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Global and blow-up solutions for nonlinear parabolic problems with a gradient term under Robin boundary conditions

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Abstract

In this paper, we study the global and blow-up solutions of the following nonlinear parabolic problems with a gradient term under Robin boundary conditions:

$$\begin{cases} (b(u))_t = \nabla \cdot (g(u)\nabla u) + f(x, u, |\nabla u|^2, t) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + \gamma u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases}$$

where $D \subset \mathbb{R}^N$ ($N \geq 2$) is a bounded domain with smooth boundary ∂D . By constructing auxiliary functions and using maximum principles, the sufficient conditions for the existence of a global solution, an upper estimate of the global solution, the sufficient conditions for the existence of a blow-up solution, an upper bound for 'blow-up time', and an upper estimate of 'blow-up rate' are specified under some appropriate assumptions on the functions f, g, b and initial value u_0 .

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Keywords: global solution; blow-up solution; parabolic problem; Robin boundary condition; gradient term

1 Introduction

In this paper, we study the global and blow-up solutions of the following nonlinear parabolic problems with a gradient term under Robin boundary conditions:

$$\begin{cases} (b(u))_t = \nabla \cdot (g(u)\nabla u) + f(x, u, q, t) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + \gamma u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = u_0(x) > 0 & \text{in } \bar{D}, \end{cases} \quad (1.1)$$

where $q := |\nabla u|^2$, $D \subset \mathbb{R}^N$ ($N \geq 2$) is a bounded domain with smooth boundary ∂D , $\partial/\partial n$ represents the outward normal derivative on ∂D , γ is a positive constant, u_0 is the initial value, T is the maximal existence time of u , and \bar{D} is the closure of D . Set $\mathbb{R}^+ := (0, +\infty)$. We assume, throughout the paper, that $b(s)$ is a $C^3(\mathbb{R}^+)$ function, $b'(s) > 0$ for any $s \in \mathbb{R}^+$, $g(s)$ is a positive $C^2(\mathbb{R}^+)$ function, $f(x, s, d, t)$ is a nonnegative $C^1(\bar{D} \times \mathbb{R}^+ \times \bar{\mathbb{R}}^+ \times \mathbb{R}^+)$ function, and $u_0(x)$ is a positive $C^2(\bar{D})$ function. Under the above assumptions, the classical theory [1] of parabolic equation assures that there exists a unique classical solution $u(x, t)$ with

some $T > 0$ for problem (1.1) and the solution is positive over $\overline{D} \times [0, T)$. Moreover, the regularity theorem [2] implies $u(x, t) \in C^3(D \times (0, T)) \cap C^2(\overline{D} \times [0, T))$.

Many papers have studied the global and blow-up solutions of parabolic problems with a gradient term (see, for instance, [3–13]). Some authors have discussed the global and blow-up solutions of parabolic problems under Robin boundary conditions and have got a lot of meaningful results (see [14–20] and the references cited therein). Some special cases of problem (1.1) have been treated already. Zhang [21] dealt with the following problem:

$$\begin{cases} u_t = \nabla \cdot (g(u)\nabla u) + f(u) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + \gamma u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = u_0(x) > 0 & \text{in } \overline{D}, \end{cases}$$

where $D \subset \mathbb{R}^N$ ($N \geq 2$) is a bounded domain with smooth boundary ∂D . By constructing auxiliary functions and using maximum principles, the sufficient conditions characterized by functions f , g and u_0 were given for the existence of a blow-up solution. Zhang [22] investigated the following problem:

$$\begin{cases} (b(u))_t = \Delta u + f(u) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + \gamma u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = u_0(x) > 0 & \text{in } \overline{D}, \end{cases}$$

where $D \subset \mathbb{R}^N$ ($N \geq 2$) is a bounded domain with smooth boundary ∂D . By constructing some auxiliary functions and using maximum principles, the sufficient conditions were obtained there for the existence of global and blow-up solutions. Meanwhile, the upper estimate of a global solution, the upper bound of ‘blow-up time’ and the upper estimate of ‘blow-up rate’ were also given. Ding [21] considered the following problem:

$$\begin{cases} (b(u))_t = \nabla \cdot (g(u)\nabla u) + f(u) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + \gamma u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = u_0(x) > 0 & \text{in } \overline{D}, \end{cases}$$

where $D \subset \mathbb{R}^N$ ($N \geq 2$) is a bounded domain with smooth boundary ∂D . By constructing some appropriate auxiliary functions and using a first-order differential inequality technique, the sufficient conditions were obtained for the existence of global and blow-up solutions. For the blow-up solution, an upper and a lower bound on blow-up time were also given.

