(\lambda+1)(1-\lambda)^{1/3}] \text{AiryBi}[\lambda(1-\lambda)^{1/3}] \\ & - \text{AiryAiPrime}[\lambda(1-\lambda)^{1/3}] (\lambda(1-\lambda)^{2/3} \text{AiryBi}[(\lambda+1)(1-\lambda)^{1/3}] \\ & + \text{AiryBiPrime}[(\lambda+1)(1-\lambda)^{1/3}])], \end{aligned} \quad (4.15)$$

where  $\text{AiryAi}[z]$  and  $\text{AiryBi}[z]$  are Airy functions  $\text{Ai}(z)$  and  $\text{Bi}(z)$ , respectively, and  $\text{AiryAiPrime}[z]$  and  $\text{AiryBiPrime}[z]$  are derivatives of Airy functions. The function  $\mathcal{K}(\lambda)$  will be

$$\mathcal{K}(\lambda) := \cos \lambda - \lambda \sin \lambda. \quad (4.16)$$

Figures 13, 14 and Tables 11, 12 illustrate the applications of the method to this problem.

**Competing interests**

The author declares that he has no competing interests.

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