In this paper, we study problem (1.1). Since the function $f(x, u, q, t)$ contains a gradient term $q = |\nabla u|^2$, it seems that the methods of [21–23] are not applicable for problem (1.1). In this paper, by constructing completely different auxiliary functions with those in [21–23] and technically using maximum principles, we obtain some existence theorems of a global solution, an upper estimate of the global solution, the existence theorems of a blow-up solution, an upper bound of ‘blow-up time’, and an upper estimates of ‘blow-up rate’. Our results extend and supplement those obtained [21–23].

We proceed as follows. In Section 2 we study the global solution of (1.1). Section 3 is devoted to the blow-up solution of (1.1). A few examples are presented in Section 4 to illustrate the applications of the abstract results.

2 Global solution

The main result for the global solution is the following theorem.

Theorem 2.1 *Let u be a solution of problem (1.1). Assume that the following conditions*

(i)-(iv) are satisfied:

(i) for any $s \in \mathbb{R}^+$,

$$\begin{aligned} (sb'(s))' &\geq 0, \quad sb'(s) - (sb'(s))' \leq 0, \quad \left(\frac{g(s)}{b'(s)}\right)' \leq 0, \\ \left[\frac{1}{g(s)}\left(\frac{g(s)}{b'(s)}\right)' + \frac{1}{b'(s)}\right]' + \frac{1}{g}\left(\frac{g(s)}{b'(s)}\right)' + \frac{1}{b'(s)} &\leq 0; \end{aligned} \quad (2.1)$$

(ii) for any $(x, s, d, t) \in D \times \mathbb{R}^+ \times \overline{\mathbb{R}^+} \times \mathbb{R}^+$,

$$\begin{aligned} f_t(x, s, d, t) &\leq 0, \quad f_d(x, s, d, t) \left[\left(\frac{1}{b'(s)}\right)' + \frac{1}{b'(s)}\right] \leq 0, \\ \left(\frac{f(x, s, d, t)b'(s)}{g(s)}\right)_s - \frac{f(x, s, d, t)b'(s)}{g(s)} &\leq 0; \end{aligned} \quad (2.2)$$

(iii)

$$\int_{m_0}^{+\infty} \frac{b'(s)}{e^s} ds = +\infty, \quad m_0 := \min_{\overline{D}} u_0(x); \quad (2.3)$$

(iv)

$$\alpha := \max_{\overline{D}} \frac{\nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0)}{e^{u_0}} > 0, \quad q_0 := |\nabla u_0|^2. \quad (2.4)$$

Then the solution u to problem (1.1) must be a global solution and

$$u(x, t) \leq H^{-1}(\alpha t + H(u_0(x, t))), \quad (x, t) \in \overline{D} \times \overline{\mathbb{R}^+}, \quad (2.5)$$

where

$$H(z) := \int_{m_0}^z \frac{b'(s)}{e^s} ds, \quad z \geq m_0, \quad (2.6)$$

and H^{-1} is the inverse function of H .

Proof Consider the auxiliary function

$$P(x, t) := b'(u)u_t - \alpha e^u. \quad (2.7)$$

Now we have

$$\nabla P = b''u_t \nabla u + b' \nabla u_t - \alpha e^u \nabla u, \quad (2.8)$$

$$\Delta P = b'''u_t |\nabla u|^2 + 2b'' \nabla u \cdot \nabla u_t + b''u_t \Delta u + b' \Delta u_t - \alpha e^u |\nabla u|^2 - \alpha e^u \Delta u, \quad (2.9)$$

and

$$\begin{aligned}
 P_t &= b''(u_t)^2 + b'(u_t)_t - \alpha e^u u_t \\
 &= b''(u_t)^2 + b' \left(\frac{g}{b'} \Delta u + \frac{g'}{b'} |\nabla u|^2 + \frac{f}{b'} \right)_t - \alpha e^u u_t \\
 &= b''(u_t)^2 + \left(g' - \frac{b''g}{b'} \right) u_t \Delta u + g \Delta u_t + \left(g'' - \frac{b''g'}{b'} \right) u_t |\nabla u|^2 \\
 &\quad + (2g' + 2f_q) \nabla u \cdot \nabla u_t + \left(f_u - \frac{b''f}{b'} - \alpha e^u \right) u_t + f_t.
 \end{aligned} \tag{2.10}$$

It follows from (2.9) and (2.10) that

$$\begin{aligned}
 \frac{g}{b'} \Delta P - P_t &= \left(\frac{b'''g}{b'} + \frac{b''g'}{b'} - g'' \right) u_t |\nabla u|^2 + \left(2 \frac{b''g}{b'} - 2g' - 2f_q \right) \nabla u \cdot \nabla u_t \\
 &\quad + \left(2 \frac{b''g}{b'} - g' \right) u_t \Delta u - \alpha \frac{g}{b'} e^u |\nabla u|^2 - \alpha \frac{g}{b'} e^u \Delta u - b''(u_t)^2 \\
 &\quad + \left(\frac{b''f}{b'} - f_u + \alpha e^u \right) u_t - f_t.
 \end{aligned} \tag{2.11}$$

By (1.1), we have

$$\Delta u = \frac{b'}{g} u_t - \frac{g'}{g} |\nabla u|^2 - \frac{f}{g}. \tag{2.12}$$

Substitute (2.12) into (2.11), to get

$$\begin{aligned}
 \frac{g}{b'} \Delta P - P_t &= \left(\frac{b'''g}{b'} - \frac{b''g'}{b'} - g'' + \frac{(g')^2}{g} \right) u_t |\nabla u|^2 + \left(2 \frac{b''g}{b'} - 2g' - 2f_q \right) \nabla u \cdot \nabla u_t \\
 &\quad - \frac{(b')^2}{g} \left(\frac{g}{b'} \right)' (u_t)^2 + \left(\frac{fg'}{g} - \frac{b''f}{b'} - f_u \right) u_t + \left(\alpha \frac{g'}{b'} e^u - \alpha \frac{g}{b'} e^u \right) |\nabla u|^2 \\
 &\quad + \alpha \frac{f}{b'} e^u - f_t.
 \end{aligned} \tag{2.13}$$

With (2.8), we have

$$\nabla u_t = \frac{1}{b'} \nabla P - \frac{b''}{b'} u_t \nabla u + \alpha \frac{e^u}{b'} \nabla u. \tag{2.14}$$

Next, we substitute (2.14) into (2.13) to obtain

$$\begin{aligned}
 \frac{g}{b'} \Delta P + \left[2 \left(\frac{g}{b'} \right)' + 2 \frac{f_q}{b'} \right] \nabla u \cdot \nabla P - P_t &= \left(\frac{b'''g}{b'} + \frac{b''g'}{b'} - g'' + \frac{(g')^2}{g} - 2 \frac{(b'')^2 g}{(b')^2} + 2 \frac{b''f_q}{b'} \right) u_t |\nabla u|^2 \\
 &\quad + \left(2\alpha \frac{b''g}{(b')^2} e^u - \alpha \frac{g'}{b'} e^u - \alpha \frac{g}{b'} e^u - 2\alpha \frac{f_q}{b'} e^u \right) |\nabla u|^2 \\
 &\quad - \frac{(b')^2}{g} \left(\frac{g}{b'} \right)' (u_t)^2 + \left(\frac{fg'}{g} - \frac{b''f}{b'} - f_u \right) u_t + \alpha \frac{f}{b'} e^u - f_t.
 \end{aligned} \tag{2.15}$$

In view of (2.7), we have

$$u_t = \frac{1}{b'}P + \alpha \frac{e^u}{b'}. \quad (2.16)$$

Substituting (2.16) into (2.15), we get

$$\begin{aligned} & \frac{g}{b'} \Delta P + \left[2 \left(\frac{g}{b'} \right)' + 2 \frac{f_q}{b'} \right] \nabla u \cdot \nabla P \\ & + \left\{ \left[g \left(\frac{1}{g} \left(\frac{g}{b'} \right)' \right)' + 2 f_q \left(\frac{1}{b'} \right)' \right] |\nabla u|^2 + \frac{g}{(b')^2} \left(\frac{fb'}{g} \right)_u \right\} P - P_t \\ & = -\alpha e^u \left\{ g \left[\left(\frac{1}{g} \left(\frac{g}{b'} \right)' + \frac{1}{b'} \right)' + \frac{1}{g} \left(\frac{g}{b'} \right)' + \frac{1}{b'} \right] + 2 f_q \left[\left(\frac{1}{b'} \right)' + \frac{1}{b'} \right] \right\} |\nabla u|^2 \\ & - \frac{(b')^2}{g} \left(\frac{g}{b'} \right)' (u_t)^2 - \alpha \frac{g e^u}{(b')^2} \left[\left(\frac{fb'}{g} \right)_u - \frac{fb'}{g} \right] - f_t. \end{aligned} \quad (2.17)$$

The assumptions (2.1) and (2.2) guarantee that the right-hand side of (2.17) is nonnegative, *i.e.*,

$$\begin{aligned} & \frac{g}{b'} \Delta P + \left[2 \left(\frac{g}{b'} \right)' + 2 \frac{f_q}{b'} \right] \nabla u \cdot \nabla P \\ & + \left\{ \left[g \left(\frac{1}{g} \left(\frac{g}{b'} \right)' \right)' + 2 f_q \left(\frac{1}{b'} \right)' \right] |\nabla u|^2 + \frac{g}{(b')^2} \left(\frac{fb'}{g} \right)_u \right\} P - P_t \\ & \geq 0 \quad \text{in } D \times (0, T). \end{aligned} \quad (2.18)$$

By applying the maximum principle [24], it follows from (2.18) that P can attain its non-negative maximum only for $\overline{D} \times \{0\}$ or $\partial D \times (0, T)$. For $\overline{D} \times \{0\}$, by (2.4), we have

$$\begin{aligned} \max_{\overline{D}} P(x, 0) &= \max_{\overline{D}} \{ b'(u_0)(u_0)_t - \alpha e^{u_0} \} = \max_{\overline{D}} \{ \nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0) - \alpha e^{u_0} \} \\ &= \max_{\overline{D}} \left\{ e^{u_0} \left[\frac{\nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0)}{e^{u_0}} - \alpha \right] \right\} = 0. \end{aligned}$$

We claim that P cannot take a positive maximum at any point $(x, t) \in \partial D \times (0, T)$. In fact, suppose that P takes a positive maximum at a point $(x_0, t_0) \in \partial D \times (0, T)$, then

$$P(x_0, t_0) > 0 \quad \text{and} \quad \frac{\partial P}{\partial n} \Big|_{(x_0, t_0)} > 0. \quad (2.19)$$

With (1.1) and (2.16), we have

$$\begin{aligned} \frac{\partial P}{\partial n} &= b'' u_t \frac{\partial u}{\partial n} + b' \frac{\partial u_t}{\partial n} - \alpha e^u \frac{\partial u}{\partial n} = -\gamma b'' u u_t + b' \left(\frac{\partial u}{\partial n} \right)_t + \gamma \alpha e^u \\ &= -\gamma b'' u u_t + b' (-\gamma u)_t + \gamma \alpha e^u = -\gamma (ub')' u_t + \gamma \alpha e^u \\ &= -\gamma (ub')' \left(\frac{1}{b'} P + \alpha \frac{1}{b'} e^u \right) + \gamma \alpha e^u \\ &= -\gamma \frac{(ub')'}{b'} P + \gamma \alpha e^u \frac{ub' - (ub')'}{b'} \quad \text{on } \partial D \times (0, T). \end{aligned} \quad (2.20)$$

Next, by using the fact that $(sb'(s))' \geq 0$, $sb'(s) - (sb'(s))' \leq 0$ for any $s \in \mathbb{R}^+$, it follows from (2.20) that

$$\frac{\partial P}{\partial n} \Big|_{(x_0, t_0)} \leq 0,$$

which contradicts with inequality (2.19). Thus we know that the maximum of P in $\bar{D} \times [0, T)$ is zero, *i.e.*,

$$P \leq 0 \quad \text{in } \bar{D} \times [0, T),$$

and

$$\frac{b'(u)}{e^u} u_t \leq \alpha. \quad (2.21)$$

For each fixed $x \in \bar{D}$, integration of (2.21) from 0 to t yields

$$\int_0^t \frac{b'(u)}{e^u} u_t \, dt = \int_{u_0(x)}^{u(x,t)} \frac{b'(s)}{e^s} \, ds \leq \alpha t, \quad (2.22)$$

which implies that u must be a global solution. Actually, if that u blows up at finite time T , then

$$\lim_{t \rightarrow T^-} u(x, t) = +\infty.$$

Passing to the limit as $t \rightarrow T^-$ in (2.22) yields

$$\int_{u_0(x)}^{+\infty} \frac{b'(s)}{e^s} \, ds \leq \alpha T$$

and

$$\int_{m_0}^{+\infty} \frac{b'(s)}{e^s} \, ds = \int_{m_0}^{u_0(x)} \frac{b'(s)}{e^s} \, ds + \int_{u_0(x)}^{+\infty} \frac{b'(s)}{e^s} \, ds \leq \int_{m_0}^{u_0(x)} \frac{b'(s)}{e^s} \, ds + \alpha T < +\infty,$$

which contradicts with assumption (2.3). This shows that u is global. Moreover, it follows from (2.22) that

$$\int_{u_0(x)}^{u(x,t)} \frac{b'(s)}{e^s} \, ds = \int_{m_0}^{u(x,t)} \frac{b'(s)}{e^s} \, ds - \int_{m_0}^{u_0(x)} \frac{b'(s)}{e^s} \, ds = H(u(x, t)) - H(u_0(x)) \leq \alpha t.$$

Since H is an increasing function, we have

$$u(x, t) \leq H^{-1}(\alpha t + H(u_0(x))).$$

The proof is complete. \square

3 Blow-up solution

The following theorem is the main result for the blow-up solution.

Theorem 3.1 *Let u be a solution of problem (1.1). Assume that the following conditions*

(i)-(iv) are fulfilled:

(i) *for any $s \in \mathbb{R}^+$,*

$$\begin{aligned} (sb'(s))' &\geq 0, \quad sb'(s) - (sb'(s))' \geq 0, \quad \left(\frac{g(s)}{b'(s)}\right)' \geq 0, \\ \left[\frac{1}{g(s)}\left(\frac{g(s)}{b'(s)}\right)' + \frac{1}{b'(s)}\right]' + \frac{1}{g}\left(\frac{g(s)}{b'(s)}\right)' + \frac{1}{b'(s)} &\geq 0; \end{aligned} \quad (3.1)$$

(ii) *for any $(x, s, d, t) \in D \times \mathbb{R}^+ \times \overline{\mathbb{R}^+} \times \mathbb{R}^+$,*

$$\begin{aligned} f_t(x, s, d, t) &\geq 0, \quad f_d(x, s, d, t) \left[\left(\frac{1}{b'(s)}\right)' + \frac{1}{b'(s)}\right] \geq 0, \\ \left(\frac{f(x, s, d, t)b'(s)}{g(s)}\right)_s - \frac{f(x, s, d, t)b'(s)}{g(s)} &\geq 0; \end{aligned} \quad (3.2)$$

(iii)

$$\int_{M_0}^{+\infty} \frac{b'(s)}{e^s} ds < +\infty, \quad M_0 := \max_{\overline{D}} u_0(x); \quad (3.3)$$

(iv)

$$\beta := \min_{\overline{D}} \frac{\nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0)}{e^{u_0}} > 0, \quad q_0 := |\nabla u_0|^2. \quad (3.4)$$

Then the solution u of problem (1.1) must blow up in finite time T , and

$$T \leq \frac{1}{\beta} \int_{M_0}^{+\infty} \frac{b'(s)}{e^s} ds, \quad (3.5)$$

$$u(x, t) \leq G^{-1}(\beta(T - t)), \quad (x, t) \in \overline{D} \times [0, T), \quad (3.6)$$

where

$$G(z) := \int_z^{+\infty} \frac{b'(s)}{e^s} ds, \quad z > 0, \quad (3.7)$$

and G^{-1} is the inverse function of G .

Proof Construct the following auxiliary function:

$$Q(x, t) := b'(u)u_t - \beta e^u. \quad (3.8)$$

Replacing P and α with Q and β in (2.17), respectively, we get

$$\begin{aligned} & \frac{g}{b'} \Delta Q + \left[2 \left(\frac{g}{b'} \right)' + 2 \frac{f_q}{b'} \right] \nabla u \cdot \nabla Q \\ & + \left\{ \left[g \left(\frac{1}{g} \left(\frac{g}{b'} \right)' \right)' + 2 f_q \left(\frac{1}{b'} \right)' \right] |\nabla u|^2 + \frac{g}{(b')^2} \left(\frac{fb'}{g} \right)_u \right\} Q - Q_t \\ & = -\beta e^u \left\{ g \left[\left(\frac{1}{g} \left(\frac{g}{b'} \right)' + \frac{1}{b'} \right)' + \frac{1}{g} \left(\frac{g}{b'} \right)' + \frac{1}{b'} \right] + 2 f_q \left[\left(\frac{1}{b'} \right)' + \frac{1}{b'} \right] \right\} |\nabla u|^2 \\ & - \frac{(b')^2}{g} \left(\frac{g}{b'} \right)' (u_t)^2 - \beta \frac{g e^u}{(b')^2} \left[\left(\frac{fb'}{g} \right)_u - \frac{fb'}{g} \right] - f_t. \end{aligned} \quad (3.9)$$

Assumptions (3.1) and (3.2) imply that the right-hand side in equality (3.9) is nonpositive, *i.e.*,

$$\begin{aligned} & \frac{g}{b'} \Delta Q + \left[2 \left(\frac{g}{b'} \right)' + 2 \frac{f_q}{b'} \right] \nabla u \cdot \nabla Q \\ & + \left\{ \left[g \left(\frac{1}{g} \left(\frac{g}{b'} \right)' \right)' + 2 f_q \left(\frac{1}{b'} \right)' \right] |\nabla u|^2 + \frac{g}{(b')^2} \left(\frac{fb'}{g} \right)_u \right\} Q - Q_t \\ & \leq 0 \quad \text{in } D \times (0, T). \end{aligned} \quad (3.10)$$

With (3.4), we have

$$\begin{aligned} \min_{\bar{D}} Q(x, 0) &= \min_{\bar{D}} \{ b'(u_0)(u_0)_t - \beta e^{u_0} \} = \min_{\bar{D}} \{ \nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0) - \beta e^{u_0} \} \\ &= \min_{\bar{D}} \left\{ e^{u_0} \left[\frac{\nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0)}{e^{u_0}} - \beta \right] \right\} = 0. \end{aligned} \quad (3.11)$$

Substituting P and α with Q and β in (2.20), respectively, we have

$$\frac{\partial Q}{\partial n} = -\gamma \frac{(ub')'}{b'} Q + \gamma \beta e^u \frac{ub' - (ub')'}{b'} \quad \text{on } \partial D \times (0, T). \quad (3.12)$$

Combining (3.10)-(3.12) with the fact that $(sb'(s))' \geq 0$, $sb'(s) - (sb'(s))' \geq 0$ for any $s \in \mathbb{R}^+$, and applying the maximum principles again, it follows that the minimum of Q in $\bar{D} \times [0, T)$ is zero. Thus

$$Q \geq 0 \quad \text{in } \bar{D} \times [0, T),$$

and

$$\frac{b'(u)}{e^u} u_t \geq \beta. \quad (3.13)$$

At the point $x^* \in \bar{D}$, where $u_0(x^*) = M_0$, integrate (3.13) over $[0, t]$ to get

$$\int_0^t \frac{b'(u)}{e^u} u_t \, dt = \int_{M_0}^{u(x^*, t)} \frac{b'(s)}{e^s} \, ds \geq \beta t, \quad (3.14)$$

which implies that u must blow up in finite time. Actually, if u is a global solution of (1.1), then for any $t > 0$, (3.14) shows

$$\int_{M_0}^{+\infty} \frac{b'(s)}{e^s} ds \geq \int_{M_0}^{u(x^*,t)} \frac{b'(s)}{e^s} ds \geq \beta t. \quad (3.15)$$

Letting $t \rightarrow +\infty$ in (3.15), we have

$$\int_{M_0}^{+\infty} \frac{b'(s)}{e^s} ds = +\infty,$$

which contradicts with assumption (3.3). This shows that u must blow up in finite time $t = T$. Furthermore, letting $t \rightarrow T$ in (3.14), we get

$$T \leq \frac{1}{\beta} \int_{M_0}^{+\infty} \frac{b'(s)}{e^s} ds.$$

By integrating inequality (3.13) over $[t, s]$ ($0 < t < s < T$), for each fixed x , we obtain

$$\begin{aligned} G(u(x, t)) &\geq G(u(x, t)) - G(u(x, s)) = \int_{u(x,t)}^{+\infty} \frac{b'(s)}{e^s} ds - \int_{u(x,s)}^{+\infty} \frac{b'(s)}{e^s} ds \\ &= \int_{u(x,t)}^{u(x,s)} \frac{b'(s)}{e^s} ds = \int_t^s \frac{b'(u)}{e^u} u_t dt \geq \beta(s - t). \end{aligned}$$

Hence, by letting $s \rightarrow T$, we have

$$G(u(x, t)) \geq \beta(T - t).$$

Since G is a decreasing function, we obtain

$$u(x, t) \leq G^{-1}(\beta(T - t)).$$

The proof is complete. \square

4 Applications

When $b(u) \equiv u$ and $f(x, u, q, t) \equiv f(u)$, the results stated in Theorem 3.1 are valid. When $g(u) \equiv 1$ and $f(x, u, q, t) \equiv f(u)$ or $f(x, u, q, t) \equiv f(u)$, the conclusions of Theorems 2.1 and 3.1 still hold true. In this sense, our results extend and supplement the results of [21–23].

In what follows, we present several examples to demonstrate the applications of the abstract results.

Example 4.1 Let u be a solution of the following problem:

$$\begin{cases} u_t = \Delta u + \frac{2+u}{1+u} |\nabla u|^2 + \frac{e^{-u}(e^{-u}+e^q)}{1+u} (e^{-t} + |x|^2) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + 2u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = 2 - |x|^2 & \text{in } \bar{D}, \end{cases}$$

where $q = |\nabla u|^2$, $D = \{x = (x_1, x_2, x_3) \mid |x|^2 < 1\}$ is the unit ball of \mathbb{R}^3 . The above problem can be transformed into the following problem:

$$\begin{cases} (ue^u)_t = \nabla \cdot ((1+u)e^u \nabla u) + (e^{-u} + e^q)(e^{-t} + |x|^2) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + 2u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = 2 - |x|^2 & \text{in } \bar{D}. \end{cases}$$

Now

$$\begin{aligned} b(u) &= ue^u, & g(u) &= (1+u)e^u, & f(x, u, q, t) &= (e^{-u} + e^q)(e^{-t} + |x|^2), \\ u_0(x) &= 2 - |x|^2, & \gamma &= 2. \end{aligned}$$

In order to determine the constant α , we assume

$$s := |x|^2,$$

then $0 \leq s \leq 1$ and

$$\begin{aligned} \alpha &= \max_{\bar{D}} \frac{\nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0)}{e^{u_0}} \\ &= \max_{\bar{D}} \{32|x|^2 - 4|x|^4 - 18 + (1 + |x|^2)[\exp(-4 + 2|x|^2) + \exp(-2 + 5|x|^2)]\} \\ &= \max_{0 \leq s \leq 1} \{32s - 4s^2 - 18 + (1 + s)[\exp(-4 + 2s) + \exp(-2 + 5s)]\} \\ &= 50.4417. \end{aligned}$$

It is easy to check that (2.1)-(2.3) hold. By Theorem 2.1, u must be a global solution, and

$$\begin{aligned} u(x, t) &\leq H^{-1}(\alpha t + H(u_0(x))) = -1 + \sqrt{50.4417t + (1 + u_0(x))^2} \\ &= -1 + \sqrt{50.4417t + (3 - |x|^2)^2}. \end{aligned}$$

Example 4.2 Let u be a solution of the following problem:

$$\begin{cases} u_t = \Delta u - \frac{1}{u(1+u)} |\nabla u|^2 + \frac{u(e^u - e^{-q})}{1+u} (6 + t|x|^2) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + 2u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = 2 - |x|^2 & \text{in } \bar{D}, \end{cases}$$

where $q = |\nabla u|^2$, $D = \{x = (x_1, x_2, x_3) \mid |x|^2 < 1\}$ is the unit ball of \mathbb{R}^3 . The above problem may be turned into the following problem:

$$\begin{cases} (u + \ln u)_t = \nabla \cdot ((1 + \frac{1}{u}) \nabla u) + (e^u - e^{-q})(6 + t|x|^2) & \text{in } D \times (0, T), \\ \frac{\partial u}{\partial n} + 2u = 0 & \text{on } \partial D \times (0, T), \\ u(x, 0) = 2 - |x|^2 & \text{in } \bar{D}. \end{cases}$$

Now we have

$$\begin{aligned} b(u) &= u + \ln u, & g(u) &= 1 + \frac{1}{u}, & f(x, u, q, t) &= (e^u - e^{-q})(6 + t|x|^2), \\ u_0(x) &= 2 - |x|^2, & \gamma &= 2. \end{aligned}$$

By setting

$$s := |x|^2,$$

we have $0 \leq s \leq 1$ and

$$\begin{aligned} \beta &= \min_{\bar{D}} \frac{\nabla \cdot (g(u_0) \nabla u_0) + f(x, u_0, q_0, 0)}{e^{u_0}} \\ &= \min_{\bar{D}} \left\{ \frac{-6|x|^4 + 26|x|^2 - 36}{(2 - |x|^2)^2 \exp(2 - |x|^2)} + 6[1 - \exp(-3|x|^2 - 2)] \right\} \\ &= \min_{0 \leq s \leq 1} \left\{ \frac{-6s^2 + 26s - 36}{(2 - s)^2 \exp(2 - s)} + 6[1 - \exp(-3s - 2)] \right\} \\ &= 0.0735. \end{aligned}$$

Again it is easy to check that (3.1)-(3.3) hold. By Theorem 3.1, u must blow up in finite time T , and

$$\begin{aligned} T &\leq \frac{1}{\beta} \int_{M_0}^{+\infty} \frac{b'(s)}{e^s} ds = \frac{1}{0.0735} \int_2^{+\infty} \left(1 + \frac{1}{s}\right) \frac{1}{e^s} ds = 2.5066, \\ u(x, t) &\leq G^{-1}(\beta(T - t)) = G^{-1}(0.0735(T - t)), \end{aligned}$$

where

$$G(z) = \int_z^{+\infty} \frac{b'(s)}{e^s} ds = \int_z^{+\infty} \left(1 + \frac{1}{s}\right) \frac{1}{e^s} ds, \quad z \geq 0,$$

and G^{-1} is the inverse function of G .

Remark 4.1 We can see from Example 4.1 that when the equation has a gradient term with exponential increase, the functions g and b increase exponentially to ensure that the solution of (1.1) blows up. It follows from Example 4.2 that when the equation has a gradient term with exponential decay, the appropriate assumptions on the functions g and b can guarantee the solution of (1.1) to be global.

Competing interests

The author declares that he has no competing interests.

Author's contributions

All results belong to Juntang Ding.

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References

1. Amann, H: Quasilinear parabolic systems under nonlinear boundary conditions. *Arch. Ration. Mech. Anal.* **92**, 153-192 (1986)
2. Sperb, RP: *Maximum Principles and Their Applications*. Academic Press, New York (1981)
3. Ding, JT, Guo, BZ: Global existence and blow-up solutions for quasilinear reaction-diffusion equations with a gradient term. *Appl. Math. Lett.* **24**, 936-942 (2011)
4. Tersenov, A: The preventive effect of the convection and of the diffusion in the blow-up phenomenon for parabolic equations. *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* **21**, 533-541 (2004)
5. Zheng, SN, Wang, W: Effects of reactive gradient term in a multi-nonlinear parabolic problem. *J. Differ. Equ.* **247**, 1980-1992 (2009)
6. Payne, LE, Song, JC: Lower bounds for blow-up time in a nonlinear parabolic problem. *J. Math. Anal. Appl.* **354**, 394-396 (2009)
7. Chen, SH: Global existence and blowup of solutions for a parabolic equation with a gradient term. *Proc. Am. Math. Soc.* **129**, 975-981 (2001)
8. Chen, SH: Global existence and blowup for quasilinear parabolic equations not in divergence form. *J. Math. Anal. Appl.* **401**, 298-306 (2013)
9. Chipot, M, Weissler, FB: Some blowup results for a nonlinear parabolic equation with a gradient term. *SIAM J. Math. Anal.* **20**, 886-907 (1989)
10. Fila, M: Remarks on blow up for a nonlinear parabolic equation with a gradient term. *Proc. Am. Math. Soc.* **111**, 795-801 (1991)
11. Souplet, P, Weissler, FB: Poincaré's inequality and global solutions of a nonlinear parabolic equation. *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* **16**, 335-371 (1999)
12. Souplet, P: Recent results and open problems on parabolic equations with gradient nonlinearities. *Electron. J. Differ. Equ.* **2001**, 1-19 (2001)
13. Souplet, P: Finite time blow-up for a non-linear parabolic equation with a gradient term and applications. *Math. Methods Appl. Sci.* **19**, 1317-1333 (1996)
14. Friedman, A, McLeod, B: Blow-up of positive solutions of semilinear heat equations. *Indiana Univ. Math. J.* **34**, 425-447 (1985)
15. Enache, C: Blow-up phenomena for a class of quasilinear parabolic problems under Robin boundary condition. *Appl. Math. Lett.* **24**, 288-292 (2011)
16. Payne, LE, Schaefer, PW: Blow-up in parabolic problems under Robin boundary conditions. *Appl. Anal.* **87**, 699-707 (2008)
17. Rault, JF: The Fujita phenomenon in exterior domains under the Robin boundary conditions. *C. R. Math. Acad. Sci. Paris* **349**, 1059-1061 (2011)
18. Li, YF, Liu, Y, Xiao, SZ: Blow-up phenomena for some nonlinear parabolic problems under Robin boundary conditions. *Math. Comput. Model.* **54**, 3065-3069 (2011)
19. Liu, Y, Luo, SG, Ye, YH: Blow-up phenomena for a parabolic problem with a gradient nonlinearity under nonlinear boundary conditions. *Comput. Math. Appl.* **65**, 1194-1199 (2013)
20. Li, YF, Liu, Y, Lin, CH: Blow-up phenomena for some nonlinear parabolic problems under mixed boundary conditions. *Nonlinear Anal., Real World Appl.* **11**, 3815-3823 (2010)
21. Zhang, LL: Blow-up of solutions for a class of nonlinear parabolic equations. *Z. Anal. Anwend.* **25**, 479-486 (2006)
22. Zhang, HL: Blow-up solutions and global solutions for nonlinear parabolic problems. *Nonlinear Anal. TMA* **69**, 4567-4574 (2008)
23. Ding, JT: Global and blow-up solutions for nonlinear parabolic equations with Robin boundary conditions. *Comput. Math. Appl.* **65**, 1808-1822 (2013)
24. Protter, MH, Weinberger, HF: *Maximum Principles in Differential Equations*. Prentice-Hall, Englewood Cliffs (1967)

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